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MACHINE CALCULATIONS OF ENERGY TRANSFER
PHENOMENA IN A BOMBARDED LATTICE
WARREN L. GAY

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MACHINE CALCULATIONS OF ENERGY TRANSFER
PHENOMENA IN A BOMBARDED LATTICE

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Warren L. Gay

MACHINE CALCULATIONS OF ENERGY TRANSFER
PHENOMENA IN A BOMBARDED LATTICE

by

Warren L. Gay

Lieutenant, United States Navy

Submitted in partial fulfillment of
the requirements for the degree of

MASTER OF SCIENCE
IN
PHYSICS

United States Naval Postgraduate School
Monterey, California

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MACHINE CALCULATIONS OF ENERGY TRANSFER
PHENOMENA IN A BOMBARDED LATTICE

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This work is accepted as fulfilling
the thesis requirements for the degree of
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IN

PHYSICS

from the

United States Naval Postgraduate School

ABSTRACT

This computer simulation investigates the ways in which an embedding lattice influences a collision event. Low energy events were studied with a Born-Mayer potential function; high energy with a Thomas-Fermi-Firsov potential. Conclusions were: (1) the lattice will increase ranges for particle energies above 400 ev; (2) the apparent mass concept is not a valid description of events in a lattice; (3) lattice effects will significantly modify the low energy portion of the target atom energy distribution function; (4) there is no evidence that a copper atom will "rebound" from a copper lattice.

The writer wishes to express his appreciation for the assistance and encouragement given him by Associate Professor Don E. Harrison, Jr. of the U. S. Naval Postgraduate School in this investigation.

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1. History.

Radiation damage is the overall term applied to erosion, disruption, or rearrangement of a crystal lattice by bombardment of ions or atoms. Sputtering, a phenomena sometimes observed in radiation damage, applies to the process of lattice atoms being ejected or "knocked out" of the crystal. The sputtering process was first observed in 1852 by Grove¹ in the cathodes of electrical discharge tubes.

Since Grove's observation, experimental and theoretical physicists have been attempting to develop a sputtering theory that will predict measurable experimental quantities as a function of mass ratio, position in the periodic table, energy of bombarding particle and angle of incidence. To date, no such complete theory is available.

A tremendous volume of material has been written on the subject of sputtering since its discovery. In 1955, Irvine² listed 306 papers concerned with the sputtering process. These early papers disagree violently because important experimental conditions were not obtained or simply ignored; consequently many of the conclusions reached are contradictory.

During this period two important theories were proposed. Kingdon and Langmuir³, in 1923, developed a momentum transfer theory which agreed quite well with experimental data they had obtained for thoriated tungsten filaments. The "evaporation theory" was proposed in 1928 by von Hippel and Blechschmidt⁴ and improved by Townes⁵ in 1944. It assumes that the incident

¹W. R. Grove, Phil. Trans. Roy. Soc. London, 142, 87 (1842).

²M. M. Irvine, Dissertation, Lehigh University (1955).

³K. H. Kingdon and I. Langmuir, Phys. Rev. 22, 148 (1923).

⁴A. von Hippel and E. Blechschmidt, Ann. Physik 86, 1006 (1928).

⁵C. H. Townes, Phys. Rev. 65, 319 (1944).

particles heat a small volume of the crystal to a very high temperature and subsequently surface atoms evaporate. More recently (1953), Harrison⁶, following a suggestion made by Keywell⁷, presented a statistical theory of sputtering which could be placed somewhere between the two models mentioned above. We will not present details of these theories. We must mention however, that they do not provide an adequate description of the sputtering process.

The trend in sputtering theories has recently been toward the consideration of the individual collision processes within the lattice. Henschke⁸ has used this approach in his theory of sputtering and Harrison and Magnuson⁹ have also applied it, with modifications, to the theoretical study of sputtering threshold energies.

The theoretical work performed by Henschke uses the momentum transfer concept initiated by Kingdon and Langmuir³ to account for all the experimental phenomena observed at low incident particle energy. His basic treatment of collisions between atoms in a lattice is based on the assumption

...that the collisions described can be treated with the general principles of classical mechanics, using impulsive forces, in a manner similar to the well-known collisions with restitution.¹⁰

However, he did not in all cases use the masses of the individual particles involved in the equations pertaining to the two body collisions. Instead he postulated an "effective" mass. If the collision is between a surface

⁶D. E. Harrison, Jr., Phys. Rev. 102, 1473 (1956).

⁷F. Keywell, Phys. Rev. 87, 160 (1952).

⁸E. B. Henschke, Phys. Rev. 106, 737 (1957).

⁹D. E. Harrison and G. D. Magnuson, Phys. Rev. 122, 1421 (1961)

¹⁰See reference 8, p. 738.

atom and a moving ion or atom and is such that the surface atom is struck on its "inside" hemisphere, according to Henschke, an effective mass is not required and the two particles involved can be considered to have their actual masses. Should the collision of an atom or ion be directed inward from the target lattice surface, Henschke states,

The bulk of the target is behind the struck atom and produces a very large 'effective' mass compared to the mass of the ion or to the mass of the target atom.¹¹

The reason Henschke did not assume pure elastic collisions was stated as

Energy losses are due to the fact that the target atom is coupled rather strongly with the atoms of the lattice. Before the moment of highest compression is reached, the ion and the struck atom exchange energy with the neighboring atoms of the lattice. Debye waves are thus excited and dissipated irreversibly into the lattice.¹²

He goes on to state,

...the final step in each sputtering process at any angle of incidence of the ion can be generally described as a collision of the ion with an upper surface atom, in which this atom is hit on its inside hemisphere so as to obtain an impulse with a component in the direction of the outward normal to the surface. If the energy transferred in this direction to this target atom by the impact of the ion is equal to or greater than the heat of vaporization, with which the atom is assumed to be bound to the crystal lattice plane, then this atom is ejected in the collision.¹³

Early study in radiation damage was precipitated by the advent of the nuclear reactor. Recently however, with the imminent possibility of thermonuclear power production and ion propulsion engines for space vehicles, other practical applications of the processes involved in radiation damage have

¹¹E. B. Henschke, Phys. Rev. 121, 1290 (1961).

¹²See reference 8, p. 738.

¹³See reference 8, p. 737.

become extremely important.

Robinson and his co-workers¹⁴ have made theoretical studies of the ranges in solids of atoms having energies from 1 to 100 kev using digital computer techniques on the basic assumption that the moving atom loses its energy through repeated binary elastic collisions with atoms of the solid. The masses used in these calculations are the true masses of the interacting particles, but they note that their assumption is certainly not valid below 100 ev. Gibson et. al.¹⁵ have also taken advantage of the speed available with modern digital computers in the study of radiation damage. They do not assume binary collisions but instead use iteration techniques employing Newton's equation of motion to solve the complex many body problem. The initial success of these programs have encouraged their originators to explore further the possibilities of these techniques. Although the computer programs are designed to study radiation damage, the basic principles involved are also important in sputtering.

The work just discussed is very sensitive to the mathematical form of the interatomic potential. The Born-Mayer, Bohr (screened Coulomb), and the Thomas-Fermi-Firsov are some of the potentials used in these calculations. The ability of any of the potential functions to describe a physical situation is highly dependent upon the energy range under consideration and the specific atoms or ions involved. No single potential has yet been devised that satisfactorily represents the interaction under all circumstances.

¹⁴O. S. Oen, D. K. Holmes, and M. T. Robinson, Jour Appl. Phys. 34, 302 (1963).

¹⁵J. B. Gibson, A. N. Goland, M. Milgram, and G. H. Vineyard, Phys. Rev. 120, 1229 (1960)

An obvious unresolved question is inherent in these theories. The many body approach and the two body assumption, with a minimum energy limitation, appear to be reasonable approaches to the problem. Both methods describe possible events in a lattice, but are they compatible? The general use of the two body approach is desirable because of its simplicity. The problem to be studied is: when can the two body collision be assumed, if at all, and should an "apparent" mass, as proposed by Henschke⁸, be employed?

2. Objective.

We intend to investigate theoretically the energy transfer process between an incoming copper atom and a copper lattice. The many body approach, using computer techniques similar to Gibson et. al.¹⁵, will be utilized and then compared with the simple two body solution. We hope the results will provide some quantitative answers to the questions posed in the preceding section.

3. Procedure.

The representative crystal lattice consists of a cube of 63 atoms arranged in a face centered cubic structure. This arrangement gives a total of 63 atoms. These atoms are free to move, if disturbed. In an effort to simulate a larger lattice, the movable 63 atoms are surrounded on all sides by stationary immovable atoms positioned as a continuation of the crystal. This procedure also gives the bombarded face some characteristics similar to a binding energy.

In order to simplify the calculations, we defined a quantity called the "lattice unit". One lattice unit is equal to half the length of a cube side. For copper, the length of a cube side is 3.614A, therefore one lattice unit is 1.807A. A three dimensional Cartesian coordinate system is used for reference with the origin placed with a (100) face of the movable 63 atom "core" parallel to the x-z plane at $y=0$. Any penetration into the lattice core from this plane is in the + y direction. The immovable atoms surrounding the core are in the x-z plane at $y=-1$ and $y=5$, the x-y plane at $z=0$ and $z=6$, and the y-z plane at $x=0$ and $x=6$. For convenience all atoms are numbered. Number one is the "bullet", numbers 2-64 are the "core" and 65-172 are the immovable atoms.(see Fig. 1)

The program depends upon a potential function subroutine; so potential modification is accomplished without any changes in the main program. The Born-Mayer and the Thomas-Fermi-Firsov have been used in the present work (see Appendix VI for exact forms and constants). All calculations are with the potential in "eroded" form. To "erode" a potential, the value of the potential at the nearest neighbor separation is subtracted from the value of the potential at smaller distances. The eroded form makes the

potential zero at distances greater than nearest neighbors.

The interaction between atoms is represented by an eroded repulsive potential. Also, the entire lattice is initially at absolute zero since no vibrational energy is simulated. These two approximations imply that the lattice has no potential or kinetic energy before interaction with the bullet and all the energy in the lattice at any time thereafter is derived from the bombarding atom.

We realize that each atom in the lattice is contained in some type of "potential well", but we are measuring the energy from the ground state level of this "well" rather than from the true zero of potential at infinite separations. Thus, before interactions with the bullet, our lattice is not held together by the surrounding stationary atoms as in the Gibson et. al.¹⁵ model because the erosion also has removed all forces. Forces appear only when atoms move from their equilibrium positions.

The use of an eroded potential does not provide as good a model as the one used by Gibson et. al., but it reduces the computer time required for calculations.

The "binding energy" created by the layer of fixed atoms covering the front face is probably much too large, but it does approximately simulate the behavior of the next layer of atoms.

The calculations which move atoms are similar to those of Gibson et. al.¹⁵ Newton's equation of motion can be rearranged to give the change in velocity of a body acted on by an unbalanced force. ($F \cdot \Delta T / m \equiv \Delta V$) The change in velocity can be related to a change in position if an average velocity is assumed ($X_{NEW} = [V + \frac{\Delta V}{2}] \cdot \Delta T + X_{OLD}$). The unbalanced force used in the first equation is an average force calculated by a double iteration procedure as follows: (1) assume an atom at position 1 with velocity 1

(2) calculate the total force on the atom as a result of all the other atoms in the lattice (this means normally only about 8-10 nearest atoms because the potential is eroded) (3) call this calculated force, force 1, and use the equation of motion to move the atom to a temporary position, position 2. (4) now repeat the force calculations for position 2, call this force 2. (5) go back to position 1, and use the average of force 1 and force 2 to move the atom to a new position, position 3. Procedures 1 through 5 constitute one "time step". Forces are eroded in the same manner as potentials and are calculated by a subroutine based on the partial derivative of the potential function with respect to distance.

The basic value of ΔT is dependent upon the original energy of the bombarding atom. ΔT is the time in seconds required for the incoming atom to traverse one lattice unit. The atom will lose energy as it interacts with the lattice but ΔT remains unchanged. To add flexibility to the program, the basic value of ΔT can be multiplied by any desired factor (time step multiplier). ΔT is established at the beginning of the program and is constant for the duration of the calculations.

The bombarding copper atom is originally located in the plane $y = \sqrt{2}$ and given a velocity in the + y direction only. This requires the bombarding particle to approach perpendicular to the (100) face of the movable lattice core. The immovable atoms in the plane ($y = -1$) that cover the bombarded face do not interact with the bullet.

An impact area was chosen on the (100) face so that all points in this area would be representative of any point in a (100) plane. An impact point is defined as the location on the (100) face toward which the bombarding atom is directed. It may or may not actually pass through this point. By moving in 1/10 lattice unit increments, 36 impact points in the impact "triangle"

were assigned. (see Fig. 2 and 3).

For each case the two body problem is solved after the many body problem with the same double iteration scheme. Lattice atom number eight was assigned as the "target" in the many body lattice problem so that its motion could be compared with the simple two body problem. The geometrical relationship and physical constants of the two body problem correspond to those of the incoming atom and "target" in the crystal.

The program was written in FORTRAN language for use on a CDC 1604 computer. The input consists of the incoming particle mass and energy and the time step multiplier. From these inputs the bullet is assigned a location and an appropriate velocity component in the + y direction. Also calculated is the actual value of ΔT .

The next section of the program assigns coordinates to all atoms in the crystal, movable and immovable, gives each atom a number and sets their velocities equal to zero. If the distance between any two atoms is less than $\sqrt{2}$ lattice units, the force between these two atoms is calculated and stored in memory. This is done for each atom and the summation of all the forces calculated gives the resultant force on each particle.

The change in the velocities of all movable atoms as a result of unbalanced forces is calculated; then the original velocities of all particles and the calculated change in velocity are used to move atoms to temporary locations. The calculations to obtain resultant forces are repeated for the temporary locations, old and new forces are averaged and used to move all atoms to new locations. Forces and velocities are calculated as vector components using the coordinate system established for the lattice.

After each movement of all the atoms, the kinetic energy of the target is calculated and compared to its former value. The procedure is repeated

until the kinetic energy of the target reaches a maximum. The potential and kinetic energy of each movable atom is then calculated. The total energy of the entire lattice including bullet is also found and compared to the original "bullet" energy. This comparison gives a reasonable measure of the accuracy achieved. The two body problem is solved in the next section using the same iteration technique. Appendices I-V contain a detailed explanation of the program.

For each "run", the computer will solve the many body and two body problem 36 times, corresponding to the 36 different impact points. The form of the printed output for each impact point is shown in Fig. 4 and 5. The "triangular" output as seen in Fig. 6 indicates how two chosen parameters vary over the impact area.

The computer running time for each of the 36 impact points was 2-4 minutes, the time varying with the value chosen for ΔT . We originally assumed that the accuracy of the results would increase for smaller and smaller values of ΔT . We found that the value of ΔT which will give the minimum error in total energy appears to be a complicated function of the time step multiplier, the incoming bullet energy, and the impact point. By trial and error methods, we found that below 100 ev for the Born-Mayer potential, a time step multiplier of around 0.09 would produce reasonable errors at all impact points.

4. Potentials

The Born-Mayer potential is an exponential function of the internuclear separation with constants which may be obtained from experimentally measurable elastic moduli. The constants used in the program are the same as those of Gibson et. al.¹⁶, with special emphasis on their potential number two. Near equilibrium separations, the Born-Mayer potential is thought to be an adequate approximation. We expect the potential to fail for energies above a few hundred electron volts.

The Thomas-Fermi-Firsov (TFF) potential is the result of theoretical work by Firsov¹⁷ based on the Thomas-Fermi model of the atom. It is a screened Coulomb potential with a more complicated screening function than that used in the Bohr potential. The TFF is satisfactory for intermediate separations. This places it between the Bohr potential used for small separations and the Born-Mayer potential. Abrahamson and Hatcher¹⁸ state that the Thomas-Fermi approximation becomes unreliable when the internuclear distance exceeds $\sim 1\text{\AA}$. The form of the TFF potential and the appropriate constants used in the program are the same as those used in computer programs at the Oak Ridge National Laboratory¹⁹.

The mathematical forms of the potential and force functions as they appear in the program are derived in Appendix VI.

¹⁶See reference 15, p. 1233.

¹⁷O. B. Firsov, J. Exptl. Theoret. Phys. (U.S.S.R.) 32, 1464 (1957). translation: Soviet Phys. -JETP 5, 1192 (1957).

¹⁸A. A. Abrahamson and R. D. Hatcher, Phys. Rev. 121, 159 (1961).

¹⁹R. T. Robinson (private communication). The author would like to thank Dr. Robinson for the use of his results prior to their publication.

5. Results

A. General description of events in the lattice.

The number of atoms disturbed in the lattice is a function of two variables, the number of time steps the computer performs and the impact point of the bullet. In general, the lattice is more disrupted by impact points in the area near $x = 3.0$, $z = 4.0$, and the lattice is least affected by head on collisions with the target. No more than one half of the atoms were ever disturbed before the target reached a maximum in kinetic energy.

The impact area is not symmetrically located with respect to the surrounding immovable atoms but for appropriate points in the impact area, excellent symmetry of displacement was observed for the movable atoms. This indicates that the size of the model is adequate for present purposes.

Because of the limitation placed on the number of time steps, only a few atoms have any significant kinetic energy when the interaction is stopped. These are normally the target, bullet, and the atoms directly behind the target and bullet in the lattice.

Although our primary purpose was to compare lattice interactions with the two body interaction, we altered the program so that it would run a predetermined number of time steps in order to observe "chains" and "channels". Impact points at $x = 3.0$, $z = 4.0$ and $x = 3.0$, $z = 3.0$ correspond to the beginning of (100)chains. These are not "close packed" chains but they transfer energy reasonably well, especially above 100 ev. Impact point $x = 2.5$, $z = 3.5$ is the beginning of a (100) channel. The (100) channel is not as "wide open" as a (110) channel and a 100 ev bullet does not travel more than $\sim 3A$ into the lattice.

B. Comparison of lattice interactions with the two body interaction.

A strict comparison of interactions in the lattice to the two body problem is not possible, but certain limited correlations can be made. The difficulty arises because the end of the interaction in the lattice cannot be defined explicitly. There are several criteria available that may be used in an attempt to compare the interactions. If the two body interaction was allowed to continue until the potential between atoms was zero, for comparison purposes the lattice interaction could proceed until: 1) the potential between target and bullet is equal to zero 2) the total energy (kinetic plus potential) of the target reaches a maximum or 3) the kinetic energy of the target reaches a maximum. All of the criteria apply to the same physical situation in the two body problem but are not equivalent in the lattice interaction.

We chose the maximum kinetic energy of the target as the comparison stopping point because this represents the time when the target and bullet have almost ceased to interact and are just beginning to interact with other lattice atoms.

A fourth criteria was suggested after the present results had been obtained. When the kinetic energy of the target in the lattice reaches a maximum, find the distance between the target and bullet and allow the two body interaction to proceed until the distance between atoms is equal to the separation in the lattice. This criteria is possibly better than the one chosen but at the present time it has not been utilized.

The geometrical relationships of the target, bullet, and lattice introduce complications into the comparison scheme. The motion of the target in the lattice is not restricted isotropically (see Fig. 7). If the interaction with the bullet moves the target toward atom number six (shaded area),

the movement of the target is restricted more than if the target motion is initially toward the cross-hatched area. Restriction of the target motion lengthens the interaction time between bullet and target, therefore the energy transferred to the target is greater than in the two body problem. This effect is noticable only at impact parameters of 0.5 lattice units or greater. At impact distances smaller than this the interaction occurs so rapidly that the target does not move a significant amount and the effect is reduced. The kinetic energy transferred at distances smaller than 0.5 lattice units is less than that transferred in the two body case because there is absorption of energy by the remainder of the lattice. The geometry of the bombarded face is not the only factor that effects energy transfer. The atoms behind the target in the lattice also play an important role but their effect is not as immediately obvious as those just discussed.

Figure 8 illustrates some of the points just mentioned. Along the line A B ($x = 3.0$, $z = 3.0-4.0$) the target behavior in the lattice is similar to the two body interaction. These impact points cause the target to move in the negative z direction toward a "hole" in the lattice face. At point D on the line A C , the effect of restricted target motion becomes obvious, and also in the area enclosed by the dashed line. The geometrical effects of the lattice are fairly predictable and no unusual phenomena are observed.

For a static system of two equal mass atoms with a conservative repulsive force between them, the total potential energy of the system can be halved and the result assigned as the energy of the atom. If the two atoms are allowed to move, the total energy (potential plus kinetic) of each will be constant and equal to the original potential energy. In a system composed of three or more atoms, this is generally not true. For any system where the atoms are not static, the process of assigning an energy to any one atom is

no longer possible (there is one exception). This is the situation in the lattice.

The potential energy of an atom in the lattice is defined in the program as half the potential associated with its position. This definition is merely a convenience.

The area of the impact triangle from point A to the arc E-F (see Fig. 8) is least effected by the geometry of the lattice and closer inspection of these interactions is justified. The kinetic energy transferred to the n-body target is less than that of the two body target. In percentages, the kinetic energy transferred becomes less as the bullet energy is decreased. (see Figs 8 through 11). In each case the target also acquires some potential energy (see Figs. 12 through 15). As stated earlier, it is difficult to assign a specific potential energy to the target. We do know the limits involved, i.e. the target could eventually receive none of it, or twice the indicated value. Either limiting situation is unlikely.

If we make the assumption that one-half of the potential "belongs" to the target, then the total energy of the target is approximately equal to the energy transferred in the two body problem (within 2%) for energies greater than 50 ev.

Bullet behavior for the area under discussion is also very similar to the two body problem (see Figs. 16 through 19). For the remainder of the impact triangle, the bullet is strongly effected by other atoms in the lattice.

Three runs were made with the TFF potential at energies of 5, 10, and 30 kev (Figs. 20 through 25). Since the TFF potential is unreliable beyond 1A, the section ABC in Figs. 20 through 25 is the only part of the impact area that can be considered. This section includes all impact points that

result in a closest point of approach of 1A or less. We noted the same general agreement with the two body interaction as found for the Born-Mayer potential. The only apparent difference is a shift in the energy scale.

For a head-on collision in the lattice there is very little geometrical effect, and an "effective" mass can be described. At 25 ev, the mass of the target is apparently about 2.5 times the bullet mass, but the energy transfer is still approximately 80%. For a 500 ev bullet, the effective mass increase is negligible. We used the program, as modified to observe chains and channels, and could find no evidence of bullet recoil which could be attributed to an effectively heavy target as proposed by Henschke⁸. Although an effective mass concept for the head-on collision is possible, it does not properly describe the subsequent motion of either target or bullet. If the bullet motion results in a glancing hit with the target, no single effective mass can be assigned to the target because it is a function of the impact parameter. An average effective mass is not applicable because the geometrical effects of the lattice on the direction and energy of the recoil atom are far more significant than the mass of the target in a glancing hit.

These results apply to a collision anywhere in the lattice and are not limited to the interaction of surface atoms with incoming particles. The results indicate that Henschke's⁸ apparent mass concept and the rebound phenomena associated with it are not a good description of collision events in a lattice.

6. Conclusions

A. The Bullet.

For energies above 40 ev in the impact area AEF (see Figs 16 and 17), the n-body bullet's kinetic energy after interaction is essentially the same as the two body bullet. Between 40 and 400 ev for impact points outside the area AEF, the n-body bullet has considerably less energy after interaction than the two body bullet. Above 400 ev, the n-body and two body energy transfers agree within 3% for all points in the impact area.

The angular behavior of a bullet in the lattice is much more complicated. Even at energies above 400 ev, the scattering angle in the n-body problem is affected by the geometry of the lattice. In general, the scattering angle in the lattice is smaller than that of the two body problem. It is possible that the lattice is attempting to focus the energy into preferred directions. Our bullets in the (100) direction appear to focus in the (100) direction, especially at higher energies.

The cumulative effects of energy transfer and lattice geometry will affect the range of energetic atoms. The reduced scattering angle in the n-body model for atoms with energies above 400 ev should lead to ranges that exceed those found by two body approximation methods. Below 400 ev, the bullet loses more energy than in the two body case; so the effects may cancel, or perhaps ranges calculated by two body methods may be too large.

B. The Target.

In the area AEF (see Figs. 8, 9, 12, 13), the n-body energy transferred to the target is the same (within 2%) as that in the two body problem for energies above 50 ev. For the remainder of the impact area, the n-body target receives more energy than the two body target for all energies above

25 ev. Isolated points at various energies may not conform to this rule but the effect is always present for impact points near the apex of the impact area ($x=2.5$, $z=3.5$). Near the apex, if the bullet energy is large (300 ev or greater), the energy transferred to the target is a very small fraction of the total for both the n-body and two body problems. However, the energy transferred to the target in the two cases can differ by as much as 75% of the transferred energy.

As a consequence of this phenomena, the low energy portion of the n-body target atom energy distribution (assuming more than one collision has occurred) will be much higher than that expected with a two body collision assumption for lattice interactions. Precise measurements of energy transfer for large impact parameters do not appear to be warranted in view of the results obtained.

Unfortunately, the face centered cubic lattice structure was incorporated in the program before we learned of Veksler's²⁰ experimental work with molybdenum targets (bcc), but certain qualitative comparisons are still possible. Veksler has interpreted the lattice behavior in terms of the effective mass concept but his general conclusions are consistent with our work. Our model gives good evidence to support Veksler's position that the pair collision model using elastic spheres is not acceptable.

Two body approximations inherently imply that certain information is available from preliminary n-body calculations. This work attempts to answer certain questions about interactions and cross sections. As anticipated, the results are not definitive but they do indicate sensitive areas which require further examination.

²⁰V. I. Veksler, J. Exptl. Theoret. Phys. (U.S.S.R.) 42, 325 (1962). translation: Soviet Physics - JETP 15, 222 (1962).

