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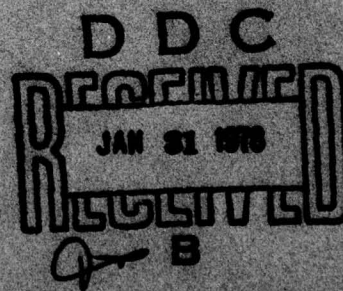
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HARMONIC GENERATION IN Nd:LASER-PRODUCED PLASMAS

There have been many observations of second harmonic light¹⁻⁸ emitted from laser-produced plasma. However, there have been no reported observations of third or higher harmonic emission in plasmas generated by a Nd:glass laser, although third harmonic light⁹⁻¹⁰ and higher harmonics¹¹ have been observed from CO₂ laser-produced plasmas. We report here the first observations of the third, fourth, and fifth harmonics lines produced when a high irradiance Nd:laser beam (1.06 μm) is focused onto a target. It is generally thought that integral harmonic emission is generated at or near the critical surface (where fundamental and local plasma frequencies are equal) and, therefore, can yield information on density gradients and electric fields in this region of the plasma.^{2,12} (A slower fall-off in the intensity of the harmonics would be an indication of a steeper density gradient or a higher field strength.) Thus, information on the harmonic emission is important in understanding the laser-target absorption process. These emissions may also interfere with diagnostic light probe studies such as shadowgraphy, interferometry, Faraday rotation, etc. which use light at the harmonic frequencies of the incident laser beam. As an example, a time-integrated photograph of third harmonic emitted light (3546 Å) taken through an interference filter (3546 $\pm$ 50 Å) is shown in Fig. 1a. The small spatial extent

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of the emission ($\sim 5 \mu\text{m}$), even though averaged over time and spread by refraction effects, is indicative of a very localized emission region and, perhaps, of steep density gradients near the critical density.

The experimental arrangement is shown in Fig. 1b. One beam (6-10 J) from the NRL Pharos II Nd:laser ($1.06 \mu\text{m}$) operating at 75-psec pulse duration (FWHM) was focused at normal incidence onto a polystyrene (CH) slab target in an evacuated ($\leq 10^{-4}$ Torr) chamber using an f/1.9 aspheric lens. Prepulses were monitored on each shot and did not exceed $1 \mu\text{J}$ in these experiments. The half-energy focal diameter was measured¹³ to be $30 \mu\text{m}$ yielding an average irradiance of $\sim 10^{16} \text{ W/cm}^2$. A focal-shift monitor¹⁴ ensured that the target was in focus on each shot. The fractional absorption of laser light was inferred to be $\sim 50\%$ by measuring the light backscattered through the focusing lens and outside the lens with calorimeters.¹⁵ Simultaneous observations of the harmonic light were made with a 3/4-m Czerny-Turner mount spectrograph (resolution in first order $\sim 1 \text{ \AA}$ for a slit width of $100 \mu\text{m}$) and an interference filter - photovoltaic silicon photodiode combination (E.G.G. UV-100 B biased to $-22\frac{1}{2} \text{ V}$). These detectors viewed the target at 45° to the target normal and in a (horizontal) plane at 45° to the plane of the incident laser's electric field. The rise time of the photodiode and recording oscilloscope was $\sim 2 \text{ nsec}$, and the photodiode was proved linear in the range discussed here. Care was taken to block the higher order spectra from entering the spectrograph. An achromatic optical train was achieved by using

a spherical and a flat mirror to focus the plasma onto the spectrograph slit. Three types of film were used in the spectrograph to cover the range of wavelengths and sensitivities encountered in this experiment, i.e., Kodak I-Z plates, Kodak Royal-X Pan, and Polaroid Type 57. (All spectra required only one shot.)

Microdensitometer scans of representative time-integrated spectral profiles for the fundamental emission, ω_0 , and the integral harmonics through the fifth, $5\omega_0$, are shown in Fig. 2. The wavelength of each harmonic is indicated by a vertical dashed line. The peak of the fundamental line is blue-shifted by about $13 \text{ \AA}$. If the shift ($13 \text{ \AA}$) is a Doppler shift, this would correspond to an average outward velocity of the reflecting surface of $3 \times 10^7 \text{ cm/sec}$. A blue shift of the fundamental has also been observed by others.¹⁰ However, the centroids of the rest of the harmonic lines are red-shifted. It appears that the scattered fundamental line is produced by a different mechanism or spectral region than that which produces the harmonic lines. The $5\omega_0$ line at $2128 \text{ \AA}$ is just resolved above the plasma continuum. To observe this harmonic line, the spectrograph slit was opened to $500 \text{ }\mu\text{m}$ so that the resolution was reduced to $\sim 5 \text{ \AA}$. Because of the expected low signal level, we did not attempt to observe $6\omega_0$ at $1773 \text{ \AA}$. All of the observed harmonic spectral widths are large compared to that of the incident laserwidth ($\sim 1 \text{ \AA}$).

