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Refinement of plume models in
the infrared spectral region

By

Kenneth E. Harwell

University of Tennessee Space Institute
30 June 1978

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Final Report for
Contract DAAK40-77-C-0032

REFINEMENT OF PLUME MODELING IN
THE INFRARED SPECTRAL REGION

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Tullahoma, Tennessee 37388

30 June 1978

Approved for public release; distribution unlimited

Prepared for
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ABSTRACT

A research program was conducted to develop a validated numerical model for the prediction of radiation emitted in the infrared wavelength spectral band from turbojet and rocket exhaust plumes. Two existing gas dynamics models (LAPP and REP3) were employed to calculate the flow field properties (temperature and species concentrations) required as input for the radiation computer codes. A new finite difference computer code, which was developed as part of this project, and includes the calculation of imbedded shock waves in the exhaust plume along with finite-rate chemistry and turbulent mixing, was only partially successful. A refined version of the MIRADCOM Radiation Model was developed and was used successfully to predict CO_2 and H_2O radiation emitted from several cases for which experimental data were available.

1. INTRODUCTION

During the past few years, considerable effort has been expended to design infrared seeker systems for ground-to-air and air-to-air missile guidance systems. Since the next generation of seeker systems must possess head-on capability, infrared radiation from the missile exhaust plume is an important component of the overall infrared signature of a missile system. In order to characterize completely the missile seeker-target interaction, it is necessary to develop infrared radiation predictive models which have the capability of predicting all the pertinent physical phenomena associated with the plume infrared signature.

Infrared radiation emitted from the hot gases in an exhaust plume is characterized by the types of radiating species and their number densities and the environmental conditions (temperature and pressure). All radiative predictive models require a knowledge of these parameters, so that the first step in developing a model is to develop or to obtain a gas dynamics model which yields the spatial distribution of radiating specie concentrations, temperature, and pressure.

Unfortunately, the exhaust plumes of turbojet aircraft and rockets are characterized by the existence of a complex shock wave structure and the presence of a turbulent viscous mixing layer between the inner hot jet flow and the external freestream flow. The inner flow usually consists of a high-temperature, high-velocity, chemically-reacting gas produced by the combustion of a hydrocarbon fuel while the external flow usually is low-temperature, low-velocity air that flows around the exterior of the aircraft or missile body.

Even though considerable progress had been made in developing workable gas dynamics computational models,[1-17] at the time this project was initiated, there was no available model which adequately predicted the mixed supersonic and subsonic flow regions present in the central core of the exhaust plumes. This report describes the results of a research program which was initiated to develop a finite difference computer code which would include the calculation of imbedded shock waves and Mach discs along with finite-rate chemistry and turbulent mixing phenomena.

During the course of this project, a gas dynamics computer code developed by Dr. Brian Spalding for the Rocket Propulsion Establishment in the United Kingdom was made available to the U.S. Army MIRADCOM as part of the TTCP program with Great Britain. This program was modified by the Principal Investigator to operate on the MIRADCOM computer. The results of a few limited calculations are described in this report.

The MIRADCOM Radiation Model (developed by H. Tracy Jackson, Jr. and extended by K. E. Harwell) was refined as part of the research effort described in this report. Results obtained using this model will be described and a comparison made between the theoretical predictions and experimental data.

2. THEORETICAL FOUNDATIONS

2.1 Gas Dynamics Theoretical Models

A complete review of the theory for turbulent exhaust plumes is beyond the scope of this report. A description of the flow theories is given by Jackson, Poslajko, and Harwell in Ref. 18,

The appropriate equations of motion for the gas are the conservation equations:

Conservation of Mass:
$$\frac{\partial \rho}{\partial t} + \nabla \cdot \rho \vec{V} = 0$$

Conservation of Momentum:

$$\rho \left[\frac{\partial \vec{V}}{\partial t} + (\vec{V} \cdot \nabla) \vec{V} \right] = - \nabla p + \nabla \cdot \underline{\underline{\tau}}$$

Conservation of Energy:

$$\rho \frac{DE}{Dt} = \rho \frac{D \left[e + \frac{\vec{V} \cdot \vec{V}}{2} \right]}{Dt} = - \nabla \cdot (p \vec{V}) + \nabla \cdot \underline{\underline{\tau}} \cdot \vec{V} - \nabla \cdot \vec{q}$$

After supplementing these conservation equations with an Equation of State

$$e = e(T, p)$$

and constitutive relations relating the shearing stress τ to strain of the fluid element and ultimately to the velocity gradients in the flow, a complete system of flow equations is obtained. For example, the laminar shear stress can be expressed as

$$\tau_{ij} = \lambda \nabla \cdot \vec{V} \delta_{ij} + 2\mu \left(\frac{\partial V_j}{\partial X_i} + \frac{\partial V_i}{\partial X_j} \right)$$

The theoretical approach to the development of the turbulent shear stress model is not described here, but the model employed was the eddy diffusivity model proposed by Donaldson and Gray [19] for the compressible free mixing of a primary jet with quiescent air:

$$\tau_t = -\bar{\rho} \overline{u'v'} = K \bar{\rho} L U_s \frac{\partial \bar{u}}{\partial r}$$

where K is the mixing rate factor, L is a typical local scale length, $\bar{\rho}$ is the local density, and U_s is a typical velocity difference. These parameters are defined in Refs. [12 and 19].

It had been intended to incorporate the Turbulent Kinetic Energy (TKE) turbulence model approach into the computer model. However, due to problems with the basic computer code, this was not accomplished. In this model [see Ref.11], an additional conservation equation for the turbulent kinetic energy is added to the previous set of conservation equations. This equation for an axisymmetric flow becomes

$$\rho \bar{u} \frac{\partial \bar{k}}{\partial x} + \rho \bar{v} \frac{\partial \bar{k}}{\partial r} = \left[\frac{\rho}{P_{rk}} \epsilon_t \frac{\partial \bar{k}}{\partial r} \right] + \rho \epsilon_t \left(\frac{\partial \bar{u}}{\partial r} \right)^2 - a_2 \frac{\bar{k}^{3/2}}{\ell_k}$$

where

$$\bar{k} = \frac{1}{2} \left[\overline{(u')^2} + \overline{(v')^2} \right]$$

$$\epsilon_t = \tau_t / \bar{\rho} \frac{\partial \bar{u}}{\partial r}$$

$$\tau_t = a_1 \bar{\rho} \bar{k}$$

and

$$P_{rk} = \bar{\rho} C_p \frac{\epsilon_t}{\bar{k}}$$

a_1 and a_2 are Universal empirical relationships and ℓ_k is a length scale for the TKE. The reader is referred to Harsha [20] for a detailed formulation of this model. Suffice it to state that the addition of the TKE equation in the matrix formulation of the computer model is straightforward and can be accomplished without major difficulty.

The chemical reaction relations employed in the gas dynamics model was identical to those employed in the Low Altitude Plume Program (LAPP) developed for the U.S. Air Force Rocket Propulsion Laboratory [5]. The reader is referred to the LAPP report [5] for a detailed description of the formalism.

2.1.1 Gas Dynamics Finite Difference Computer Model

A finite difference numerical scheme was used to approximate the gas dynamics conservation equations described in the previous section.

A method developed by Daniel Matuska of the USAF Weapons Laboratory for computing hydrodynamic properties of fluids was modified to include the viscous stress tensor.

The finite differencing technique used was the standard one in which the derivative of the function F is defined as $[F(X + H) - F(X)]/H$ where H is finite.

The degree of success in using the finite difference technique to solve the differential equations depends on the "stability," "convergence" and "compatibility" of the method. The method is stable if any bounded starting procedure yields a uniformly bounded solution of the difference equation as H goes to zero. A method is convergent if the values of $Y(N)$, the solution of the finite difference expression approximating the difference equation, tend to the values of $Y(T)$ of the exact solution. The method is compatible

if the truncation error goes to zero as H goes to zero. The difference between the true solution of the differential equation and the finite difference equation method is influenced by the step size H and may require very small values of H to be stable and convergent.

One of the major advantages of the finite difference approach is that the problem is reduced to an initial value problem with well defined boundary conditions which results in a unique time-dependent solution. The major advantage in terms of the present problem of interest, is that the mixed subsonic/transonic/supersonic flow problem can be solved without using a patching solution.

The main disadvantages of the approach are that the complete boundary must be defined as a function of time, that small steps may be required which results in a fine computational mesh which in turn requires a large computer memory if realistic resolution is desired, and that computational times may become prohibitive.

The calculation procedure used in the present program is to treat the problem in two phases. In the first phase, the "Lagrangian terms" (coordinate system fixed to the fluid element) are integrated. In the second phase, the convective terms are integrated (mass momentum and energy is transported across boundaries). The details of the calculation procedure are described in Ref. [21].

2.2 Infrared Radiation Theoretical Model

A numerical infrared radiation band model was developed to predict the infrared spectral and spatial radiation intensity distributions in exhaust plumes. The model essentially extends the models of Jackson [22,23], which were developed for CO₂ radiation from aircraft plumes, to include water vapor radiation and to include the capability of treating higher temperatures than the Jackson code.

The statistical band model and spectral parameters used in the code were developed for application to rocket plume studies [24]. The transmissivity τ_R , in the model becomes

$$\tau_R = \exp [-P/(1 + P^2/4Q)^{1/2}]$$

where $P = \int k du$

$$Q = \int \frac{k\gamma}{d} du$$

$k(\omega, T)$ = absorption coefficient

u = optical depth

γ = average line width, and

$d(\omega, T)$ = average line spacing

The absorption coefficients k and reciprocal line spacings, $1/d$, are tabulated as a function of frequency, ω and temperature, T in Ref. [24].

The average line width was computed using the relation

$$\gamma = [120.1 P_{H_2O}/T] + [.09(P_{H_2O} + P_{N_2}) + .04 P_{O_2} + .12 P_{CO_2}] [273/T]^{1/2}$$

2.2.1 Infrared Radiation Computer Model

The infrared radiation computational model combined a water vapor emissivity formalism developed by Reardon [25] with the plume integration

scheme developed by Jackson [23] for carbon dioxide emission from aircraft plumes.

The computation of radiance from a jet exhaust plume was accomplished by dividing the plume into sufficiently small incremental blocks so that average conditions could be assigned and the blocks treated as homogeneous. The approach used in the geometric division of the plume is described in detail by Jackson [23]. The plume is subdivided into slabs parallel to the line of sight. Each slab is then cut into horizontal strips which lie along the line-of-sight vector.

The radiation from each of these strips is computed by integration using the modified Curtis-Godson approximation [24].

Integration is accomplished over the length of the strip to obtain the parameters P and Q as shown earlier,

$$P = \int k du$$

$$Q = \int \frac{kY}{d} du$$

The transmittance $\tau_R = \exp [-P/(1 + P^2/4Q)^{1/2}]$ is then computed.

Radiance is then calculated by integrating over the strip emissivity. For a homogeneous block, Radiance = (Blackbody) (emissivity) so that the inhomogeneous strip radiance is

$$\text{STRIP RADIANCE} = \int (\text{Blackbody}) d\epsilon$$

Since $d\epsilon = - d\tau_R$ the computational algorithm

$$\text{STRIP RADIANCE} = - \sum_i (\text{Blackbody})_i \Delta\tau_{R_i}.$$

The algorithms used for calculating the radiation from inhomogeneous strips were tested by applying them to homogeneous strips and comparing with direct radiation calculations. The differences in calculated values were less than

one percent.

A listing of the infrared computer program is given in Appendix I.

A list of the input requirements is given in Appendix II. Appendix III includes a listing of a sample case with output.

3. DISCUSSION OF RESULTS

This section contains a discussion of the results obtained as part of the contract effort. The procedure used to validate the infrared radiation computer code is described in Section 3.1. Section 3.2 presents the results of the infrared radiation code predictions for a turbojet flight test for which there are experimental data. The prediction of the infrared radiation emitted from a small kerosene/gaseous oxygen rocket engine exhaust plume is presented in Section 3.3. A discussion of the progress in developing the Gas Dynamics Finite Difference Computer Model is described in Section 3.4. Some preliminary results obtained using the REP3 Computer Code are described in Section 3.5.

3.1 Validation of the UTSI Infrared Radiation Code

In order to validate the Infrared Radiation Code, the spectral radiance of water vapor and carbon dioxide mixtures experimental data are available [Refs.24,27-29]. Sukanek and Davis [26] performed a similar assessment of the NASA band model formulation developed by Reardon [25]. It was decided that several selected cases used by Sukanek and Davis would be used for the evaluation since this would provide a comparison of the present model calculations with the experimental data and with the predictions of the NASA/Reardon Radiation Code. The cases selected for the comparison are given in Table 3.1.

The spectral radiance of 60-cm thick slabs of hot H₂O vapor and CO₂ were computed for the 2.7- μ and 4.3- μ wavelength bands. Figures 3.1-3.8 present the results of the computations.

The comparison between the measured and predicted values of the spectral radiance in the 2.7- μ band for an isothermal slab of water vapor is presented

TABLE 3.1 CASES CONSIDERED FOR COMPARISON

Subject	Fig. No.	Particle Pressure (mmHg)					Temperature Profile	Length (m)
		H ₂ O	CO ₂	N ₂	CO	H ₂		
Isothermal Hot H ₂ O	1	38.0		38.7			1215	0.6
Isothermal Hot CO ₂	2		38.0	38.0			1202	0.6
Isothermal Hot CO ₂	3		7.58	68.43			1200	0.6
Isothermal Mixture	4	35.26	7.58	27.09	1.44	4.64	1200	0.6
Isothermal Mixture	5	35.26	7.58	27.06	1.47	4.62	1200	0.6
Nonisothermal H ₂ O	6	670					382/537/723/953/1128/1160/990/751/558/389	0.6
Nonisothermal CO ₂	7		760				386/528/719/953/1130/1160/975/737/541/387	0.6
Nonisothermal Mixture	8	57	28	675			378/537/723/958/1127/1158/990/752/555/383	0.6

NOTE: The temperature of the gas was either homogenous, in which case one temperature is reported, or was inhomogeneous, in which case the temperatures are given at the center of each 6 cm wide cell.

in Fig. 3.1. At the center of the spectral band the computed values are slightly lower than the measured ones. The NASA code predicts slightly higher values than the UTSI code, but in general the agreement is fairly good between the calculated and measured values. The calculated values are within 20% of the measured ones.

The measured and predicted values of spectral radiance in the 2.7- μ and 4.3- μ wavelength bands for an isothermal slab of CO_2 are presented in Figs. 3.2 and 3.3. In contrast to the H_2O predictions, the predictions of the spectral radiance for CO_2 do not agree very well with the measurements. Errors of 30 to 50 percent are observed. Similar disagreement was found by Lindquist and others [27-29] who attributed to the probable cause of the disagreement to the lack of knowledge of the CO_2 broadening parameters at low pressures. Ludwig and others [24] show better agreement between the predictions and measurements at higher pressures ($P_{\text{CO}_2} = 1-2$ atmospheres).

Figures 3.4 and 3.5 present the results for the spectral radiance of an isothermal mixture of CO_2 and H_2O in the 4.3 and 2.7- μ bands, respectively. The calculated values are higher than the measurements in the center of the spectral bands. Both codes yield results which are in fair agreement with the measured values.

Figure 3.6 presents the results for a nonisothermal slab of water vapor. There is fairly good agreement between the predicted and measured values, but the present code appears to predict values which are below the measured and calculated values using the NASA code.

Figure 3.7 presents the comparison between the predicted and measured values of spectral radiance in the 2.7- μ spectral band of a nonisothermal slab of CO_2 . The agreement between the predicted and measured values is quite good.

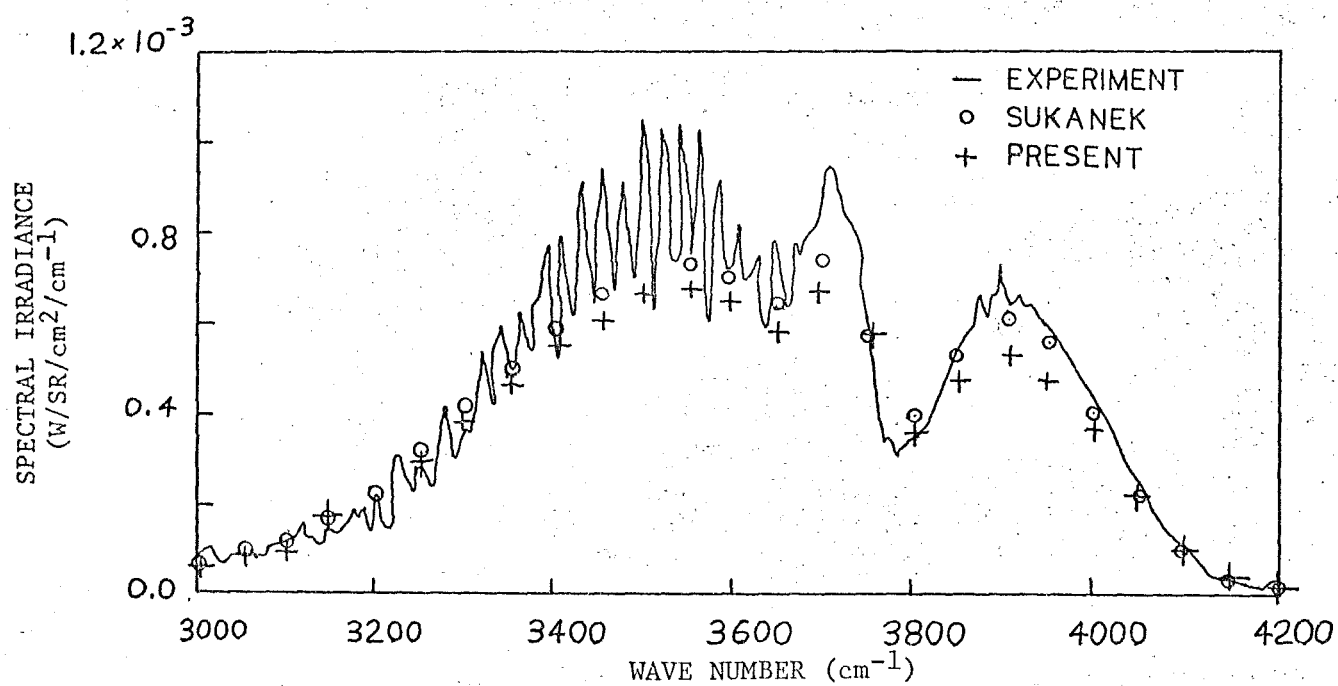


Fig. 3.1 SPECTRAL RADIANCE IN THE 2.7- μ m BAND FOR AN ISOTHERMAL SLAB OF H₂O VAPOR.

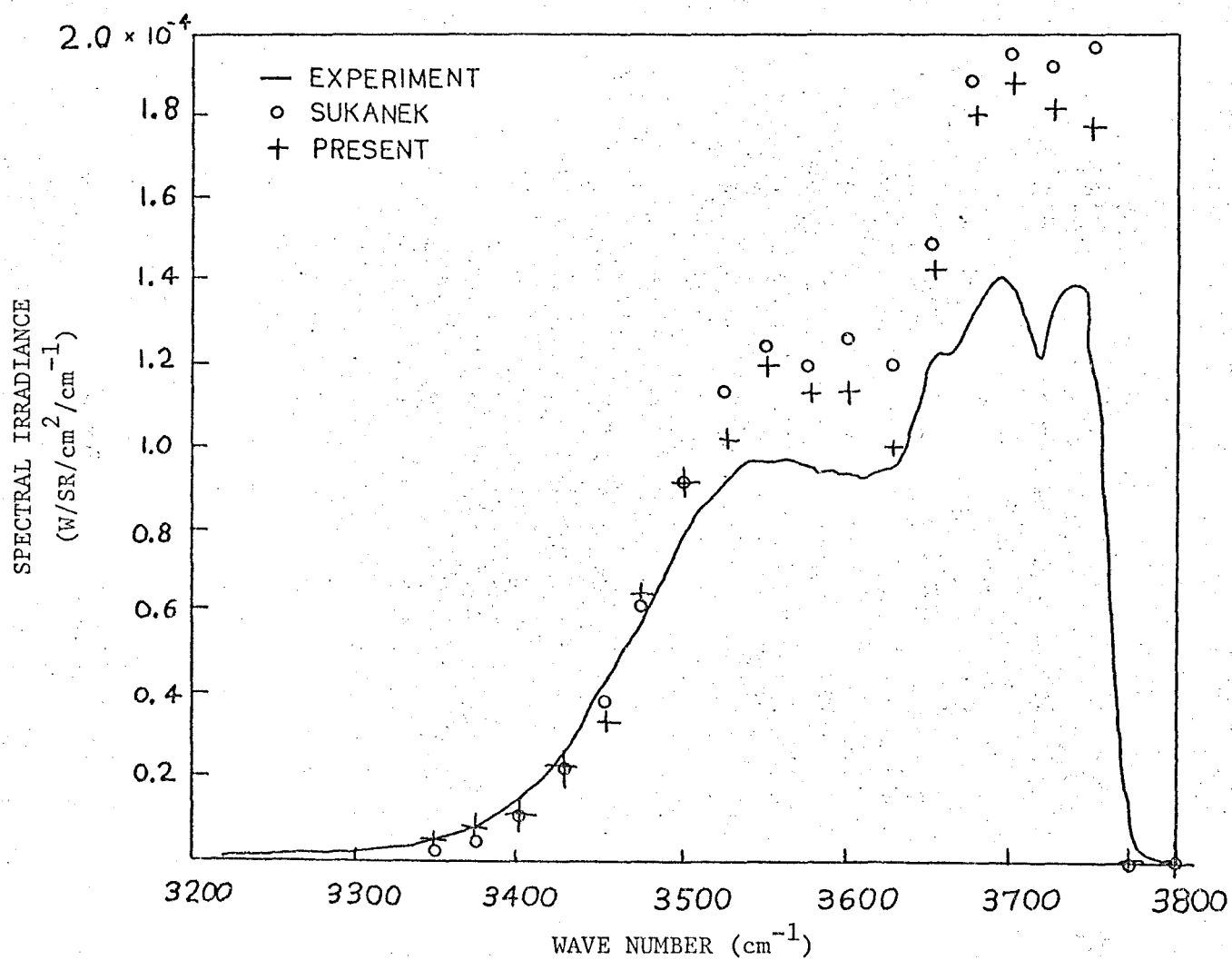


Fig. 3.2 SPECTRAL RADIANCE IN THE 2.7- μ m BAND FOR AN ISOTHERMAL SLAB OF CO₂.

