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THE ENERGY BUDGET AT THE EARTH'S SURFACE:
ESTIMATES OF THE DIFFUSION RESISTANCE
OF SOME LARGE SUNFLOWER LEAVES IN THE FIELD

Contribution by:

L. A. Hunt, Ivan I. Impens and E. R. Lemon

INTERIM REPORT

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UNITED STATES ARMY ELECTRONICS COMMAND

Atmospheric Sciences Laboratory, Research Division
Fort Huachuca, Arizona

Cross Service Order 2-67

Microclimate Investigations, SWCRD-ARS-U. S. Department of Agriculture
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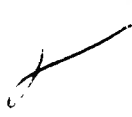
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INTERIM REPORT

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Prepared by

L. A. Hunt, Ivan I. Impens and E. R. Lemon

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For

U. S. Army Electronics Command
Atmospheric Sciences Laboratory, Research Division
Fort Huachuca, Arizona

ESTIMATES OF THE DIFFUSION RESISTANCE
OF SOME LARGE SUNFLOWER LEAVES IN THE FIELD

L. A. Hunt, Ivan I. Impens and E. R. Lemon

Measurements of the resistance to water vapor movement from the intercellular spaces to the leaf surface (leaf resistance, r_l) were made in a one-hectare field crop of sunflowers (Helianthus annuus L.) at Ellis Hollow (near Ithaca, N. Y.). Some of these measurements, supplemented with determinations of net radiation above and below the leaves, leaf temperature, and both air temperature and humidity, were used to estimate the resistance to mass and sensible heat transfer from the surface of the leaf to the bulk air outside the leaf boundary layer (boundary layer resistance, r_a).

The sunflower crop was sown early in June 1966, and thinned to 60,000 plants per hectare and irrigated in July. Measurements to determine both leaf and boundary layer resistances were made during clear intervals on days with scattered clouds (1 and 3 August); measurements of leaf resistance alone were made during the afternoon of a day that was cloudless from early morning (4 August). One large, representative leaf lamina approximately 27 cm wide along both major axes and located 1.5-1.8 m from ground level, was selected on each of ten plants. The experimental leaf was supported roughly parallel to the soil surface by cellophane tape strips joining the petiole and stem. All remaining leaves on each test plant were wrapped around the stem to avoid their interference in direct solar radiation on the test leaf.

Leaf resistances were measured with a small diffusion porometer of the type described by van Bavel, Nakayama and Ehrler (19). Five measurements were made on both upper and lower leaf surfaces, with each individual determination being made on a separate leaf. The leaves were examined in random order. Net radiation above and below the leaf was determined with a miniature Funk net radiometer which has a 12-mm-diameter sensing head (Middleton & Co., Adelaide, Australia). Leaf temperature was measured with a bolometer (Barnes Eng. Co. infrared thermometer, Model IT-3) with a 7° field of view. The indicated temperature was generally 0.5 °C higher when viewing the leaf from below than from above. The temperature recorded when viewing from below was considered more correct, and thus was used for calculation (7). Temperature and humidity of the bulk air in the community at the leaf level were measured with an aspirated wet and dry bulb thermistor psychrometer (Atkins Technical Inc.).

The leaf resistance was determined from diffusion meter measurements for the upper and lower leaf surfaces by.

$$\frac{1}{r_l} = \frac{1}{r_{lb}} + \frac{1}{r_{ld}}$$

where r_{lb} and r_{ld} are the measured resistances for abaxial and adaxial leaf surfaces respectively. The boundary layer resistance, r_a , was calculated by solving the following energy balance equation:

$$R_N = \frac{L (\Delta\rho)}{r_a + r_l} + \frac{C_p \rho_a (\Delta T)}{r_a}$$

where R_N is the net radiation intercepted by the leaf, in $\text{cal cm}^{-2} \text{sec}^{-1}$; L is the heat of vaporization of water (590 cal g^{-1}); C_p is the specific heat of air ($0.24 \text{ cal g}^{-1} \text{ }^\circ\text{C}^{-1}$); ρ_a is the air density ($1.205 \times 10^{-3} \text{ g cm}^{-3}$); $\Delta\rho$ is the difference between the saturation water vapor density at leaf temperature and water vapor density of bulk air, in g cm^{-3} ; and ΔT is the difference between leaf and bulk air temperature, in $^\circ\text{C}$. The bulk air is considered as that having representative properties of the air in the plant community at the level of the leaf measurements but outside the leaf boundary layer.