The spectral widths (FWHM) for the second, third, and fourth harmonics in wavelength and frequency units (to facilitate a comparison with theory) are shown in Table 1. It is seen that the spectral half-

widths in frequency units change at most by 60%, whereas the half-widths in wavelength units change by almost a factor of 3. The observed spectral broadening thus involves mechanisms other than Doppler broadening (either due to random or non-local mass motion) such as instabilities and magnetic fields.¹⁶

Assuming shot-to-shot reproducibility of the over 150 data shots, relative (total) intensities of the harmonics measured with the silicon photodiode-interference filter are shown in Fig. 3. These values of relative intensity are corrected for the transmission and the band pass of the filters and the spectral responsivity (given by the manufacturer) of the silicon photodiode. The error bars indicate the estimated error in the measurement due to the spectral responsivity uncertainty and the signal-to-noise ratio. There is a large decrease ($\sim 2 \times 10^2$) in the intensity of the second harmonic relative to the reflected or scattered fundamental; however, the next three harmonics decrease in steps of less than a factor of 10. These ratios of harmonic line intensities are the same (within a factor of 5 over a range of 5 orders of magnitude) as the corresponding harmonic intensities reported by the National Research Council Group¹¹ in Ottawa, Canada, using a CO_2 laser at an irradiance of $I = 10^{14} \text{ W/cm}^2$. This fact suggests that a comparable interaction takes place in Nd: and CO_2 : laser irradiation which scales as $I\lambda^2$.

In summary, we report emission spectra and relative intensities of the fundamental and second through the fifth harmonics of the incident laser line when a Nd:laser (10^{16} W/cm^2) is focused onto a

CH slab target. The emission is sufficiently intense that the contrast in diagnostic probe studies at these harmonic wavelengths would be adversely affected. (This difficulty has been obviated at NRL by shifting the probe wavelength using a Raman cell to a nonharmonic wavelength¹⁷.) Although the present theories of optical absorption of intense light in plasma have not yet predicted the ratios, widths, and shifts at the higher harmonics, these harmonics should be directly related to the absorption process and would aid in estimating density gradients and electric fields produced in the absorbing layer. Hopefully, these experimental results will stimulate the appropriate theoretical explanations. An interesting result of this study is that the spectra and relative intensities of the harmonics produced in Nd:laser-target interactions at 10^{16} W/cm² are strikingly similar to recent results¹¹ found in CO₂ laser-produced plasmas at 10^{14} W/cm².

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- a. Work sponsored by the U.S. ERDA.
- b. Permanent address: Dept. of Physics and Astronomy, University of Maryland, College Park, MD.
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Table 1 — Spectral half-widths (FWHM) of the harmonics in frequency and wavelength units. The estimated error from the photographic plate calibration and a small number of representative shots is $\pm 10\%$.

HARMONIC	$\Delta\omega_{\text{FWHM}}$ [rad sec ⁻¹]	$\Delta\lambda_{\text{FWHM}}$ [Å]
$2\omega_0$ (5320 Å)	1.9×10^{13}	29
$3\omega_0$ (3547 Å)	1.7×10^{13}	11
$4\omega_0$ (2660 Å)	2.7×10^{13}	10

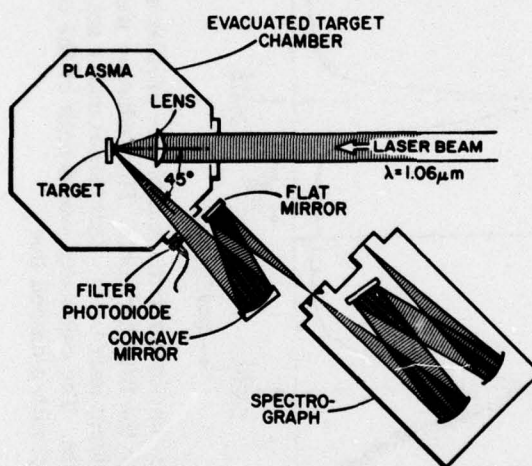
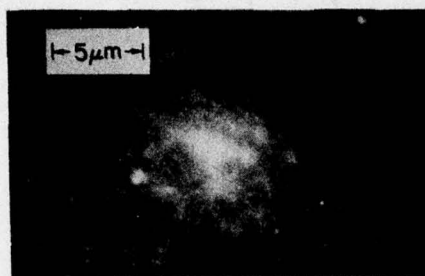


Fig. 1 — a) A time-integrated photograph of the third harmonic light ($3546 \text{ \AA}$) emitted tangential to the target surface when a 75-psec, $\sim 10^{16} \text{ W/cm}^2$ laser pulse strikes a CH_2 plane target. The laser beam was incident from the left. A $3547 \text{ \AA}$ interference filter with a bandwidth of $50 \text{ \AA}$ was used in front of the camera.
b) The experimental arrangement.