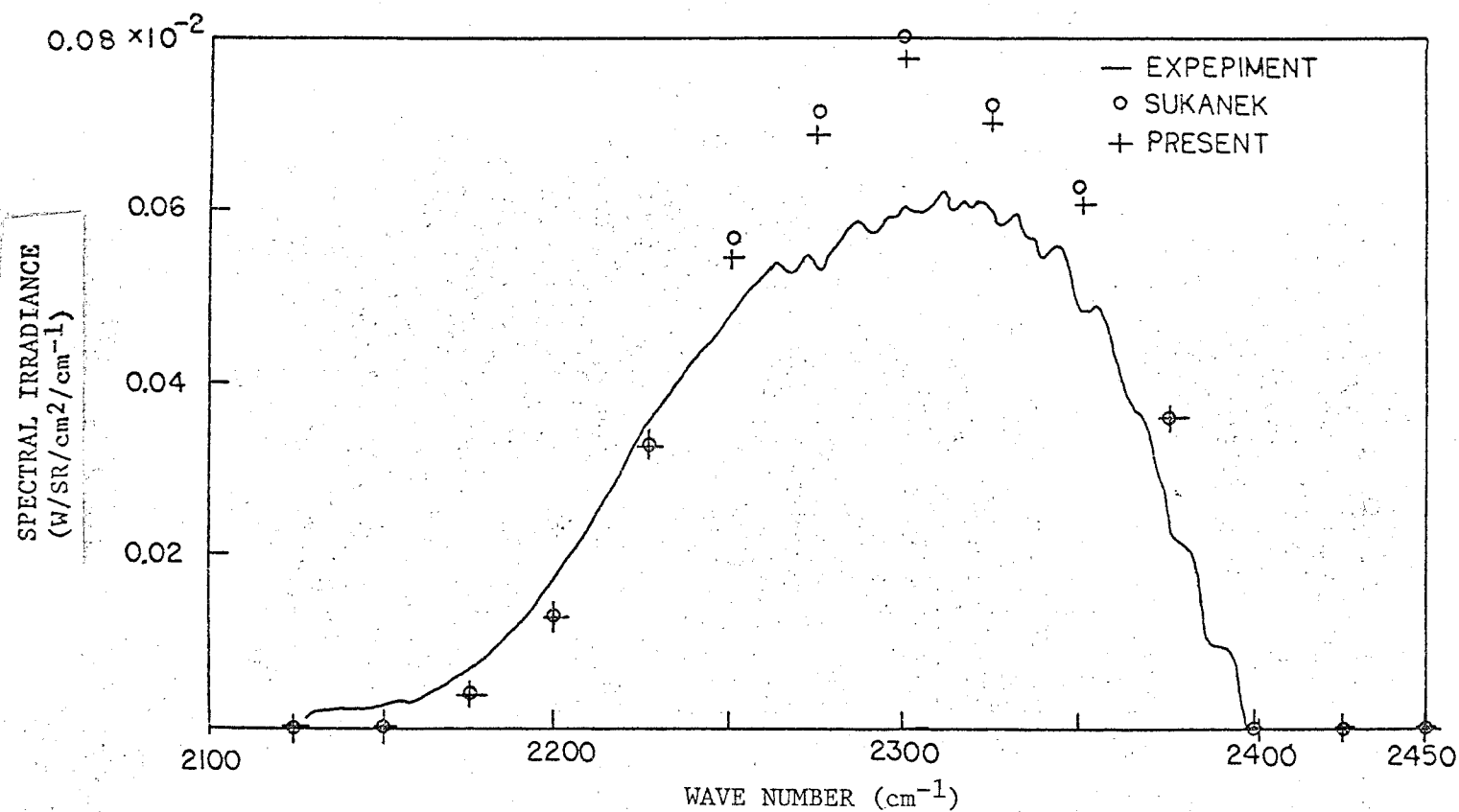


Fig. 3.3 SPECTRAL RADIANCE IN THE 4.3- μ m BAND FOR AN ISOTHERMAL SLAB OF CO₂.

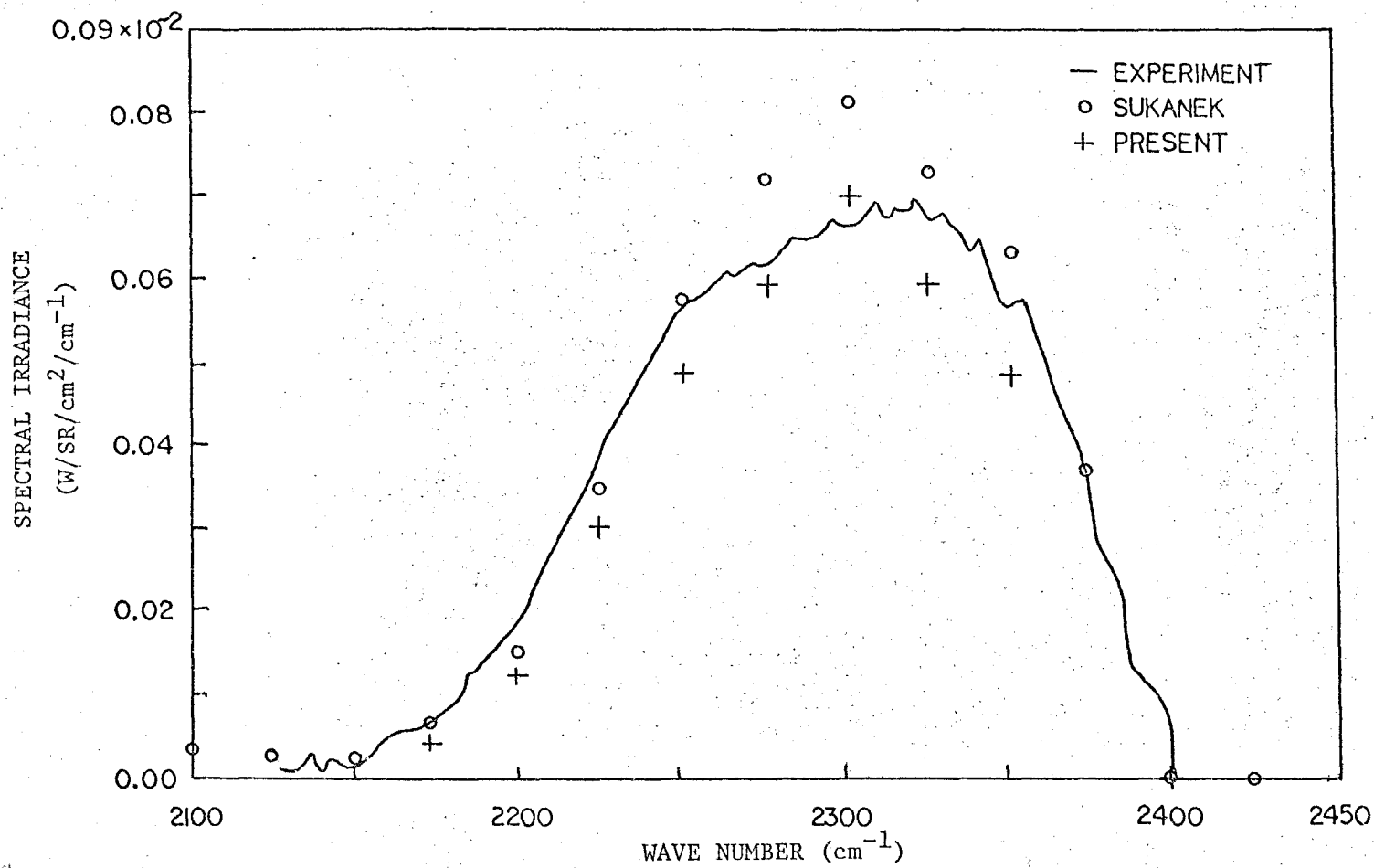


Fig. 3.4 SPECTRAL RADIANCE IN THE 4.3- μ m BAND FOR AN ISOTHERMAL MIXTURE OF CO₂ and H₂O.

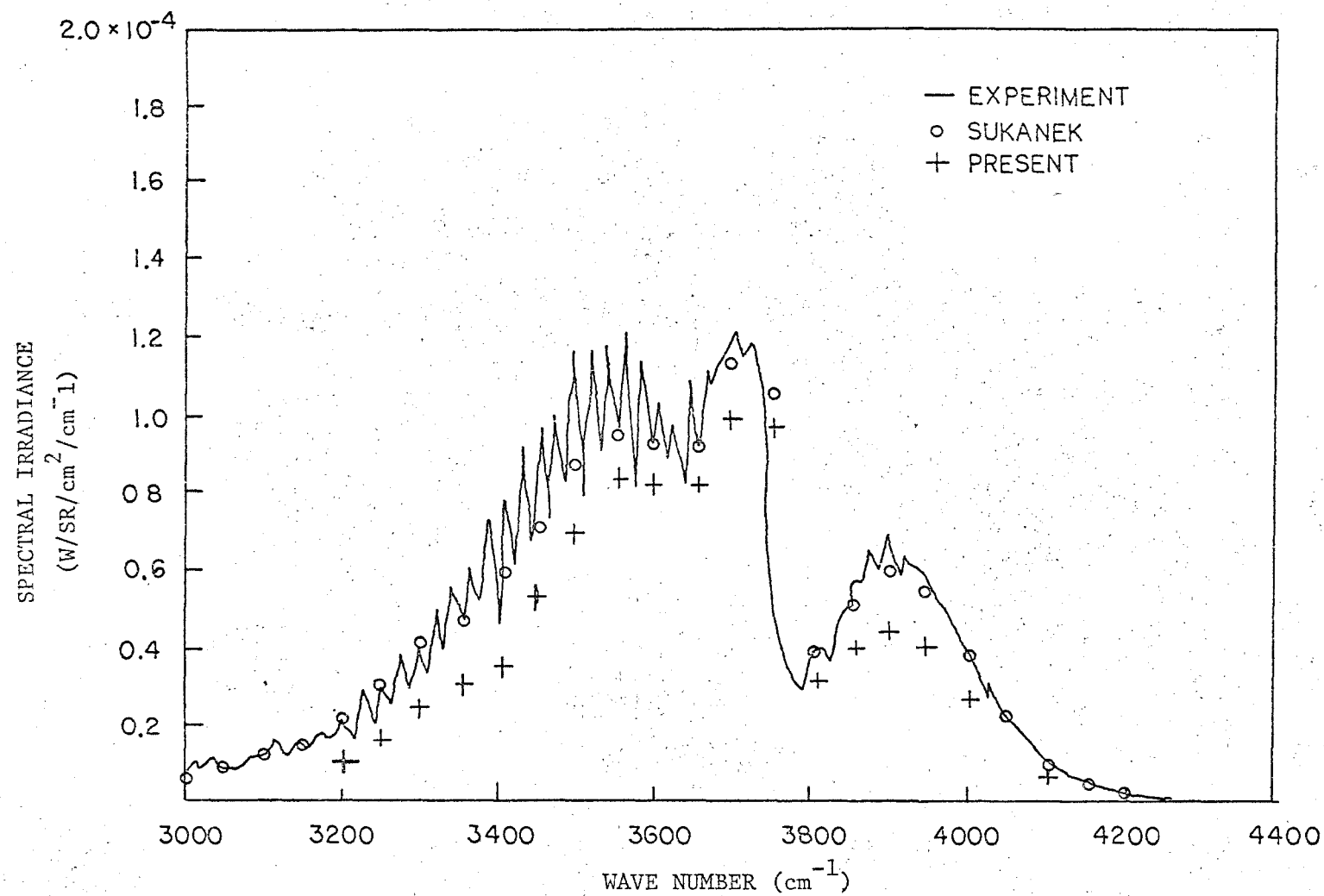


Fig. 3.5 SPECTRAL RADIANCE IN THE 2.7- μ m BAND FOR AN ISOTHERMAL MIXTURE OF CO₂ and H₂O.

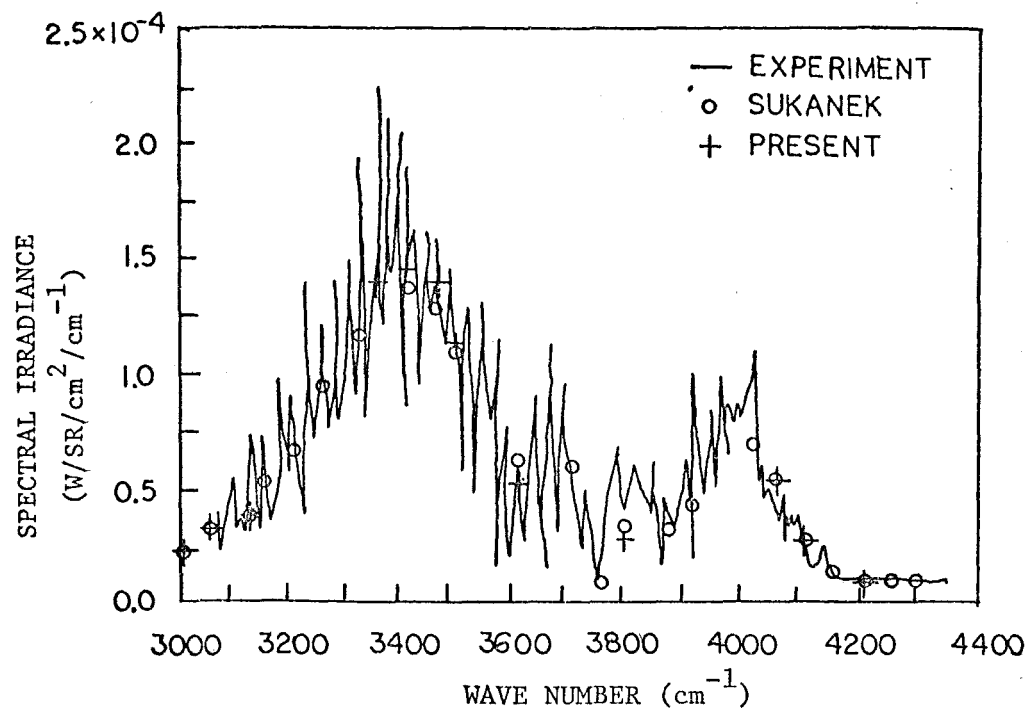


Fig. 3.6 SPECTRAL RADIANCE IN THE 2.7- μm BAND FOR A NONISOTHERMAL SLAB OF H_2O VAPOR.

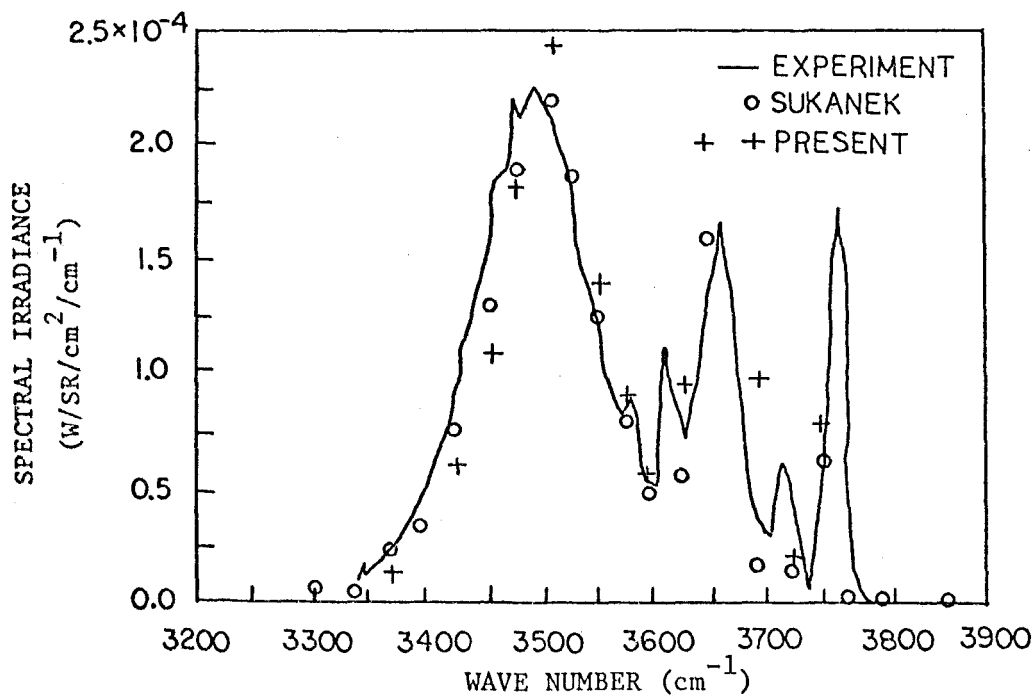


Fig. 3.7 SPECTRAL RADIANCE IN THE 2.7- μm BAND FOR A NONISOTHERMAL SLAB OF CO_2 .

Figure 3.8 presents the results for a nonisothermal mixture of CO_2 and H_2O . There is fairly good agreement over the entire $2.7\text{-}\mu$ wavelength band.

In summary, the UTSI radiation code appears to yield results which agree fairly well with the NASA code even though there are some differences. With the exception of the low pressure CO_2 cases, the predicted values are within $\pm 20\%$ of the measured values.

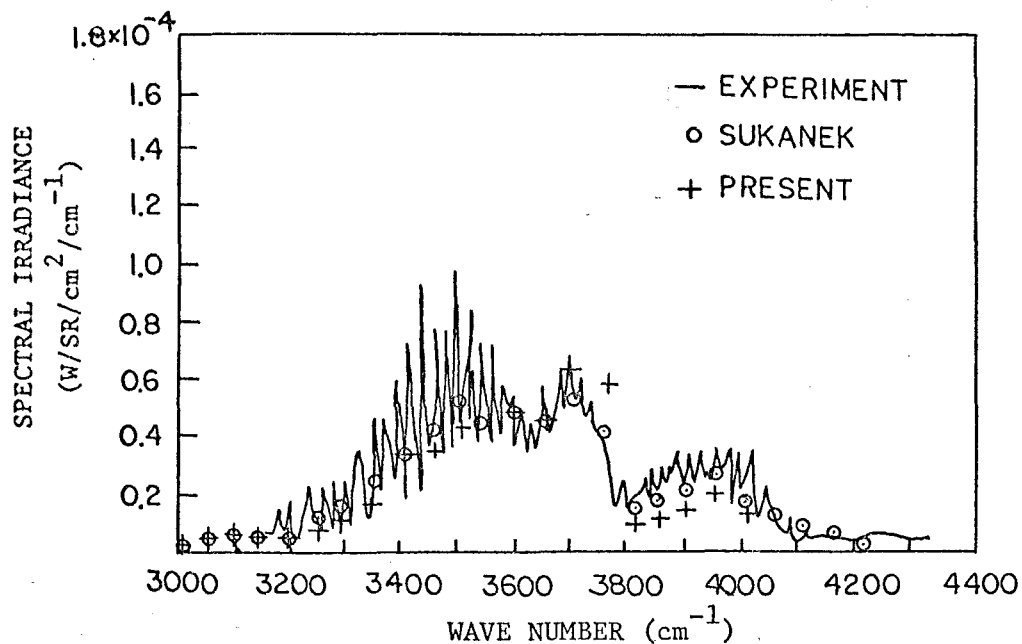


Fig. 3.8 SPECTRAL RADIANCE IN THE $2.7\text{-}\mu\text{m}$ BAND FOR A NONISOTHERMAL MIXTURE OF CO_2 AND H_2O .