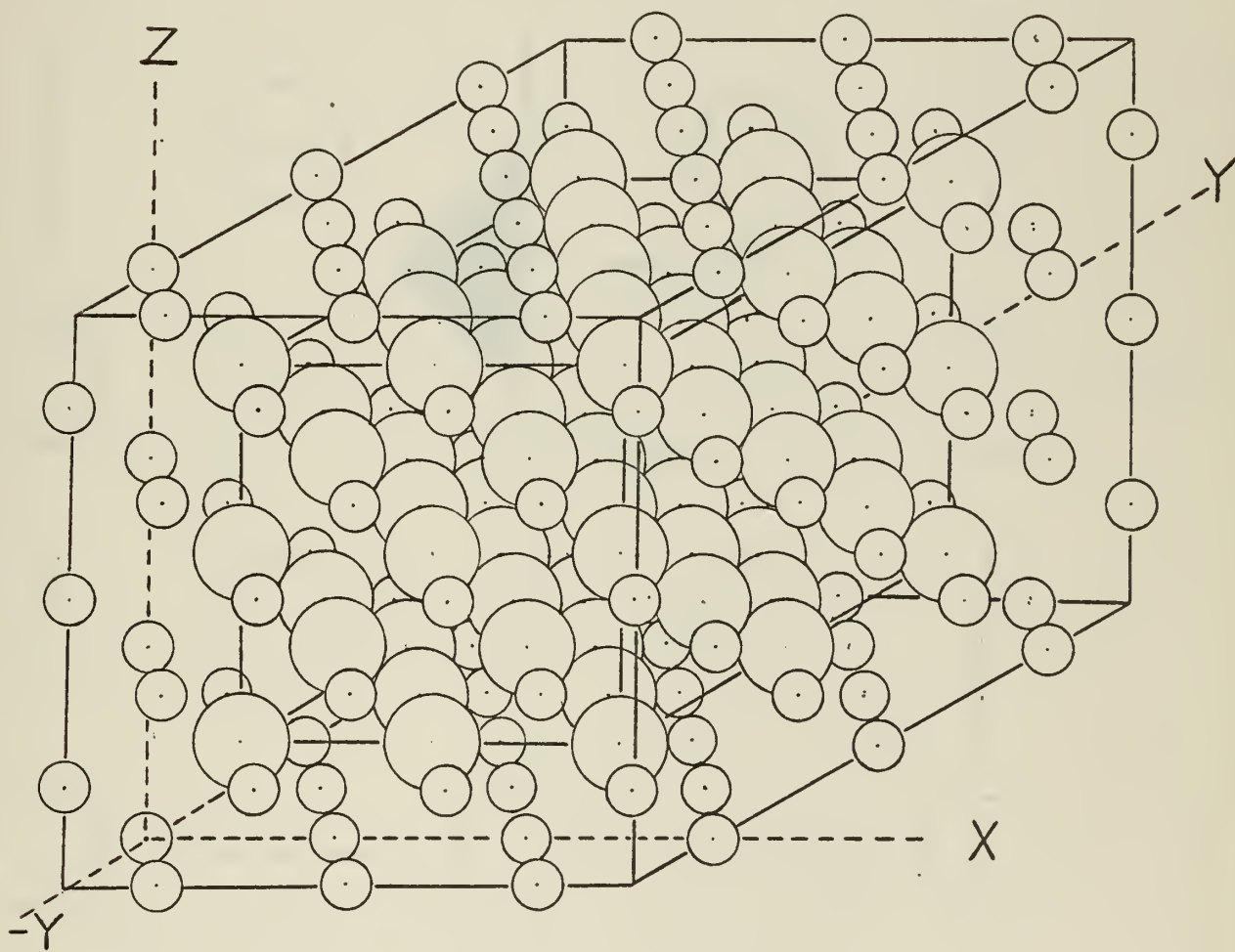


Fig. 1. THE MODEL
Large circles indicate movable atoms;
small circles indicate fixed atoms.

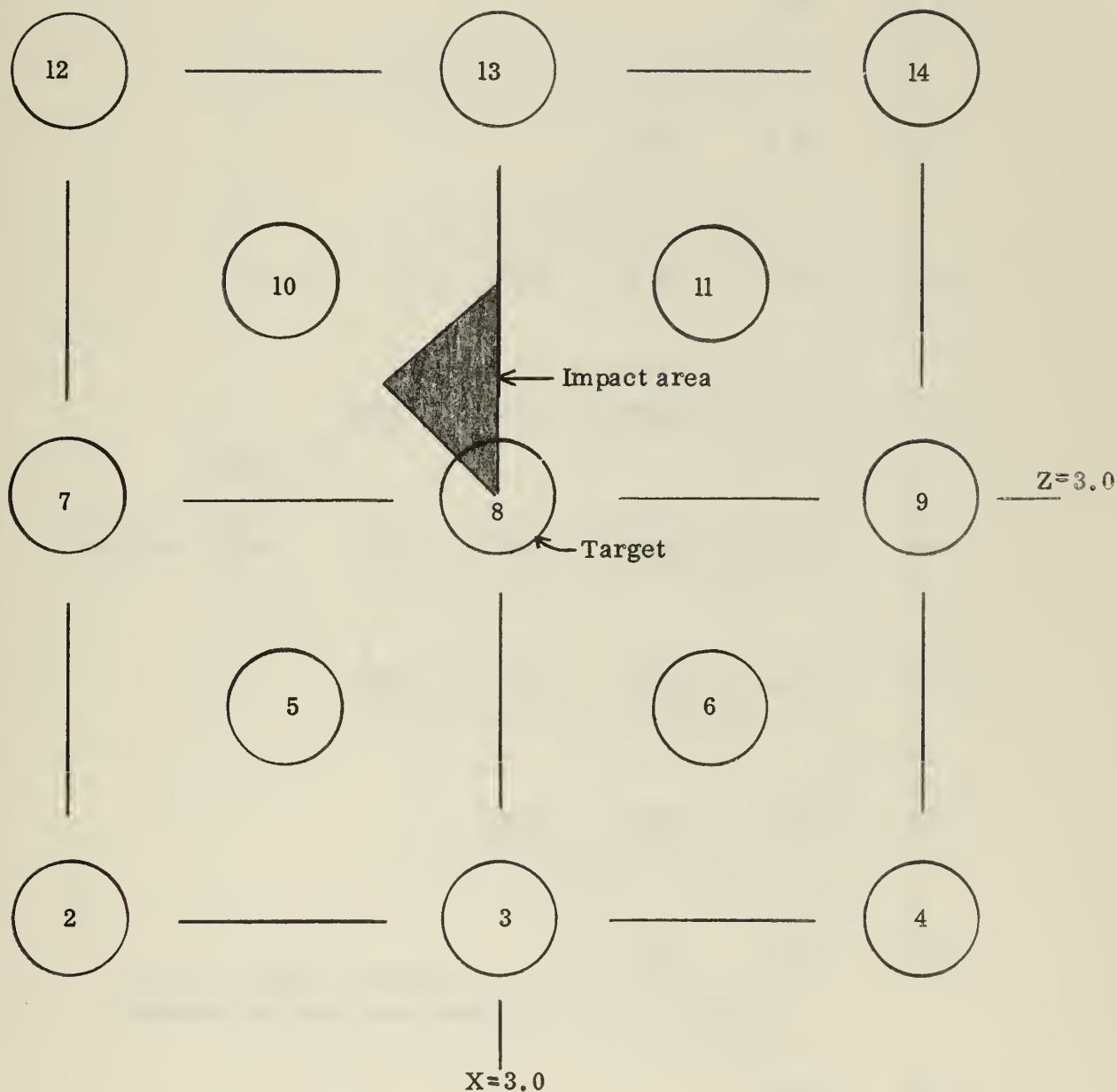
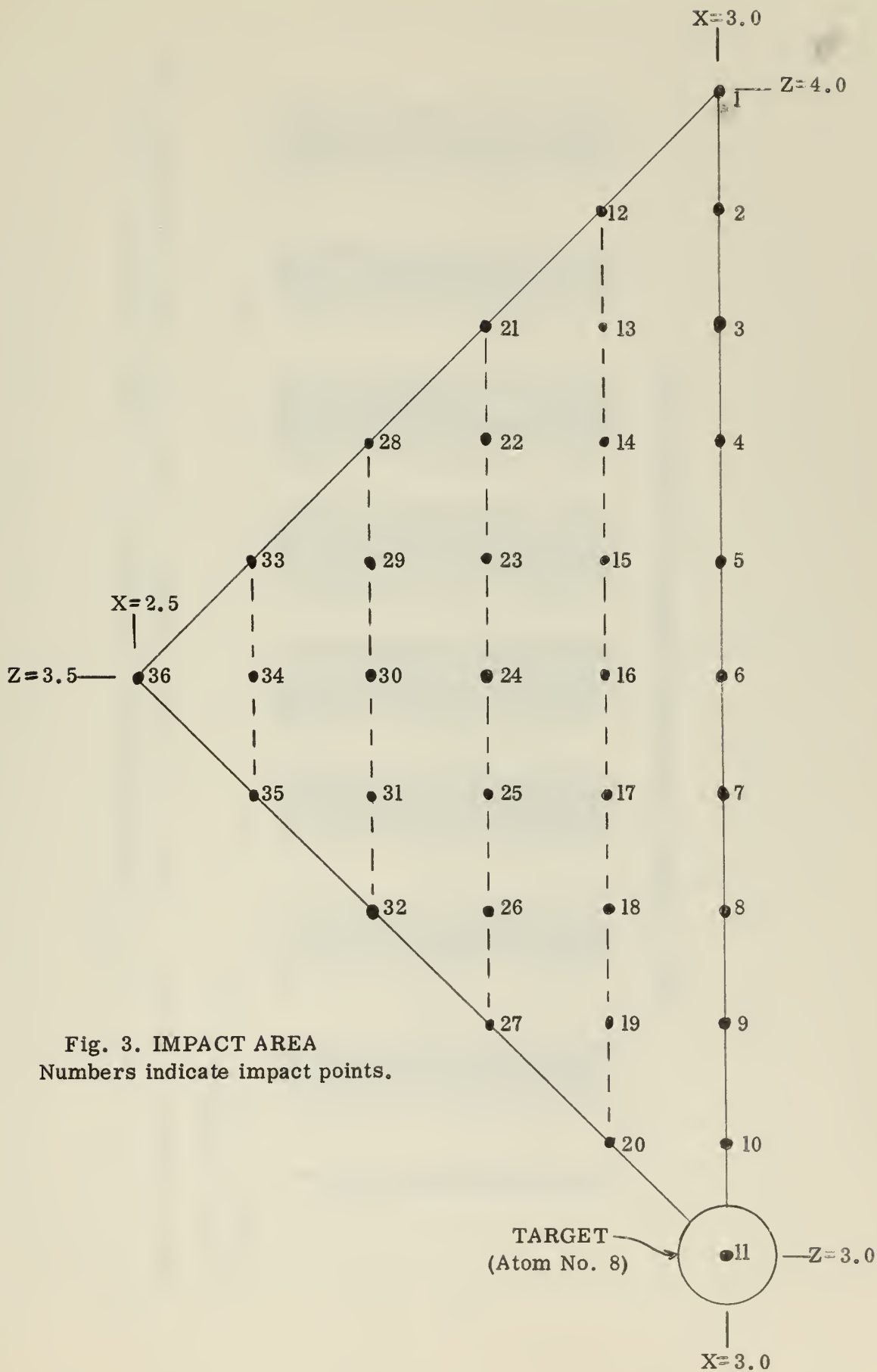


Fig. 2. FACE (100) of LATTICE CORE.
Numbers indicate atoms; shaded portion is impact area



BORN-MAYER ERODED POTENTIAL				COPPER LATTICE AND BULLET				
BULLET INCOMING LOCATION	X= 3.000	Z= 3.900	ENERGY= 100.0	TIME STEP MULTIPLIER= .080				
LATTICE LOCATIONS								
ATOM	DX	DY	DZ	VX	VY	VZ	KE	PE
1	.0000	1.30093	.00094	.00E+00	.49E+03	.16E+03	.08939	1.50805
6	.00056	-.00001	-.00052	.25E+02	.59E+00	.22E+02	.00037	.041628
8	.00000	-.00094	-.15781	.00E+00	.21E+03	.25E+04	.00737	.30768
9	.00019	-.00000	-.00020	.84E+01	.67E-01	.88E+01	.00005	.014597
10	-.07887	.00832	.00659	.14E+04	.19E+02	.10E+03	.62477	.250877
11	-.00024	.00000	.00025	.10E+02	.83E-01	.11E+02	.00007	.018227
12	.00000	.00152	.04014	.10E+00	.72E+02	.78E+03	.20369	.218278
13	.00024	.00000	.00025	.10E+02	.83E-01	.11E+02	.00007	.018227
14	.00000	.00079	.00074	.00E+00	.35E+02	.32E+02	.00072	.056989
18	-.00013	.00016	-.00002	.38E+01	.44E+01	.30E+00	.00001	.002889
20	.00013	.00016	-.00002	.38E+01	.44E+01	.30E+00	.00001	.002889
21	.00028	.00029	.04273	.11E+02	.17E+05	.51E-01	.00009	.018895
22	.00000	.00028	-.00000	.00E+00	.12E+02	.13E-01	.52631	2.408845
23	.00028	.00003	-.00000	.11E+02	.12E+02	.13E-01	.00009	.018844
24	.00000	.00003	.00007	.52E+00	.25E+01	.23E+01	.00000	.002133
25	.00000	.00055	.00077	.00E+00	.48E+02	.78E+02	.00000	.002133
26	.00000	.00068	-.00002	.00E+00	.63E+02	.32E+01	.00000	.26470
33	.00096	.00068	-.00002	.10E+03	.63E+02	.32E+01	.00000	.38456
35	.00096	.00068	-.00002	.10E+03	.63E+02	.32E+01	.00000	.38456
36	.00000	.00086	.00119	.00E+00	.85E+02	.14E+03	.00846	.56089

Fig. 4. TYPICAL COMPUTER OUTPUT for an impact point

BORN-MAYER ERODED POTENTIAL COPPER LATTICE AND BULLET
 BULLET INCOMING LOCATION X= 3.000 Z= 3.900 ENERGY= 100.0 TIME STEP MULTIPLIER= .080

RESULTS OF SIMPLE TWO BODY INTERACTION

KIN ENRG PROJ= 98.293220 KIN ENRG TARG= 1.70688 TOT ENRG= 100.00010 ENRG IN= 100.00000

RESULTS OF LATTICE INTERACTION

TARGET VALUES

TIME STEPS POT E TARGET W/R TO BULLET KE TARGET POT E TARGET W/R TO LAT-BULLET
 32. 269015 2.007372 .00000

BULLET VALUES

KIN ENERGY = .089 POT E W/R TO LAT-TARGET = 1.369 VELOCITY IN M/SEC X= .00E-01 Y= 4.95E+02 Z= 1.64E+02

LATTICE VALUES

POT ENERGY 6.87605 KIN ENERGY 94.09908 TOT ENERGY 100.97512 PER ERROR .97512

COMPARISON VALUES OF LATTICE REACTION AGAINST SIMPLE TWO BODY PROBLEM BOTH INTERACTIONS STOPPED WHEN TARGET KINETIC ENERGY MAXIMUM

PERCENT KE TRANSFERRED TWO BODY= 1.707 LATTICE= 2.007

RATIO KE TRANSFER IN LATTICE / KE TRANSFER TWO BODY = 1.17605

RATIO KE PROJ IN LATTICE / KE PROJ TWO BODY = .00091

RATIO BULLET MASS / TARGET MASS = 1.000

POT + KE ABSORBED BY LATTICE = 97.06764

SCATTERING ANGLE TWO BODY = 7.51 LATTICE = 18.29

RECOIL ANGLE TWO BODY = 82.49 LATTICE = 85.02

RATIO SCATTERING ANGLE TWO BODY / SCATTERING ANGLE LATTICE = .41035

RATIO RECOIL ANGLE TWO BODY / RECOIL ANGLE LATTICE = .97033

IMPACT PARAMETER = .9000 RATIO KE BULLET AT END OF RUN / ORIGINAL ENERGY = 8.93871E-04

Fig. 5. TYPICAL COMPUTER OUTPUT for an impact point

BORN-MAYER ERODED POTENTIAL
CCPPE LATTICE AND BULLET

BULLET INCOMING ENERGY = 100.0 EV TIME STEP MULTIPLIER = .080 BULLET INCIDENT ON (100) FACE

TOP NO. IS KINETIC ENERGY IN EV OF BULLET IN LATTICE

BOTTOM NO. IS KINETIC ENERGY IN EV OF TARGET IN LATTICE

B		C	
(X=3.0)	.124 * .710	(Z=3.5)	41.205 * 53.160
	.089 * 2.007	(Z=3.2)	25.710 * 68.855
			3.241 * 91.836
			(3.0, 3.0) TARGET ATOM
(X=2.9)	.067 * 1.972		
	8.557 * 5.279		14.870 * 79.645
			6.298 * 88.624
(X=2.8)	4.380 * 5.104		
	45.648 * 11.263		21.749 * 71.643
(X=2.7)	23.435 * 11.135		33.076 * 50.190
(X=2.6)	45.033 * 19.221		36.110 * 32.090
(X=2.5)	43.295 * 23.701		

Fig. 6. TRIANGULAR COMPUTER OUTPUT; (*) indicates impact point

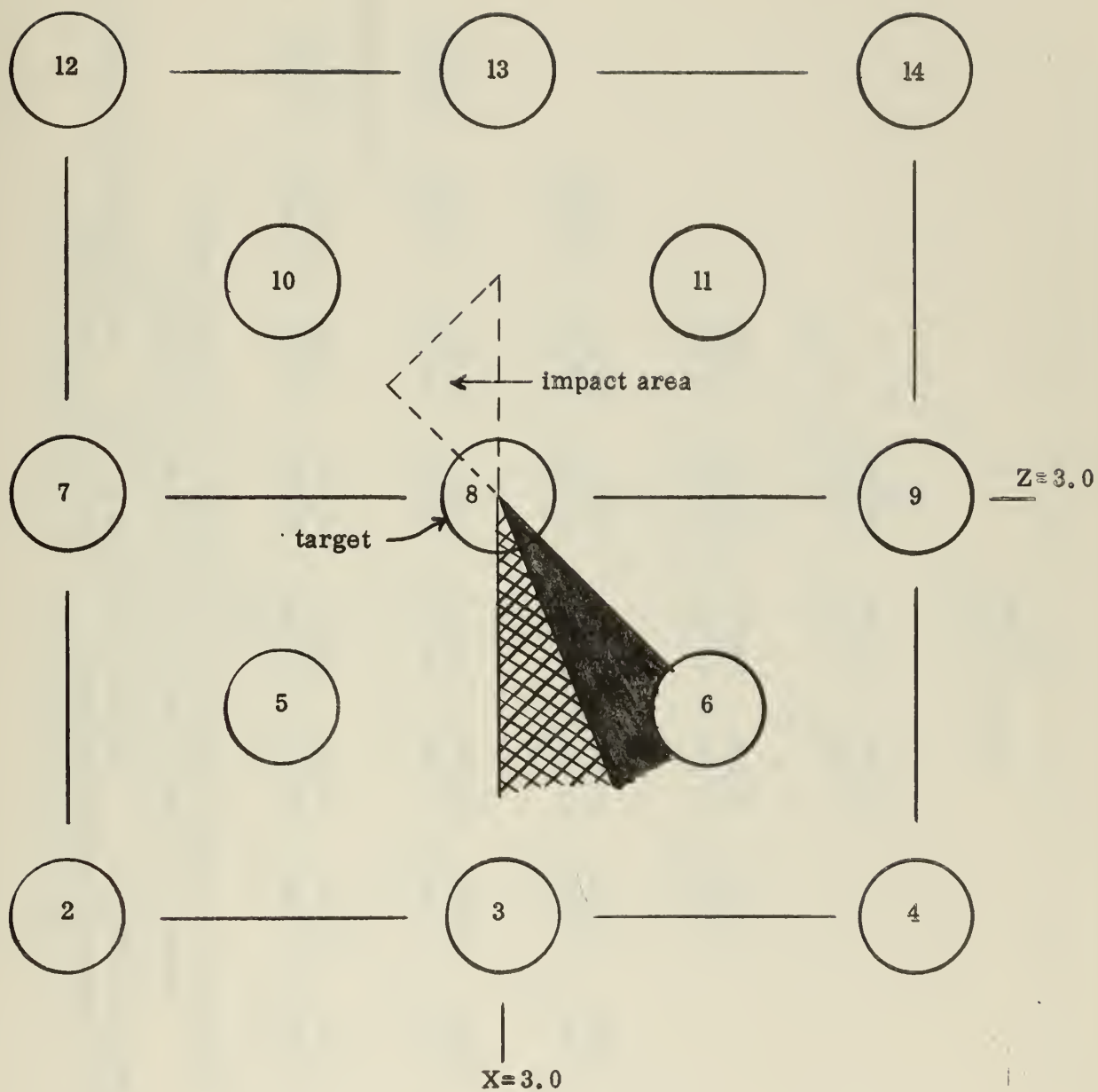


Fig. 7. GEOMETRICAL EFFECTS OF THE LATTICE FACE.
Target movement is more restricted if initial motion is in shaded area rather than in crosshatched area.

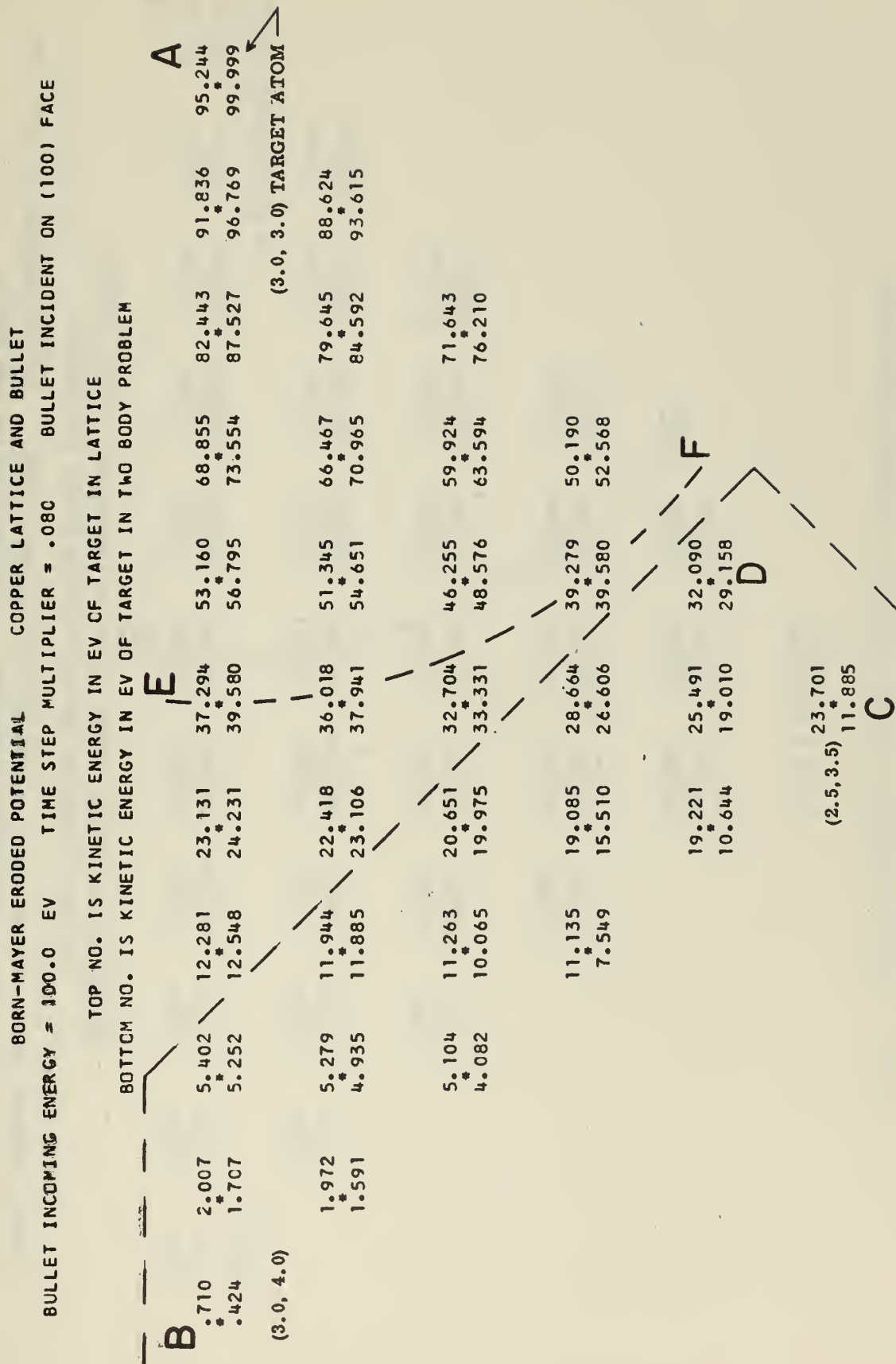


Fig. 8. GEOMETRICAL EFFECTS on 100 ev bullet; (*) indicates impact point

BORN-AYER EXCESS POTENTIAL

COPPER LATTICE AND BULLET

BULLET INCOMING ENERGY = 80.0 EV TIME STEP MULTIPLIER = .C90 BULLET INCIDENT ON (100) FACE

TOP NO. IS KINETIC ENERGY IN EV OF TARGET IN LATTICE

BOTTOM NO. IS KINETIC ENERGY IN EV CF TARGET IN TWO BODY PROBLEM

	E				A					
.810	2.122	5.364	11.451	20.442	31.687	43.878	55.795	65.780	72.806	75.383
.485	1.837	5.277	11.834	21.705	34.091	47.525	60.294	70.764	77.617	79.999
(3.0, 4.0)									(3.0, 3.0) TARGET ATOM	
	2.094	5.272	11.188	19.900	30.673	42.429	54.086	63.757	70.410	
	1.719	4.979	11.255	20.776	32.791	45.872	58.338	68.576	75.284	

(3.0, 4.0)