The leaf and boundary layer resistances for clear periods during two days with scattered clouds are given in Table 1. The leaf resistances, which varied from 3.3 to 4.7 sec cm^{-1} , were higher than the minimum values (0.3 sec cm^{-1}) reported by Holmgren, Jarvis and Jarvis (9) for sunflowers. This may be attributable either to the environment and/or to the leaves themselves, because (a) intermittent clouds may have caused partial stomatal closure, and (b) minimum stomatal resistance is not constant throughout the life of a leaf (9, 16).

Values of boundary layer resistances calculated by the energy balance equation ($0.05 - 0.13 \text{ sec cm}^{-1}$) are lower than values calculated from formulas based either on: (a) forced convection laminar flow theory applicable to isolated, rigid, smooth bodies that are flat and thin, oriented parallel to a relatively uniform airstream; or (b) empirical data dealing with water vapor loss from isolated small leaf models studied under laboratory conditions. Combining both theoretical and empirical formulas, Monteith (12) suggested that the following formula was most appropriate for describing boundary resistance to water vapor transfer of leaves under field conditions:

$$r_a = 0.65 \left(\frac{x}{u}\right)^{0.5} \text{ sec cm}^{-1} \text{ per unit area of leaf}$$

where x is the width of the leaf (cm) and u is the velocity of the bulk air stream (cm sec^{-1}). This formula, when adjusted for the difference in both thickness of boundary layer and diffusion coefficient, yields the following expression for boundary resistance to heat transfer:

$$r_a = 0.75 \left(\frac{x}{u}\right)^{0.5}$$

Such formulas when applied to leaves of the size examined in the present work yield the following estimates of r_a :

u (cm sec ⁻¹)	r _a (sec cm ⁻¹)	
	<u>Water vapor</u>	<u>Heat</u>
10	1.07	1.23
100	0.33	0.41
1,000	0.10	0.12

Thus, exceptionally high windspeeds would be required to yield r_a values of the magnitude found in Table 1.

Another method of determining r_a of sunflower leaves in the field was tried. We used a transient heat pulse method described by Linacre (11):

$$r_a = 4.05 \times 10^{-6} t_{0.5} S_A(F) B$$

where $t_{0.5}$ is the time required after shading for the temperature to fall from its initial value to that halfway between the initial and final states, $S_A(F)$ is the leaf surface area per unit of fresh weight (specific leaf area on a fresh weight basis), and

$$B = \frac{r_l + r_a (\Delta + \gamma)/\gamma}{r_l + r_a}$$

in which γ is the psychrometric constant and Δ is defined by:

$$\Delta = \frac{e_s - e_{as}}{T} \left[\frac{d(e_{as})}{dT} \right]_{T_a}$$

where e_s is the saturated water vapor pressure at leaf temperature, e_a is the saturated water vapor pressure at the bulk air temperature, and ΔT is the difference between leaf and bulk air temperature. T_a is the absolute temperature of the air. Four separate temperature determinations were made using a bolometer during clear weather on the morning of 4 August, when the windspeed immediately above the crop was approximately 100 cm sec^{-1} . Intermittent shading was accomplished with a $1 \times 1 \text{ m}$ plywood board held normal to the incoming sunlight with its central point approximately 50 cm from the leaf. This arrangement permitted sharp delineation between the unshaded and shaded regimes. The $t_{0.5}$ values were determined with the bolometer coupled to a sensitive mV strip chart recorder (Leeds & Northrup Co.) when ΔT was approximately 2.5°C ; actual values (8 sec, 8 sec, 10 sec, 7 sec) were read from curves drawn by hand through the recorder trace. The mean r_a estimate was $0.06 \pm 0.01 \text{ sec cm}^{-1}$, a value substantially less than those calculated on the basis of Monteith's equation. Evidently, therefore, the present findings reflect other than the commonly accepted ideas of the magnitude of leaf boundary layer resistances. Results of this method thus agree with those determined by the energy balance method.