3.2 Infrared Radiation Predictions for Turbojet Flight Test

The UTSI Infrared Radiation Code was utilized to calculate the infrared radiation emitted in the wavelength interval from 4 to 5 microns from a turbojet aircraft exhaust for which experimental data are available.

H. Tracy Jackson, Jr. at MIRADCOM provided the data given in Table 3.2.

The Low Altitude Plume Program (LAPP) was exercised for the input data given in Table 3.2. The first portion of the LAPP code output for this case, labeled "FLT1-PASS 7," is included as Figs. 3.9-3.11. The output data from the LAPP code was then used as input data for the Infrared Radiation Code. Refer to Fig. 3.12 for a typical set of input data.

The results of the computations are presented in Figs. 3.13-3.15. The CO_2 contribution to the radiation per length of plume (Watts/SR/cm) is given in Fig. 3.13 while the H_2O contribution is given in Fig. 3.14. The H_2O radiation is very small compared to the CO_2 radiation.

The spatial distribution of the radiation is given in Fig. 3.15. Also, shown on the figure is the centerline temperature distribution. The total radiation emitted from the plume in the 4.0 to 5.0 micron wavelength interval was 165.7 Watts/SR. This radiation is at the plume and is uncorrected for detector response characteristics or for atmospheric attenuation between the plume and the detector.

It had been anticipated that the Jackson MIRADCOM Radiation Code would be used to calculate the radiation for the same test flight so that a comparison could be made between the two codes. Unfortunately, the results from the MIRADCOM Code did not become available prior to the completion of the project.

TABLE 3.2

TURBOJET FLIGHT TEST DATA

Flight 10, Pass 7

TPT = 790° K

Vel = 512 Kts. = 263 m/sec.

Asp α = 19°

O/F = 62

Range - 2212 meters

Too = 526° R

ρ_{∞} = 1955 lb/ft²

C^* = 1124 ft/sec.

Air Flow 182 lb/sec., 99% RPM

V jet = 1801.6 ft/sec.

ARMSTRONG RESEARCH LABORATORIES PRINCETON N.J.
A FAST COMPUTER PROGRAM FOR NONEQUILIBRIUM ROCKET PLUME PREDICTIONS
THIS VERSION INCLUDES SHOOT AND MODEL 6 MODIFICATIONS

F L11- PASS 7

PRESSURE(CONSTANT) = 1.0000000E 00 ATMOSPHERES

NOZZLE RADIUS= 8.9999998E-01 FEET

LEWIS NUMBER(CONSTANT)= 1.0000000E 00 PRANDTL NUMBER(CONSTANT)= 6.9999999E-01

X INITIAL(FEET)= 0.0

X FINAL(FEET)= 3.0000000E 01

PRINT INCREMENT= 1.0000000E 00

MINIMUM STEP SIZE= 9.9999994E-11

DDALDSON/GRAY VISCOSITY MODEL

		JET	EDGE
TEMPERATURE(DEG. KELVIN)		7.9000000E 02	2.9200000E 02
VELOCITY (FEET/SECOND)		1.8000000E 03	1.0000000E 02
MOLE FRACTION	H	2.8829817E-07	2.8849433E-07
MOLE FRACTION	H2	2.8829817E-07	2.8849433E-07
MOLE FRACTION	H2O	3.3150621E-02	2.8849433E-07
MOLE FRACTION	CO	2.8829817E-07	2.8849433E-07
MOLE FRACTION	CO2	3.2630604E-02	3.2879249E-04
MOLE FRACTION	O	2.8829817E-07	2.8849433E-07
MOLE FRACTION	OH	2.8829817E-07	2.8849433E-07
MOLE FRACTION	O2	1.5670282E-01	2.0846844E-01
MOLE FRACTION	N2	7.7751452E-01	7.9120100E-01

	REACTIONS BEING CONSIDERED	KK=A*EXP(B/RT)/T**N	A	N	B	(MOLECULE-MU-SEC UN)
1	O + O + M = O2 + M		1.000E-29	1.0	0.0	
2	H + H + M = H2 + M		1.000E-29	1.0	0.0	
3	O + H + M = OH + M		1.000E-29	1.0	0.0	
4	H + OH + M = H2O + M		1.000E-29	1.0	0.0	
5	CO + O + M = CO2 + M		5.000E-29	1.0	-4000.0	
6	OH + O + M = H2O + O		1.000E-11	0.0	-1000.0	
7	OH + H2 = H2O + H		4.000E-11	0.0	-5500.0	
8	O + H2 = OH + H		3.000E-11	0.0	-8200.0	
9	H + O2 = OH + O		3.000E-10	0.0	-16500.0	
10	CO + OH = CO2 + H		5.000E-13	0.0	-600.0	

Fig. 3.9 Initial Conditions for LAPP Computer Program

X/R DELTA X FEE1 PRESS(ATM)
0.0 8.999991E-02 1.000000E 00

HADE INNER MIXING MACH NUMBER MIXING RATE
RADIUS/R ZONE RADIUS/R AT HALF RADIUS COEFFICIENT
1.137226E 00 1.013722E 00 6.269119E-01 4.209092E-02

Y/R	VELOCITY FEET/SEC	TEMPERATURE K	DENSITY GM/CC	MACH NO.	ENTHALPY CAL/GM	VISCOSITY LB/FT/SEC	ELECTRON DENSITY(1/ML)	PSI
0.0	1.800000E 03	7.900000E 02	4.448628E-04	9.920356E-01	-4.688548E 01	1.104238E-01	0.0	
0.1000	1.800000E 03	7.900000E 02	4.448628E-04	9.920356E-01	-4.688548E 01	1.104238E-01	1.121742E-01	
0.2000	1.800000E 03	7.900000E 02	4.448628E-04	9.920356E-01	-4.688548E 01	1.104238E-01	2.243484E-01	
0.3000	1.800000E 03	7.900000E 02	4.448628E-04	9.920356E-01	-4.688548E 01	1.104238E-01	3.365226E-01	
0.4000	1.800000E 03	7.900000E 02	4.448628E-04	9.920356E-01	-4.688548E 01	1.104238E-01	4.486969E-01	
0.5000	1.800000E 03	7.900000E 02	4.448628E-04	9.920356E-01	-4.688548E 01	1.104238E-01	5.608711E-01	
0.6000	1.800000E 03	7.900000E 02	4.448628E-04	9.920356E-01	-4.688548E 01	1.104238E-01	6.730453E-01	
0.7000	1.800000E 03	7.900000E 02	4.448628E-04	9.920356E-01	-4.688548E 01	1.104238E-01	7.852195E-01	
0.8000	1.800000E 03	7.900000E 02	4.448628E-04	9.920356E-01	-4.688548E 01	1.104238E-01	8.973937E-01	
0.9000	1.800000E 03	7.900000E 02	4.448628E-04	9.920356E-01	-4.688548E 01	1.104238E-01	1.009567E 00	
1.0000	1.800000E 03	7.900000E 02	4.448628E-04	9.920356E-01	-4.688548E 01	1.104238E-01	1.121741E 00	
1.2745	1.800000E 02	2.920000E 02	1.204386E-03	8.888811E-02	-2.495791E 00	2.989527E-01	1.233916E 00	
1.8438	1.800000E 02	2.920000E 02	1.204386E-03	8.888811E-02	-2.495791E 00	2.989527E-01	1.345090E 00	

Fig. 3.10 Sample Output of LAPP Computer Program

MODE FRACTIONS

A	H	H2	H2O	CO	CO2	O	OH	PT
2.88298E-07	2.88298E-07	3.31506E-02	2.88298E-07	3.26306E-02	2.88298E-07	2.88298E-07	2.88298E-07	1
2.88298E-07	2.88298E-07	3.31506E-02	2.88298E-07	3.26306E-02	2.88298E-07	2.88298E-07	2.88298E-07	2
2.88298E-07	2.88298E-07	3.31506E-02	2.88298E-07	3.26306E-02	2.88298E-07	2.88298E-07	2.88298E-07	3
2.88298E-07	2.88298E-07	3.31506E-02	2.88298E-07	3.26306E-02	2.88298E-07	2.88298E-07	2.88298E-07	4
2.88298E-07	2.88298E-07	3.31506E-02	2.88298E-07	3.26306E-02	2.88298E-07	2.88298E-07	2.88298E-07	5
2.88298E-07	2.88298E-07	3.31506E-02	2.88298E-07	3.26306E-02	2.88298E-07	2.88298E-07	2.88298E-07	6
2.88298E-07	2.88298E-07	3.31506E-02	2.88298E-07	3.26306E-02	2.88298E-07	2.88298E-07	2.88298E-07	7
2.88298E-07	2.88298E-07	3.31506E-02	2.88298E-07	3.26306E-02	2.88298E-07	2.88298E-07	2.88298E-07	8
2.88298E-07	2.88298E-07	3.31506E-02	2.88298E-07	3.26306E-02	2.88298E-07	2.88298E-07	2.88298E-07	9
2.88298E-07	2.88298E-07	3.31506E-02	2.88298E-07	3.26306E-02	2.88298E-07	2.88298E-07	2.88298E-07	10
2.88298E-07	2.88298E-07	3.31506E-02	2.88298E-07	3.26306E-02	2.88298E-07	2.88298E-07	2.88298E-07	11
2.88298E-07	2.88298E-07	3.31506E-02	2.88298E-07	3.26306E-02	2.88298E-07	2.88298E-07	2.88298E-07	12
2.88298E-07	2.88298E-07	3.31506E-02	2.88298E-07	3.26306E-02	2.88298E-07	2.88298E-07	2.88298E-07	13

NET RATE OF PRODUCTION (W-DOT/RHJ*U)

A	H	H2	CO	CO2	O	OH	PT
5.31415E-10	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	6.40717E-08	6.87579E-08	1
5.31415E-10	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	6.40717E-08	6.87579E-08	2
5.31415E-10	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	6.40717E-08	6.87579E-08	3
5.31415E-10	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	6.40717E-08	6.87579E-08	4
5.31415E-10	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	6.40717E-08	6.87579E-08	5
5.31415E-10	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	6.40717E-08	6.87579E-08	6
5.31415E-10	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	6.40717E-08	6.87579E-08	7
5.31415E-10	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	6.40717E-08	6.87579E-08	8
5.31415E-10	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	6.40717E-08	6.87579E-08	9
5.31415E-10	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	6.40717E-08	6.87579E-08	10
5.31415E-10	5.31415E-10	-1.89146E-09	2.77323E-11	-2.77323E-11	6.40717E-08	6.87579E-08	11

Fig. 3.11 Sample Output of LAPP Computer Program

MOLE FRACTIONS

U/F	U2	U2	B	SRMO	MED	2	PT
0.0	1.56703E-01	7.77515E-01	0.0	0.0	0.0	0.0	1
0.1000	1.56703E-01	7.77515E-01	0.0	0.0	0.0	0.0	2
0.2000	1.56703E-01	7.77515E-01	0.0	0.0	0.0	0.0	3
0.3000	1.56703E-01	7.77515E-01	0.0	0.0	0.0	0.0	4
0.4000	1.56703E-01	7.77515E-01	0.0	0.0	0.0	0.0	5
0.5000	1.56703E-01	7.77515E-01	0.0	0.0	0.0	0.0	6
0.6000	1.56703E-01	7.77515E-01	0.0	0.0	0.0	0.0	7
0.7000	1.56703E-01	7.77515E-01	0.0	0.0	0.0	0.0	8
0.8000	1.56703E-01	7.77515E-01	0.0	0.0	0.0	0.0	9
0.9000	1.56703E-01	7.77515E-01	0.0	0.0	0.0	0.0	10
1.0000	1.56703E-01	7.77515E-01	0.0	0.0	0.0	0.0	11
1.27445	2.06464E-01	7.91201E-01	0.0	0.0	0.0	0.0	12
1.63363	2.06464E-01	7.91201E-01	0.0	0.0	0.0	0.0	13

NET RATE OF PRODUCTION (W-DOT/RHO*U)

U2	U2	B	SRMO	MED	2	PT
-6.51553E-08	0.0	0.0	0.0	0.0	0.0	1
-6.51553E-08	0.0	0.0	0.0	0.0	0.0	2
-6.51553E-08	0.0	0.0	0.0	0.0	0.0	3
-6.51553E-08	0.0	0.0	0.0	0.0	0.0	4
-6.51553E-08	0.0	0.0	0.0	0.0	0.0	5
-6.51553E-08	0.0	0.0	0.0	0.0	0.0	6
-6.51553E-08	0.0	0.0	0.0	0.0	0.0	7
-6.51553E-08	0.0	0.0	0.0	0.0	0.0	8
-6.51553E-08	0.0	0.0	0.0	0.0	0.0	9
-6.51553E-08	0.0	0.0	0.0	0.0	0.0	10
-6.51553E-08	0.0	0.0	0.0	0.0	0.0	11

END INLET

Fig. 3.11 (cont.)

PRESSURE		EXIT RADIUS		P _{N2}		P _{CO}		P _{CO2}		P _{H2O}		P _{O2}	
.1000E 01 0.2720E 02 120													
0.0 = X POSITION (cm)		T (K)											
0.0	R (cm)	0.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.2720E	010	.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.5440E	010	.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.8160E	010	.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.1088E	020	.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.1360E	020	.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.1632E	020	.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.1904E	020	.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.2176E	020	.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.2448E	020	.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.2720E	020	.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.3460E	020	.2920E	030	.1000E-190	.1000E-190	.3288E-030	.2885E-060	.1000E-19					
0.382E	02 = X (cm)												
0.0		0.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.2720E	010	.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.5440E	010	.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.8160E	010	.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.1088E	020	.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.1360E	020	.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.1632E	020	.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.1904E	020	.7899E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.2176E	020	.7876E	030	.1000E-190	.1000E-190	.3245E-010	.3296E-010	.1000E-19					
0.2448E	020	.7640E	030	.1000E-190	.1000E-190	.3061E-010	.3108E-010	.1000E-19					
0.2728E	020	.6661E	030	.1000E-190	.1000E-190	.2319E-010	.2346E-010	.1000E-19					
0.3123E	020	.4492E	030	.1000E-190	.1000E-190	.9216E-020	.9121E-020	.1000E-19					
0.3968E	020	.3143E	030	.1000E-190	.1000E-190	.1604E-020	.1309E-020	.1000E-19					
0.5252E	020	.2935E	030	.1000E-190	.1000E-190	.4128E-030	.8655E-040	.1000E-19					
0.6351E	020	.2920E	030	.1000E-190	.1000E-190	.3288E-030	.2885E-060	.1000E-19					
0.720E	02												
0.0		0.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.2720E	010	.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.5440E	010	.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.8160E	010	.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.1088E	020	.7900E	030	.1000E-190	.1000E-190	.3263E-010	.3315E-010	.1000E-19					
0.1360E	020	.7899E	030	.1000E-190	.1000E-190	.3263E-010	.3314E-010	.1000E-19					
0.1632E	020	.7892E	030	.1000E-190	.1000E-190	.3257E-010	.3309E-010	.1000E-19					
0.1903E	020	.7848E	030	.1000E-190	.1000E-190	.3222E-010	.3273E-010	.1000E-19					
0.2175E	020	.7667E	030	.1000E-190	.1000E-190	.3078E-010	.3125E-010	.1000E-19					
0.2449E	020	.7178E	030	.1000E-190	.1000E-190	.2694E-010	.2731E-010	.1000E-19					
0.2744E	020	.6269E	030	.1000E-190	.1000E-190	.2022E-010	.2041E-010	.1000E-19					

Fig. 3.12 Typical Set of Input Data for Radiation Code

FL1.PASS7 X(cm)	W/SR/CM	STATION RADIATION (W/SR)
0.0	0.8030024E 00	0.1062637E 01
38.19	0.6650596E 00	0.2803171E 02
71.97	0.5855711E 00	0.2112463E 02
102.27	0.5204300E 00	0.1675532E 02
135.10	0.4548181E 00	0.1600888E 02
165.70	0.3968118E 00	0.1302993E 02
195.02	0.3400359E 00	0.1190510E 02
228.89	0.2897733E 00	0.9721726E 01
259.24	0.2418252E 00	0.8068358E 01
298.38	0.1936032E 00	0.6521466E 01
329.39	0.1511664E 00	0.5345317E 01
359.80	0.1196913E 00	0.4125719E 01
395.49	0.9243929E-01	0.3779300E 01
429.76	0.7282054E-01	0.2831890E 01
462.94	0.5788064E-01	0.2166581E 01
494.77	0.4560884E-01	0.1662321E 01
525.78	0.3838268E-01	0.1317710E 01
560.59	0.3459402E-01	0.1270374E 01
594.32	0.2598157E-01	0.1021555E 01
626.96	0.2198615E-01	0.7828305E 00
658.51	0.1888854E-01	0.6446391E 00
689.25	0.1645325E-01	0.5431341E 00
723.25	0.1424051E-01	0.5217939E 00
756.43	0.1251272E-01	0.4438905E 00
788.80	0.1109931E-01	0.3821360E 00
820.08	0.9954128E-02	0.3292734E 00
850.82	0.1471464E-01	0.3791106E 00
884.27	0.8120896E-02	0.3819935E 00
909.84	0.0	0.1038170E 00

----- * RADIATION VALUE WAS CALCULATED BASED ON THE SLOPE BETWEEN THE PREVIOUS TWO POINTS
TOTAL RADIATION EMITTED=0.1612E 03 WATTS/STER
CENTROID= 0.1921E 03CM

Fig. 3.13 CO₂ Contribution to Infrared Radiation per Length of Plume

FL1.PASS7

		STATION RADIATION (W/SR)
0.0	0.2072786E-01	0.1062637E 01
38.19	0.1714092E-01	0.7230808E 00
71.97	0.1537959E-01	0.5493103E 00
102.27	0.1386401E-01	0.4430514E 00
135.10	0.1225346E-01	0.4287230E 00
165.70	0.1080529E-01	0.3527986E 00
198.02	0.9366962E-02	0.3259194E 00
228.89	0.8083675E-02	0.2693681E 00
259.24	0.6870915E-02	0.2269739E 00
298.38	0.5615387E-02	0.2443608E 00
329.39	0.4496639E-02	0.1567802E 00
359.86	0.3643198E-02	0.1240556E 00
395.49	0.2894438E-02	0.1165630E 00
429.76	0.2336578E-02	0.8953859E-01
462.94	0.1902269E-02	0.7033050E-01
494.77	0.1566058E-02	0.5518754E-01
525.78	0.1313538E-02	0.4464533E-01
560.59	0.1114177E-02	0.4226155E-01
594.32	0.9213609E-03	0.3432758E-01
626.96	0.7910719E-03	0.2794681E-01
658.51	0.6881264E-03	0.2333586E-01
689.25	0.6061282E-03	0.1989019E-01
723.25	0.5306695E-03	0.1932555E-01
756.43	0.4710052E-03	0.1661982E-01
788.80	0.4216461E-03	0.1444663E-01
820.08	0.3811687E-03	0.1255593E-01
850.82	0.3897771E-03	0.1184792E-01
884.27	0.3154266E-03	0.1179666E-01
909.84	0.0	0.4032392E-02

* RADIATION VALUE WAS CALCULATED BASED ON THE SLOPE BETWEEN THE PREVIOUS TWO POINTS

TOTAL RADIATION EMITTED=0.4459E 01 WATTS/STER

CENTRJD= 0.2050E 03CM

Fig. 3.14 H₂O Contribution to Infrared Radiation per Length of Plume

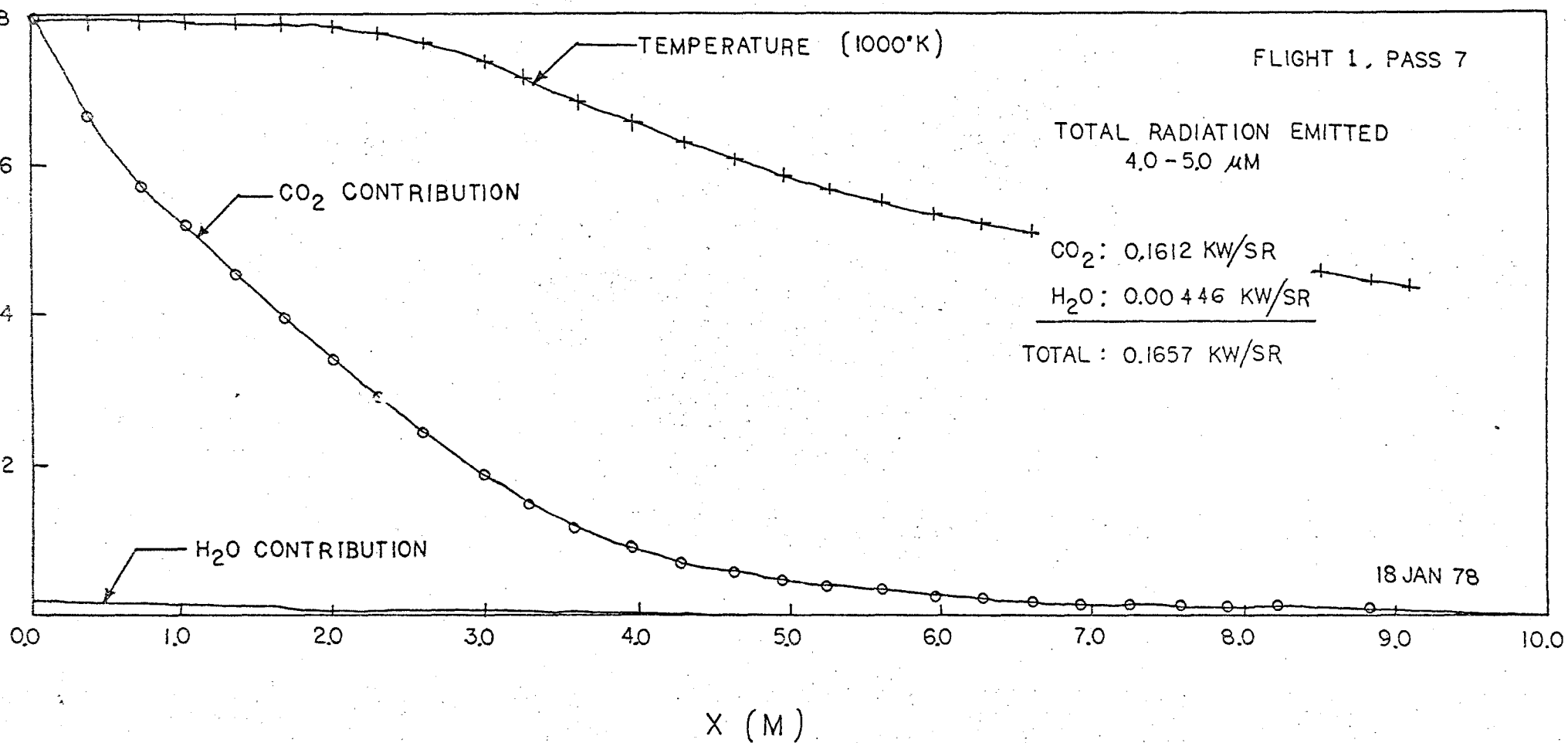


Fig. 3.15 SPATIAL DISTRIBUTION OF TEMPERATURE AND RADIATION IN THE TURBOJET EXHAUST.