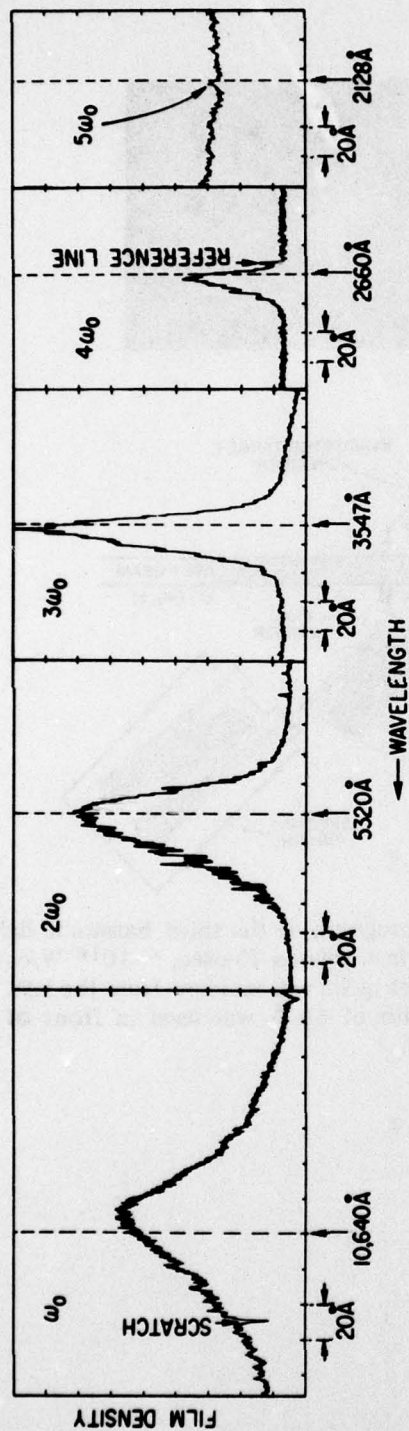


Fig. 2 — Representative time-integrated spectral line profiles are shown for the fundamental through the fifth harmonic. This radiation was observed at 45 degrees from the incident laser beam. Extraneous marks are indicated on the ω_0 and $4\omega_0$ profiles. The calculated wavelength of the center of each spectral line is indicated with a dashed line.

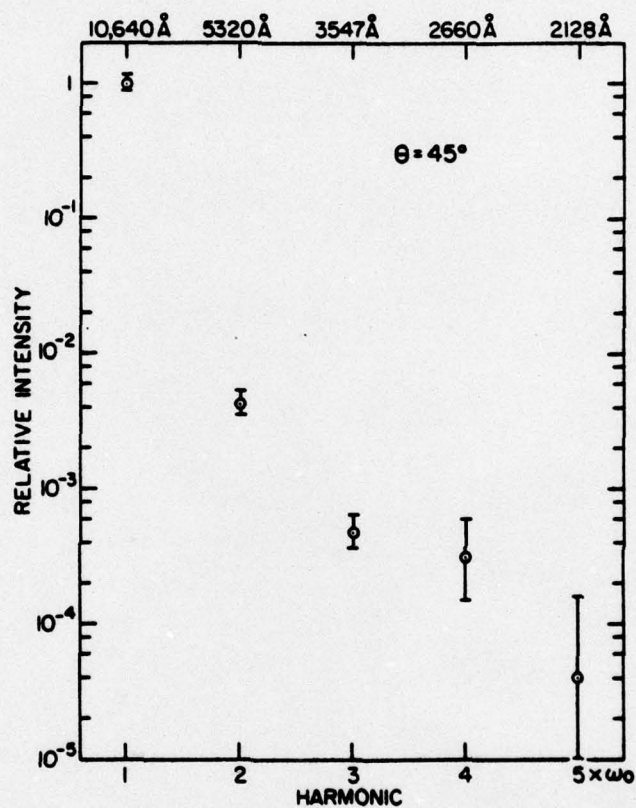


Fig. 3 — Intensities of the harmonics relative to the fundamental emission at 45° to the incident laser beam. The error bars are estimated from the spectral responsivity uncertainty and the signal-to-noise ratio.