(3.0, 3.0) TARGET ATOM

28

10.860	17.547	25.166	33.394	41.790
7.393	14.390	23.657	34.091	44.261
				F
	18.124	22.998	27.936	
	10.164	17.359	25.736	

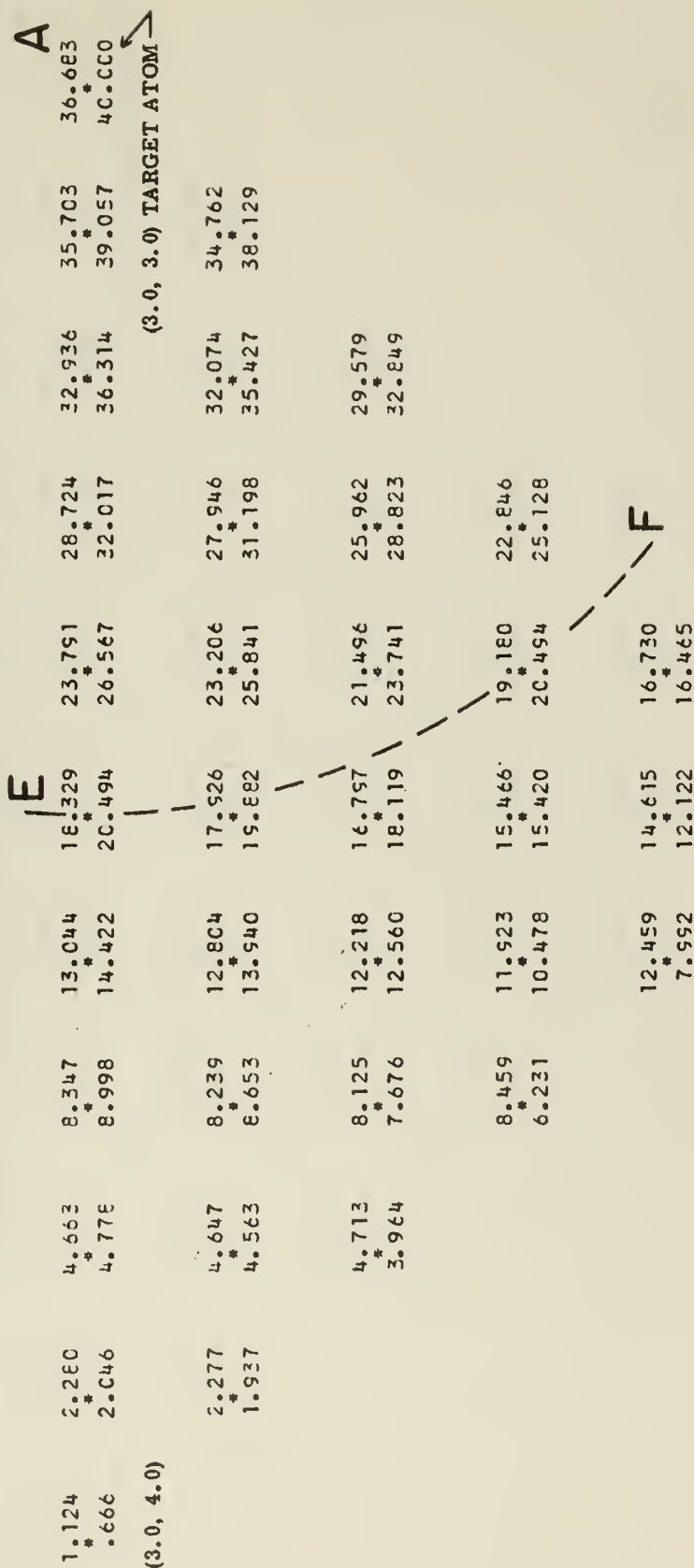
W

(2.5, 3.5) 22.045
11.255

Fig. 9. OUTPUT TRIANGLE for 80 ev bullet; (*) indicates impact point

BOERN-MAYER ERODED POTENTIAL COPPER LATTICE AND BULLET
 BULLET INCOMING ENERGY = 40.0 EV TIME STEP MULTIPLIER = .090 BULLET INCIDENT ON (100) FACE

TOP AC. IS KINETIC ENERGY IN EV OF TARGET IN LATTICE
 BOTTOM NO. IS KINETIC ENERGY IN EV OF TARGET IN TWO BODY PROBLEM



¹⁴356
 (2.5, 3.5) * 8.653

Fig. 10. OUTPUT TRIANGLE for 40 ev bullet; (*) indicates impact point

BORN-MAYER ERODED POTENTIAL

BULLET INCOMING ENERGY = 25.0 EV TIME STEP MULTIPLIER = .090 BULLET INCIDENT ON (100) FACE

TOP NO. IS KINETIC ENERGY IN EV OF TARGET IN LATTICE

BOTTOM NO. IS KINETIC ENERGY IN EV OF TARGET IN TWO BODY PROBLEM

[illegible]

Fig. 11. OUTPUT TRIANGLE for 25 ev bullet; (*) indicates impact point

BORN-MAYER, ERODED POTENTIAL

COPPER LATTICE AND BULLET

BULLET INCOMING ENERGY = 100.0 EV

TIME STEP MULTIPLIER = .C8C

BULLET INCIDENT ON (100) FACE

BULLET INCOMING ENERGY = 100.0 EV TIME STEP MULTIPLIER = .C8C BULLET INCIDENT ON (100) FACE

TOP NO. IS POTENTIAL ENERGY IN EV OF TARGET IN LATTICE

BOTTOM NO. IS POTENTIAL ENERGY IN EV OF BULLET IN LATTICE

[illegible]

Fig. 12. OUTPUT TRIANGLE for 100 ev bullet; (*) indicates impact point

COPPER LATTICE AND BULLET

BULLET INCOMING ENERGY =: 80.0 EV TIME STEP MULTIPLIER = .090 BULLET INCIDENT ON (100) FACE

TOP NO. IS. POTENTIAL ENERGY IN EV OF: TARGET IN LATTICE

BOTTOM NO. IS POTENTIAL ENERGY IN EV OF BULLET IN LATTICE

[illegible]

Fig. 13. OUTPUT TRIANGLE for 80 ev bullet; (*) indicates impact point

BORN-MAYER ERODED POTENTIAL

COPPER LATTICE AND BULLET

BULLET INCOMING ENERGY = 40.0 EV

TIME STEP MULTIPLIER = .09C

BULLET INCIDENT ON (100) FACE

TOP NO. IS POTENTIAL ENERGY IN EV OF TARGET IN LATTICE

BOTTOM NO. IS POTENTIAL ENERGY IN EV OF BULLET IN LATTICE

[illegible]

Fig. 14. OUTPUT TRIANGLE for 40 ev bullet; (*) indicates impact point

BORN-MAYER ERODED POTENTIAL COPPER LATTICE AND BULLET

BULLET INCOMING ENERGY = 25.0 EV TIME STEP MULTIPLIER = .C90 BULLET INCIDENT ON (100) FACE

JOP NO. IS POTENTIAL ENERGY IN EV OF TARGET IN LATTICE

BOTTOM NC. IS POTENTIAL ENERGY IN EV OF BULLET IN LATTICE

.353	.400	.559	.719	.996	1.159	1.315	1.476	1.470	1.464	1.460
*1.954	*5.453	*3.603	*2.554	*1.821	*1.187	*.872	*.916	*.674	*.566	*.541
(3.0, 4.0)										(3.0, 3.0) TARGET ATOM
.408	.543	.543	.732	.938	1.138	1.295	1.449	1.469	1.467	
*5.347	*3.878	*2.442	*2.442	*1.991	*1.535	*1.116	*.674	*.754	*.609	
.585	.795	.959	.959	1.115	1.268	1.364	1.469			
*3.091	*2.100	*2.233	*2.337	*1.860	*1.174	*.998				
.855	1.025	1.144	1.213	1.290						
*1.469	*1.604	*2.367	*2.760	*2.014						
1.139	1.228	1.228	1.228	1.228						
*.989	*1.560	*2.517								
(2.5, 3.5)	1.288									
	*.962									

Fig. 15. OUTPUT TRIANGLE for 25 ev bullet; (*) indicates impact point

COPPER LATTICE AND BULLET

BULLET INCOMING ENERGY = 100.0 EV TIME STEP MULTIPLIER = .080 BULLET INCIDENT ON (100) FACE

TOP NO. IS KINETIC ENERGY IN EV OF BULLET IN LATTICE

BOTTOM NO. IS KINETIC ENERGY IN EV OF BULLET IN TWOBODY PROBLEM

[illegible]

Fig. 16. OUTPUT TRIANGLE for 100 ev bullet; (*) indicates impact point

COPPER LATTICE AND BULLET

BULLET INCOMING ENERGY = 8C.0 EV TIME STEP MULTIPLIER = .C9C BULLET INCIDENT ON (100) FACE

TOP NC. IS KINETIC ENERGY IN EV OF BULLET IN LATTICE

BOTTOM NO. IS KINETIC ENERGY IN EV OF BULLET IN TWOBODY PROBLEM

[illegible]

Fig. 17. OUTPUT TRIANGLE for 80 ev bullet; (*) indicates impact point

BORN-MAYER ERODED POTENTIAL COPPER LATTICE AND BULLET

BULLET INCOMING ENERGY = 25.0 EV TIME STEP MULTIPLIER = .090 BULLET INCIDENT ON (100) FACE

TOP NO. IS KINETIC ENERGY IN EV OF BULLET IN LATTICE

BOTTOM NO. IS KINETIC ENERGY IN EV OF BULLET IN TWOBODY PROBLEM

		E		A		(3.0, 3.0) TARGET ATOM					
.195	3.755	12.597	12.848	12.014	9.442	6.848	4.212	1.989	.519	.011	
24.251	23.012	20.905	18.001	14.559	10.918	7.418	4.363	2.000	.509	.000	
(3.0, 4.0)											

Fig. 19. OUTPUT TRIANGLE for 25 ev bullet; (*) indicates impact point

THOMAS-FERMI-FIRSOV POTENTIAL, ERODED, COPPER-COPPER

BULLET INCOMING ENERGY = 5000.0 EV TIME STEP MULTIPLIER = .090 BULLET INCIDENT ON (100) FACE

TOP NO. IS KINETIC ENERGY IN EV OF TARGET IN LATTICE

BOTTOM NO. IS KINETIC ENERGY IN EV OF TARGET IN TWO BODY PROBLEM

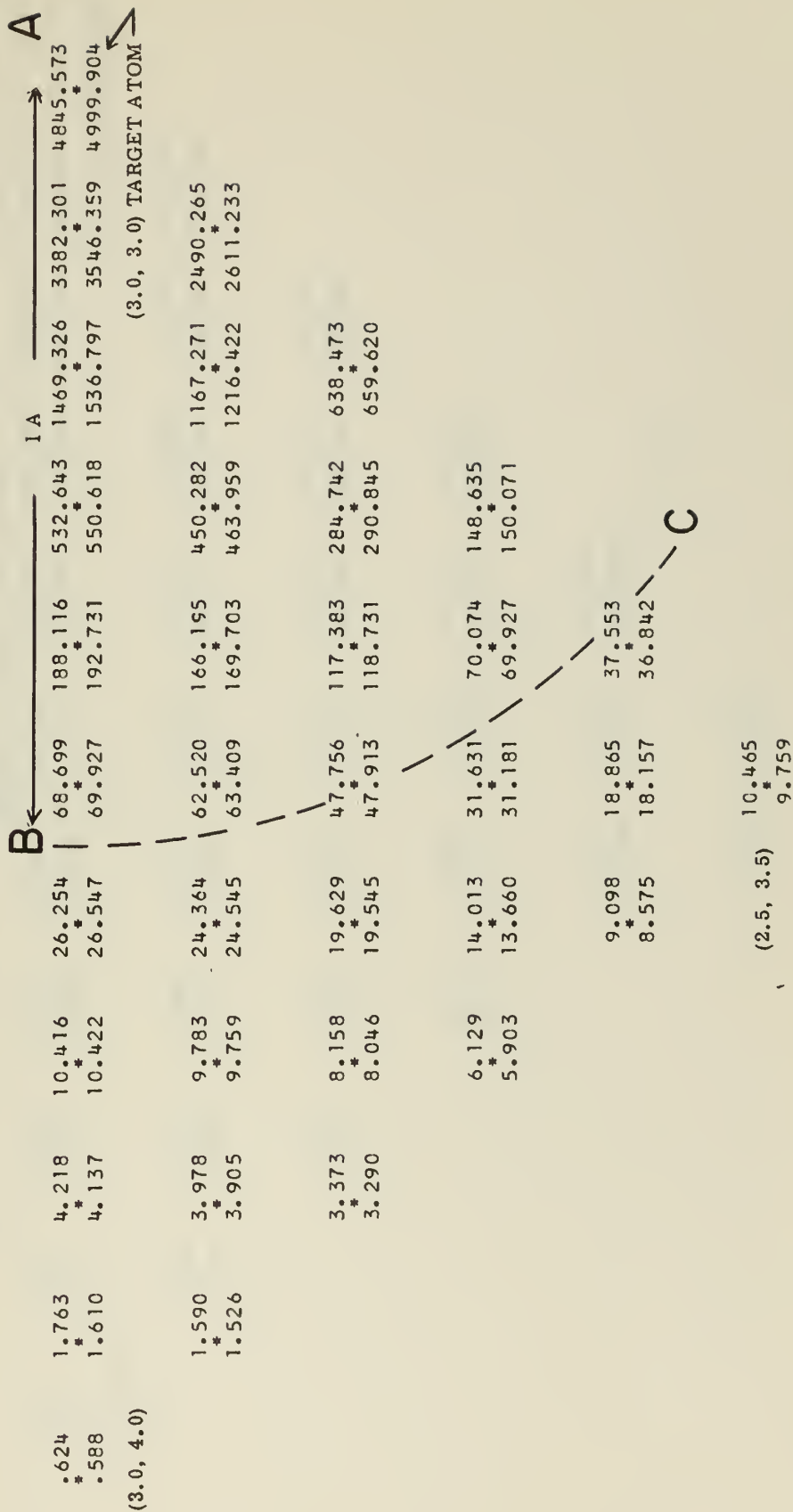


Fig. 20. OUTPUT TRIANGLES for 5 kev bullet; (*) indicates impact point

THOMAS-FERMI-FIRSOV POTENTIAL, ERODED, CUPPER-COPPER

BULLET INCOMING ENERGY = 10000 EV TIME STEP MULTIPLIER = .090 BULLET INCIDENT ON (100) FACE

TOP NO. IS KINETIC ENERGY IN EV OF TARGET IN LATTICE

BOTTOM NO. IS KINETIC ENERGY IN EV OF TARGET IN TWO BODY PROBLEM

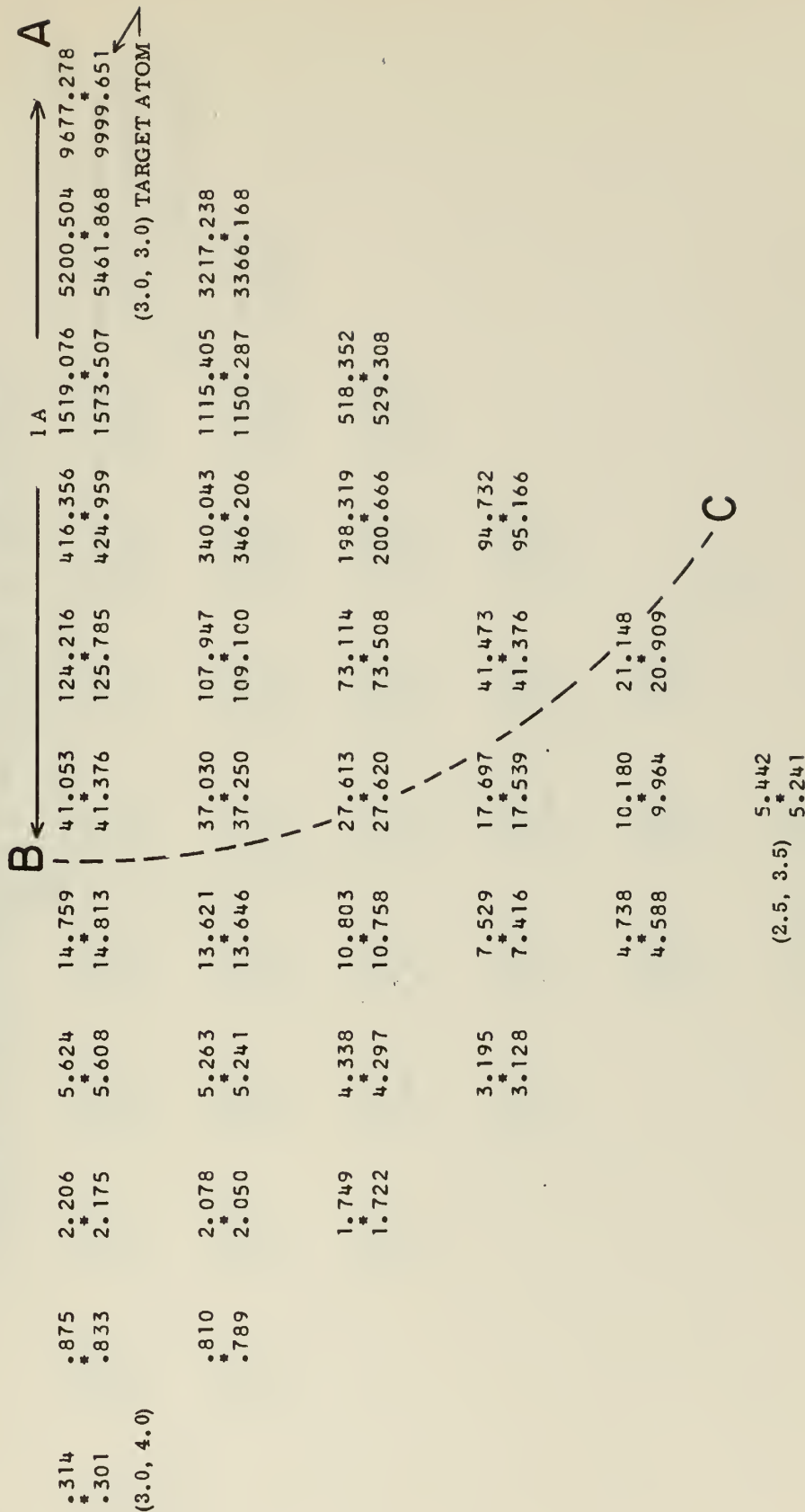


Fig. 21. OUTPUT TRIANGLE for 10 kev bullet; (*) indicates impact point

THOMAS-FERMI-FIRSOV POTENTIAL, ERODED, COPPER-COPPER

BULLET INCOMING ENERGY = 5000.0 EV TIME STEP MULTIPLIER = .090 BULLET INCIDENT ON (100) FACE

TOP NO. IS POTENTIAL ENERGY IN EV OF TARGET IN LATTICE

BOTTOM NO. IS POTENTIAL ENERGY IN EV OF BULLET IN LATTICE

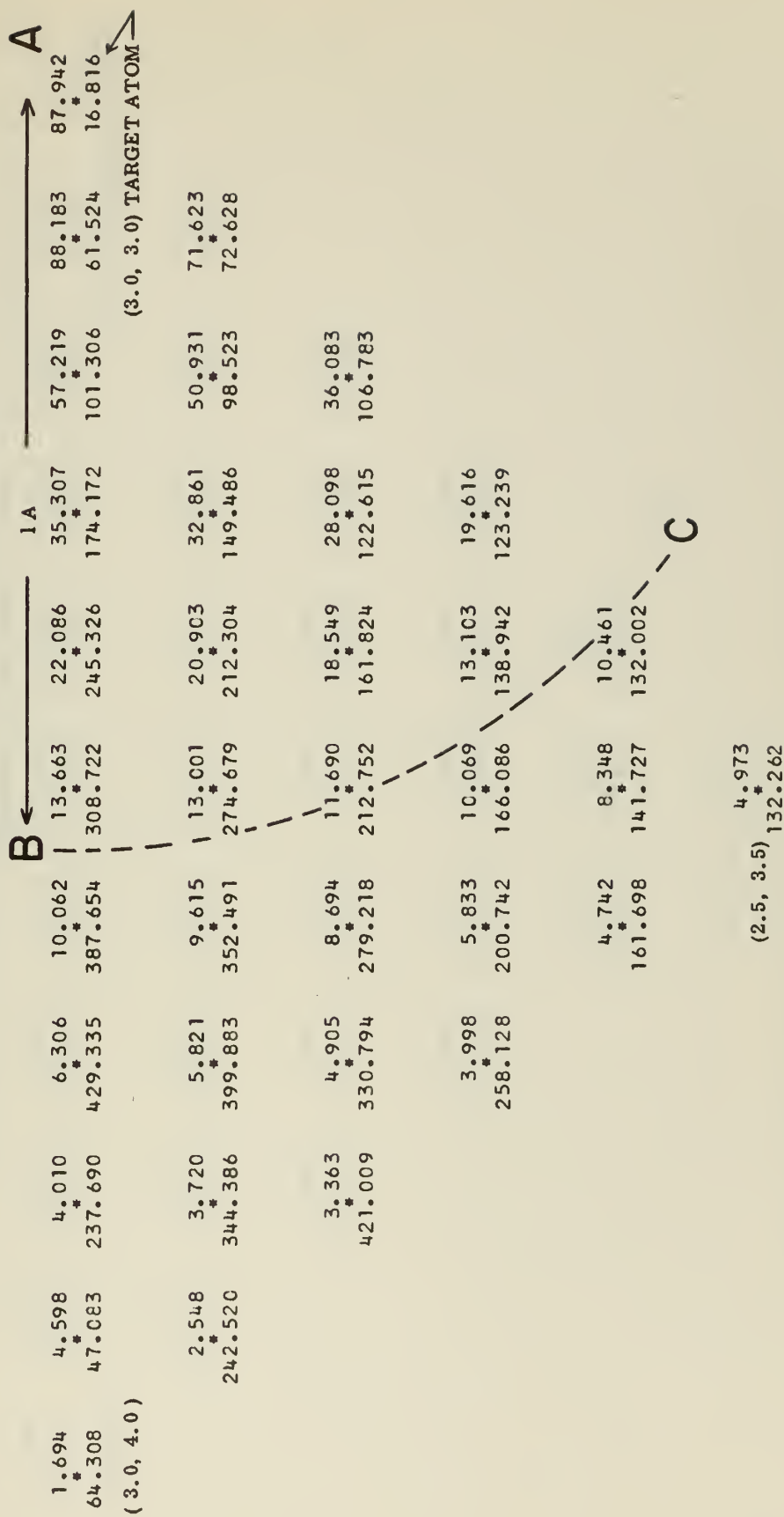


Fig. 23. OUTPUT TRIANGLE for 5 keV bullet; (*) indicates impact point

THOMAS-FERMI-FIRSOV POTENTIAL, ERODED, COPPER-COPPER

BULLET INCOMING ENERGY = 10000 EV TIME STEP MULTIPLIER = .090 BULLET INCIDENT ON (100) FACE

TOP NO. IS POTENTIAL ENERGY IN EV OF TARGET IN LATTICE

BOTTOM NO. IS POTENTIAL ENERGY IN EV OF BULLET IN LATTICE

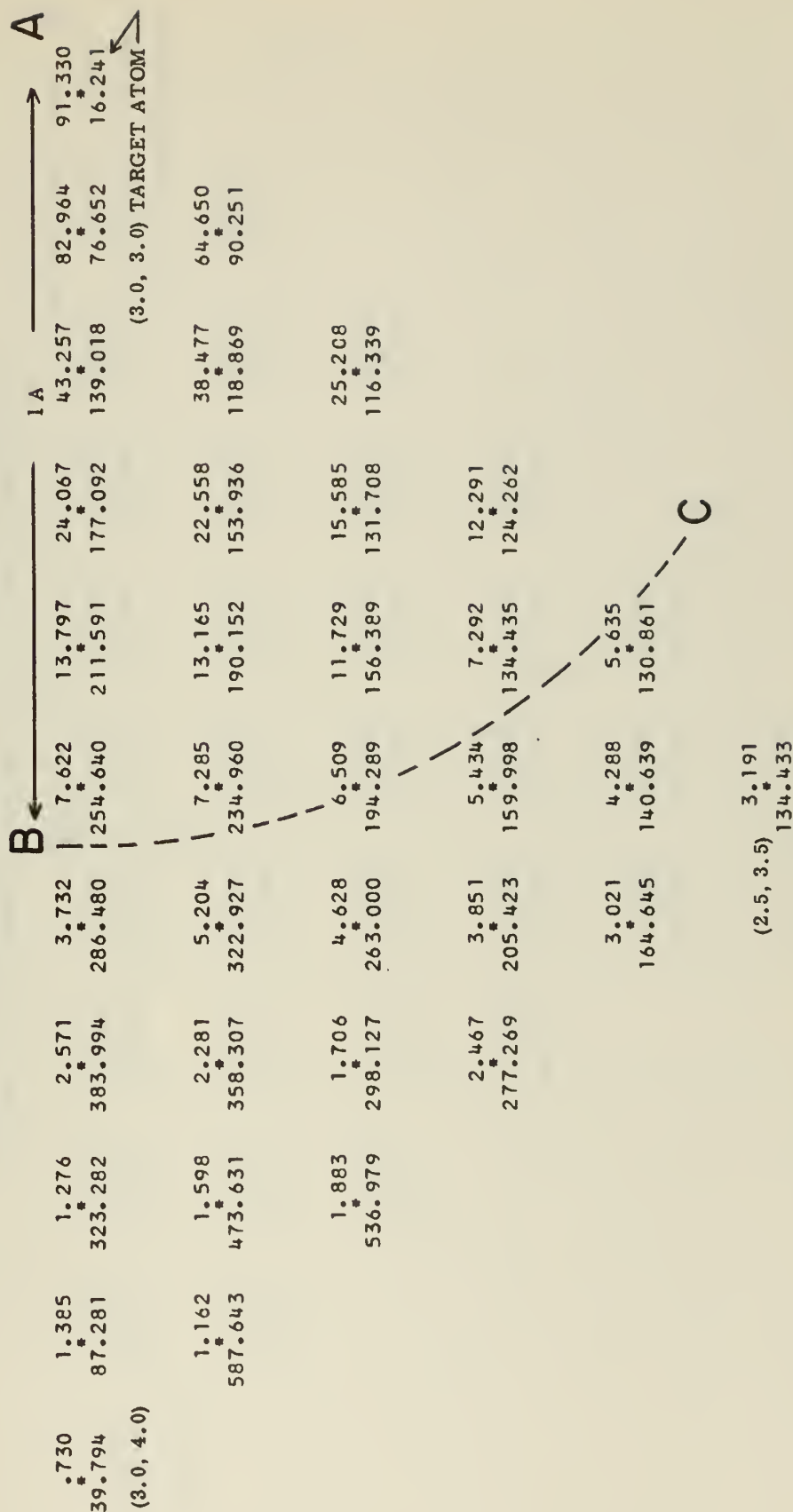


Fig. 24. OUTPUT TRIANGLE for 10 kev bullet; (*) indicates impact point

APPENDIX I

GENERAL DISCUSSION OF THE PROGRAM

The masses of the lattice atoms are equal to that of copper and are an integral part of the program. The bullet mass however, can be varied by changing the input data. The potentials and forces are calculated by function subroutines at the end of the program. The main program is not disrupted when a different potential function is substituted.