3.3 Infrared Radiation Predictions for a Small Kerosene/Gaseous Oxygen Rocket Engine

Two gas dynamics and the UTSI infrared radiation computational models were used to predict the physical properties in hot rocket exhaust plumes. Experimental data were obtained in the exhaust plume of a small kerosene/gaseous oxygen rocket engine using a three-dimensional laser Doppler velocimeter system to measure the velocity distribution. The spatial distribution of radiated energy in the 4-5 μ wavelength band was measured using an infrared radiometer. Detailed comparisons between the computational and experimental results are described in this section.

3.3.1 Description of Experimental Equipment

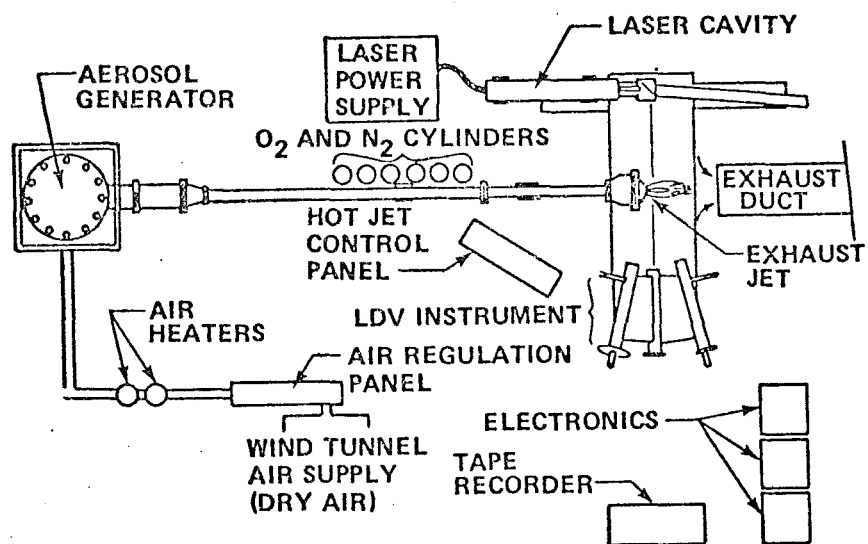
Small Laboratory Rocket Engine: The exhaust plumes studied in these measurements were produced by a small kerosene/gaseous oxygen rocket engine [30] operating in the NASA Laser Doppler Velocimeter Facility [31]. Engine design and performance parameters are listed in Table 3.3. Data were taken at a variety of oxidizer/fuel (O/F) ratios and at two chamber pressures. The experimental data used to compare with the theoretical calculations were obtained at a chamber pressure of 11.1 atm. and an O/F ratio of 2.25.

NASA Laser Doppler Velocimeter Facility: The NASA Marshall three-dimensional laser Doppler velocimeter [31] was used to measure the velocity distribution in the exhaust plume and mixing layer produced by the inner hot rocket exhaust jet mixing with an outer subsonic airflow. A schematic diagram of the LDV experimental arrangement showing the installation of the small rocket engine is shown in Fig. 3.16. As described in Ref. [18] the simultaneous measurement of the three Doppler velocity components yielded the three velocity components u, v, and w.

TABLE 3.3

ENGINE DESIGN AND PERFORMANCE PARAMETERS

Fuel	Kerosene
Oxidizer	Gaseous oxygen
Purge	Nitrogen 100 psig
Coolant	Water at 3 gal/min
Chamber pressure	50 and 148 psig
Fuel flow rate	
50 psig	1.38 g/sec
148 psig	3.80 g/sec
Oxygen flow rate	
50 psig	3.05 g/sec
148 psig	8.70 g/sec
Throat diameter	0.200 in.
Exit diameter	0.310 in.
Throat area	0.0314 in. ²
Exit area	0.0755 in. ²
Area ratio (A_E/A_t)	2.40
O/F ratio	2.25
Chamber diameter	0.780 in.
Chamber length	1.620 in.
L^*	25.20 in.
C^* at 148 psig	6045 fps



3.16 SCHEMATIC DIAGRAM OF LASER VELOCIMETER ARRANGEMENT

In order to simulate external airflow around the small rocket engine, an outer flow nozzle was used to produce an external flow velocity of 61 m/sec. A schematic diagram of the rocket engine and outer flow nozzle is presented in Fig. 3.17.

Infrared Radiation Diagnostic Equipment: The primary instrument used to obtain the radiation intensity data was an Electro-Optical Industries Model 470 Radiometer. The data were obtained using a fixed filter, which transmitted radiation in the 4-5 μ band.

A large metal shield with a 1-mm aperture was moved in the axial direction to obtain the variation of radiant intensity along the exhaust plume centerline. A second set of measurements was made using a 1-mm vertical slot in the large metal shield. The slot length was greater than the diameter of the exhaust plume. The experimental setup is shown in Fig. 3.18.

3.3.2 Typical Experimental Results

Velocity Distribution Data: Experimental data for mean flow velocities $(\bar{u}, \bar{v}, \bar{w})$, turbulent intensities $[(\bar{u}')^2, (\bar{v}')^2, (\bar{w}')^2]$, and turbulent velocity correlations were obtained in the radial direction at various axial locations downstream from the nozzle exit plane. Data were obtained at X/D locations of 2.4, 4.8, 8.4, 11.3, and 14.2, where X is the distance measured from the nozzle exit plane and D is the diameter of the exit nozzle. Typical data obtained at X/D = 11.3 for the mean axial flow velocity are presented in Fig. 3.19. Other data are described in Ref. [18].

Infrared Radiation Data: The infrared radiation data have been described previously [30]. Typical data obtained at an O/F ratio of 2.25 and an external flow velocity of 210 fps are presented in Figs. 3.20 and 3.21 for the 1-mm aperture and slot, respectively. The centerline radiation intensity

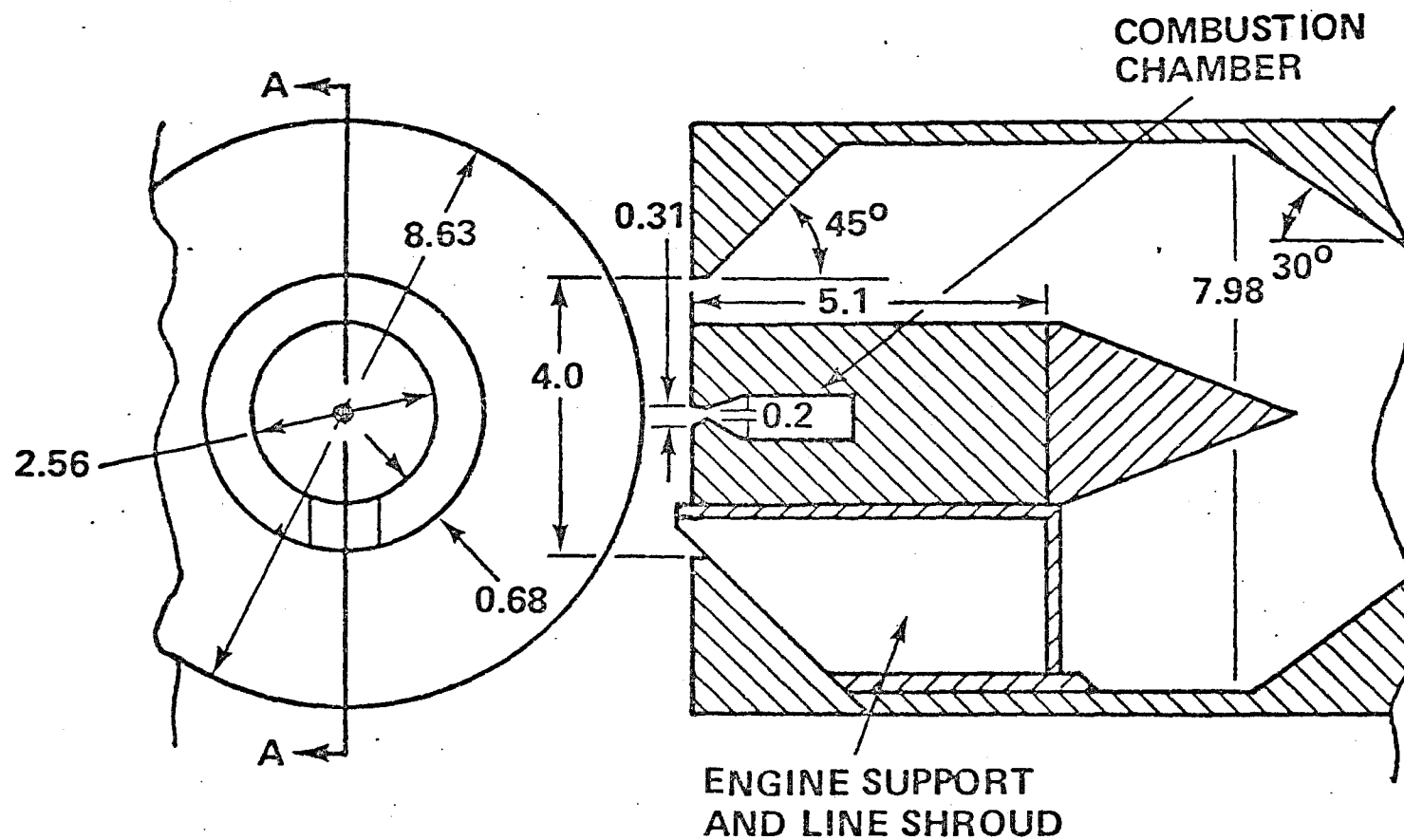


Fig. 3.17 Schematic of the exhaust plume and external flow apparatus.

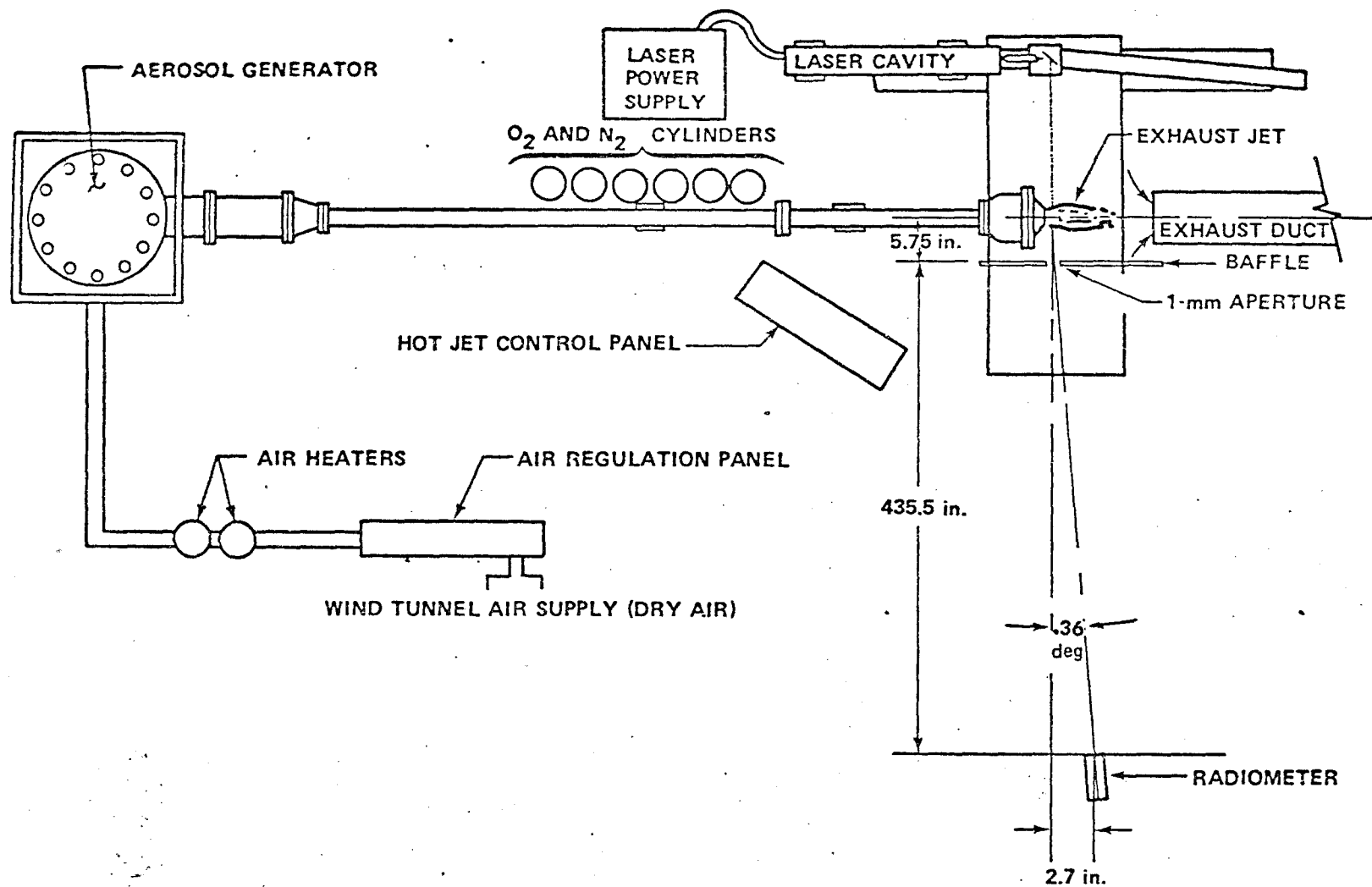


Fig. 3.18 Exhaust plume test arrangement in NASA LDV facility.

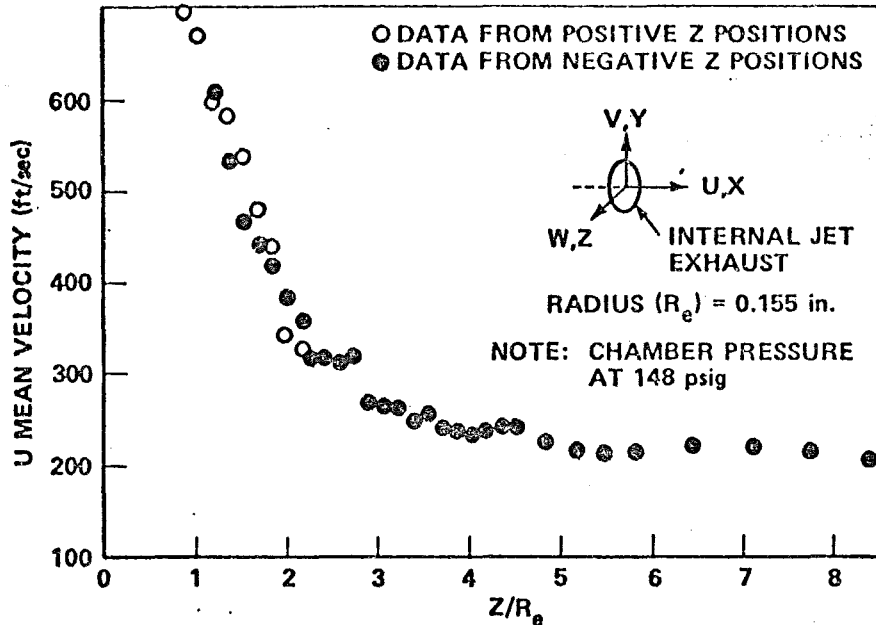


Fig. 3.19 Profile of u mean velocity component at $X/D = 11.3$

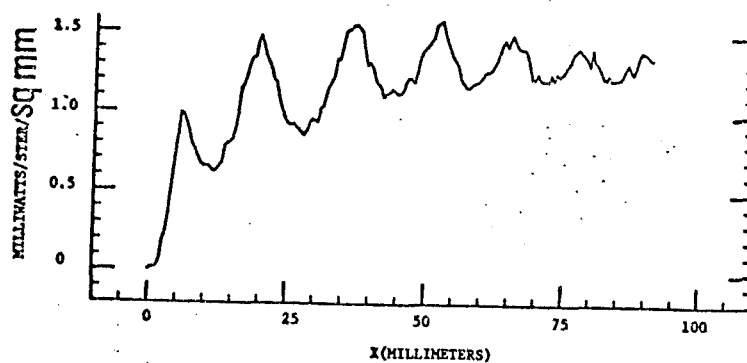


Fig. 3.20 Spatial distribution of radiation along axis of plume for aperture diameter = 1-mm; $O/F = 2.25$, velocity = 150 fps.

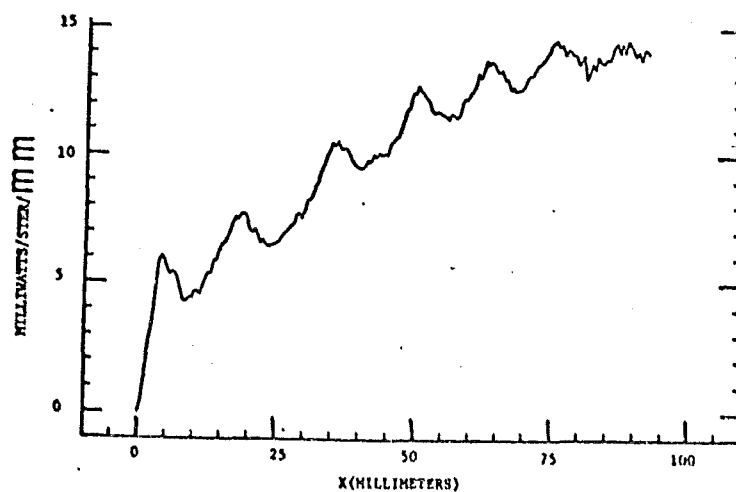


Fig. 3.21 Spatial distribution of radiation along axis of plume for 1-mm slot width; $O/F = 2.25$, velocity = 210 fps.

distribution in Fig. 3.20 clearly shows the presence and influence of shocks and/or Mach disks in the exhaust plume. The effect of afterburning in the mixing layers on the infrared signature is shown clearly in Fig. 3.21 where the shock-induced radiation peaks are superimposed on an axially increasing radiation intensity. The data presented in Figs. 3.20 and 3.21 were obtained in the near field of the exhaust plume (X/D less than 12) and indicate that the peak radiation intensity has not been reached. Typically, the axial distribution of radiation is similar to that shown in Fig. 3.22. Even though the data presented in Fig. 3.22 are for a lower mass flow and chamber pressure, the axial variation is similar to that obtained for the high-chamber-pressure cases. The data presented here are only samples selected to indicate the measured trends in preparation for a discussion of the comparison between theory and experiment.