Such results could stem from hairs present on both the upper and lower surfaces of the sunflower leaves. This follows, as pointed out by Wolpert (20), because heat transfer through a laminar boundary layer is by conduction. Thus, when certain roughness elements lie within a laminar boundary layer, the distance through which heat must travel by conduction through the fluid medium is reduced, and the rate of heat transfer is increased, relative to a surface without roughness projections. In certain cases, the reduction in the boundary layer resistance of rough, relative to smooth, surfaces may equal the percentage increase of surface area due to the roughness elements (20).

However, the present findings could also stem from the turbulence of the bulk air in the plant community. The laminar, forced convection equation was derived from studies under conditions of relatively steady air movement. When the airstream is turbulent, the edge of a laminar boundary layer is irregular and its resistance is reduced by penetrating eddies. Hinze (8) indicates that for turbulent intensities of 10%, the resistance value is 10 - 50% less than that given by Monteith's equation. The actual reduction depends on the geometry of the object considered and the scale of turbulence of the airstream. Under field conditions the turbulent intensity is often greater than 10% (18), thus reductions of more than 50% in the laminar flow resistance values should be expected.

Furthermore, the results could reflect a complete change in the nature of the boundary layer. With flat, smooth objects that are thin and rigid, a transition from a laminar to a turbulent boundary layer occurs when the Reynolds number is $3-5 \times 10^5$. Transition will occur at the lower value when (a) the free stream is turbulent, or (b) the free stream is not uniform (2). With rough, irregular, flexible objects, the boundary layer probably is turbulent from the leading edge (2). In the present study a Reynolds number of 3×10^5 would require windspeeds of 1500 cm sec^{-1} . Transition at windspeeds as low as $100-500 \text{ cm sec}^{-1}$ could arise from hairs and venal undulations making the surface aerodynamically rough, as well as from elastic properties permitting flexing and bending of the leaf surface under the influence of larger eddies in the bulk air stream. However, only a small amount of leaf movement was observed; thus it is likely that the geometric properties of the surface would have played the more important role among the various leaf characteristics in bringing about any complete change in the nature of the boundary layer. The scale of turbulence experienced by the test leaves being in an array of other plant parts under field conditions could be of the greatest importance of all.

Our results reported here for sunflower leaves agree with those calculated for leaves in a stand of corn at the same location (17, 3). These latter calculations involved: (a) the application of an energy budget approach to evaluate both the eddy diffusivity (K) and the boundary layer resistance (r_a) at different levels within the crop canopy, (b) mathematical approximation of the attenuation of both of these components, and (c) comparison of the attenuation characteristics of K and r_a . Interpretation was based on the implicit assumption that windspeed was attenuated within the canopy in the same way as K .

Our results, by indicating that the r_a values of large sunflower leaves under field conditions are less than those estimated on the basis of laminar flow theory, also add weight to the contentions of Philip (14) and Budagovsky (4). Philip (14), commenting on Raschke's (15) generalization that leaf boundary layers are laminar and that narrow leaves have the smallest r_a values, has argued that large leaves may have a greater scale of turbulence and may in consequence have smaller r_a values than small leaves. Budagovsky (4) has also commented on the possibility of leaf boundary layers being turbulent under outdoor conditions.