A coordinate system was established with an atom located at the origin. For a face centered cubic lattice in this coordinate system, the sum of the coordinates of an atom is always an even number. Fixed point numbers were used for lattice units, and the volume which contains all the atoms (movable and fixed) was scanned in the x, y, and z directions. If the sum of the coordinates of a point was even, an atom was placed there, otherwise a space was left. Truncation of quotients in fixed point arithmetic allows a simple test for odd or even numbers. The fixed point coordinates of the atoms were stored in memory as floating point numbers for later calculations. The "core" atoms were assigned numbers from 2-64 and the fixed atoms were numbered from 65 to 172.

In the calculation of the forces or potentials associated with an atom, the atoms closer than nearest neighbors in an undisturbed lattice must be considered. A "nested" Do loop was used for these calculations. The outside Do loop was indexed from 1 to 64 and the inner Do loops initial index was always one greater than the outer Do and ended at 172. This arrangement avoids the calculation of forces or potentials twice for the same pair of atoms. The x, y and z coordinate separations, in that order, between the atom under consideration (index of outer loop) and any other atom were tested. If any of these separations were found to be greater than ROE,

the inner Do index was advanced by one. If the coordinates passed these three tests, the square of the distance between atoms was compared to the square of ROE (ROE²). The true distance of any atom that passed this test was used in the force or potential function and the result applied to both atoms. (In potential calculations, half the potential was given to each atom). If the square of the distance between atoms was equal to or greater than ROE², the inner Do index was advanced by one. Most of the atoms will be eliminated in the first three tests. The procedure avoids time consuming square root calculations by the computer.

The double iteration method used to move atoms requires two calculations of resultant forces for each "time step". The same section of the program was used for both force calculations, but the equations required to move the atoms twice in the time step are not identical. The fixed point variable INDEX was used as a switching device to jump over the motion equations that locate the atoms at temporary positions. The jump is executed after the second calculation of forces.

The movement of atoms is continued until the kinetic energy of the target is a maximum. It is necessary to advance one time step beyond the maximum to establish the maximum. The particle locations and velocities at the time of the maximum are still in memory, and these are used in kinetic and potential energy calculations or as output.

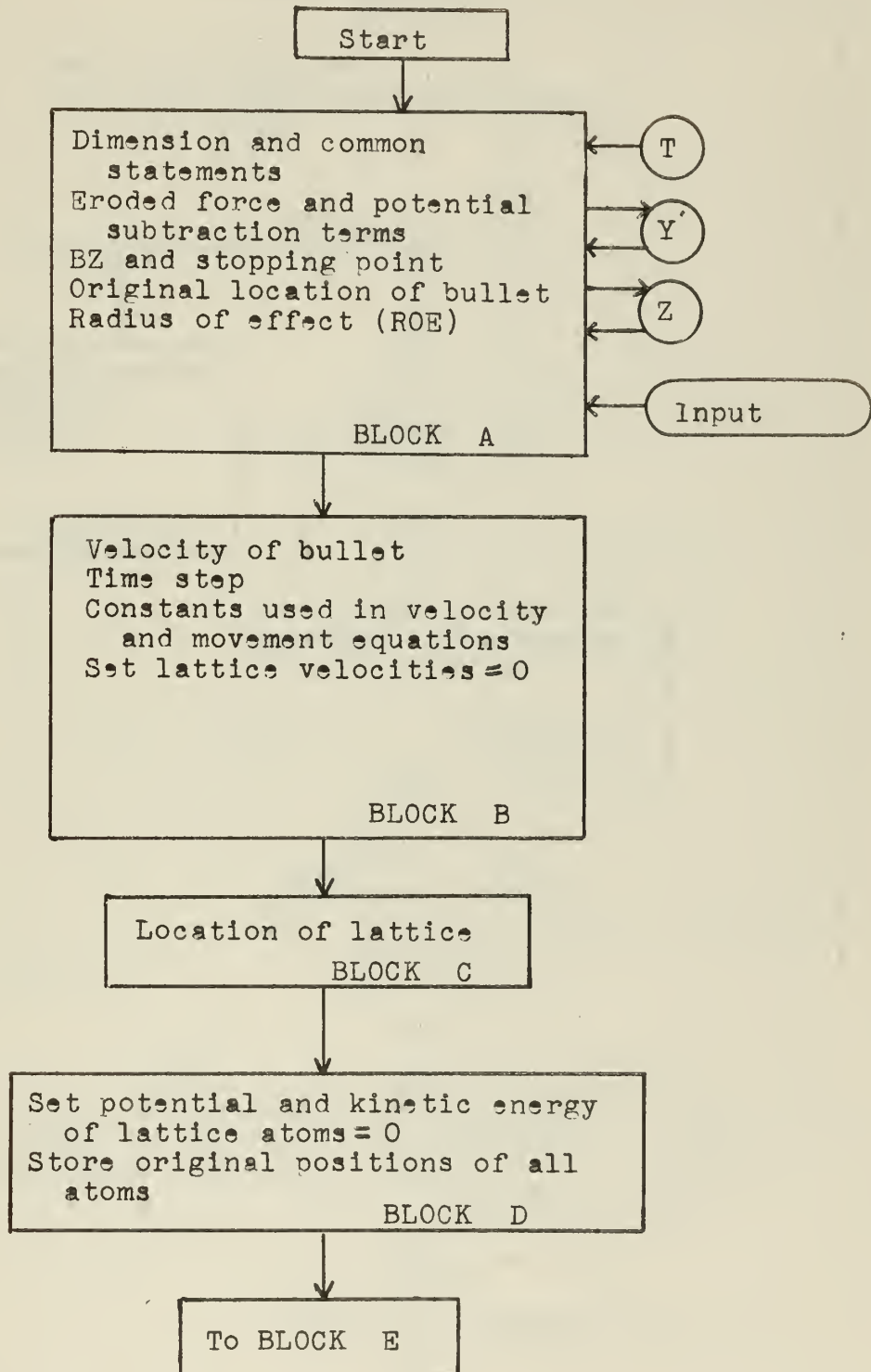
Although the two body problem is in two dimensions, the impact parameter is the same as that in the lattice. The two body problem is solved in the same manner as the lattice except that the interaction is stopped when the atoms are separated by a distance equal to ROE which means the potential and force between atoms is zero. The interaction is initially started with the separation of atoms slightly less than ROE. The ΔT for

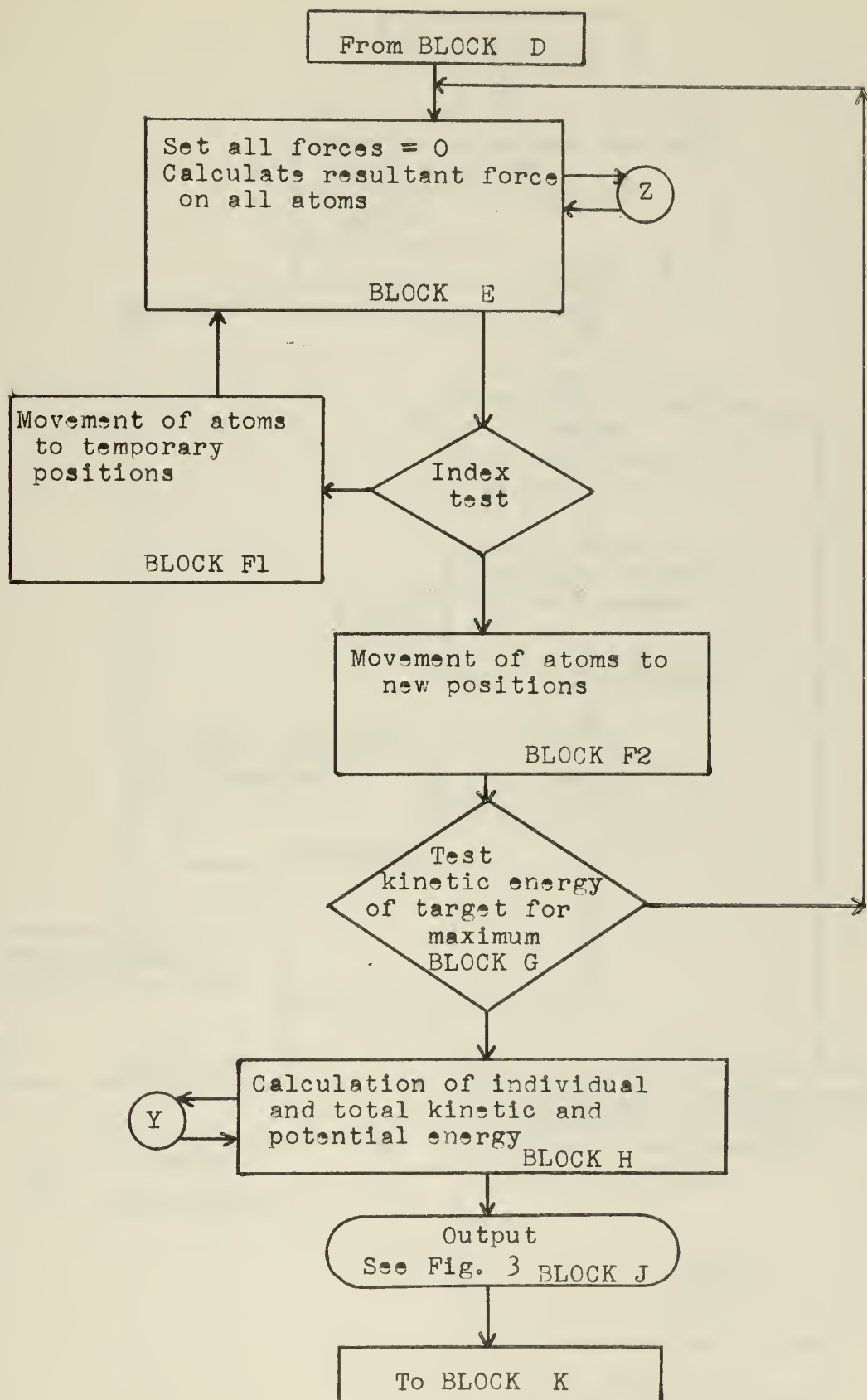
the two body problem is much smaller than the one used for lattice calculations. The small ΔT consistently produces very much more accurate results in the two body problem than in the lattice.

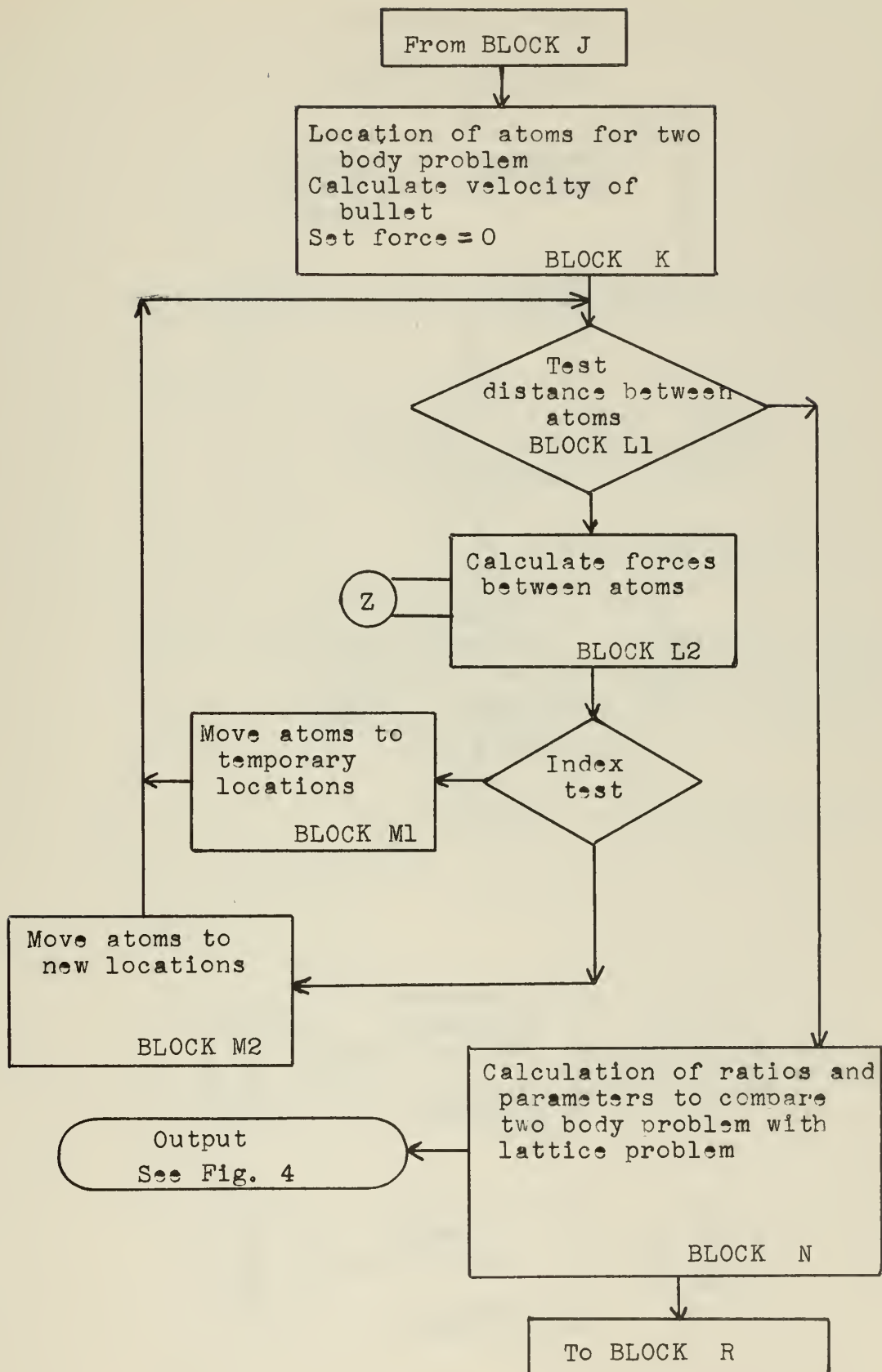
The input consists of the bullet energy and mass, its x coordinate, and the time step multiplier. In order to calculate data for all 36 impact points, six input cards are required; each is identical except for the x coordinate of the bullet. The x coordinates assigned are 3.0, 2.9, 2.8, 2.7, 2.6, and 2.5 in that order. After receiving the input from one card, the corresponding z coordinate on the line AB is calculated (see Fig. 6). The computer proceeds through the program for this point, then subtracts 0.1 lattice units from the z coordinate and repeats. It proceeds down a column until all calculations for impact points on the line BC have been made. The computer pauses at this point to await instructions. Here it may receive a new data card and start another column. After the calculations for the point on the apex of the triangle ($x=2.5$, $z=3.5$) are complete, the computer can be directed to print the output in the triangular form shown by Fig. 6.

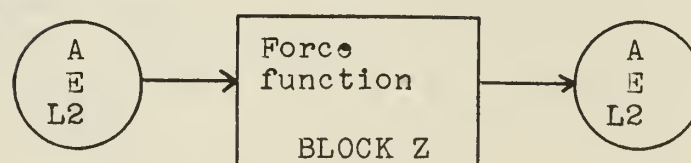
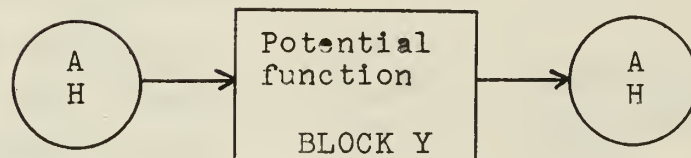
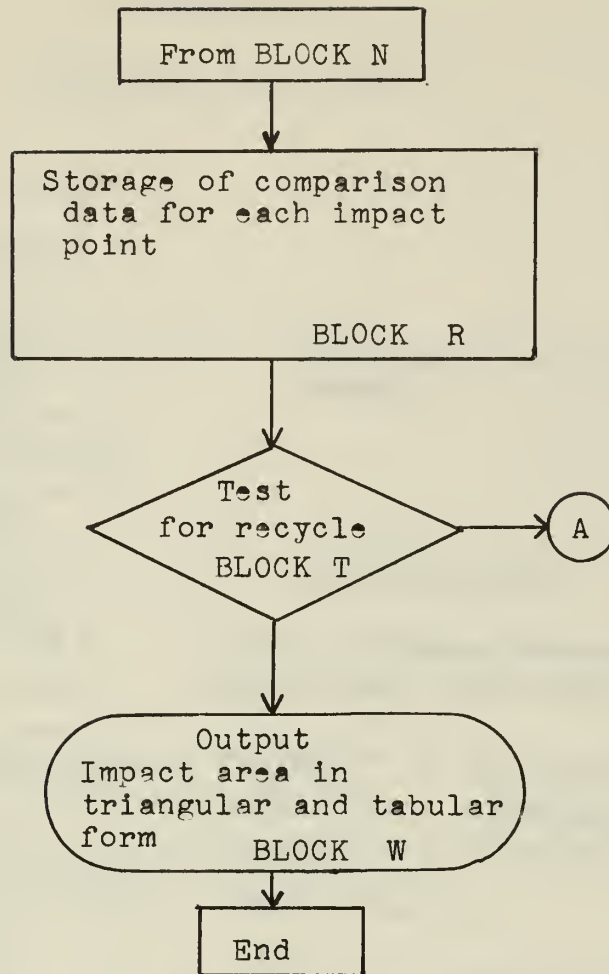
APPENDIX II

BLOCK DIAGRAM





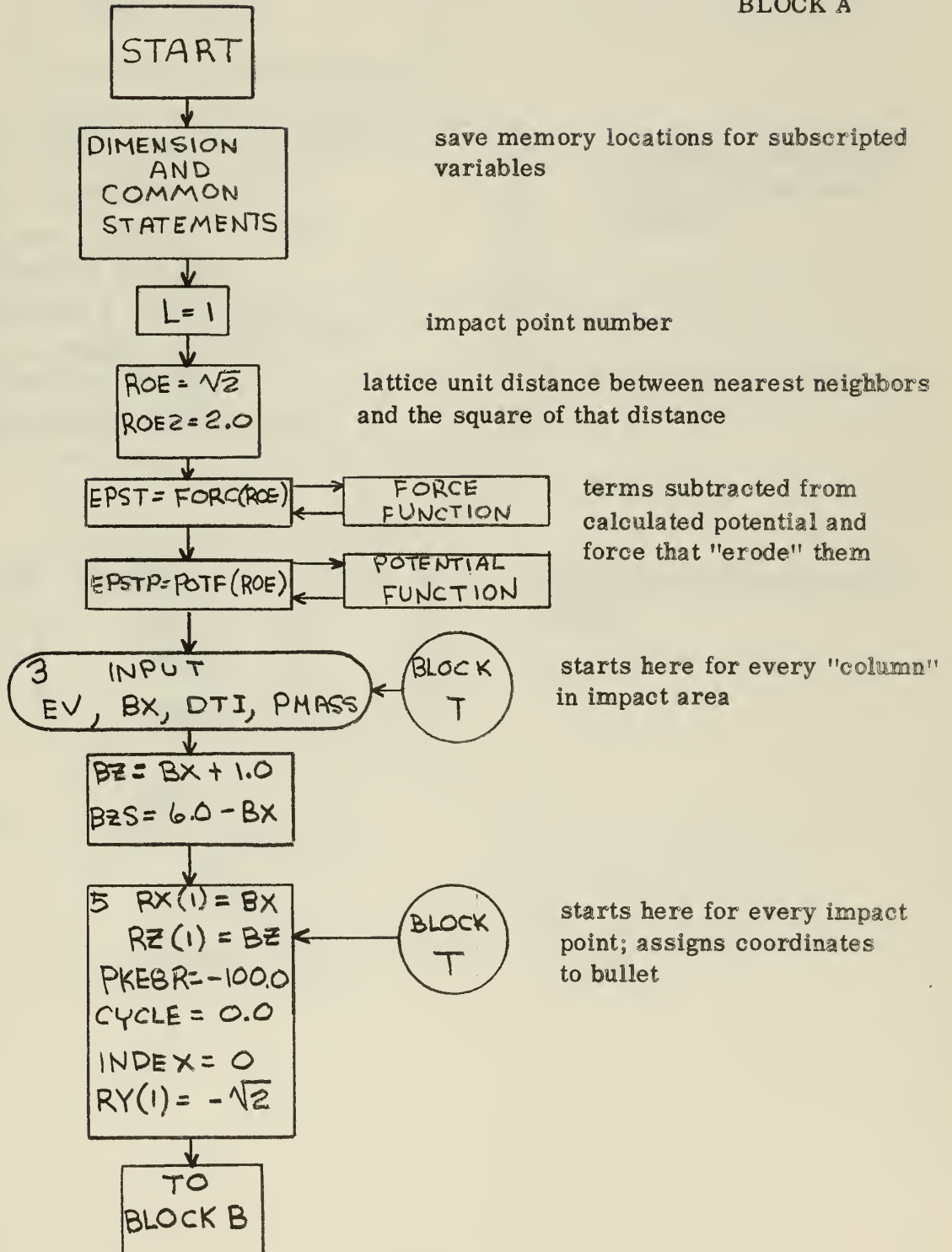




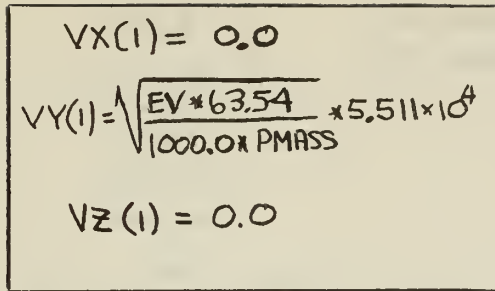
APPENDIX III

DETAILED FLOW CHART

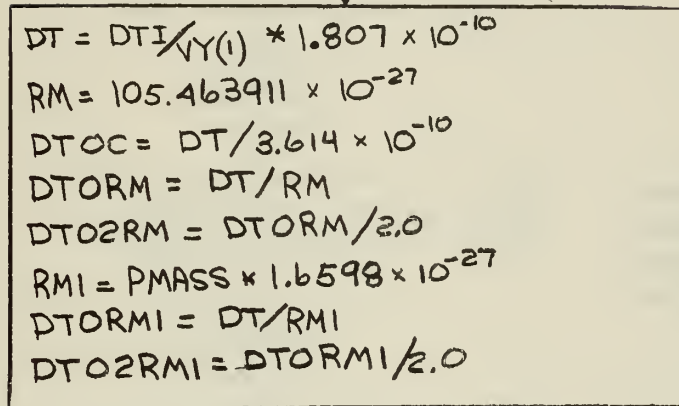
BLOCK A



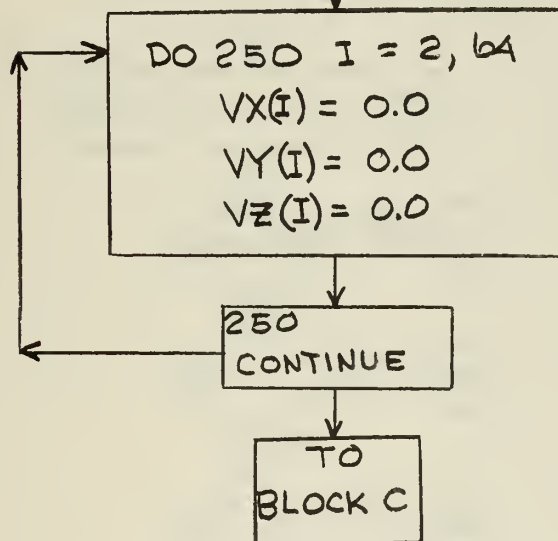
BLOCK B



velocity components of bullet



calculates constants used in force and movement equations



sets velocity of lattice atoms to zero

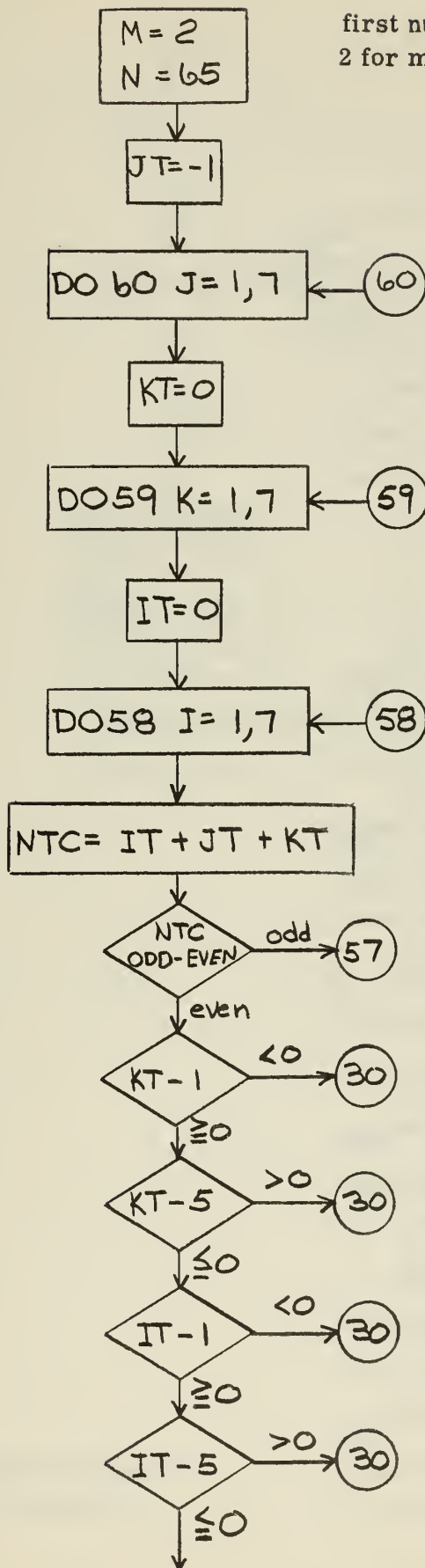
BLOCK C

first numbers to be assigned to lattice atoms
2 for movable core, 65 for fixed atoms