3.3.3 Comparison of Calculated and Measured Results

The LAPP and TKE computer codes were exercised for the rocket performance parameters listed in Table 3.3. The NASA Lewis Chemical Equilibrium Computer Program was used to calculate the exit plane gas properties. Because of the short expansion nozzle length, the chemical species were assumed to be frozen at the throat values. These exit plane conditions then were used as initial conditions for the LAPP and TKE computer codes.

The results predicted by the LAPP and TKE computer codes are presented in Figs. 3.23-3.28. The variation of the centerline velocity with distance from the nozzle exit plane is presented in Fig. 3.23. The LAPP code predicts a constant-velocity core having a length approximately equal to $X/D = 5$, which is followed by an exponential velocity decay due to turbulent mixing predicted by the Donaldson-Gray viscosity model. The TKE (REP3) code predicts a sharp decrease in velocity at approximately an S/D value of one. There are also

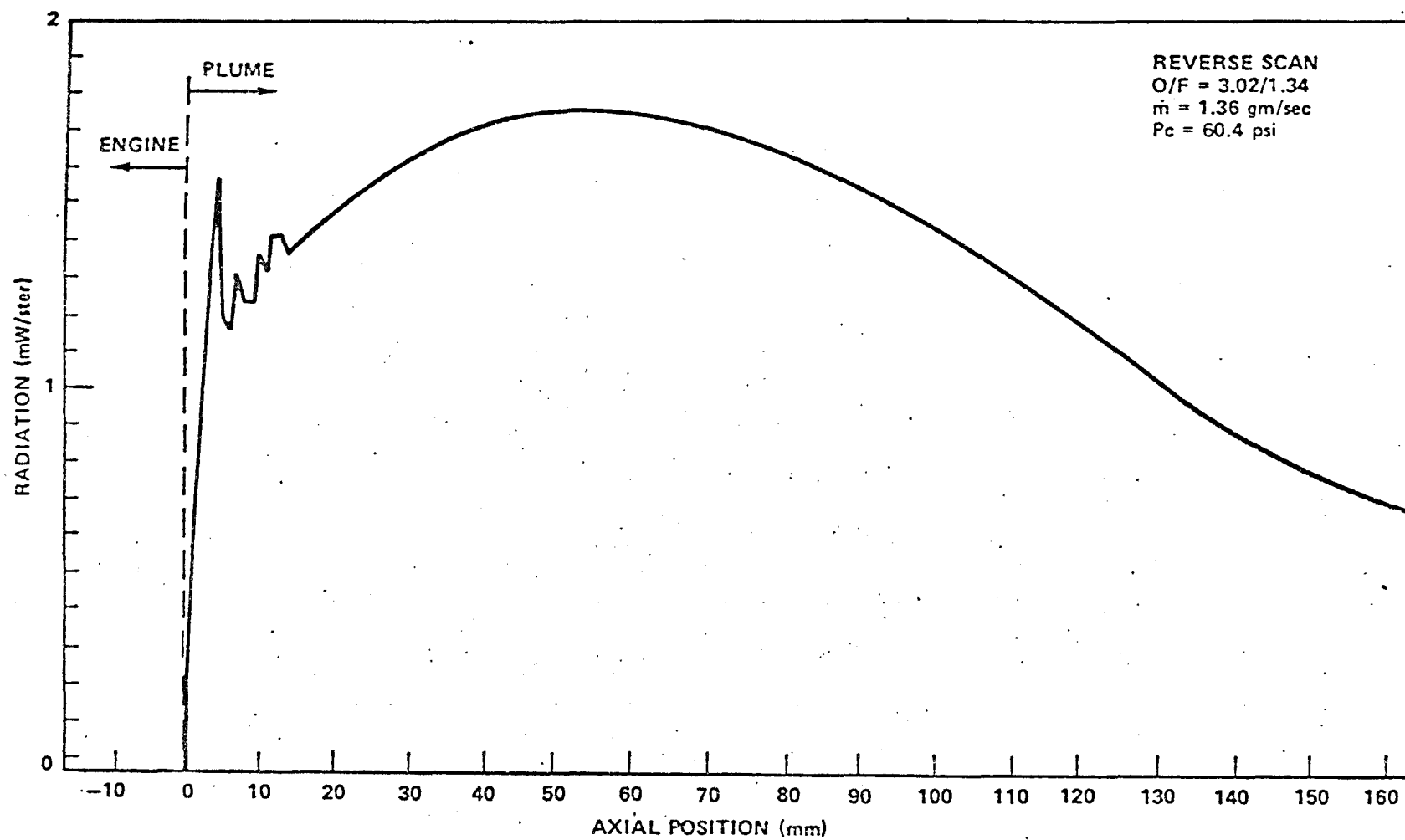


Fig. 3.22 Axial variation of radiation at O/F ratio of 2.27.

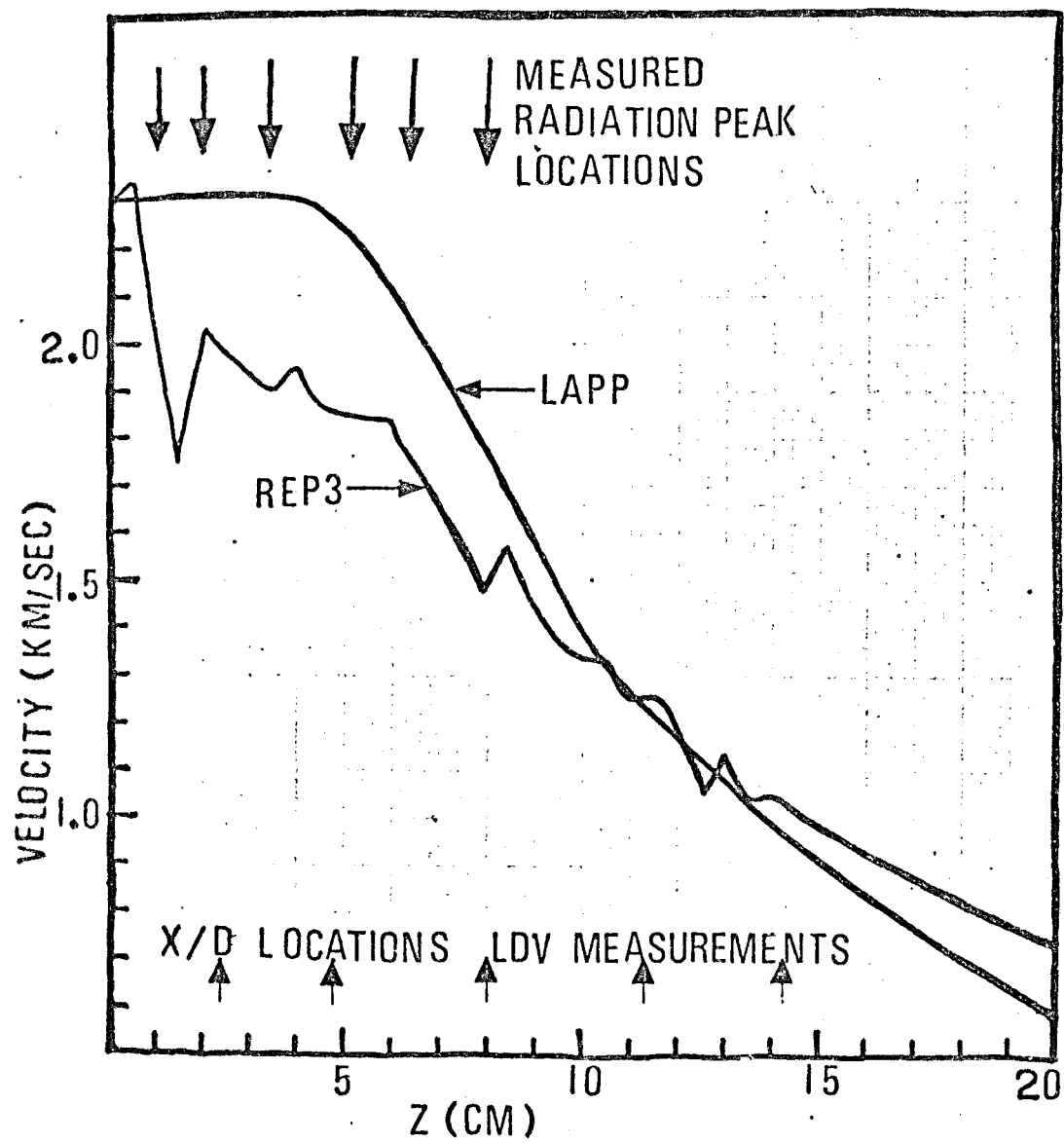


Fig. 3.23 Calculated axial velocity distribution.

several other oscillations in velocity, as shown by the jagged nature of the curve. A particularly large decrease in velocity occurs near an X/D value of 8. Also noted in Fig. 3.23 are the locations of the measured radiation peaks shown previously in Fig. 3.20. The X/D locations for the LDV measurements also are indicated on Fig. 3.23.

The predicted variations of temperature at the jet centerline with axial distance are given in Fig. 3.24. The LAPP code predicts a small temperature rise due to afterburning in the exhaust plume. The TKE code predicts a sharp rise in temperature at an X/D value of one, which then is followed by a series of oscillations in temperature. The predicted oscillations do not agree with the measured locations of radiation peaks (denoted by arrows on the Figs.). It is quite obvious that neither of the computer codes does an adequate job of predicting the spatial variation of velocity and temperature produced by shock waves and Mach disks in the near field of the exhaust plume. This is not surprising since LAPP has no provision for treating shocks.

Figures 3.25-3.38 present the variation of the axial component of the velocity with radial distance from the jet centerline. As shown in Fig. 3.25 for an X/D value of 2.4, the predicted results (curves) do not agree with the experimental measurements (symbols). Because of the limited number of particles in the hot core of the jet, the NASA LDV system was not able to make measurements closer to the jet centerline than about one exit nozzle radius. The variation of the radial velocity component also is shown. The TKE code predicts an increasing value of v with radius, whereas the measured value decreases with radial distance. The external flow is an axial flow at 61 m/sec.

The lack of agreement between theory and experiment also is indicated in Fig. 3.26 for the velocity profile at an X/D value of 8.0. The LAPP and TKE

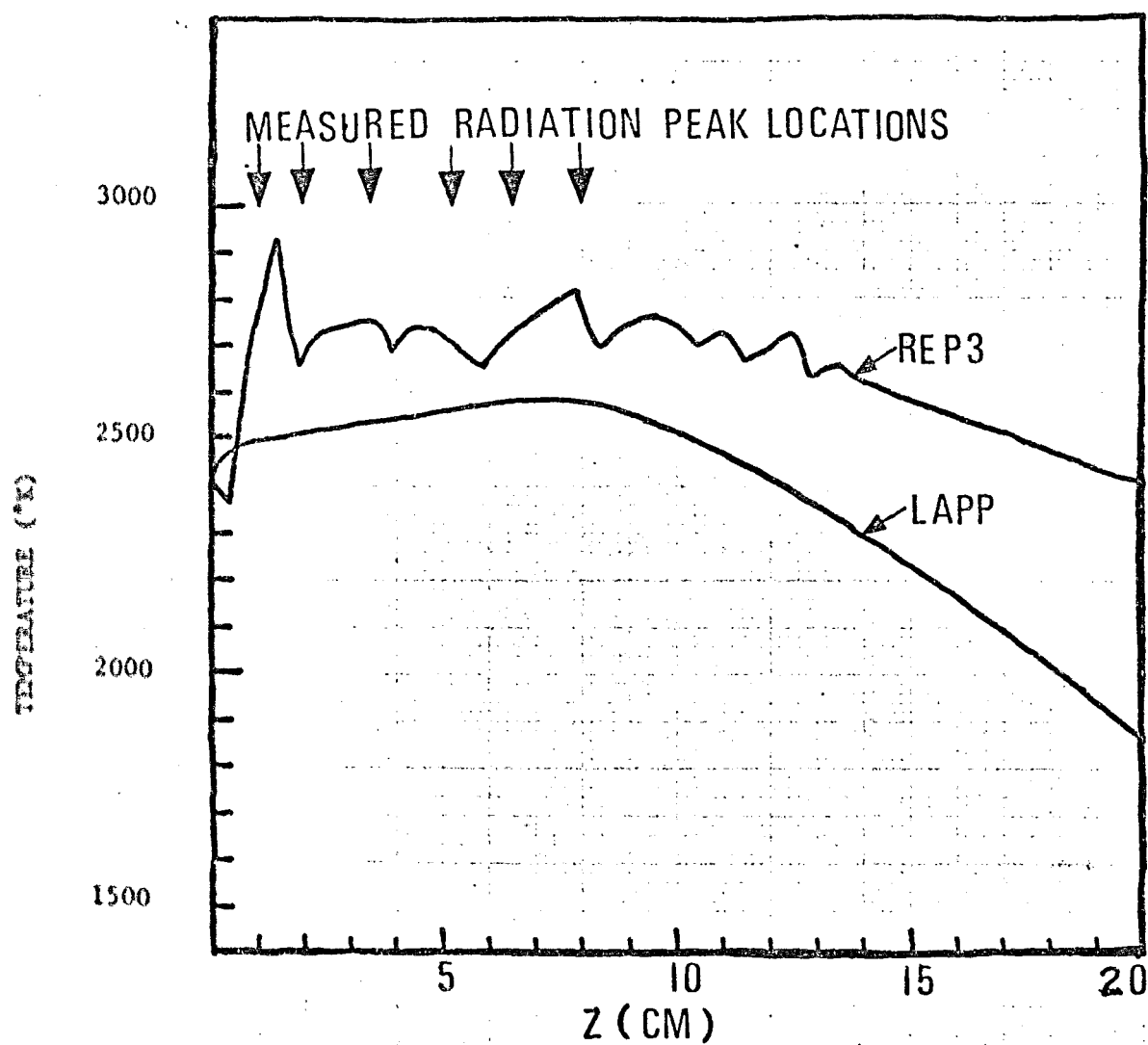


Fig. 3.24 Calculated temperature distribution.

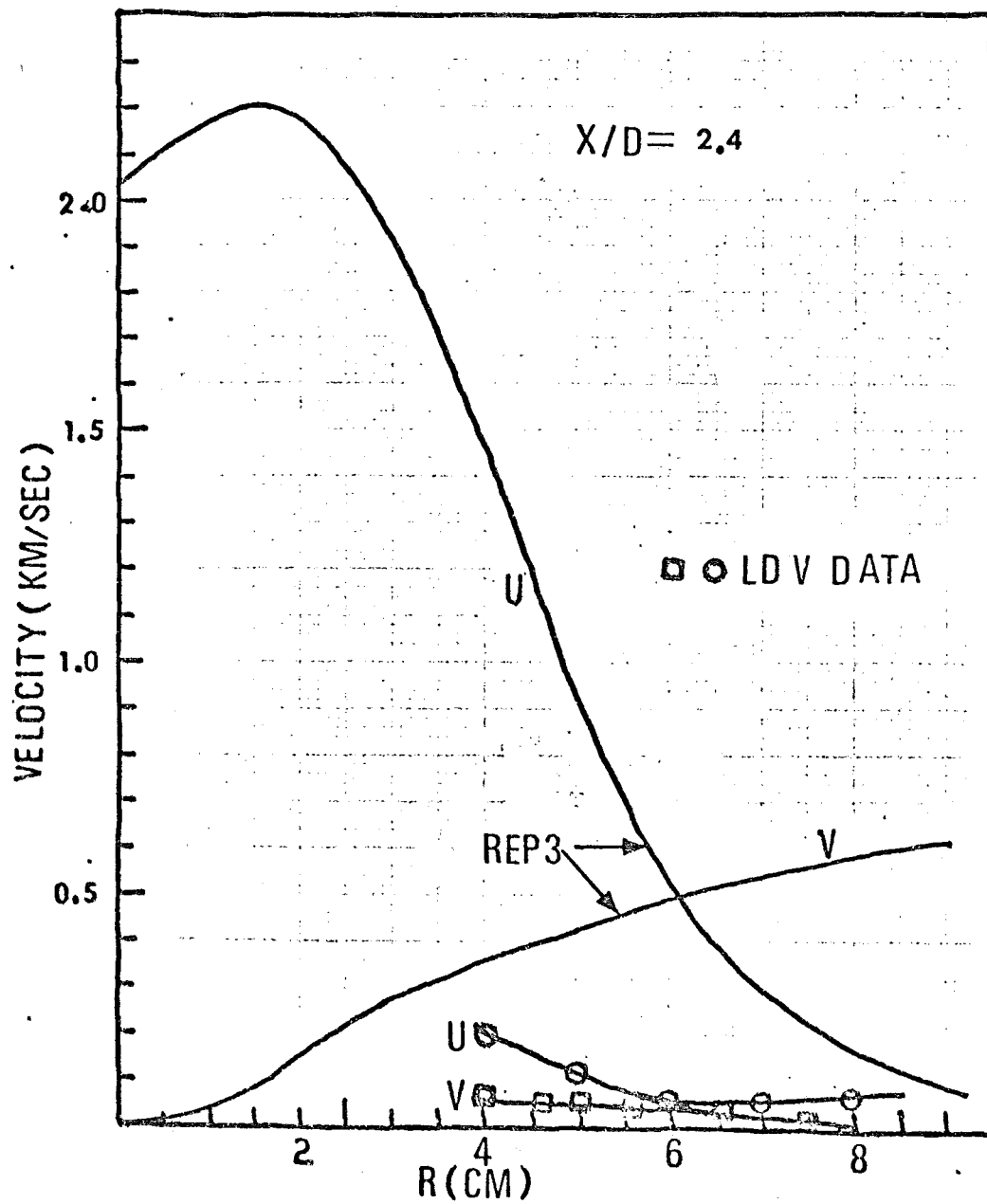


Fig. 3.25 Velocity profiles at $X/D = 2.4$

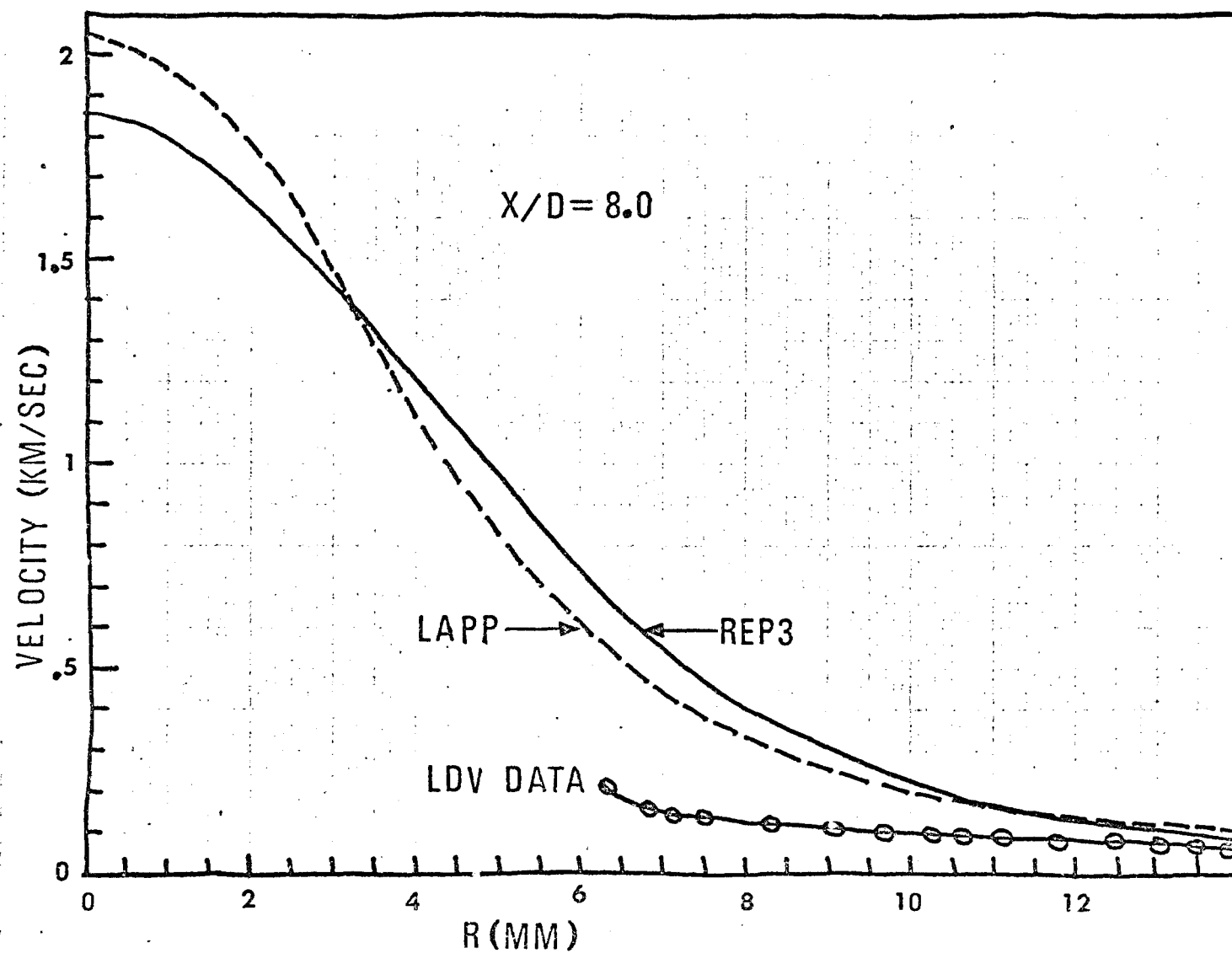


Fig. 3.26 Velocity profiles at $X/D = 8.0$

calculated results are similar but are in excess of the experimental values. Similar results are presented in Figs. 3.27 and 3.28 for X/D values of 11.3 and 14.2.