The realization that some form of turbulent boundary layer may exist around leaves under field conditions is of secondary importance to water vapor movement from the intercellular spaces to the bulk air in the plant community. This follows, as pointed out by Impens et al (10), because r_a under outdoor conditions is usually an order of magnitude smaller than r_l . If one assumes a laminar boundary layer and a windspeed of 100 cm sec^{-1} , and then calculates the air resistance for the leaves tested, the sum of r_a and r_l gives a mean of 4.3 sec cm^{-1} , whereas the measured mean of r_a and r_l is 4.1 sec cm^{-1} (mean values from Table 1). The boundary layer resistance is of primary importance, however, in controlling sensible heat losses from leaves. This fact should not be minimized.

The leaf resistances recorded on the clear day varied from 3.2 to 4.8 sec cm^{-1} (Table 2). These values, when expressed on an hourly basis, were higher in the periods 1100 to 1300 and 1400 to 1500 than during 1300 to 1345. This latter pattern, which agrees in general with that reported for both bulrush millet (1) and corn (3) indicates stomatal closure during the middle of the day and stomatal opening in the afternoon. By contrast to the millet and corn, however, stomatal opening in the sunflower apparently occurred in the period 1300 to 1345 rather than later in the afternoon around 1500. These differences presumably stem from differences either in the soil-plant water relations, or in the intrinsic pattern of stomatal behavior of the species.

By contrast to the hourly mean r_L values, the 15-min means exhibited a shorter cycle, with minima occurring at 1100, 1300 and 1430, and maxima at 1215, 1415 and 1500. Such a pattern conforms with that reported by Ehrler et al. (6) for cotton and by Parkhurst and Gates (13) for cottonwood, under standard conditions. In both of the latter cases, the cycling was explained in terms of an "overshoot" mechanism, in which partial stomatal closure, occurring in response to an increased leaf water deficit, resulted in rehydration of leaf tissue and in opening of the stomates. Increased water loss following stomatal opening ultimately resulted in an increased leaf water deficit and stomatal closure. The present data suggest that a similar mechanism could be involved in the field. Such a mechanism could be associated with the cyclical pattern reported for the rate of evaporation of tamarisk during the afternoon hours of days of high solar radiation (5).

The present results thus indicate that erroneous conclusions may result from both: (a) computation of leaf boundary layer resistance on the basis of laminar flow theory, and (b) analysis of diurnal changes in evapotranspiration on the assumption that leaf resistance exhibits only one maximum value at midday or thereabouts.

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Table 1. Leaf and boundary layer resistances during clear periods on two days with scattered clouds

Time (EST) Hour	Leaf resistances (r_l) (diffusion meter method) (sec cm ⁻¹)	Boundary layer resistance (r_g) (energy balance method) (sec cm ⁻¹)
0900 *	4.7	0.06
1000	3.6	0.12
1030	3.3	0.06
1130	3.7	0.07
1200	3.4	0.05
1230	4.2	0.06
1300	3.9	0.07
1330	4.3	0.13
1400	4.4	0.12
1500	4.7	0.06
Standard error of mean (n = 5)	0.4	0.02

* 3 August. Remainder on 1 August 1966.

Table 2. Leaf resistances on a clear day. 4 August 1966.

Time (EST) Hour	Leaf resistance (r_l) (diffusion meter method) (sec cm^{-1})	
1100	3.7	
1115	4.1	4.2
1130	4.5	
1145	4.6	
1200	4.7	
1215	4.8	4.4
1230	4.4	
1245	3.6	
1300	3.2	
1315	3.9	3.8
1330	3.9	
1345	4.3	
1400	4.3	
1415	4.9	
1430	3.6	4.4
1445	4.7	
1500	4.8	
Standard error of Mean (n - 5)	0.3	

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13 ABSTRACT Studies of resistance to gaseous exchange between large sunflower leaves and the bulk air in a crop canopy were made. Two components of the diffusive pathway for mass and sensible heat were evaluated; (a) the resistance from the interior of the leaf to the leaf surface, and (b) the resistance from the surface of the leaf through the leaf boundary air layer to the bulk air. It was found that: (a) leaf resistance not only displays diurnal trends but shorter fluctuations, and (b) boundary air layer resistance was significantly smaller than predicted from classical boundary layer formulae.			

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