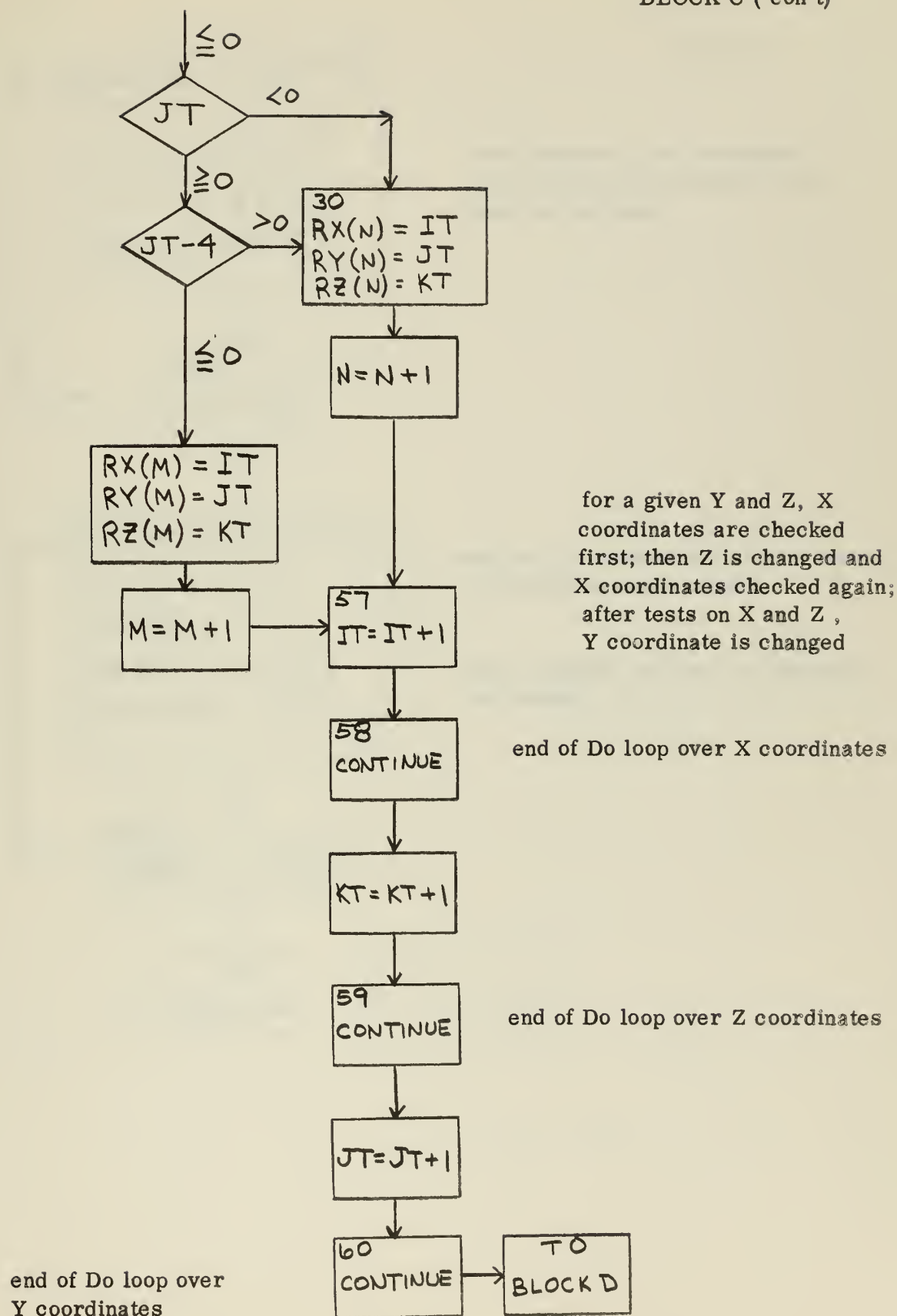
IT, JT, KT are X, Y, Z coordinates
in the system used for reference

each point with integer values for
coordinates is tested; if the sum of
the coordinates is an even number,
an atom is assigned to that position

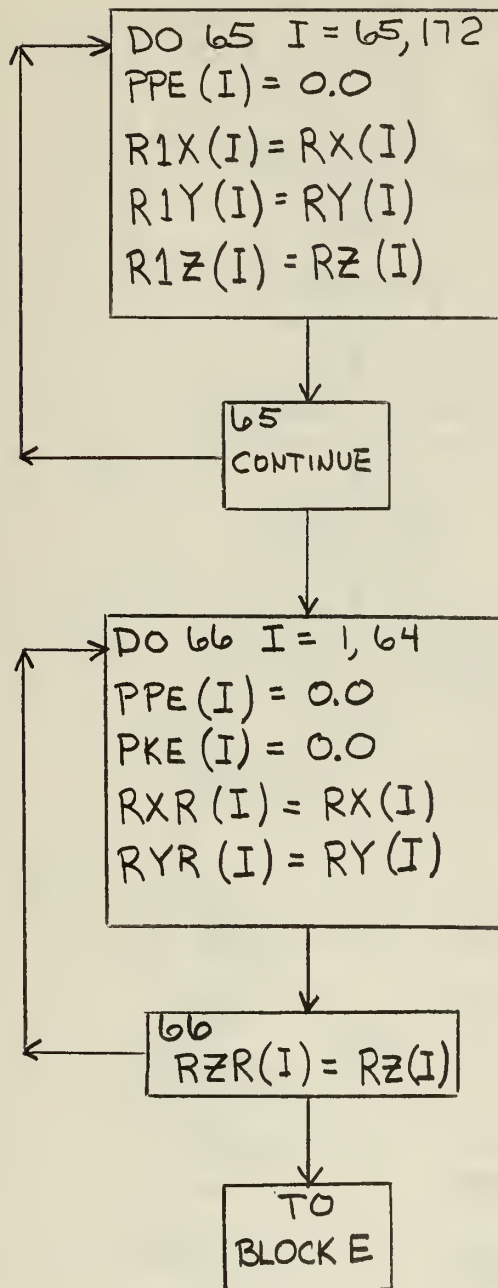
these decision processes
determine if an atom is in the
"core"; this is necessary in
order to have core atoms with
consecutive numbers 2-64



BLOCK C (con't)



BLOCK D

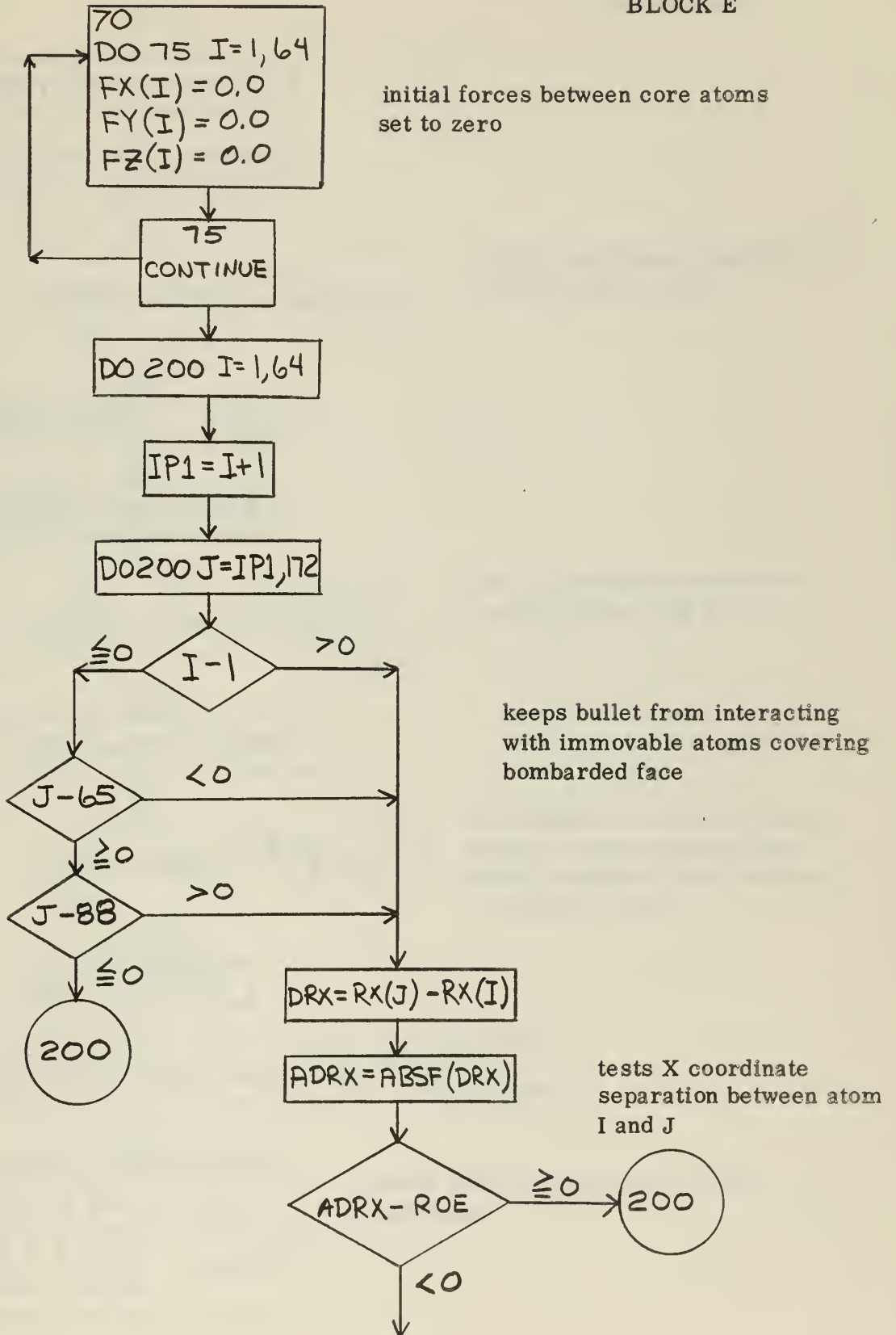


sets temporary and permanent
coordinates of immovable atoms
equal to each other

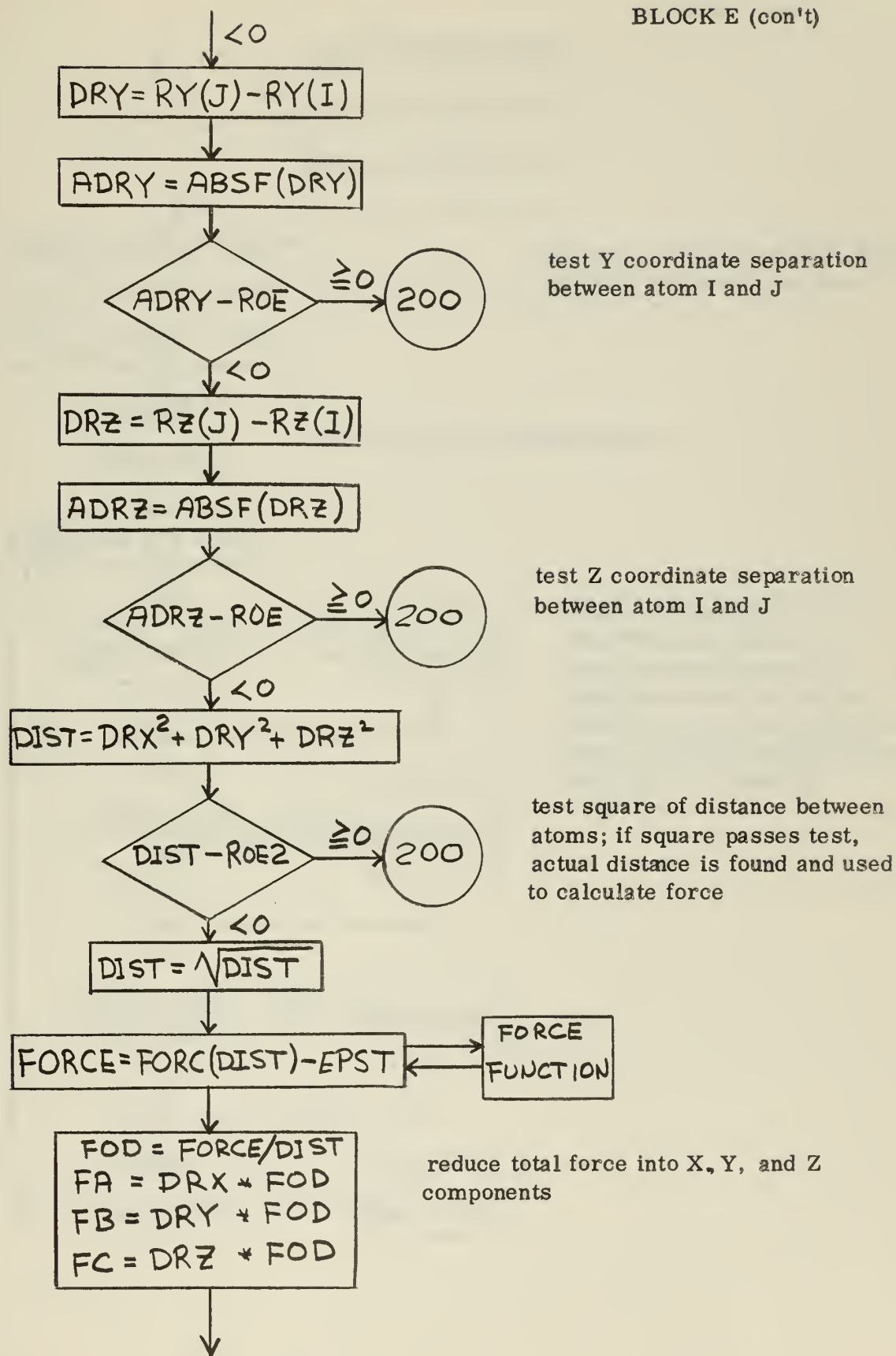
sets initial potential and kinetic energy
of core atoms equal to zero

stores original positions of bullet and
core atoms

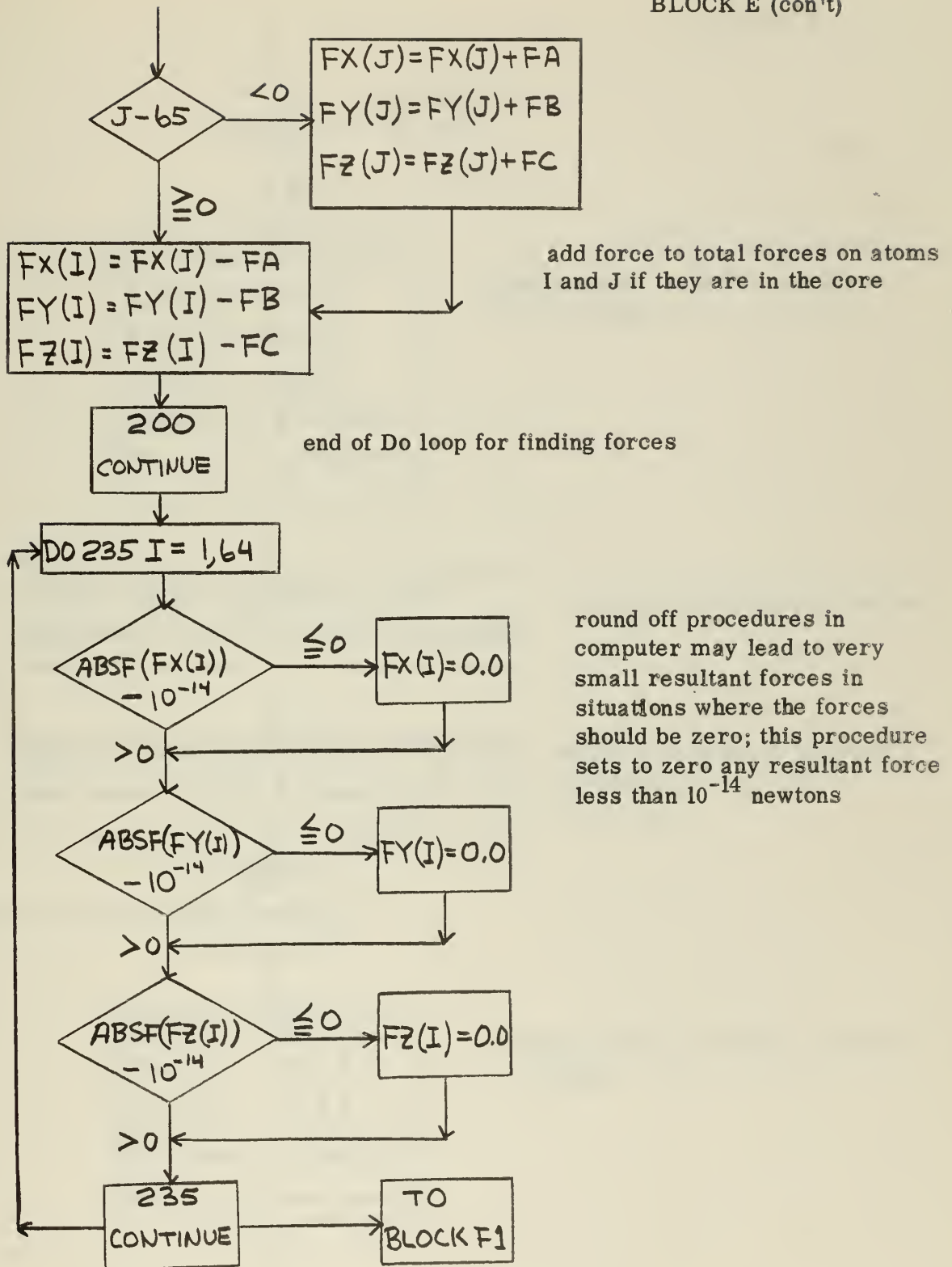
BLOCK E



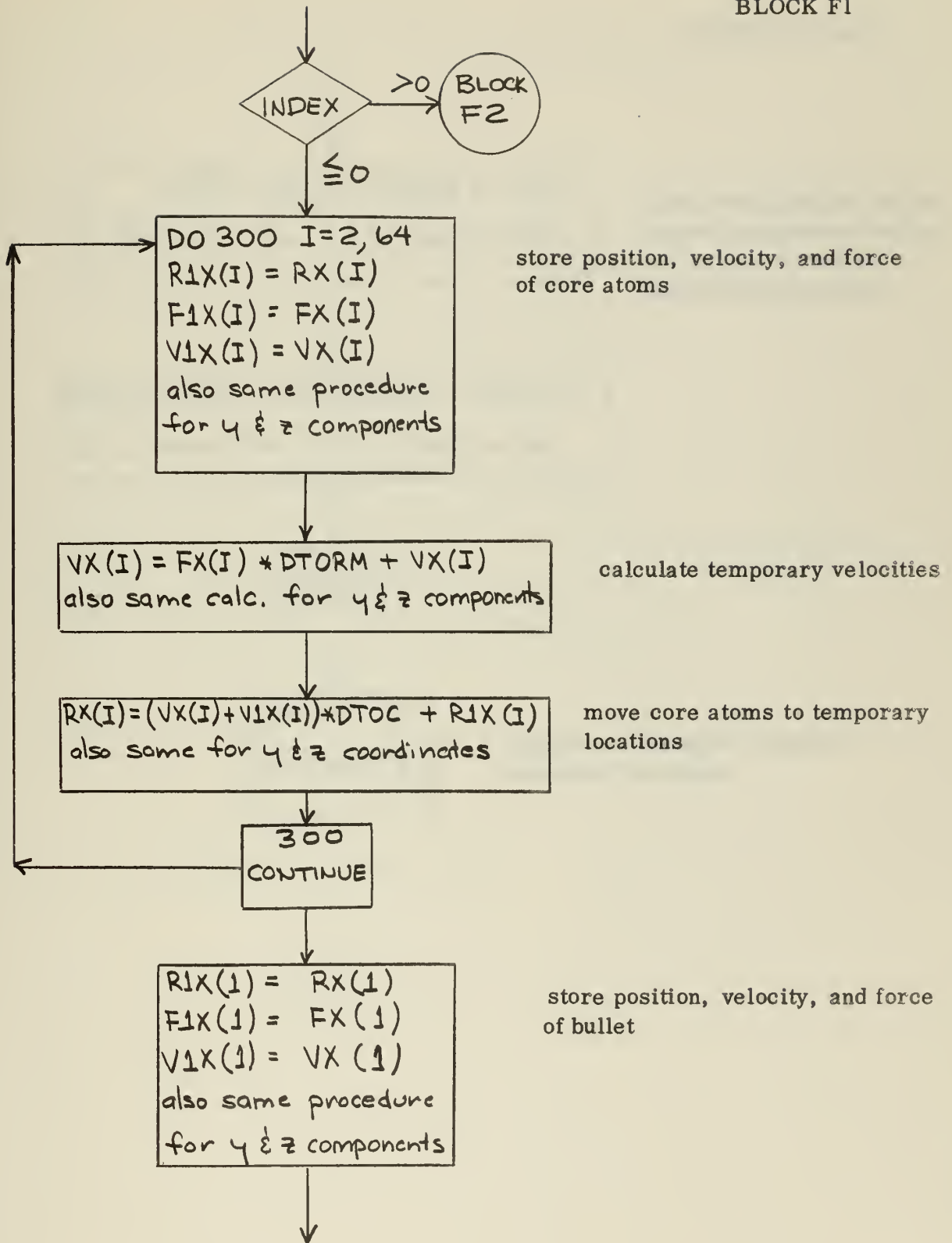
BLOCK E (con't)



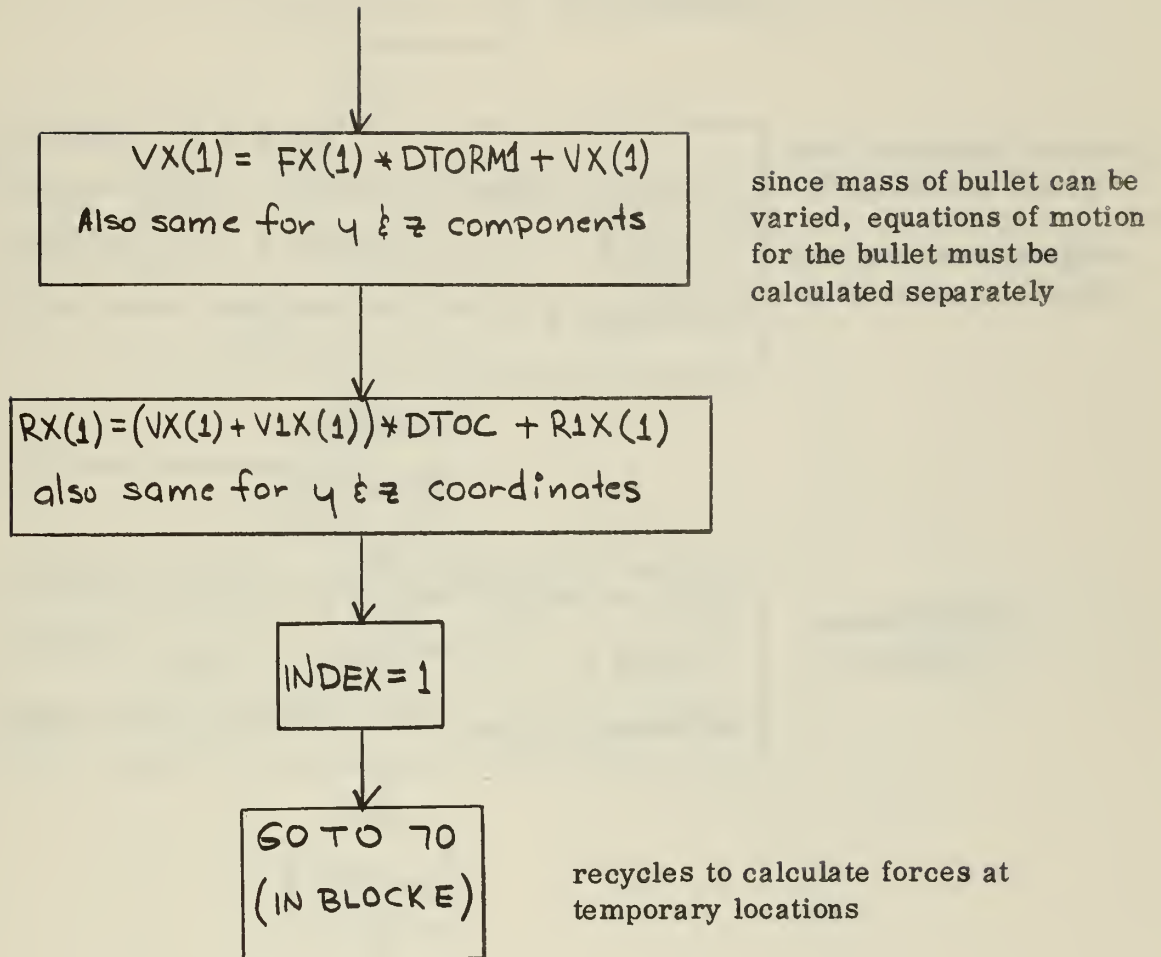
BLOCK E (con't)