The LDV data indicate that the exhaust plume is much narrower than predicted by the LAPP and TKE gas dynamics models. Since both the LAPP and TKE models assumed that turbulent mixing began at the nozzle exit plane and ignored the large recirculation (base flow) region, it is expected that turbulent mixing occurs too fast compared to the physical situation where turbulent mixing begins at the end of the recirculation zone, which is several nozzle diameters downstream from the exit. It is known that the base flow region can have a pronounced effect on exhaust plume spread and initial turbulent mixing layer growth. Because of time and budget constraints, we were unable to use a base flow model to assess its effect on the calculated temperature and velocity profiles.

The infrared radiation band model program was used to calculate the variation of infrared radiation along the exhaust plume for the temperature and species concentration profiles predicted by the LAPP and TKE gas dynamics computer codes. The computed results are shown in Fig. 3.29. The lack of agreement between theory and experiments was not unexpected, since the calculated temperature distributions did not yield the spatial structure due to shock waves that are evident in the measured infrared data and are visible in photographs of the exhaust plume.

It is not clear what causes the difference between the REP3 and LAPP calculations. The LAPP code plus the radiation code yield results that are within a factor of 1.5 of the measured results while the factor is approximately 2.5 for the REP3 plus radiation code results.

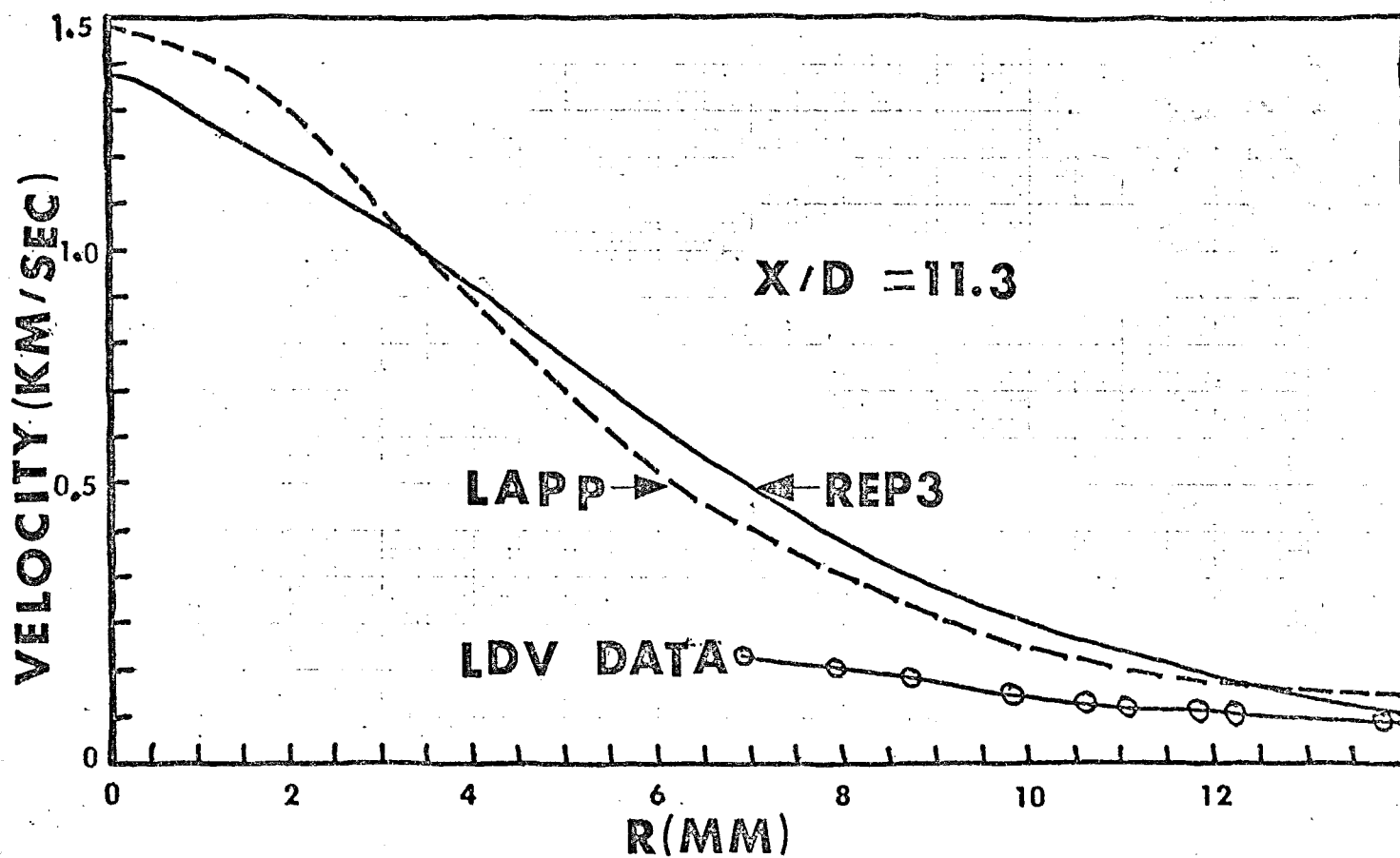


Fig. 3.27 Velocity profiles at $X/D = 11.3$

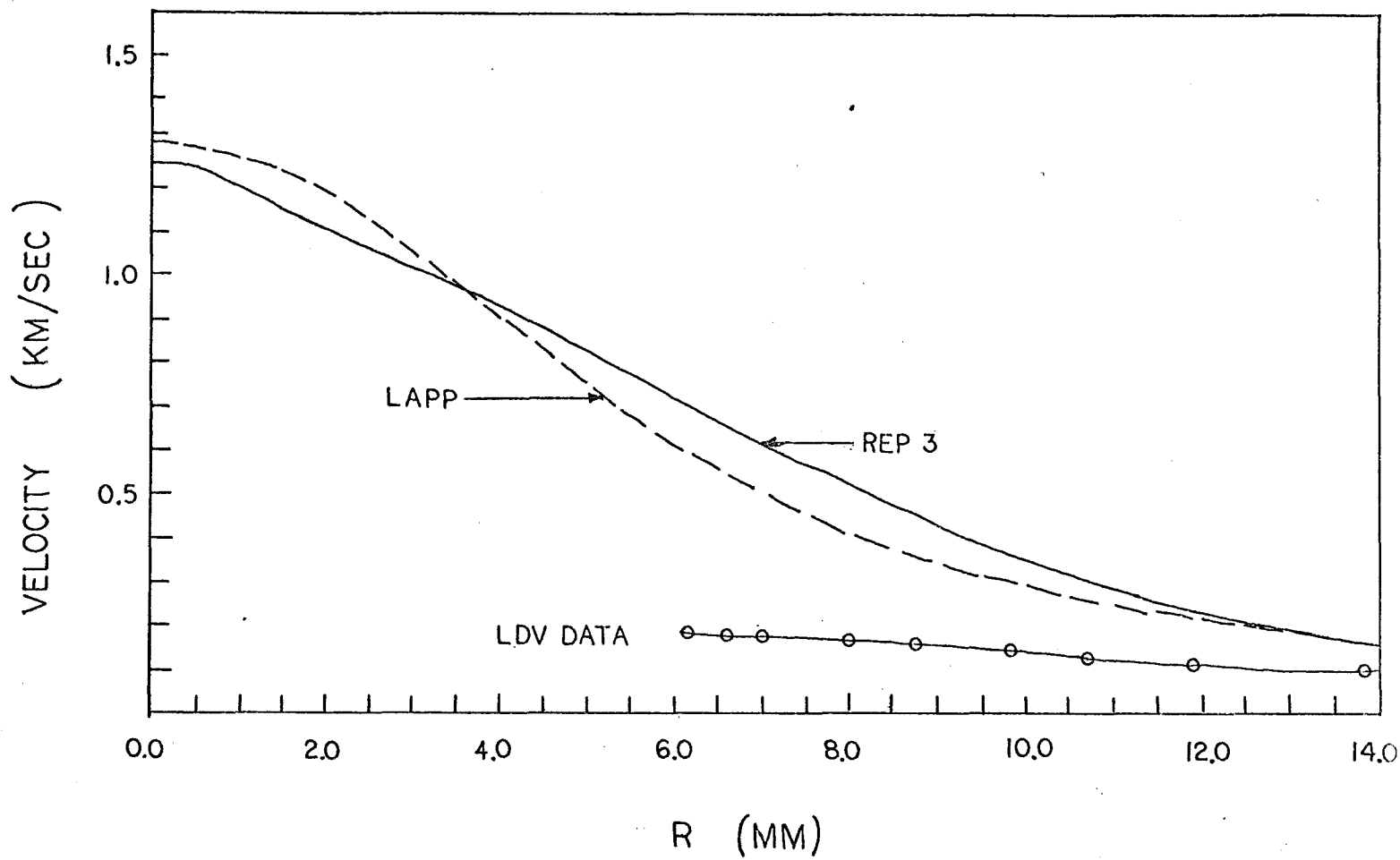


Fig. 3.28 VELOCITY PROFILES AT $X/D = 14.2$

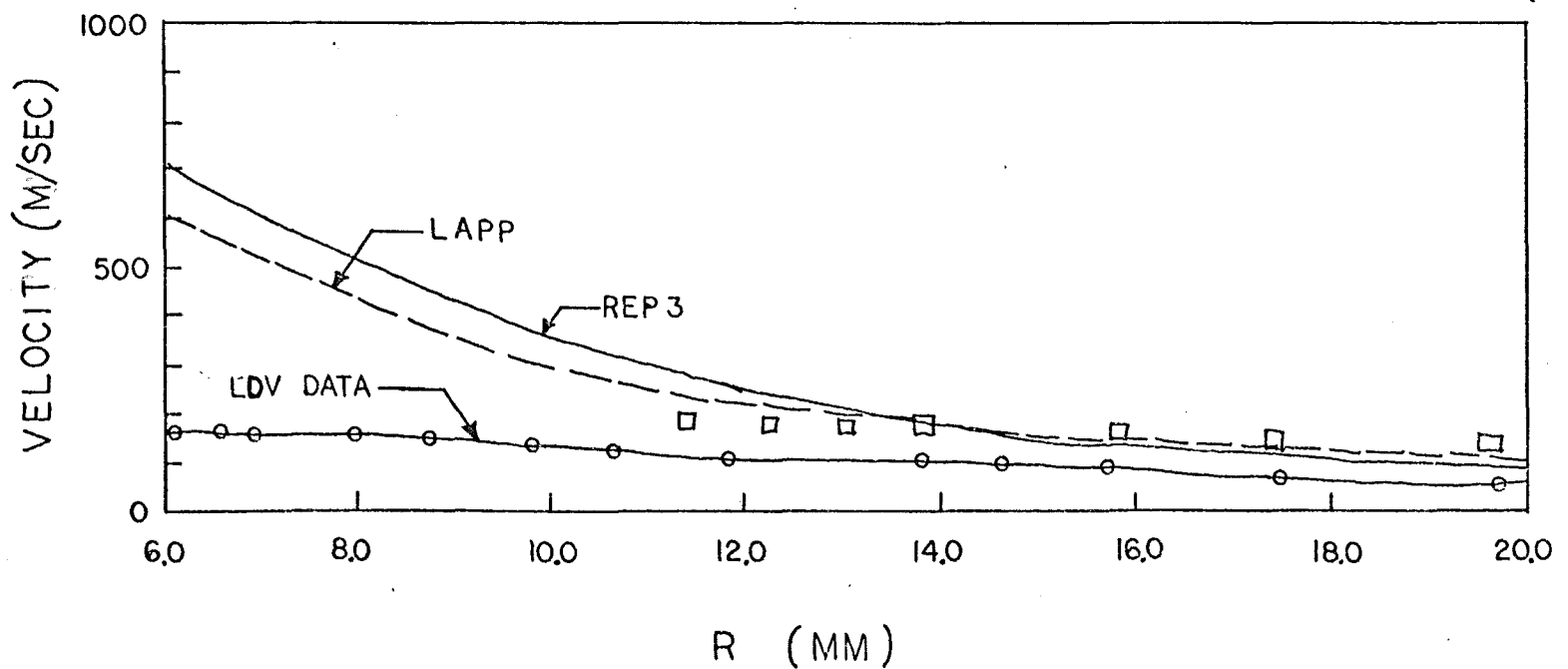


Fig. 3.28 (continued)

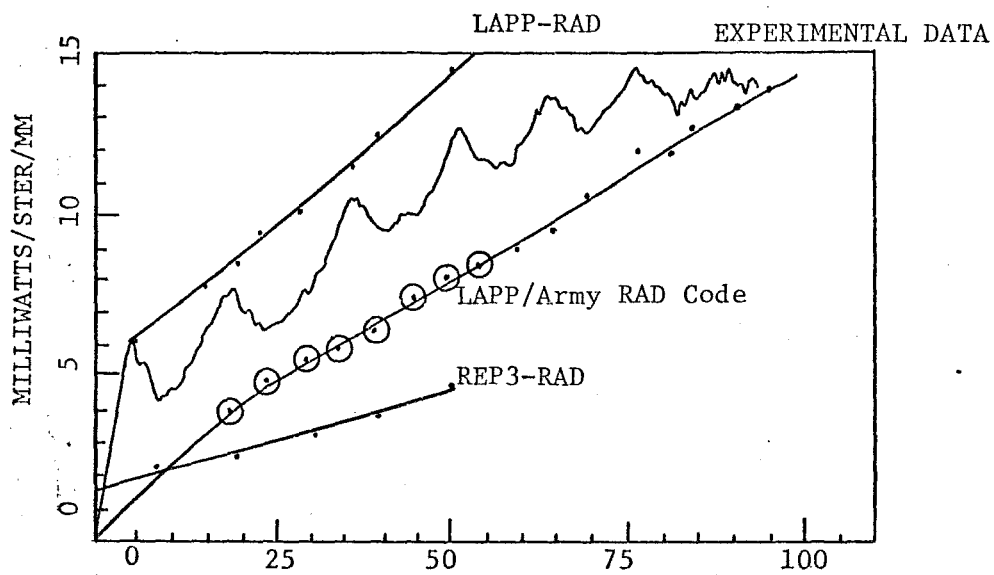


Fig. 3.29 Comparison of calculated and measured IR radiation.

3.4 Gas Dynamics Finite Difference Computer Model

A Gas Dynamics Finite Difference Computer Model was developed as part of the contract effort. Numerous difficulties and problems were encountered during the development process, but the program was operational at the conclusion of the effort. Unfortunately, the present version of the code is not entirely satisfactory due to the fact that the computational times are large. Approximately one hour on the IMB 360 computer system was required to march the solution 360 cycles (which corresponds to a flow time of 0.2 msec.). While this time would be approximately 6-10 minutes on the CDC 660 computer, it is still felt to be quite excessive since the test flow being analyzed was a relatively simple expanding flow. In its present form, the finite difference code successfully calculates shock waves on wedges and blunt bodies and is moderately successful in calculating the shock-Mach disc structure in a cold gas nozzled expansion.

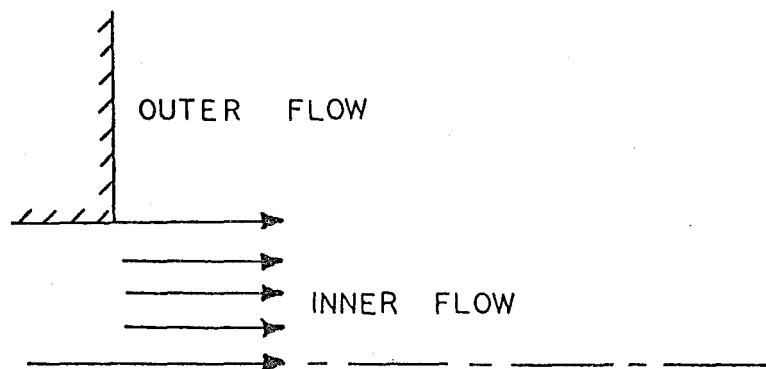
Several typical nozzle expansion test cases were computed to assess whether the numerical results were reasonable. These results are described in the next section.

3.4.1 Overexpanded Nozzle Exhaust

In order to assess the FD model's ability to handle strong shock waves which result in the exhaust plume from an overexpanded nozzle, the flow having the initial characteristics given in Fig. 3.30.

The computed Mach number variation along the jet centerline is given in Fig. 3.31. The Mach number decreased to a subsonic value in a very short distance downstream from the exit plane and remained subsonic for 80-cm in the flow direction.

An analytical estimate of the flow conditions was made by assuming that an oblique shock wave would be formed at the nozzle exit and would raise the



<u>INNER FLOW</u>	<u>OUTER FLOW</u>
$M_{\text{exit}} = 1.133$	$M = 0$
$T_{\text{exit}} = 300^{\circ}\text{K}$	$T = 300^{\circ}\text{K}$
$P_{\text{exit}} = 0.8 \text{ atmospheres}$	$P = 1.0 \text{ atmospheres}$

Fig. 3.30 INPUT FLOW PARAMETERS

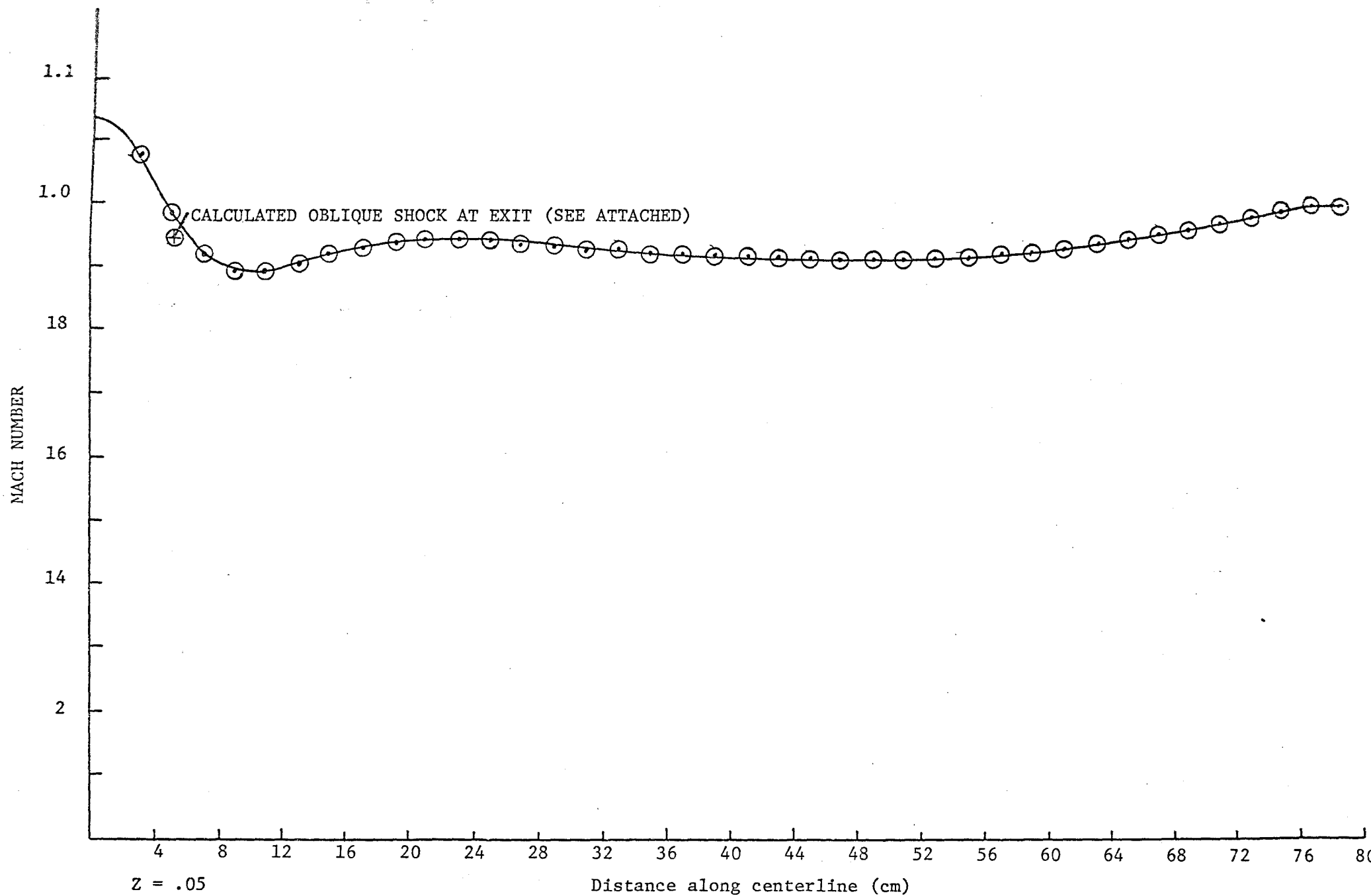


Fig. 3.31 CENTERLINE MACH NUMBER VARIATION

pressure behind the shock wave to the ambient pressure of 1.0 atm. The oblique shock angle was then determined to be approximately 76.6° as shown in Fig. 3.32. The shock wave would intersect the centerline axis at a distance of 0.2323-cm downstream of the nozzle exit plane. Gas dynamics calculations yielded a value of 0.94 for the Mach number in region (2) behind the shock wave. This calculated value is indicated on Fig. 3.31 and is in good agreement with the finite difference code result.

As shown in Fig. 3.32, the shock wave would reflect from the centerline and would process the flow back parallel to the centerline axis. In doing so, the pressure in region (3) would increase while the Mach number M_3 would be less than M_2 . This agrees qualitatively with the FD code results.

At the intersection of the reflected shock with the free boundary, an expansion wave would be generated to reduce the pressure at the edge of the plume back to one atm. The expansion fan would accelerate and expand the flow in regions (4-6) until a new shock pattern would be formed between regions (7 and 8). Then the entire pattern would repeat.

While the FD code calculations appear to be quantitatively correct, there were no available numerical data to compare with the FD code results in the subsonic flow region. In order to determine whether the FD code predictions are reasonable in an expanding flow, another test case was designed in which a supersonic flow would expand from a nozzle. The results for this flow are described in the next section.

3.4.2 Underexpanded Nozzle Exhaust

As stated above, a test case involving the supersonic expansion of air from a nozzle was employed so that the FD code predictions could be compared with the results obtained using the Method of Characteristics.

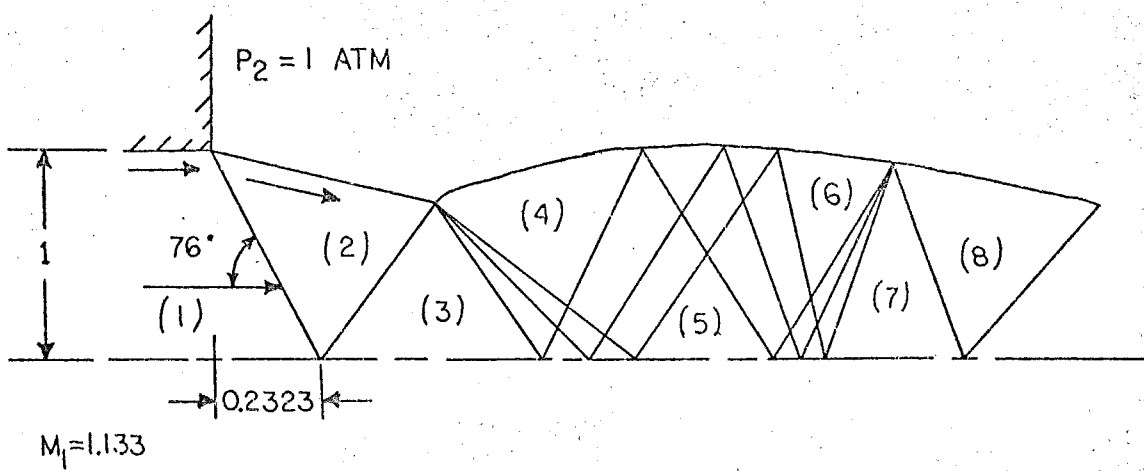


Fig. 3.32 ASSUMED SHOCK PATTERN FOR HAND CALCULATION.

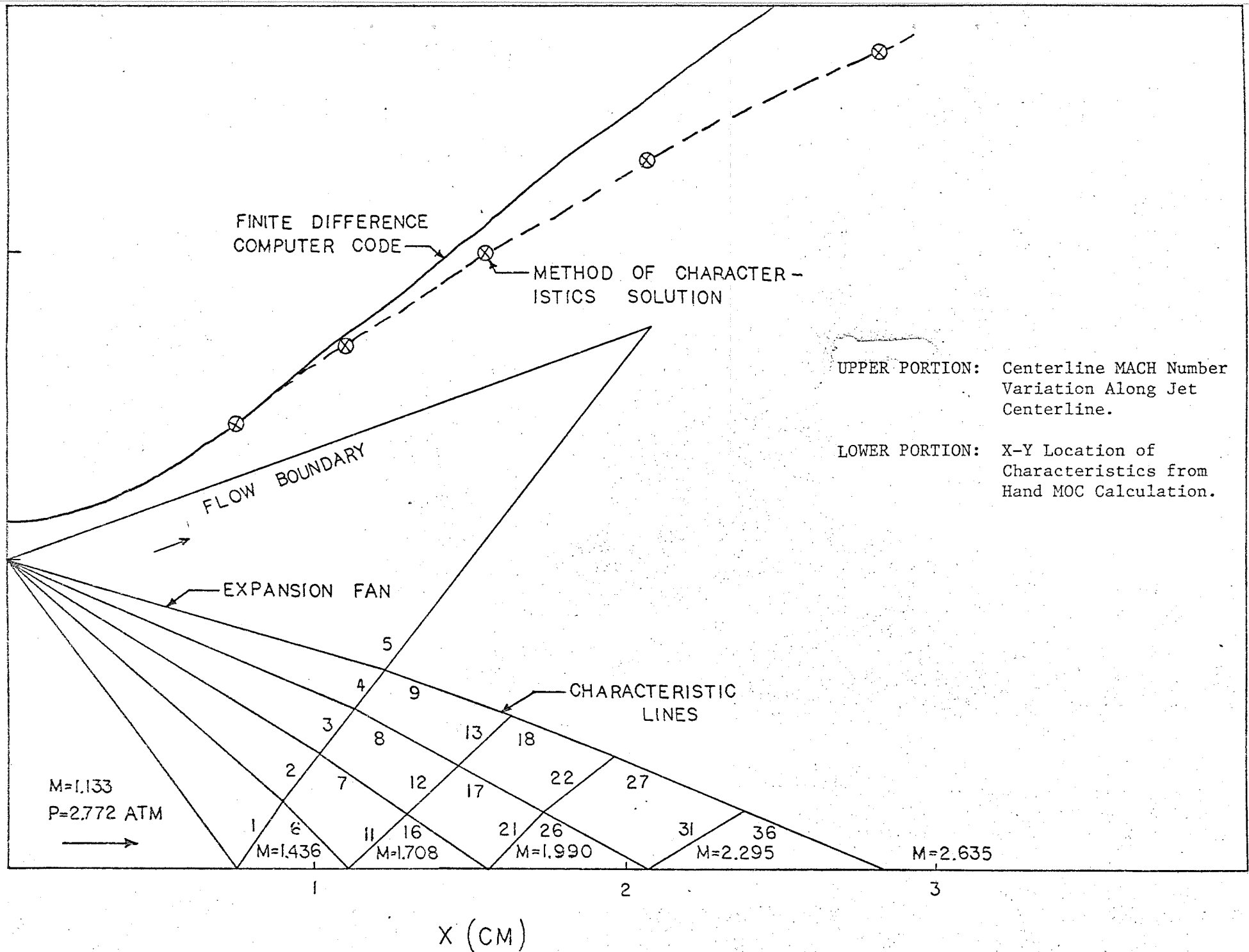


Fig. 3.33 COMPARISON OF METHOD OF CHARACTERISTICS AND FINITE DIFFERENCE MODEL RESULTS

The method of Characteristic technique used was Temple's Method as taught by the Principal Investigator in a first-year graduate course in gas dynamics. The characteristics net produced in the X-Y physical plane is shown in the lower portion of Fig. 3.33. A comparison of the MOC results (using 5 characteristics to start the calculation) and the Finite Difference Code results is presented in the upper portion of Fig. 3.33. The results are in good agreement. The lack of agreement at large distances from the nozzle exit plane is to be expected since the hand-calculated MOC solution used only five characteristics in the initial fan and is therefore expected to be more approximate than the Finite Difference Computer Code which uses a more complete set of differential equations.

Even though we were encouraged that the computer code appears to yield reasonable numerical results for an expanding plume, we are discouraged by the fact that the computer program required approximately one hour on the IBM 360 computer system to march the solution 0.3 msec. (or 360 cycles). While this time would be approximately 6-10 minutes on the CEC 6600, it still is felt to be quite excessive in view of the simple flow calculated.

3.4.3 Other Test Cases

In earlier work [21], the finite difference code was used to calculate the viscous flow between two parallel walls with one of the walls moving relative to the other. This is the so-called classical Couette Flow problem whose theoretical solution is well known [32]. The computed numerical results agreed well with the theoretical model results. In addition, the computer solution converged in 10 to 20 iteration cycles and the total computational time on the IBM 370 computer was approximately 30 seconds.

The code also successfully predicted the flow over a two-dimensional wedge in which a strong bow shock wave is formed.

In both of these cases, we have solid boundaries which restrict the flow and appear to assist in the convergence of the solution to a stable value.

3.4.4 Chemical Reactions and Turbulent Mixing Computations

Due to problems with the basic finite difference computer code, the chemical reaction and turbulent mixing portions of the code were not fully operational at the end of the program.

3.5 Results Obtained Using the REP 3 TKE Computer Model

The REP3 TKE Computer Model described earlier was revised for use on the CDC computer at MIRADCOM. The test case supplied by Dr. David Jensen of the Rocket Propulsion Establishment of Great Britain was successfully run on the MIRADCOM computer.

After the REP3 program was operational, it was utilized to calculate the flow field from a small kerosene/gaseous oxygen engine similar to the one used as the test case by Jensen. The results of these calculations are presented in Figs. 3.23-3.28. As discussed in Section 3.3, the REP 3 predictions do not agree with the experimental data. Attempts to locate any problems with the calculations or with the experimental data were unsuccessful.

One computer run was devoted to rerunning the case described in the previous paragraph. The program was rerun using a much smaller grid spacing. The results are shown in Fig. 3.34. The top curve shows the oscillations in temperature while the lower curve indicates the results for the velocity. Also shown on the graph are the locations of the radiation minima (two experimental runs are shown: Runs 65 and 85). As can be seen, the first three locations of the calculated temperature minima agree with the experimental locations. Even though the fourth through seventh minima do not agree, the results are reasonable.

ARROWS LOCATE RADIATION PEAK LOCATIONS

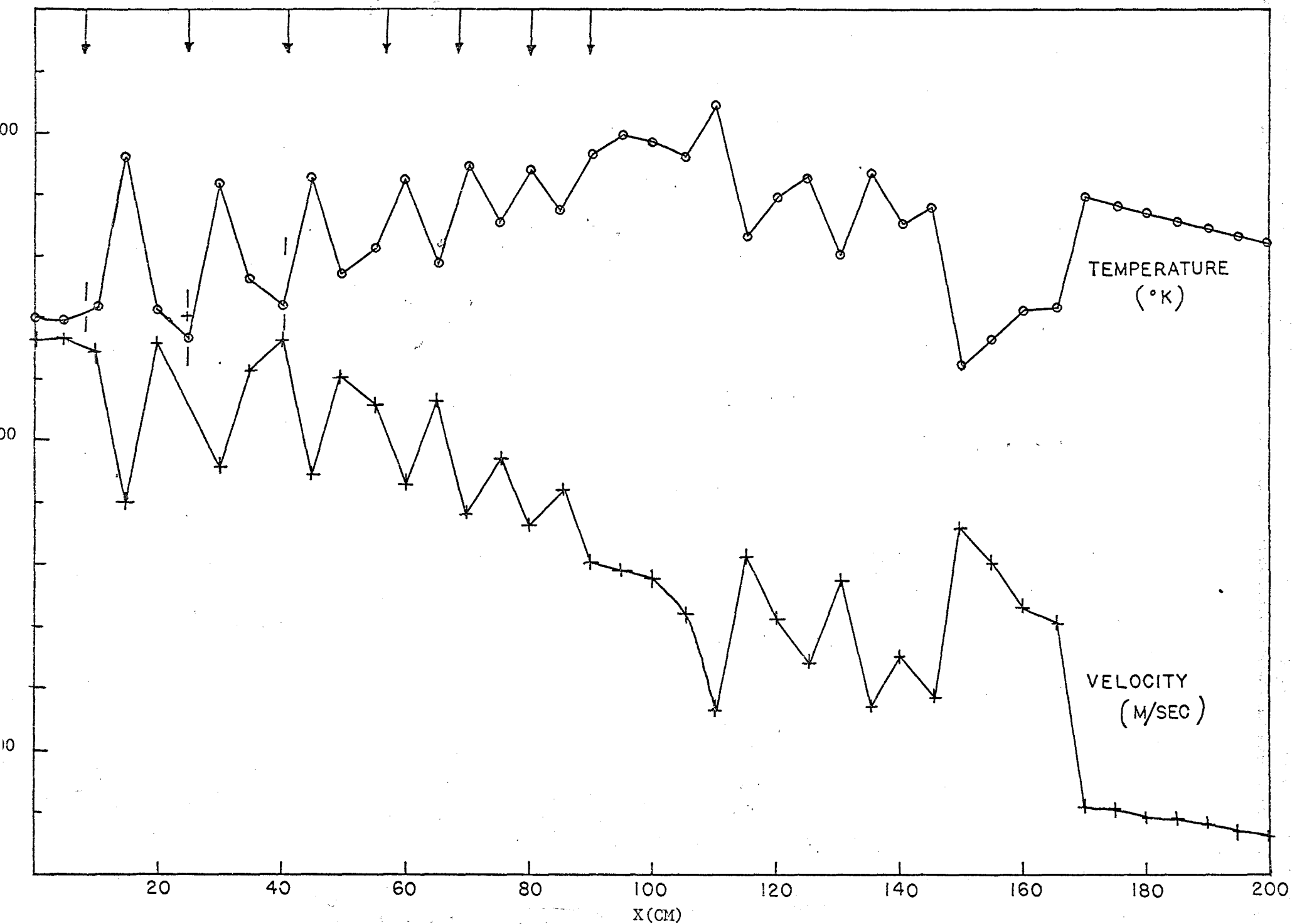


Fig. 3.34 CALCULATED RESULTS USING THE REP3 COMPUTER CODE

However, the behavior of the calculated profiles for X greater than 90-mm are disturbing in that there are a series of temperature and velocity oscillations which do not agree with any experimental or physical evidence. Furthermore, the predicted sharp drop (rise) in temperature (velocity) at 145-mm and the sharp rise (drop) in temperature (velocity) at 165-mm do not correspond to observed physical phenomena.

Clearly, further efforts are needed to ascertain whether or not the REP3 Program is correct. Unfortunately, there was no available time or funds to consider this further.

It was also discouraging that over 30 minutes of computer time on the CDC 6600 was required to obtain the small step size required to resolve the near field shock waves.

4. SUMMARY AND CONCLUSIONS

A research project was conducted which resulted in the development of an extended version of the MIRADCOM (Jackson) Infrared Radiation Computer Model.

The extended computer code was validated by comparing the predicted spectral radiance of H_2O and CO_2 in the 2.7- μm and 4.3- μm wavelength bands with experimental data and with predictions of the NASA (Reardon) Radiation Computer Code. The predicted results agree fairly well with the NASA code even though there were some slight differences. With the exception of a few results for low pressure CO_2 , the predicted values were within $\pm 20\%$ of the measured values.

The infrared computer code was used to predict the infrared radiation emitted in the wavelength interval 4 to 5 microns from the exhaust of a full scale turbojet aircraft. The predicted results appeared to be reasonable when compared with the available test data.

The infrared radiation band model was used to calculate the variation of infrared radiation along the exhaust plume of a small kerosene/gaseous oxygen rocket engine. The temperature and species concentration profiles were predicted using the LAPP and REP3 computer codes. The agreement between the theoretical calculations and the experimental measurements was poor. The lack of agreement was due to the fact that the calculated temperature distributions did not exhibit the spatial structure due to shock waves that are evident in the measured infrared data and are visible in photographs of the exhaust plume. There appears to be short comings in the ability of both the LAPP and REP3 codes to predict the correct spatial distributions of temperature and specie concentrations.

The Gas Dynamics Finite Difference Computer Model developed as part of this project was only partially completed and was able to produce reasonable numerical results at the expense of long computer runs and high computer memory requirements. The code successfully calculated the flow field for flow geometries having solid boundaries, but had problems with the free boundaries encountered in nozzle exhaust plumes. The results obtained for the near field of the nozzle appeared to be correct, but the downstream results were not validated.

The preliminary calculations obtained using the REP3 and TKE computer codes did not agree with the experimental data.

Due to the partial success of the finite difference and REP3 computer codes, further development and assessment are required before these codes are fully operational in a validated model.