BLOCK F1



BLOCK F1 (con't)



BLOCK F2

FROM INDEX check
in
BLOCK F1

forces at temporary locations have
been calculated

new velocities calculated
using average forces

new positions assigned
using average velocities

```
DO 425 I = 2, 64
VX(I) = [FX(I) + F1X(I)] * DT02RM + V1X(I)
RX(I) = [VX(I) + V1X(I)] * DTOC + R1X(I)
also same calculations for y & z components
```

425
CONTINUE

separate bullet
calculations

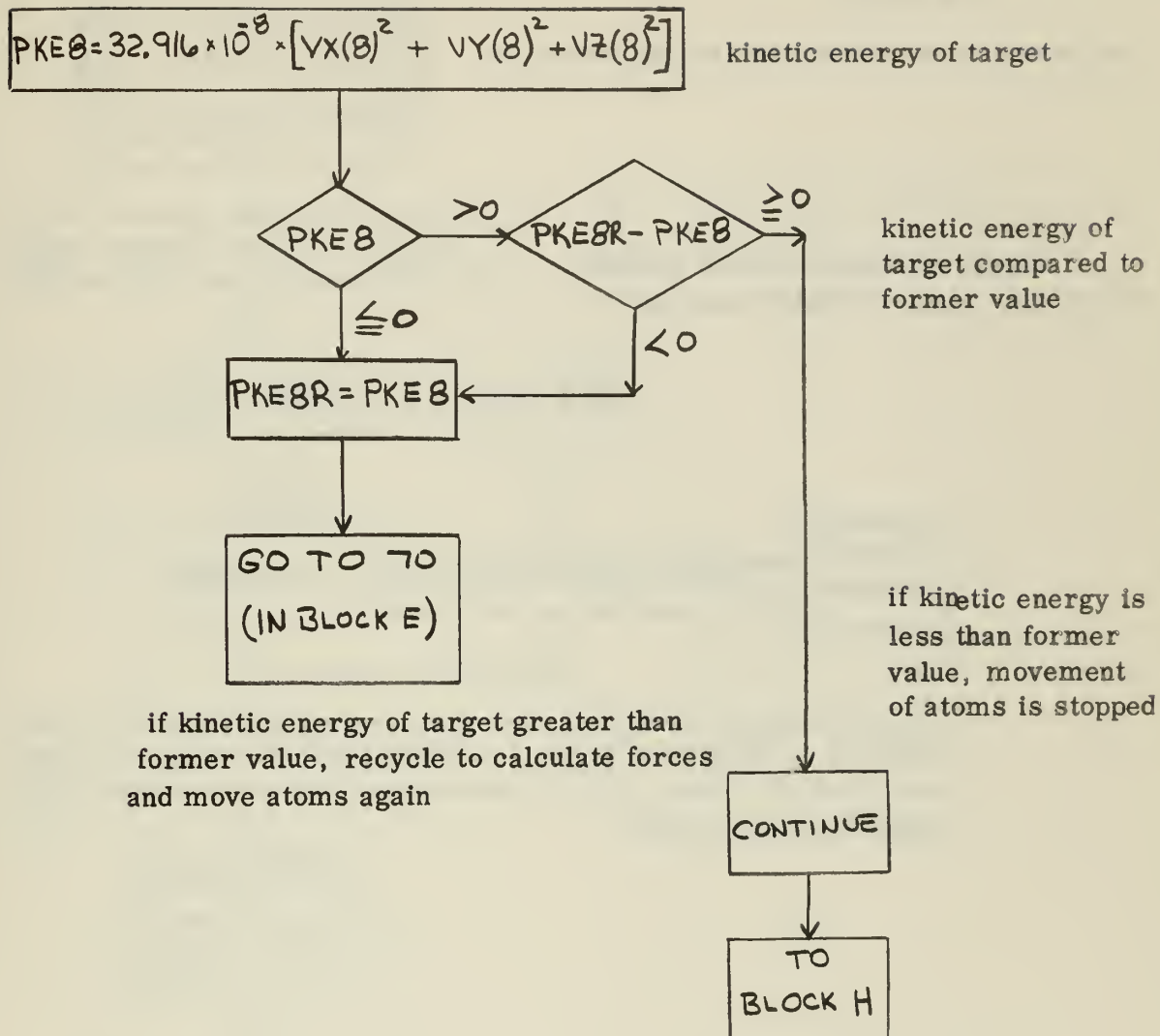
```
VX(1) = [FX(1) + F1X(1)] * DT02RM1 + V1X(1)
RX(1) = [VX(1) + V1X(1)] * DTOC + R1X(1)
also same calculations for y & z components
```

INDEX=0

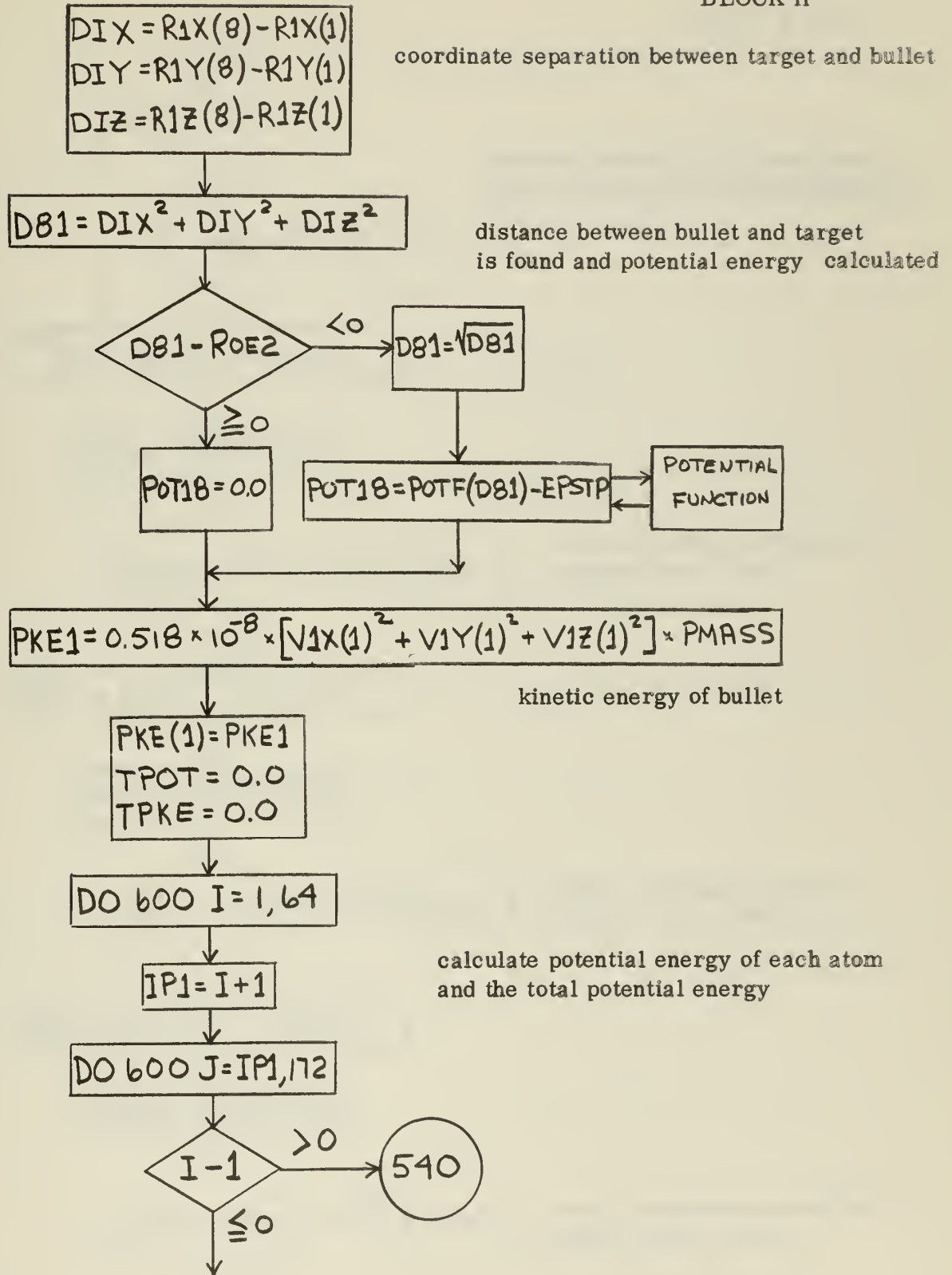
CYCLE = CYCLE + 1.0

TO
BLOCK G

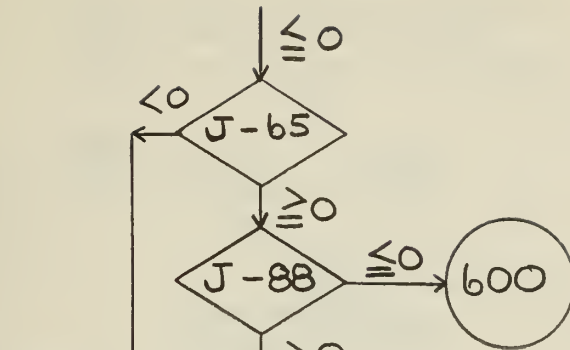
BLOCK G



BLOCK H



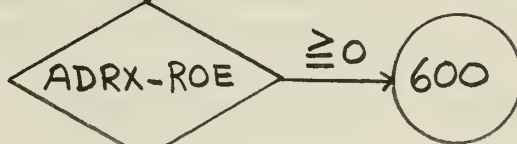
BLOCK H (con't)



prevents interaction of bullet with
immovable atoms covering bombarded
face

540
DRX = R1X(J) - R1X(I)

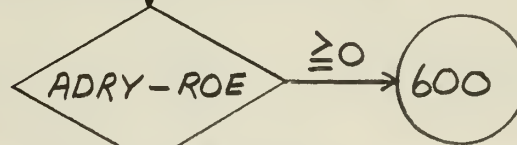
ADRX = ABSF(DRX)



test X coordinate separation
between atom I and J

DRY = R1Y(J) - R1Y(I)

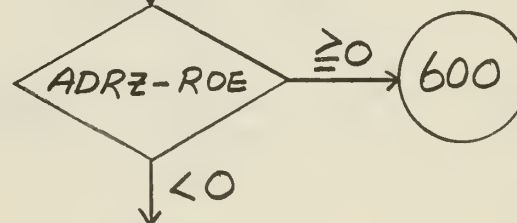
ADRY = ABSF(DRY)



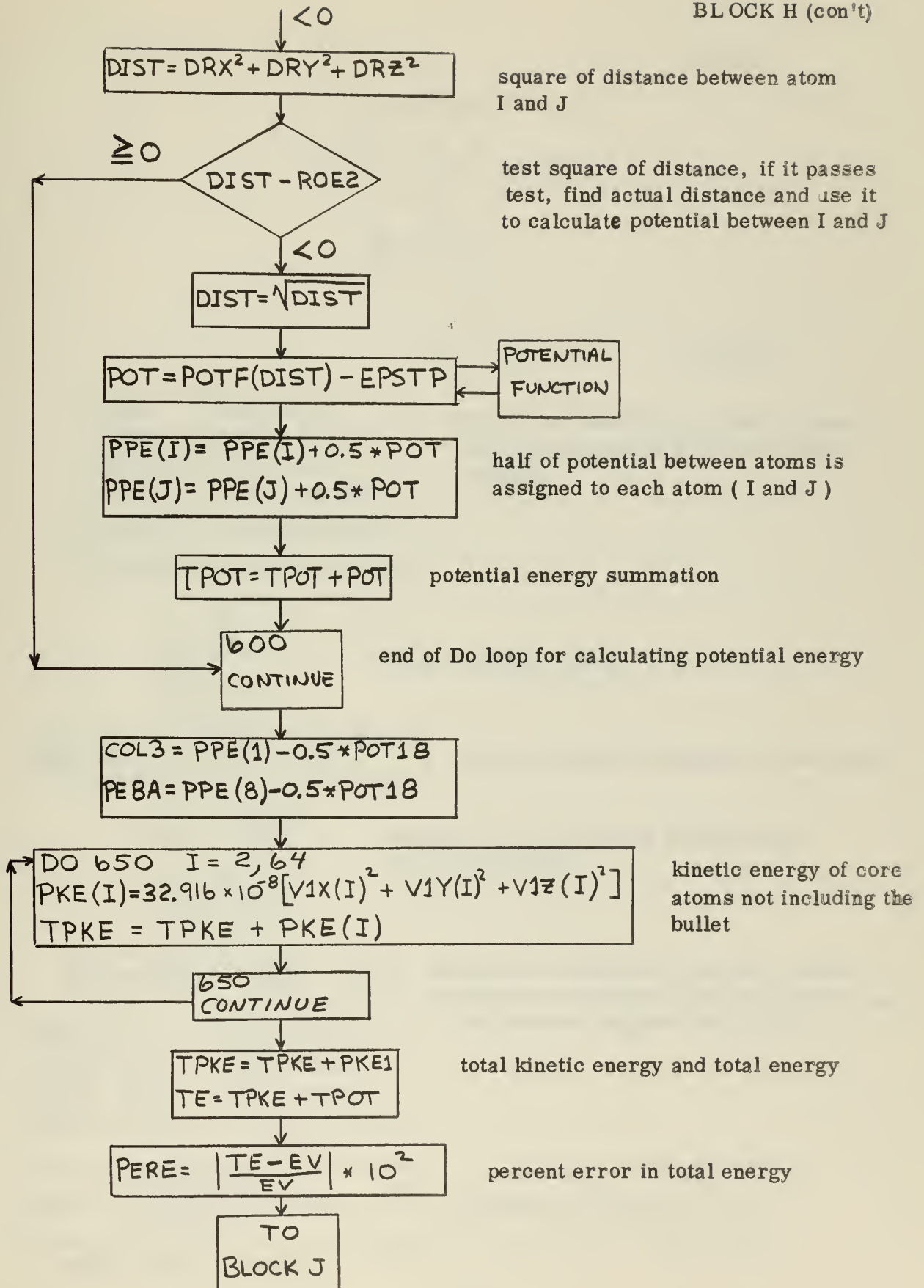
test Y coordinate separation
between atom I and J

DRZ = R1Z(J) - R1Z(I)

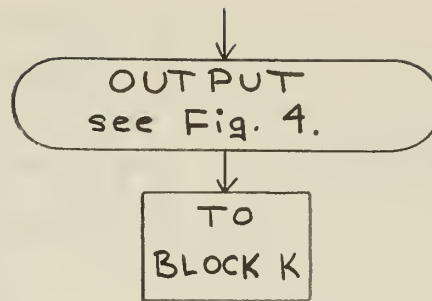
ADRZ = ABSF(DRZ)



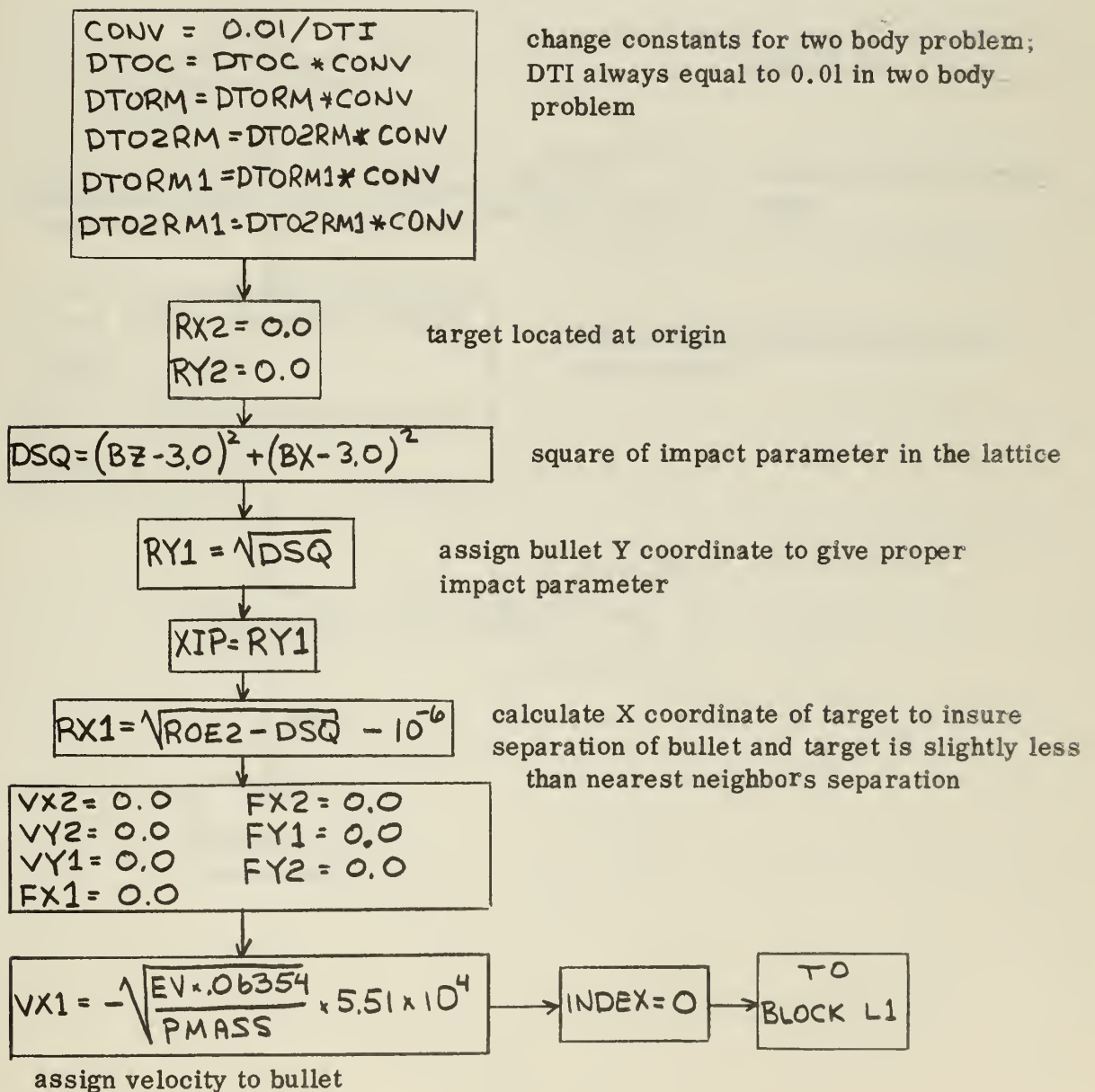
test Z coordinate separation
between atom I and J



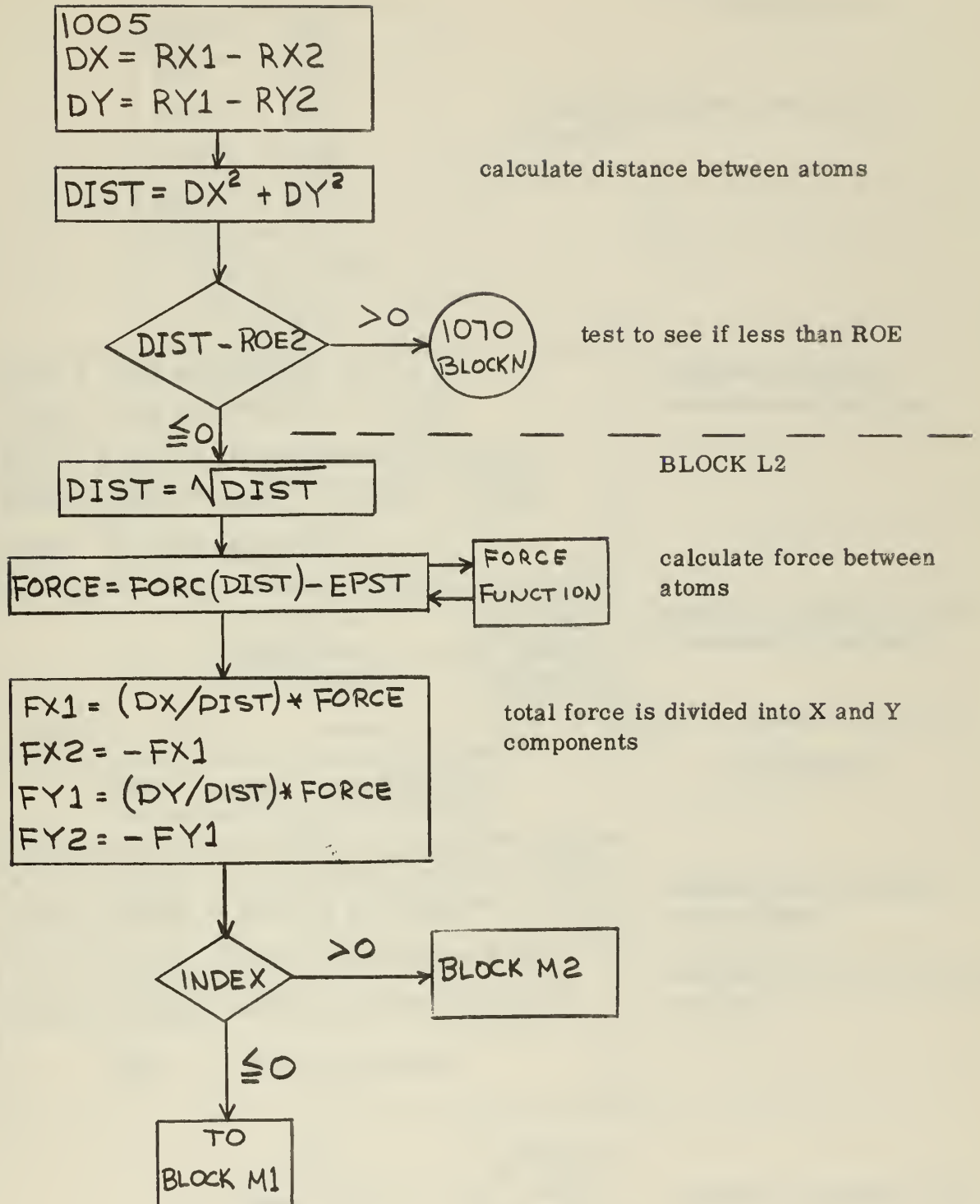
BLOCK J



BLOCK K



BLOCK L1



BLOCK M1

$RX1T = RX1$
 $RX2T = RX2$
 $FX1T = FX1$
 $VX1T = VX1$
 $VX2T = VX2$
 also same for y
 components

store old positions, forces, and velocities

$VX1 = FX1 * DTORM1 + VX1$
 $VX2 = FX2 * DTORM + VX2$
 $RX1 = [VX1 + VX1T] * DTOC + RX1T$
 $RX2 = [VX2 + VX2T] * DTOC + RX2T$
 also Y components

calculate temporary velocities and positions

INDEX = 1 → GO TO 1005 BLOCK L1

recycles to calculate forces at temporary positions

FROM INDEX TEST - BLOCK L2

BLOCK M2

$VX1 = (FX1 + FX1T) * DTO2RM1 + VX1T$
 $VX2 = (FX2 + FX2T) * DTO2RM + VX2T$
 $RX1 = (VX1 + VX1T) * DTOC + RX1T$
 $RX2 = (VX2 + VX2T) * DTOC + RX2T$
 also Y components

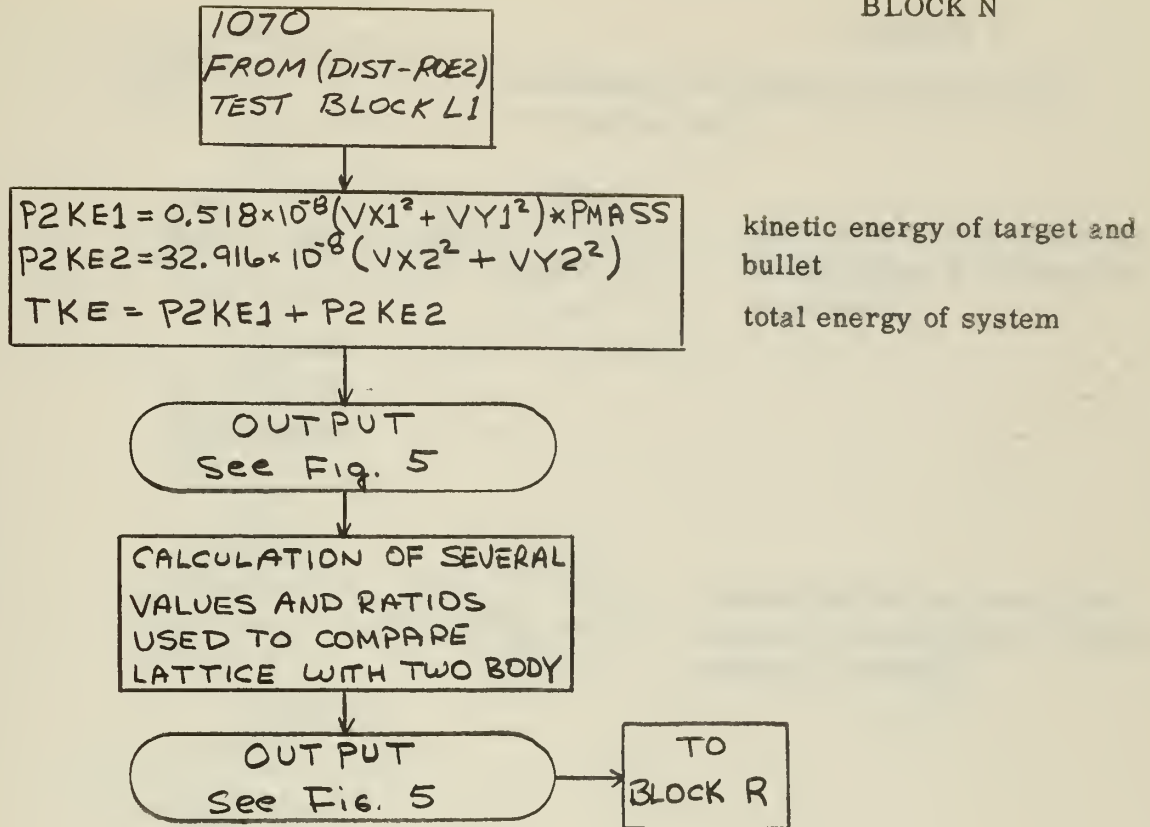
calculate final velocities and positions

average forces and velocities are used

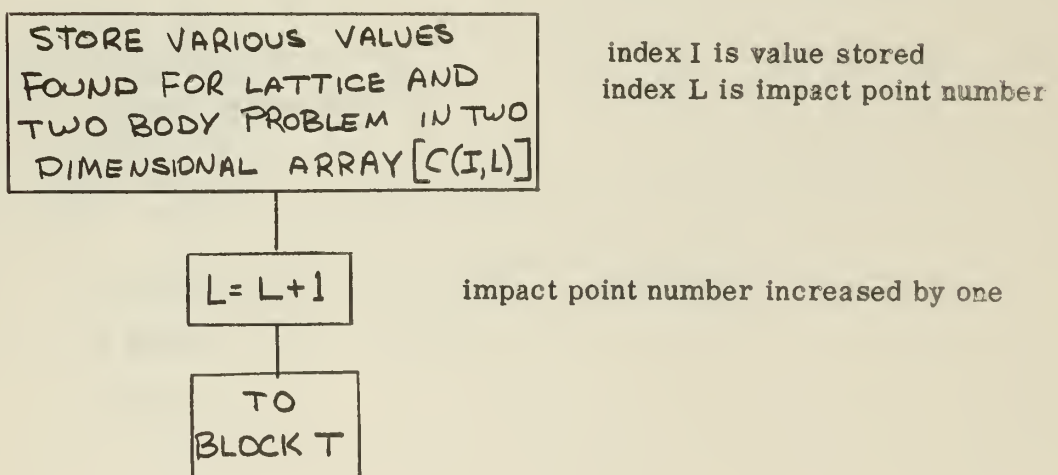
INDEX = 0 → GO TO 1005 BLOCK L1

recycles to distance test

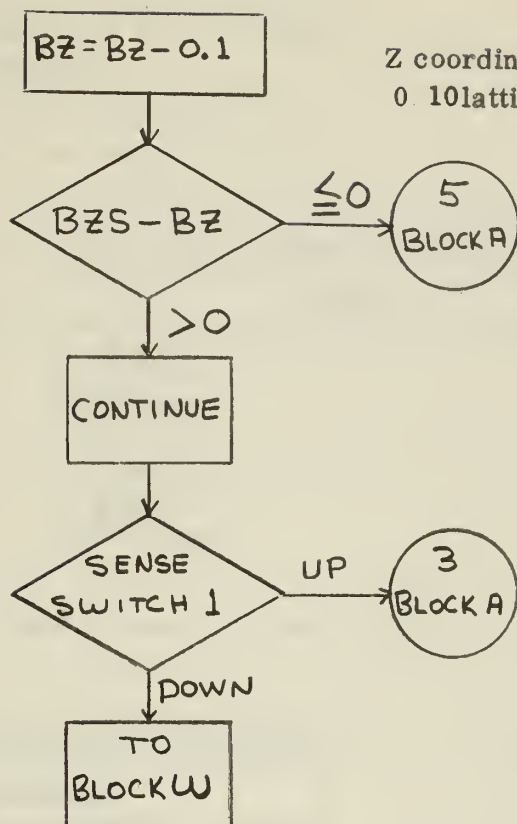
BLOCK N



BLOCK R



BLOCK T

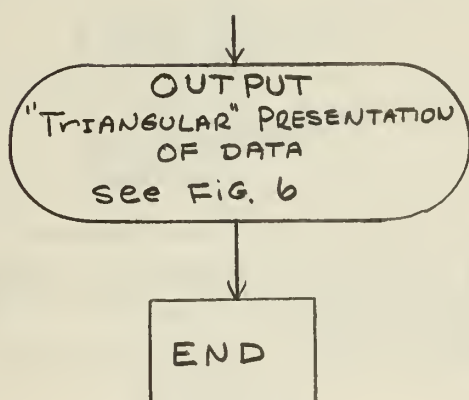


Z coordinate of impact point lowered by 0.10 lattice units

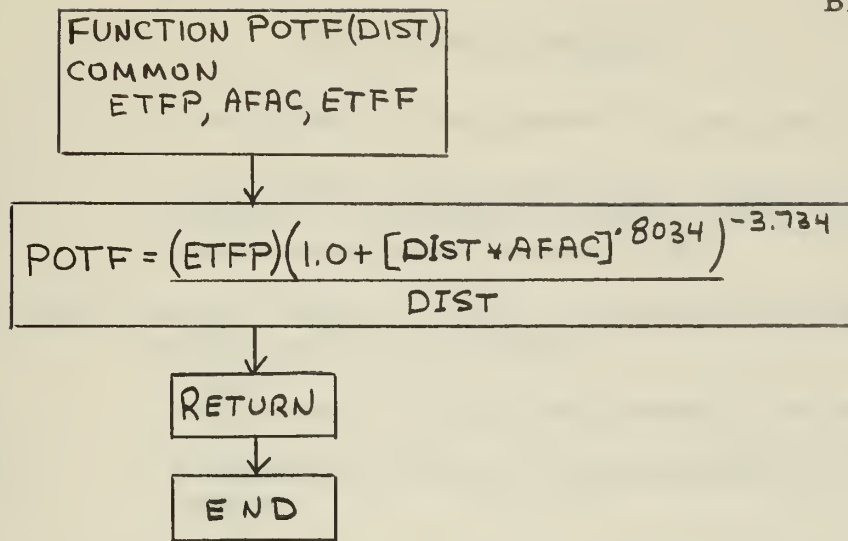
Z coordinate of impact point tested to see if it is the end of a column

if sense switch one is up, new data card is read to start another column in "triangle"

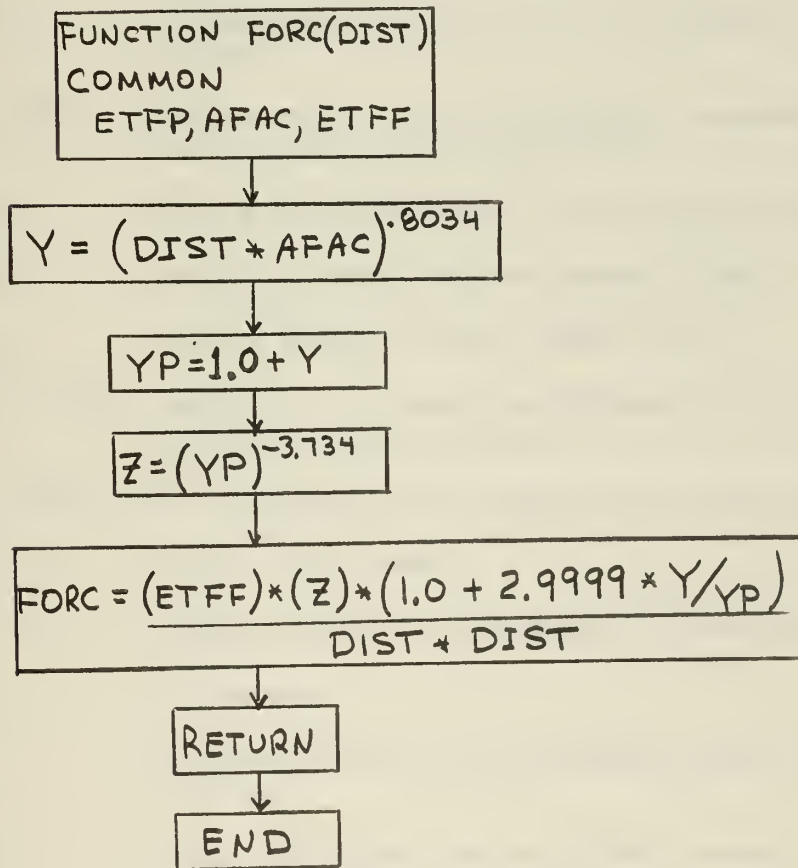
BLOCK W



BLOCK Y



BLOCK Z



APPENDIX IV

DEFINITION OF VARIABLES

RX(I)	x coordinate of atom number I at any time (lattice units)
R1X(I)	x coordinate of atom number I at former position (same as position 1 in procedure section), stored while average forces are computed (lattice units)
RY(I), RZ(I), R1Y(I), R1Z(I)	y and z coordinates defined in the same manner as RX(I) and R1X(I)
FX(I)	x component of force on atom number I at any time (newtons)
F1X(I)	x component of force at former position of atom number I (newtons)
FY(I), FZ(I), F1Y(I), F1Z(I)	y and z components of force defined in the same manner as FX(I) and F1X(I)
VX(I)	x component of velocity at any time of atom number I (m/sec)
V1X(I)	x component of velocity of atom number I at former position (m/sec)
VY(I), VZ(I), V1Y(I), V1Z(I)	y and z velocity components defined in the same manner as VX(I) and V1X(I)
C(I,L)	matrix for storing data I from impact point L
PPE(I)	potential energy of atom number I (ev)
PKE(I)	kinetic energy of atom number I (ev)
RXR(I), RYR(I), RZR(I)	x, y, and z coordinates of all atoms in initial positions (lattice units)
ETFP, AFAC, ETFF	constants used in Thomas-Fermi-Firsov potential
L	impact point number
ROE	nearest neighbors separation (lattice units)
ROE2	ROE squared
EPST	value subtracted from calculated force to erode the force (newtons)
EPSTP	value subtracted from calculated potential to erode the potential (ev)

FORC	function subroutine that calculates force; argument is the separation of atoms (newtons)
POTF	function subroutine that calculates potential; argument is the separation of atoms (ev)
EV	original kinetic energy of bullet (ev)
BX	original x coordinate of bullet (lattice units)
DTI	time step multiplier
PMASS	mass of the bullet (amu)
BZ	original z coordinate of bullet (lattice units)
BZS	z coordinate of bullet at end of column in impact triangle, used as stopping point (lattice units)
PKE8R	kinetic energy of target at the end of previous time step (ev)
CYCLE	number of time steps
INDEX	switching device
DT	length of basic time step (seconds)
RM	mass of copper atom (kg)
RM1	mass of bullet atom (kg)
DTOC	a constant used in movement equations; when multiplied by the sum of two velocities, it computes the average, multiplies it by DT and gives the result in lattice units (sec x lattice units/m)
DTORM	a constant used in velocity equations; it combines the quotient DT/RM into one constant to avoid repeated divisions (sec/kg)
DTO2RM	same as DTORM except it averages the force
DTORM1	same as DTORM but used in bullet calculations only, since the bullet mass may be different from that of the lattice atoms
DTO2RM1	same as DTO2RM but used in bullet calculations only since bullet mass may be different from that of the lattice atoms
M,N	atom numbers
JT	y coordinate of an x-z plane

KT	z coordinate of an x-y plane
IT	x coordinate of a y-z plane
NTC	sum of x,y, and z coordinates of a point in the lattice
DRX,DRY,DRZ	difference between x,y, and z coordinates respectively of any two atoms (lattice units)
ADRX, ADRY, ADRZ	absolute value of DRX, DRY, and DRZ
DIST	distance between two atoms or the square of the distance (lattice units)
FORCE	eroded force between two atoms (newtons)
FOD	FORCE divided by the distance between two atoms (newtons/lattice unit)
FA,FB,FC	x,y, and z components respectively of the force between two atoms (newtons)
PKE8	kinetic energy of target atom in the lattice (ev)
DIX,DIY,DIZ	x,y, and z coordinate separation between target and bullet (lattice units)
D81	distance between target and bullet in lattice or the square of the distance (lattice units)
POT18	potential between target and bullet (ev)
PKE1	kinetic energy of bullet at end of time step (ev)
TPOT	total potential energy of lattice (ev)
TPKE	total kinetic energy of lattice (ev)
POT	potential between any two atoms (ev)
COL3	potential of the bullet with respect to the lattice minus the target (ev)
PE8A	potential energy of the target with respect to the lattice minus the bullet (ev)
TE	total energy of the lattice (ev)
PERE	percent error in total energy (percent)
CONV	time step multiplier conversion factor for two body problem; makes two body time step multiplier equal to 0.01 regardless of the value used for the lattice problem

DSQ	square of the impact parameter in the two body problem (lattice units squared)
XIP	impact parameter in the two body problem (lattice units)
RX2,RY2	x and y coordinates of target at any time in two body problem (lattice units)
RX2T,RY2T	x and y coordinates of target at former position, stored while average force is computed (lattice units)
RX1,RY1,RX1T,RY1T	x and y coordinates of the bullet in the two body problem defined in the same manner as RX2,RY2,RX2T, and RY2T
VX2,VY2	x and y components of target velocity in two body problem (m/sec)
VX2T,VY2T	x and y components of target velocity in two body problem at former position (m/sec)
VX1,VY1,VX1T,VY1T	x and y components of velocity of the bullet in the two body problem defined in the same manner as VX2,VY2,VX2T, and VY2T
FX1,FY1	x and y components of force on the bullet in the two body problem (newtons)
FX1T,FY1T	x and y components of force on the bullet in the two body problem at its former position (newtons)
FX2,FY2,FX2T,FY2T	x and y components of force on the target in the two body problem defined in the same manner as FX1, FY1,FX1T, and FY1T
DX	x coordinate separation between atoms in the two body problem (lattice units)
DY	y coordinate separation between atoms in the two body problem (lattice units)
P2KE1	kinetic energy of the bullet in the two body problem after interaction is complete (ev)
P2KE2	kinetic energy of the target in the two body problem after interaction is complete (ev)
TKE	total kinetic energy in the two body problem (ev)
T2RA	recoil angle in the two body problem (degrees)
B2SA	scattering angle in the two body problem (degrees)

SRE8	resultant velocity of target in lattice (m/sec)
SRE1	resultant velocity of bullet in lattice (m/sec)
ATLRA	quotient of two velocities; used to find TLRA
TLRA	recoil angle in the lattice (degrees)
ABLSA	quotient of two velocities; used to find BLSA
BLSA	scattering angle in the lattice (degrees)
RBMT	bullet to target mass ratio
EABL	energy absorbed by the lattice; does not include potential or kinetic energy of target and bullet (ev)
ET2,ETL,RET,RKEP,RSA2L,RRA2L,RBFI	ratios used to compare lattice interaction with two body problem; the ratios are in terms of variables already defined.

APPENDIX V

PROGRAM LISTING

BLOCK A

```

PROGRAM TFF
DIMENSION RX(200), RY(200), RZ(200)
DIMENSION FX(200), FY(200), FZ(200)
DIMENSION VX(200), VY(200), VZ(200)
0 DIMENSION R1X(200), R1Y(200), R1Z(200), F1X(200), F1Y(200),
1 F1Z(200), V1X(200), V1Y(200), V1Z(200)
DIMENSION C(22,36), PPE(200), PKE(200), RXR(200), RYR(200), RZR(200)
COMMON ETFF, AFAC, ETFF
AFAC=3.5906
ETFF=5.9360E-6
ETFP=6.6980E+3
L=1
ROE = SQRTF (2.0)
ROE2 = 2.0
EPST = FORC ( ROE )
EPSTP = POTF(ROE)
10 FORMAT(34X54H THOMAS-FERMI-FIRSOV POTENTIAL, ERCEDED, CCPPER-COPPER
1
3 READ INPUT TAPE 2,4, EV,BX,DTI, PMASS
4 FORMAT(4E10.3)
BZ = BX + 1.0
BZS = 6.0-BX
5 RX(1) = BX
RZ(1) = BZ
PKE8R = -100.0
CYCLE = 0.0
INDEX = 0
RY(1) = -SQRTF(2.0)
12 OFORMAT(/,29H BULLET INCOMING LOCATION X= F6.3, 5H Z= F6.3,
110H ENERGY= F8.1, 24H TIME STEP MULTIPLIER= F5.3, / )
15 FORMAT(1H1)

```

BLOCK B

```

VX(1) = 0.0
VY(1)= SQRTF( EV * 63.54/(1.0E3 * PMASS) ) * 5.511811 E04
VZ(1) = 0.0
DT = DTI/ VY(1) * 1.807E-10
RM= 105.463911 E-27
DIOC = DT/3.614 E-10
DTORM = DT / RM
DIO2RM = DTORM/2.0
RM1 = PMASS *1.6598 E-27
DTORM1 = DT/RM1
DIO2RM1 = DTORM1/ 2.0
DO 250 I = 2, 64
VX(I) = 0.0
VY(I) = 0.0
250 CONTINUE

```

BLOCK C

```

M = 2
N = 65
JT = -1
DO 60 J = 1,7
KT = C
DO 59 K = 1,7
IT = C
DO 58 I = 1,7
NTC=IT+J1+KT
IF(NTC-NTC/2*2) 57,20,57
20 IF(KT - 1) 30, 21, 21
21 IF(KT - 5) 22, 22, 30
22 IF(IT - 1) 30, 23, 23
23 IF(IT - 5) 24, 24, 30
24 IF(JT-C) 30, 25, 25
25 IF(JT-4) 40, 40, 30
30 RX(N) = IT
RY(N) = JT
RZ(N) = KT
N = N + 1
GO TO 57
40 RX(M)=IT
RY(M) = JT
RZ(M) = KT
M = M + 1
57 IT=IT+1
58 CONTINUE
KT = KT + 1
59 CONTINUE
JT = JT + 1
60 CONTINUE

```

BLOCK D

```

DO 65 I = 65, 172
PPE(I)=0.0
RIX(I) = RX(I)
RIY(I) = RY(I)
RIZ(I) = RZ(I)
65 CONTINUE
DO 66 I= 1,64
PPE(I)=0.0
PKE(I)=C.C
RXR(I) =RX(I)
RYR(I)=RY(I)
66 RZR(I) = RZ(I)

```

BLOCK E

```

70 DO 75 I = 1, 64
   FX(I) = 0.0
   FY(I) = 0.0
   FZ(I) = 0.0
75 CONTINUE
   DO 200 I = 1, 64
     IP1 = I + 1
     DO 200 J = IP1, 172
       IF (I-1) 130, 130, 120
130 IF (J-65) 120, 135, 135
135 IF (J-68) 200, 200, 120
120 DRX = RX(J) - RX(I)
     ADRX = ABSF(DRX)
     IF(ADRX - ROE) 140, 200, 200
140 DRY = RY(J) - RY(I)
     ADRY = ABSF(DRY)
     IF(ADRY - ROE) 145, 200, 200
145 DRZ = RZ(J) - RZ(I)
     ADRZ = ABSF(DRZ)
     IF(ADRZ - ROE) 150, 200, 200
150 DIST = (DRX * DRX + DRY * DRY + DRZ * DRZ)
     IF (DIST - ROE2) 155, 200, 200
155 DIST = SQRTF(DIST)
     FORCE = FCRC(DIST) - EPST
     FOD = FORCE / DIST
     FA = DRX * FOD
     FB = DRY * FOD
     FC = DRZ * FOD
     IF (J-65) 160, 165, 165
160 FX(J) = FX(J) + FA
     FY(J) = FY(J) + FB
     FZ(J) = FZ(J) + FC
165 FX(I) = FX(I) - FA
     FY(I) = FY(I) - FB
     FZ(I) = FZ(I) - FC
200 CONTINUE
     DO 235 I = 1, 64
       IF( ABSF(FX(I)) - 1.0E-14) 210, 210, 215
210 FX(I) = 0.0
       IF( ABSF(FY(I)) - 1.0E-14) 220, 220, 225
220 FY(I) = 0.0
       IF( ABSF(FZ(I)) - 1.0E-14) 230, 230, 235
230 FZ(I) = 0.0
235 CONTINUE
     IF (INDEX) 240, 240, 420

```

BLOCK F1

```

240 DO 300 I = 2, 64
    R1X(I) = RX(I)
    R1Y(I) = RY(I)
    R1Z(I) = RZ(I)
    F1X(I) = FX(I)
    F1Y(I) = FY(I)
    F1Z(I) = FZ(I)
    V1X(I) = VX(I)
    V1Y(I) = VY(I)
    V1Z(I) = VZ(I)
    VX(I) = FX(I) * DTCRM + VX(I)
    VY(I) = FY(I) * DTCRM + VY(I)
    VZ(I) = FZ(I) * DTCRM + VZ(I)
    RX(I) = ( VX(I) + V1X(I) ) * DTOC + R1X(I)
    RY(I) = ( VY(I) + V1Y(I) ) * DTOC + R1Y(I)
    RZ(I) = ( VZ(I) + V1Z(I) ) * DTOC + R1Z(I)
300 CONTINUE
    R1X(I) = RX(I)
    R1Y(I) = RY(I)
    R1Z(I) = RZ(I)
    F1X(I) = FX(I)
    F1Y(I) = FY(I)
    F1Z(I) = FZ(I)
    V1X(I) = VX(I)
    V1Y(I) = VY(I)
    V1Z(I) = VZ(I)
    VX(I) = FX(I) * DTCRM1 + VX(I)
    VY(I) = FY(I) * DTCRM1 + VY(I)
    VZ(I) = FZ(I) * DTCRM1 + VZ(I)
    RX(I) = ( VX(I) + V1X(I) ) * DTOC + R1X(I)
    RY(I) = ( VY(I) + V1Y(I) ) * DTOC + R1Y(I)
    RZ(I) = ( VZ(I) + V1Z(I) ) * DTOC + R1Z(I)
    INDEX = 1
    GO TO 70

```

BLOCK F2

```

420 DO 425 I = 2, 64
    VX(I) = ( FX(I) + F1X(I) ) * DT02RM + V1X(I)
    VY(I) = ( FY(I) + F1Y(I) ) * DT02RM + V1Y(I)
    VZ(I) = ( FZ(I) + F1Z(I) ) * DT02RM + V1Z(I)
    RX(I) = ( VX(I) + V1X(I) ) * DTOC + R1X(I)
    RY(I) = ( VY(I) + V1Y(I) ) * DTOC + R1Y(I)
    RZ(I) = ( VZ(I) + V1Z(I) ) * DTOC + R1Z(I)
425 CONTINUE
    VX(I) = ( FX(I) + F1X(I) ) * DT02RM1 + V1X(I)
    VY(I) = ( FY(I) + F1Y(I) ) * DT02RM1 + V1Y(I)
    VZ(I) = ( FZ(I) + F1Z(I) ) * DT02RM1 + V1Z(I)
    RX(I) = ( VX(I) + V1X(I) ) * DTOC + R1X(I)
    RY(I) = ( VY(I) + V1Y(I) ) * DTOC + R1Y(I)
    RZ(I) = ( VZ(I) + V1Z(I) ) * DTOC + R1Z(I)
    INDEX = 0
    CYCLE = CYCLE + 1.0

```

BLOCK G

```

PKE8 = 32.9163268E-08 * ( VX(8) * VX(8) + VY(8) * VY(8) + VZ(8) * VZ(8) )
IF(PKE8) 460, 460, 455
455 IF(PKE8R-PKE8) 460, 465, 465
460 PKE8R = PKE8
    GO TO 70
465 CONTINUE

```

BLOCK H

```

DIX = R1X(8) - R1X(1)
DIY = R1Y(8) - R1Y(1)
DIZ = R1Z(8) - R1Z(1)
D81 = (DIX*DIX + DIY*DIY + DIZ*DIZ)
IF(D81 - ROE2) 428,435,435
428 D81 = SQRT(D81)
430 POT18 = POTF(D81) - EPSTP
GO TO 440
435 POT18 = 0.0
440 PKE1 = 0.51804102E-C8*(V1X(1)*V1X(1) + V1Y(1)*V1Y(1) + V1Z(1)*
V1Z(1))*PMASS
PKE(1) = PKE1
TPOT = 0.0
TPKE = 0.0
DO 600 I = 1,64
IP1 = I + 1
DO 600 J = IP1,172
IF(I - 1) 525,525,540
525 IF(J - 65) 540,530,530
530 IF(J - 82) 600,600,540
540 CRX = R1X(J) - R1X(I)
ADRX = ABSF(CRX)
IF(ADRX - ROE) 545,600,600
545 CRY = R1Y(J) - R1Y(I)
ADRY = ABSF(CRY)
IF(ADRY - ROE) 550,600,600
550 DRZ = R1Z(J) - R1Z(I)
ADRZ = ABSF(DRZ)
IF(ADRZ - ROE) 555,600,600
555 DIST = (CRX*CRX + CRY*CRY + DRZ*DRZ)
IF(DIST - ROE2) 556,600,600
556 DIST = SQRT(DIST)
560 POT = PCIF(DIST) - EPSTP
PPE(1) = PPE(1) + 0.5 * POT
PPE(J) = PPE(J) + 0.5 * POT
TPOT = TPOT + POT
600 CONTINUE
COL3 = PPE(1) - 0.5 * POT18
PE8A = PPE(8) - 0.5 * POT18
DO 650 I = 1,64
PKE(I) = 32.2163268E-C8*(V1X(I)*V1X(I) + V1Y(I)*V1Y(I) + V1Z(I)*V1Z(I))
TPKE = TPKE + PKE(I)
650 CONTINUE
TPKE = TPKE + PKE1
TE = TPKL + TPOT
PERE = ABSF((TE-EV)/EV) * 1.0E2

```

BLOCK J

```

WRITE OUTPUT TAPE 3,10
WRITE OUTPUT TAPE 3,12, BX, BZ, EV, DTI
WRITE OUTPUT TAPE 3,850
850 FORMAT( /, 19H LATTICE LOCATIONS, /)
WRITE OUTPUT TAPE 3,860
860 FORMAT(9CH ATOM DX DY DZ VX VY
1 VZ KE PE /)
SHIFT = 1.0E-05
DO 900 I = 1,64
ENI2(C).
DX = R1X(I) - RXR(I)
+SSK(DX),SLJ(L+2),+SCM(77777B),+FSB(SHIFT),AJP3(865),INI2(1).
865 DY = R1Y(I) - RYR(I)
+SSK(DY),SLJ(L+2),+SCM(77777B),+FSB(SHIFT),AJP3(875),INI2(1).
875 DZ = R1Z(I) - RZR(I)
+SSK(DZ),SLJ(L+2),+SCM(77777B),+FSB(SHIFT),AJP3(885),INI2(1).
885 ENA2(C),AJP(900).
890 WRITE OUTPUT TAPE 3,895, I,DX,DY,DZ,V1X(I),V1Y(I),V1Z(I),PKE(I),
IPPE(I)
900 CONTINUE
895 FORMAT(I10,3F10.5,3E10.2,2F10.2)

```

BLOCK K

```

CONV = 0.01/DT1
DIOC = DIOC * CONV
DTORM = DTORM * CONV
DIO2RM = DIO2RM * CONV
DTORM1=DTORM1*CONV
DIO2RM1 = DIO2RM1 * CONV
RX2 = 0.0
RY2 = 0.0
DSQ=(BZ-3.0)*(BZ-3.0)+(BX-3.0)*(BX-3.0)
RY1= SGRIF(DSQ)
XIP=RY1
RX1= SGRIF(RDE2-DSQ) - 1.0E-06
VX2 = 0.0
VY2 = 0.0
VY1 = 0.0
FX1 = 0.0
FX2= 0.0
FY1= 0.0
FY2= 0.0
VX1 = -SGRTF(EV * 0.06354 / PMASS) * 5.511811E04
INDEX = 0

```

BLOCK L1

```

1005 DX = RX1- RX2
    DY = RY1- RY2
    DIST = DX*DX + DY*DY
    IF(DIST-RDE2) 1010, 1010, 1070

```

BLOCK L2

```

1010 DIST = SGRTF(DIST)
    FORCE= FPRC(DIST) - EPST
    FX1 = DX/ DIST * FORCE
    FX2 = -FX1
    FY1= DY/DIST * FORCE
    FY2 = -FY1
    IF(INDEX) 1015,1015,1030

```

BLOCK M1

```

1015 RX1T = RX1
    RY1T = RY1
    RX2T = RX2
    RY2T = RY2
    FX1T = FX1
    FY1T = FY1
    FX2T = FX2
    FY2T = FY2
    VX1T = VX1
    VY1T = VY1
    VX2T = VX2
    VY2T = VY2
    VX1 = FX1T * DTORM1 + VX1
    VY1 = FY1T * DTORM1 + VY1
    VX2 = FX2T * DTORM1 + VX2
    VY2 = FY2T * DTORM1 + VY2
    RX1 = (VX1 + VX1T) * DIOC + RX1T
    RY1 = (VY1 + VY1T) * DIOC + RY1T
    RX2 = (VX2 + VX2T) * DIOC + RX2T
    RY2 = (VY2 + VY2T) * DIOC + RY2T
    INDEX = I
    GO TO 1005

```

BLOCK M2

```

1030 VX1= (FX1 + FX1T)* DT02RM1 + VX1T
      VY1= (FY1 + FY1T)* DT02RM1 + VY1T
      VX2= (FX2 + FX2T)* DT02RM1 + VX2T
      VY2= (FY2 + FY2T)* DT02RM1 + VY2T
      RX1= (VX1 + VX1T)* DTCC + RX1T
      RY1= (VY1 + VY1T)* DTCC + RY1T
      RX2= (VX2 + VX2T)* DTCC + RX2T
      RY2= (VY2 + VY2T)* DTCC + RY2T
      INDEX = 0
      GO TO 1005

```

BLOCK N

```

1070 P2KE1=0.51804102E-08*(VX1*VX1 + VY1*VY1)*PMASS
      P2KE2= 32.9163268E-08 *(VX2* VX2 + VY2* VY2)
      TKE = P2KE1 + P2KE2
      WRITE OUTPUT TAPE 3, 15
      WRITE OUTPUT TAPE 3, 10
      WRITE OUTPUT TAPE 3, 12, BX, BZ, EV, DTI
      WRITE OUTPUT TAPE 3, 1075
1075 FORMAT(////, 4CH RESULTS OF SIMPLE TWO BODY INTERACTION ,/ )
      WRITE OUTPUT TAPE 3, 1080, P2KE1, P2KE2, TKE, EV
10800FORMAT( 15H KIN ENRG PROJ= , F12.5 , 18H KIN ENRG TARG= ,
1 F12.5, 14H TOT ENRG= , F12.5, 12H ENRG IN= , F10.1, ////)
      WRITE OUTPUT TAPE 3, 654
654 FORMAT(31H RESULTS OF LATTICE INTERACTION ,/, )
      WRITE OUTPUT TAPE 3, 655
655 FORMAT(14H TARGET VALUES ,//, 90H TIME STEPS PCT E TARGET W/R TO
1 BULLET KE TARGET PCT E TARGET W/R TO LAT-BULLET )
      WRITE OUTPUT TAPE 3, 660, CYCLE, POT18, PKE8R, PEAS
660 FORMAT(F10.0, 1PE25.6, E18.6, E25.6 )
      WRITE OUTPUT TAPE 3, 665, PKE1, COL3, V1X(1), V1Y(1), V1Z(1)
6650FORMAT(/, 14H BULLET VALUES, //, 15H KIN ENERGY = F7.1 , 29H PCT
1 E W/R TO LAT-TARGET = F7.1, 25H VELOCITY IN M/SEC X= 1PE9.2,
2 4H Y= E9.2, 4H Z= E9.2 )
      WRITE OUTPUT TAPE 3, 675
675 FORMAT ( /, 16H LATTICE VALUES )
      WRITE OUTPUT TAPE 3, 800
8000FORMAT ( /, 60H PCT ENERGY KIN ENERGY TCT ENERGY
1 PER ERROR )
      WRITE OUTPUT TAPE 3, 810, TPOT , TPKE, TE, PERE
910 FORMAT(3F15.5 , F10.5 )
      WRITE OUTPUT TAPE 3, 1100
1100 FORMAT(//, 71H COMPARISON VALUES OF LATTICE REACTION AGAINST SIMPL
1E TWO BODY PROBLEM ,/, 65H BOTH INTERACTIONS STOPPED WHEN TARGET K
2INETIC ENERGY MAXIMUM )

```

```

T2RA = ATANF(ABSF(VY2/VX2)) * 57.296
B2SA = ATANF(ABSF(VY1/VX1)) * 57.296
SRE8 = SQRTF(PKE8R/32.9163268E-08)
SRE1 = SQRTF(PKE1 / (PMASS*0.51804102E-08))
ATLRA = ABSF(VIY(8))/SRE8
IF(ATLRA-1.0) 1105,1110,1110
1105 TLRA = ACOSF( ATLRA) * 57.296
GO TO 1115
1110 TLRA = 0.0
1115 ABLSA = ABSF(VIY(1))/SRE1
IF(ABLSA -1.0) 1116,1118,1118
1116 BLSA = ACOSF( ABLSA) * 57.296
GO TO 1119
1118 PLSA = 0.0
1119 RBMT = PMASS/ 63.54
ET2 = (PKE2/EV)*1.0E02
ETL=(PKE8R/EV)*1.0E02
RET=ETL/ET2
RKEP=PKE1/P2KE1
EABL = TL-PKE1-PKE8R-PPE(8)-PPE(1)
RSA2L=B2SA/BLSA
RRA2L = T2RA/TLRA
RBF1= PKE1/EV
WRITE OUTPUT TAPE 3,1120,ET2,ETL,RET,RKEP,RBMT
1120 FORMAT(/, 43H PERCENT KE TRANSFERED TWO BODY= F9.3,
118H LATTICE= F9.3,/, 56H RATIO KE TRANSFER IN LATTICE
2/ KE TRANSFER TWO BODY = F7.5,/, 48H RATIO KE PROJ IN LATTICE /
3 KE PROJ TWO BODY = F11.5,/, 36H RATIO BULLET MASS / TARGET MAS
4S = F6.3 )
WRITE OUTPUT TAPE 3,1130,EABL,R2SA,BLSA,T2RA,TLRA,RSA2L,RRA2L
1130 FORMAT(/, 33H POT + KE ABSORBED BY LATTICE = F12.3,/,
1 39H SCATTERING ANGLE TWO BODY = F7.2, 19H LAT
2TICE = F7.2,/, 35H RECOIL ANGLE TWO BODY = F7.2,
3 19H LATTICE = F7.2,/, 63H RATIO SCATTERING ANGLE TWO
4BODY / SCATTERING ANGLE LATTICE = F9.5,/, 55H RATIO RECOIL ANG
5LE TWO BODY / RECOIL ANGLE LATTICE = F9.5 )
WRITE OUTPUT TAPE 3,1140, XIP, RBF1
1140 FORMAT( /, 21H IMPACT PARAMETER = F7.4, 54H RATIO KE BULLE
1T AT END OF RUN / ORIGINAL ENERGY = 1PE13.5 )
WRITE OUTPUT TAPE 3,15

```

BLOCK R

```

C(1,L) = BX
C(2,L) = BZ
C(3,L) = XIP
IF(BZ-3.0) 1160,1160,1170
1160 C(4,L) = C.C
GO TO 1170
1170 C(4,L) = ACCSF((BZ-3.0)/XIP)* 57.296
1180 C(5,L) = ET2*1.0E-02
C(6,L) = B2SA
C(7,L) = ETL*1.0E-02
C(8,L) = BLSA
C(9,L) = RET
C(10,L) = PPE(8) + PKE8R
C(11,L) = TLRA
C(12,L) = PKE1
C(13,L) = P2KE1
C(14,L) = PKE8R
C(15,L) = P2KE2
C(16,L) = PERE
C(17,L) = CYCLE
C(18,L) = PPE(1)
C(19,L) = PPE(8)
C(20,L) = T2RA
C(21,L) = R1Y(1)
L=L+1

```

BLOCK T

```

BZ = BZ-0.1
IF(BZS-BZ)5,5,1300
1300 CONTINUE
IF(SENSE SWITCH 1) 3, 4000
4000 CONTINUE

```

BLOCK W

```

LMAX = L-1
WRITE OUTPUT TAPE 3,10
WRITE OUTPUT TAPE 3,1350
1350 FORMAT(//,110H      X      Z      B      ANGLE      TM(2)
1 SANGLE      TM(L)      SANGLE      TM(L)/TM(2)      TE TAR      RANGLE      ,/ )
DO 1360 J=1,LMAX
1360 WRITE OUTPUT TAPE 3,1400,(C(I,J), I= 1,11 )
1400 FORMAT(3F9.3,F10.2,F10.4,F10.2,F10.4,F10.2,F10.4,F13.4,F10.2)
DO 1500 M= 1,5
WRITE OUTPUT TAPE 3,15
J = 12
I = 14
WRITE OUTPUT TAPE 3,1425
WRITE OUTPUT TAPE 3,10
WRITE OUTPUT TAPE 3,1430,EV,DTI
WRITE OUTPUT TAPE 3,1450
1450 FORMAT(//,33X53H TOP NO. IS KINETIC ENERGY IN EV OF BULLET IN LAT
1 TICE,//,32X57H BOTTOM NO. IS KINETIC ENERGY IN EV OF TARGET IN LAT
2 ICE )
1500 WRITE OUTPUT TAPE 3,1850,(C(J,L), L=1,11), (C(I,L), L=1,11),
1 (C(J,L), L=12,20), (C(I,L), L=12,20), (C(J,L), L=21,27),
2 (C(I,L), L=21,27), (C(J,L), L=28,32), (C(I,L), L=28,32),
3 (C(J,L), L=33,35), (C(I,L), L=33,35), C(J,36),C(I,36)
DO 1530 M = 1,5
WRITE OUTPUT TAPE 3,15
J=12
I=13
WRITE OUTPUT TAPE 3,1425
WRITE OUTPUT TAPE 3,10
WRITE OUTPUT TAPE 3,1430,EV,DTI
WRITE OUTPUT TAPE 3,1510
1510 FORMAT(//, 23X53H TOP NO. IS KINETIC ENERGY IN EV OF BULLET IN LAT
1 TICE,//,28X64H BOTTOM NO. IS KINETIC ENERGY IN EV OF BULLET IN TWO
2 BODY PROBLEM )
1530 WRITE OUTPUT TAPE 3,1850,(C(J,L), L=1,11), (C(I,L), L=1,11),
1 (C(J,L), L=12,20), (C(I,L), L=12,20), (C(J,L), L=21,27),
2 (C(I,L), L=21,27), (C(J,L), L=28,32), (C(I,L), L=28,32),
3 (C(J,L), L=33,35), (C(I,L), L=33,35), C(J,36),C(I,36)
DO 1550 M = 1,5
WRITE OUTPUT TAPE 3,15
J = 14
I = 15
WRITE OUTPUT TAPE 3,1425
WRITE OUTPUT TAPE 3,10
WRITE OUTPUT TAPE 3,1430,EV,DTI
WRITE OUTPUT TAPE 3,1540
1540 FORMAT(//, 33X53H TOP NO. IS KINETIC ENERGY IN EV OF TARGET IN LAT
1 TICE,//, 27X66H BOTTOM NO. IS KINETIC ENERGY IN EV OF TARGET IN TW
2 BODY PROBLEM )
1550 WRITE OUTPUT TAPE 3,1850,(C(J,L), L=1,11), (C(I,L), L=1,11),
1 (C(J,L), L=12,20), (C(I,L), L=12,20), (C(J,L), L=21,27),
2 (C(I,L), L=21,27), (C(J,L), L=28,32), (C(I,L), L=28,32),
3 (C(J,L), L=33,35), (C(I,L), L=33,35), C(J,36),C(I,36)
WRITE OUTPUT TAPE 3,15
J=17
I=16
WRITE OUTPUT TAPE 3,1425
WRITE OUTPUT TAPE 3,10
WRITE OUTPUT TAPE 3,1430,EV,DTI
WRITE OUTPUT TAPE 3,1560
1560 FORMAT(//, 44X32H TOP NO. IS NUMBER OF TIME STEPS,//,46X28H BOTTOM
1 NO. IS PERCENT ERROR )
WRITE OUTPUT TAPE 3,1850,(C(J,L), L=1,11), (C(I,L), L=1,11),
1 (C(J,L), L=12,20), (C(I,L), L=12,20), (C(J,L), L=21,27),
2 (C(I,L), L=21,27), (C(J,L), L=28,32), (C(I,L), L=28,32),
3 (C(J,L), L=33,35), (C(I,L), L=33,35), C(J,36),C(I,36)

```

```

DO 1580 M=1,2
WRITE OUTPUT TAPE 3,15
J=8
I=6
WRITE OUTPUT TAPE 3,1425
WRITE OUTPUT TAPE 3,10
WRITE OUTPUT TAPE 3,1430,EV,DTI
WRITE OUTPUT TAPE 3,1570
1570 FORMAT(//, 40X30H TOP NO. IS SCATTERING ANGLE IN LATTICE, //, 34X52H
1 BOTTOM NO. IS SCATTERING ANGLE OF TWO BODY PROBLEM )
1580 WRITE OUTPUT TAPE 3,1850, (C(J,L), L=1,11), (C(I,L), L=1,11),
1(C(J,L), L=12,20), (C(I,L), L=12,20), (C(J,L), L=21,27),
2(C(I,L), L=21,27), (C(J,L), L=28,32), (C(I,L), L=28,32),
3(C(J,L), L=33,35), (C(I,L), L=33,35), C(J,36),C(I,36)
DO 1600 M=1,2
WRITE OUTPUT TAPE 3,15
J=11
I=20
WRITE OUTPUT TAPE 3,1425
WRITE OUTPUT TAPE 3,10
WRITE OUTPUT TAPE 3,1430,EV,DTI
WRITE OUTPUT TAPE 3,1590
1590 FORMAT(//, 42X35H TOP NO. IS RECOIL ANGLE IN LATTICE, //, 36X47H BCT
1 TCM NO. IS RECOIL ANGLE OF TWO BODY PROBLEM )
1600 WRITE OUTPUT TAPE 3,1850, (C(J,L), L=1,11), (C(I,L), L=1,11),
1(C(J,L), L=12,20), (C(I,L), L=12,20), (C(J,L), L=21,27),
2(C(I,L), L=21,27), (C(J,L), L=28,32), (C(I,L), L=28,32),
3(C(J,L), L=33,35), (C(I,L), L=33,35), C(J,36),C(I,36)
DO 1620 M=1,5
WRITE OUTPUT TAPE 3,15
J=19
I=18
WRITE OUTPUT TAPE 3,1425
WRITE OUTPUT TAPE 3,10
WRITE OUTPUT TAPE 3,1430,EV,DTI
WRITE OUTPUT TAPE 3,1610
1610 FORMAT(//, 32X55H TOP NO. IS POTENTIAL ENERGY IN EV OF TARGET IN LA
1 TTICE, //, 31X58H BOTTOM NO. IS POTENTIAL ENERGY IN EV OF BULLET IN
2 LATTICE )
1620 WRITE OUTPUT TAPE 3,1850, (C(J,L), L=1,11), (C(I,L), L=1,11),
1(C(J,L), L=12,20), (C(I,L), L=12,20), (C(J,L), L=21,27),
2(C(I,L), L=21,27), (C(J,L), L=28,32), (C(I,L), L=28,32),
3(C(J,L), L=33,35), (C(I,L), L=33,35), C(J,36),C(I,36)
DO 1640 M=1,2
WRITE OUTPUT TAPE 3,15
J=21
I=12
WRITE OUTPUT TAPE 3,1425
WRITE OUTPUT TAPE 3,10
WRITE OUTPUT TAPE 3,1430,EV,DTI
WRITE OUTPUT TAPE 3,1630
1630 FORMAT(//, 36X47H TOP NO. IS BULLET PENETRATION IN LATTICE UNITS, //
1, 32X56H BOTTOM NO. IS KINETIC ENERGY IN EV OF BULLET IN LATTICE )
1640 WRITE OUTPUT TAPE 3,1850, (C(J,L), L=1,11), (C(I,L), L=1,11),
1(C(J,L), L=12,20), (C(I,L), L=12,20), (C(J,L), L=21,27),
2(C(I,L), L=21,27), (C(J,L), L=28,32), (C(I,L), L=28,32),
3(C(J,L), L=33,35), (C(I,L), L=33,35), C(J,36),C(I,36)
1425 FORMAT(////////)
1430 FORMAT(//, 8X25H BULLET INCOMING ENERGY = ,F8.1,31H EV TIME ST
1EP MULTIPLIER = ,F5.3, 34H BULLET INCIDENT CN (100) FACE )
1850 FORMAT(////, 3X11F10.3, //, 111H * * * * *
1 * * * * *
2 * ,//, 3X11F10.3, //, 13X9F10.3, //, 18X82H * * * *
3 * * * * *
4 13X9F10.3, //, 23X7F10.3, //, 28X62H * * * *
5 * * * * * ,//, 23X7F10.3, //, 33X5F10.3, //, 38X
6 42H * * * * *
7 33X3F10.3, //, 48X22H * * * * * ,//, 43X3F10.3, //, 53XF10.
8 3, //, 56X2H * ,//, 53XF10.3)
END

```

BLOCK Y

```
FUNCTION POTF(DIST)
COMMON E1FP,AFAC,E1FF
POTF=E1FP*(1.0+(DIST*AFAC)**0.8034)**(-3.734)/DIST
RETURN
END
```

BLOCK Z

```
FUNCTION FORC(DIST)
COMMON E1FP,AFAC,E1FF
Y=(DIST*AFAC)**0.8034
YP=1.0+Y
Z=YP**(-3.734)
FORC=E1FF*Z*(1.0+2.9999*Y/YP)/(DIST*DIST)
RETURN
END
END
```

APPENDIX VI

POTENTIALS

1. Born-Mayer

A. Potential

$$\phi(r) = A e^{-\rho(r-r_0)/r_0}$$

r_0 = nearest neighbor separation at zero pressure and absolute zero

r_0 for copper = 2.51A

The three different sets of constants used in the potential form are identical to those used by Gibson et. al.¹⁶

<u>Potential</u>	<u>A(ev)</u>	<u>ρ</u>
1	0.0392	16.97
2	0.051	13.00
3	0.1004	10.34

Potential No. 2 was used extensively in the program. It can be reduced to a simplified form in the following manner:

$$\phi(r) = 0.051 \times e^{-13.0(r-2.551)/2.551}$$

if distances are expressed in lattice units

$$\phi(r) = 0.051 \times e^{(13.0 - 9.208 r)}$$

in the potential function subroutine this becomes

$$\text{POTF} = 0.051 * \text{EXPF}(13.0 - \text{DIST} * 9.208)$$

Potentials one and three become;

$$1) \text{ POTF} = 0.0392 * \text{EXPF}(16.97 - \text{DIST} * 8.9996)$$

$$3) \text{ POTF} = 0.1004 * \text{EXPF}(10.34 - \text{DIST} * 7.3115)$$

B. Force

$$\frac{\partial[\phi(r)]}{\partial r} = -A e^{-\rho(r-r_0)/r_0} \times \frac{\rho}{r_0}$$

substitution of constants for Potential No. 2

$$\frac{\partial[\phi(r)]}{\partial r} = - \frac{(13.0)(0.051)}{2.551} e^{-13.0(r-2.551)/2.551}$$

if distances are expressed in lattice units and electron volts changed to joules;

$$\frac{\partial[\phi(r)]}{\partial r} = -0.41614 \times 10^{-9} \times e^{(13.0-9.208r)}$$

in the force function subroutine this becomes;

$$\text{FORC} = 0.41614\text{E-}09 * \text{EXPF}(13.0-9.208*\text{DIST})$$

forces for potentials one and three become;

$$1) \text{FORC} = 0.3127\text{E-}09 * \text{EXPF}(16.97-8.9996*\text{DIST})$$

$$2) \text{FORC} = 0.6507\text{E-}09 * \text{EXPF}(10.34-7.3115*\text{DIST})$$

2. Thomas-Fermi-Firsov

A. Potential

$$V(r) = \frac{1}{2} E_{\text{TF}} \left[\left(\frac{a_{\text{TF}}}{r} \right) \phi \left(\frac{r}{a_{\text{TF}}} \right) \right]$$

$$\phi \left(\frac{r}{a_{\text{TF}}} \right) = \phi(x) \quad \text{Thomas-Fermi screening function}$$

$$\phi(x) = \left[1 + \left(\frac{x}{\alpha_1} \right)^{\alpha_2} \right]^{-\alpha_3}$$

$$\alpha_1 = 12^{2/3}$$

$$E_{TF} = 2 Z_1 Z_2 e^2 / a_{TF}$$

$$\alpha_2 = 0.8034$$

$$\alpha_3 = 3.734$$

$$a_{TF} = \frac{(9\pi^2/128)^{1/3} k a_H}{(\sqrt{Z_1} + \sqrt{Z_2})^{2/3}}$$

$$a_H = \hbar^2 / m e^2$$

first Bohr radius of hydrogen

$$K=1 \quad E_{TF} = 219233 \text{ eV} \quad a_{TF} = .96062 \times 10^{-11} \text{ m}$$

let $x = r/a_{TF}$ and $ATFR = 1/a_{TF}$ (in lattice units)

then $x = ATFR * DIST$ (DIST in lattice units)

$$\phi(x) = \left[1 + \left(\frac{x}{\alpha_1} \right)^{\alpha_2} \right]^{-\alpha_3} = \left[1 + \left(\frac{ATFR * DIST}{\alpha_1} \right)^{\alpha_2} \right]^{-\alpha_3}$$

let $ATFR/\alpha_1 = AFAC$

$$\phi(x) = \left[1 + (AFAC * DIST)^{0.8034} \right]^{-3.734}$$

$$V(r) = \left[\frac{E_{TF}}{2} \right] \frac{1}{x} \phi(x) = \frac{E_{TF} * \phi(x)}{2 * ATFR * DIST}$$

let $ETFP = E_{TF} / 2 * ATFR$

then

$$POTF = \frac{ETFP}{DIST} \left[1.0 + (AFAC * DIST)^{0.8034} \right]^{-3.734}$$

B. Force

$$\text{set } x = r/a_{TF}$$

$$\frac{\partial(V(r))}{\partial r} = \frac{\partial(V(r))}{\partial x} \times \frac{\partial x}{\partial r}$$

$$\frac{\partial(V(r))}{\partial r} = \frac{E_{TF}}{2} \left[-\frac{\phi(x)}{x^2} + \frac{1}{x} \frac{d(\phi(x))}{dx} \right] \times \frac{1}{a_{TF}}$$

$$\text{set } Y = \left(\frac{x}{\alpha_1} \right)^{\alpha_2}$$

then

$$\frac{d(\phi(x))}{dx} = -\alpha_2 \alpha_3 \frac{\phi(x)}{x} \frac{Y}{1+Y}$$

$$\frac{\partial(V(r))}{\partial r} = \frac{E_{TF} \times a_{TF}}{2r^2} \left[1 + \alpha_2 \alpha_3 \frac{Y}{Y+1} \right] \times \phi(x)$$

if the following substitutions and appropriate changes in units are made

$$\text{FORC} = \frac{E_{TF} \times Z}{\text{DIST} \times \text{DIST}} \left[1.0 + 2.9999 \times \frac{Y}{YP} \right]$$

$$\text{where } E_{TF} = E_{TF} \times a_{TF} / 2$$

$$x = \text{DIST} \times \text{AFAC} \quad (\text{DIST in lattice units})$$

$$Y = (x)^{0.8034} \quad YP = 1.0 + Y$$

$$Z = (YP)^{-3.734}$$

E_{TF} and AFAC are defined at the beginning of the program - Y , YP and Z are calculated in the subroutine each time it is used.

thesG2545

Machine calculations of energy transfer



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