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APPENDIX I

LISTING OF IR COMPUTER PROGRAM

IIQIXII II I

I II IIXIQII

//HARWE JOB (36931,51118,REMOTE3,R-192,T-35,L-10),NGUYEN

//STEP1 EXEC FORTGCLG,FORTREG=192K,GOREG=192K

//FORT.SYSPRINT DD DUMMY

//FORT.SYSIN DD *

COMMON/PL0T1/WAVEI,WAVEF,CENTR,TOTRAD

COMMON /RASP/RASP

COMMON /RANDP/ XKK(254,7),DDINV(254,7),UK(254),AKU(254)

DIMENSION I6(3),I8(10),I7(6),I2(10),I3(6),I4(4),I5(10),I1(5)

COMMON /RPN1/ NWAV,WAVNO(254),WAVLT(254)

COMMON/CHECK/IREAD

COMMON /INDATA/ X(40),R(40,30),T(40,30),PC02(40,30),
S PH20(40,30),PH2(40,30),PO2(40,30),PC0(40,30),PH2(40,30)

COMMON /CTROL/ TITLE(12),KOUNT,NXS,NORUN

COMMON /OTH/ PP,RJ,NRP(40)

COMMON /SWITCH/ KSWIT ,IXXJG

C

DIMENSION RADSTL(254),RADXT(254),TRLAST(254)

000130

DIMENSION CR(100),CX(100),CL(100)

000140

DIMENSION RADXL(254),RRN(100)

DIMENSION STRAD(40)

DIMENSION RAD(254),RADX(254),TRNS(254),P(254),O(254),RADSTR(254) 000230

DIMENSION WAVL5(254)

000240

DIMENSION RCALC(100)

000260

DIMENSION SRAD(254),SRADSR(254)

INTEGER IH20(2),IC02(2),IWRT(2)

REAL HRADXT(40),CRADXT(40)

REAL HCTOTX(40),HCRDXT(40)

REAL HRADX(254),CRADX(254),SRADX(254)

EQUIVALENCE (HRADX(1),RADX(1)),(CRADX(1),RAD(1)),

S (SRADX(1),TRNS(1))

EQUIVALENCE (HCTOTX(1),CR(1)),(HCRDXT(1),CX(1))

C

DATA I6N/10/

DATA I8N/37/

DATA I7N/23/

DATA I2N/33/

DATA I3N/25/

DATA I4N/19/

DATA I5N/49/

DATA I1N/20/

DATA ASTER,RCALC/1H*,100*1H /

000300

DATA IDISC/8/

DATA ISTOP/0/

DATA IDSCT/0/

DATA JPASS/0/

DATA IFLG/0/

ISWIT = 2

KSWIT = 2

JSWIT=2

LSWIT = 1

IREAD=0

KPZO=0

VZO=1

KNFO=0

IXXJG=0

C

C SENSE SWITCH 1 SET--YES,PRINT FLOW FIELD DATA

C SENSE SWITCH 2 SET--YES,CO2 RADIATION--NO,H2O RADIATION

C SENSE SWITCH 3 SET--YES,AEROCHAM INPUT TAPE--NO, MOTOR DATA INPUT TAPE

C SENSE SWITCH 4 SET--YES,H2O AND CO2 RADIATION COMBINED

C

IOTRAD=0.0

000590

C
C
C
C
C

SET PRINT LABEL

DETERMINE MODE OF INPUT TAPE -- MOTOR DATA OR AEROCHEM

172 CONTINUE

END FILE 9

REWIND 9

END FILE 8

REWIND 8

IF(KPZQ.EQ.1) GO TO 104

GO TO (101,102),JSVIT

102 CONTINUE

CALL READER

KPZQ=1

GO TO 104

101 CONTINUE

104 CONTINUE

C

NXS=KOUNT

DO 111 I=1,NXS

IF(NRP(I).GT.25) GO TO 111

NRP1=NRP(I)+1

ISET=NRP(I)

DO 112 KK=NRP1,25

R(I,KK)=R(I,ISET)

T(I,KK)=T(I,ISET)

PCO2(I,KK)=PCO2(I,ISET)

PH2O(I,KK)=PH2O(I,ISET)

PN2(I,KK)=PN2(I,ISET)

PO2(I,KK)=PO2(I,ISET)

112 CONTINUE

111 CONTINUE

C

READ WAVELENGTHS AND CORRESPONDING BAND PARAMETERS

C

READ IN H2O BAND DATA FIRST FOLLOWED BY CO2 BAND DATA

C

47 CONTINUE

IF(LSWIT .EQ.1.AND.ISTOP.EQ.3.AND.IFLG.LT.2) GO TO 21

IF(ISTOP.EQ.3) GO TO 23

21 CONTINUE

WRITE (6,9911)

000330

READ (5,9912) NNAV1, NNAV2

000340

WRITE (6,10035) NNAV1,NNAV2

10035 FORMAT(10H NNAV1 ,2(1X,15))

9912 FORMAT(2I10)

NNAV=1 + ((NNAV2-NNAV1)/25)

000360

DO 1 I=1,NNAV

000380

IF(KNTPQ.EQ.0) GO TO 9957

GO TO 9967

9957 READ (5,9922) WAVN7(I),(XKK(I,J),J=1,7)

9922 FORMAT (F9.1,7E9.2)

000400

GO TO 9977

9967 READ (5,9922) WAVN7(I),(XKK(I,J),J=1,7)

9977 CONTINUE

WRITE (6,10037) WAVN7(I),(XKK(I,J),J=1,7)

10037 FORMAT(10H WAV ,F9.1,7E9.2)

IF(KNTPQ.EQ.0) GO TO 9987

GO TO 9997

9987 READ (5,9901) WAVN7(I),(ODINV(I,J),J=1,7)

9901 FORMAT(F6.1,E9.2,6E10.2)

GO TO 9997

9997 READ (5,9910) WAVN7(I),(ODINV(I,J),J=1,7)

9910 FORMAT (F6.1,E9.2,6E10.2)

9947	CONTINUE	
	WRITE (6,10037) WAVNO(I), (DDINV(I,J), J=1,7)	
	WAVLT(I)=10000./WAVNO(I)	000420
9910	FORMAT (1X,8E15.7)	000450
1	WAVLS(I)=WAVLT(I)**5	000460
	WAVEI=WAVNO(I)	000470
	WAVEF=WAVNO(NWAV)	000480
	WRITE (6,9911)	000490
9911	FORMAT (1H1)	000500
	DO 9913 I=1,NWAV	000510
	DO 9913 J=1,7	000520
	IF (DDINV(I,J) .LE. 0.0) DDINV(I,J)=1.0E-20	000530
	IF (XKK(I,J) .LE. 0.) XKK(I,J)=1.0E-20	
9913	CONTINUE	000540
9000	FORMAT(F9.4,7E9.3)	
C		
C	PRINT BAND PARAMETERS IF OPTIONED SELECTED	
C		
	IF (LSWIT .EQ. 1. AND. ISTOP.EQ. 0. AND. IFLAG.LT. 1) CALL SAVE1	
	IF (ISWIT .EQ. 1) CALL BPPRT	
C		
23	CONTINUE	
	IF (LSWIT .EQ. 1. AND. ISTOP.EQ. 3. AND. IFLG.GE. 2) CALL SWAP	
	IFLG=IFLG+1	
C		
C	DETERMINE INPUT DATA PRINT OPTIONS	
C		
C	SUBROUTINE PRMD WRITES OUT INPUT FLOW FIELD DATA	
	IF (ISWIT .EQ. 1) CALL PRMD	
C		001520
C	ALPHA=ASPECT ANGLE IN DEGREES(0 DEG=NOSE-ON)	001530
C	ISTOP=1, WILL TERMINATE PROGRAM AFTER SINGLE ALPHA CALCULATION	001540
C	ISTOP=2, RETURNS TO THIS POINT FOR ANOTHER ALPHA	001550
C	ISTOP=3, RETURNS TO STMT 172 TO READ NEW SET FLOW FIELD DATA	001560
C	SENSE SWITCH 4 WILL OVER RIDE ANY ISTOP OPTION	
C		001570
103	READ(5,1010) ALPHA, ISTOP	001580
1010	FORMAT(F10.5, I5)	001590
	ALP=ALPHA/57.295780	001600
	CC=COS(ALP)	001610
	SC=SIN(ALP)	001620
	PI=3.1415926	001630
	TRADX=0.0	001640
	TOTRAD=0.0	001650
	CENTER=0.0	
C		
	DO 185 II=1,NWAV	001670
185	RADXT(II)=0.0	001680
C		
	WRITE (6,1401)	001690
	DO 3000 I=1,NXS	001700
	XZPEX=NRP(I)	
	WRITE (6,1401)	001710
1401	FORMAT (1H1)	001720
	WRITE (6,1404) (TITLE(JRR), JRR=1,12)	001730
1404	FORMAT(10X,12A4,42, /)	001740
	WRITE (6,1400) I, X(I), ALPHA	001750
1400	FORMAT (1X,12HSTATION NO.=,14,5X,17HAXIAL DIST. (CM)=,F10.3,5X,	001760
	*14HASPECT ANGLE (DEG)=,F5.1, /)	
	TRADX=0.0	001780
	IX=I	001800
	RD=R(IX,25)	001810
	IF (RD .LE. 0.0) RD=1.0E-20	001820
	X1=X(IX)+.01	001830
	IF (ALPHA.EQ.180.0) X2=X(NXS)	001840
	ZPEX=XZPEX	001850

```

C      ****CHANGE DELZ=RO/ZPEZ TO DELZ=1. FOR THE SLAP***
      DELZ=1.0
      DELZ=RO/ZPEX
C      ***** SET NZU=NZPEX FOR TOTAL PLUME *****
      ZU=RASP/DELZ
      NZU=NZPEX
C
      DO 190 II=1,NWAV
      SRADSR(II)=0.0
190    RADX(II)=0.0
      STOT=0.
C
C      THIS CARD HAS BEEN CHANGED AS
C      DO 2000 M=1,NZPEX
      DO 2000 M=1,NZU
      ZM=M
      ZO=DELZ*(ZM-1.)
      IF(ALPHA.EQ.0.0) RO=2.*R(1,25)+(ZM-1.)*(R(NXS,25)-2.*R(1,25))/ZPEX
      IF(ALPHA.EQ.180.0) RO=(ZM-1.)*(R(NXS,25))/ZPEX
C
C*****
C      CALCULATE COORDINATES X,R ALONG LINE-OF-SIGHT AS FUNCTION OF ITS
C      LENGTH
C
      DELL=10.
      IF(ALPHA.EQ.0.0.OR.ALPHA.EQ.180.0) DELL=X(NXS)/98.
200    XL=0.
C
      DO 203 J=1,100
      N=J
      APR=XL*SC-SQRT(RO*RO-ZO*ZO)
      CR(N)=SQRT(ZO*ZO+APR*APR)
      CX(N)=XO+XL*CC
      CL(N)=XL
      IF(CX(N).LT.0.) GO TO 204
      IF(CX(N).GT.X(NXS)) GO TO 204
      IF(ALPHA.EQ.0.0.OR.ALPHA.EQ.180.0) GO TO 203
      DO 201 K=1,NXS
      KI=K
      IF(CX(N).LE.X(K)) GO TO 202
201    CONTINUE
202    IF(CR(N).GT.R(KI,25) .AND. N .GT. 10) GO TO 204
203    XL=XL+DELL
C
C      DELL TOO SMALL, INCREASE AS FOLLOWS
C
      DELL=DELL*1.5
      GO TO 200
204    CONTINUE
      IF(N.GE.50) GO TO 299
C
C      DELL TOO LARGE, DECREASE AS FOLLOWS
C
205    DELL=CL(N)/50.
C
C      PLUME THICKNESS NEGLECTABLE AT THIS POINT, MOVE ON TO NEXT POINT
C      AFTER SETTING ALL SPECTRAL RADIATION VALUES EQUAL TO ZERO
C
      IF(DELL.GT..005) GO TO 200
210    CONTINUE
C
      DO 211 II=1,NWAV
      RADSTR(II)=0.0
211    RADSTR(II)=0.0
      GO TO 1990
C

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C*****002470
C   CALCULATE PRESSURE AND TEMPERATURE ALONG LINE-OF-SIGHT AT 5ACH X, R002480
C   002490
299 CONTINUE
   WRITE(6,307)DELTZ,DELL
307 FORMAT(2(1X,E10.4))
C
   DO 300 II=1,NWAV
   P(II)=0.0
   Q(II)=0.0
   TRNS(II)=1.0
300 RADSTR(II)=0.0
C
   DO 1000 J=1,N
   JCHECK=J
C
   DO 301 L=2,NXS
   IF(JCHECK.EQ.99) LCHECK=3
   NX1=L-1
   NX2=L
   IF(CX(J).LE.X(L)) GO TO 302
301 CONTINUE
C
302 CONTINUE
   DO 303 L=2,25
   IF(JCHECK.EQ.99) LCHECK=3
   NR2=L
   IF(CR(J).LE.R(NX2,L)) GO TO 304
303 CONTINUE
C
304 CONTINUE
   DO 305 L=2,25
   NR1=L
   IF(CR(J).LE.R(NX1,L)) GO TO 306
305 CONTINUE
C
306 CONTINUE
   DR1=CR(J)-R(NX1,NR1-1)
   DR2=R(NX1,NR1)-R(NX1,NR1-1)
   DR3=CR(J)-R(NX2,NR2-1)
   DR4=R(NX2,NR2)-R(NX2,NR2-1)
   DX1=CX(J)-X(NX1)
   DX2=X(NX2)-X(NX1)
C
C   DR1DR2,DR3DR4,DX1DX2 WERE CREATED TO ELIMINATE INDEFINITE FORMS
C
   DR1DR2=0.0
   DR3DR4=0.0
   DX1DX2=0.0
   IF(DR2.NE. 0.0) DR1DR2=DR1/DR2
   IF(DR4.NE. 0.0) DR3DR4=DR3/DR4
   IF(DX2.NE. 0.0) DX1DX2=DX1/DX2
   TX1=T(NX1,NR1-1)+(T(NX1,NR1)-T(NX1,NR1-1))*DR1DR2
   TX2=T(NX2,NR2-1)+(T(NX2,NR2)-T(NX2,NR2-1))*DR3DR4
   CT=TX1+(TX2-TX1)*DX1DX2
   PH2DX1=PH2D(NX1,NR1-1)+(PH2D(NX1,NR1)-PH2D(NX1,NR1-1))*DR1DR2
   PH2DX2=PH2D(NX2,NR2-1)+(PH2D(NX2,NR2)-PH2D(NX2,NR2-1))*DR3DR4
   CPH2D=PH2DX1+(PH2DX2-PH2DX1)*DX1DX2
C
   IF(JCHECK.EQ.99)JCHECK=11
   PC02X1=PC02(NX1,NR1-1)+(PC02(NX1,NR1)-PC02(NX1,NR1-1))*DR1DR2
   PC02X2=PC02(NX2,NR2-1)+(PC02(NX2,NR2)-PC02(NX2,NR2-1))*DR3DR4
   CPC02=PC02X1+(PC02X2-PC02X1)*DX1DX2
C
   PH2X1=PH2(NX1,NR1-1)+(PH2(NX1,NR1)-PH2(NX1,NR1-1))*DR1DR2
   PH2X2=PH2(NX2,NR2-1)+(PH2(NX2,NR2)-PH2(NX2,NR2-1))*DR3DR4

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CPN2=PN2X1+(PN2X2-PN2X1)*DX1DX2
PO2X1=PO2(NX1,NR1-1)+(PO2(NX1,NR1)-PO2(NX1,NR1-1))*DR1DR2
PO2X2=PO2(NX2,NR2-1)+(PO2(NX2,NR2)-PO2(NX2,NR2-1))*DR3DR4
CPO2=PO2X1+(PO2X2-PO2X1)*DX1DX2
C
PH2X1=PH2(NX1,NR1)+(PH2(NX1,NR1)-PH2(NX1,NR1-1))*DR1DR2
PH2X2=PH2(NX2,NR2-1)+(PH2(NX2,NR2)-PH2(NX2,NR2-1))*DR3DR4
CPH2=PH2X1+(PH2X2-PH2X1)*DX1DX2
PCOX1=PCO(NX1,NR1)+(PCO(NX1,NR1-1)-PCO(NX1,NR1-1))*DR1DR2
PCOX2=PCO(NX2,NR2-1)+(PCO(NX2,NR2)-PCO(NX2,NR2-1))*DR3DR4
CPCO=PCOX1+(PCOX2-PCOX1)*DX1DX2
IF(CT.LT.300.) CT=300. 002940
IF(CPH20.LT.0.) CPH20=1.E-20
IF(CPCO2.LT.0.) CPCO2=1.E-20
IF(CPN2.LT.0.) CPN2=1.E-20
IF(CPO2.LT.0.) CPO2=1.E-20
IF(CPH2.LT.0.0)CPH2=1.E-20
IF(CPCO.LT.0.0)CPCO=1.E-20
AVERAGE TEMPERATURES AND PRESSURES FOR BLOCK 002980
IF(J.GT.1) GO TO 20
CPH20L=CPH20
CPCO2L=CPCO2
CPN2L=CPN2
CPO2L=CPO2
CPH2L=CPH2
CPCOL=CPCO
CTL=CT
20 CONTINUE
C
PCO2AV=(CPCO2+CPCO2L)/2. 003000
PH20AV=(CPH20+CPH20L)/2. 002990
PN2AV=(CPN2+CPN2L)*.5
PO2AV=(CPO2+CPO2L)*.5
PCOAV=(CPCO+CPCOL)/2.
PH2AV=(CPH2+CPH2L)*0.5
C
TAV=(CT+CTL)/2. 003020
CPH20L=CPH20 003030
CPCO2L=CPCO2 003040
CPN2L=CPN2
CPO2L=CPO2
CPH2L=CPH2
CPCOL=CPCO
CTL=CT 003070
IF(J.EQ.1) GO TO 1000 003080
CALL PARAM(TAV,DEGL,PH20AV,PN2AV,PO2AV,PCO2AV,PCOAV,PH2AV)
1002 CONTINUE
C
ON 2005 II=1,NRAV 003100
TRLAST(II)=TRNS(II) 003110
P(II)=P(II)+OK(II)
Q(II)=Q(II)+AKU(II) 003130
IF (P(II) .EQ. 0.0) P(II)=1.0E-20 003140
IF (Q(II) .EQ. 0.0) Q(II)=1.0E-20 003150
ARG=-P(II)/ SQRT(1+.25*P(II)**2/Q(II))
TRNS(II)=EXP(ARG) 003180
RAD(II)=(TRLAST(II)-TRNS(II))*37413./WAVG5(II)/(EXP(1.439*WAVNO(II)/TAV)-1.)/3.1415926536 003200
RADSTR(II)=RADSTR(II)+RAD(II)
SRAD(II)=SRAD(II)*10000./WAVNO(II) **2
SRADSR(II)=SRADSR(II)+SRAD(II)
2005 CONTINUE
1000 CONTINUE 003230
C

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WRITE (6,2023)
2023 FORMAT (20X,'SPECTRAL RADIATION',/,10X,'WAVNO',10X,'SRADSR')
DO 2022 II=1,NWAV
2001 STOT = STOT+SRADSR(II)
C WRITE(6,2024)WAVNO(II),SRADSR(II)
C ***TAKE OUT WRITE STATEMENT OF SRADSR(II)***
2024 FORMAT(2(10X,E10.4))
2022 CONTINUE
CHECK=1, 003260
C
C RAD(II)=RADIATION FOR EACH BLOCK AT EACH WAVELENGTH (WAVLT(II))
C RAD(II) IN (WATTS/CM2-SR-MICRON))
C RADX(II) IN (WATTS/CM-SR-MICRON)
C TRADX IN (WATTS/CM-SR)
C RADXT(I) IN (WATTS/SR/MICRON)
C RADSTR(II)=RADIATION AT EACH X AND Z SUMMED OVER EACH PATH LENGTH
C RADX(II)=RADIATION TOTAL AT EACH X-STATION (SUMMED OVER Z)
C TRADX=RADX(I) INTEGRATED OVER WAVELENGTH 003290
C RADXT(I)=RADX(I) INTEGRATED OVER ALL X-STATIONS 003300
C TOTRAD=TRADX INTEGRATED OVER X 003310
C
1980 CONTINUE
1990 IF (M.EQ.1) GO TO 1998 003320
CHECK=2. 003330
IF(ALPHA.NE.0..AND.ALPHA.NE.180.) GO TO 1996
1981 DO 1982 II=1,NWAV 003360
IJ=II-1
1982 RADXT(II)=RADXT(II)+(RADSTR(II)+RADSTL(II))/2.*PI*DELZ**2*((ZM-1.) 003370
1**2-(ZM-2.)**2) 003380
CHECK=3. 003390
C THIS CARD HAS BEEN CHANGED AS
C IF (M.EQ.NZPEX) GO TO 3001
IF (M.EQ.NZU) GO TO 3001
GO TO 1998 003410
1996 DO 1997 II=1,NWAV 003420
C INTEGRATED RADIATION IS DOUBLED TO INCLUDE BOTH HALVES OF PLUME 003430
RADX(II)=RADX(II)+(RADSTR(II)+RADSTL(II))
C THIS CARD HAS BEEN CHANGED AS
C IF (M.EQ.NZPEX) RAD(II)=RADX(II)+RADSTR(II)
IF (M.EQ.NZU) RAD(II)=RADX(II)+RADSTR(II)
1997 CONTINUE
1998 DO 1999 II=1,NWAV 003450
IJ=II-1
1999 RADSTL(II)=RADSTR(II) 003460
2000 CONTINUE 003470
WRITE (6,2002)STOT
2002 FORMAT (1X,'SUM UP SRADSR AT X STATION',5X,E10.4)
WRITE(6,1402)
1402 FORMAT(8X,'WAVE NUMBER',5X,'WAVE LENGTH',5X,'RADIATION',
S 23X,'WAVE NUMBER',5X,'WAVE LENGTH',6X,'RADIATION'/)
WRITE(6,14021)
14021 FORMAT(42X,'W/SR/CM/MICRON')
C
LINCT=0
DO 12 II=1,NWAV,2
RADX(II)=RADX(II)*DELZ
IF(II.LE.NWAV) RADX(II+1)=RADX(II+1)*DELZ
IF(RADX(II).LT.0.) WRITE(6,18) RADX(II)
18 FORMAT(' NEGATIVE RADIATION IS ILLOGICAL ',E14.4)
IF((II+1).GT.NWAV) GO TO 15
WRITE(6,16) WAVNO(II),WAVLT(II),RADX(II),
S WAVNO(II+1),WAVLT(II+1),RADX(II+1)
16 FORMAT(5X,E10.0,E16.4,7X,E12.5,20X,E10.0,E16.4,7X,E12.5)
LINCT=LINCT+1
IF(LINCT.LE.54) GO TO 12
LINCT=0

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12 CONTINUE
GO TO 17
15 CONTINUE
WRITE(6,16) WAVNO(II),WAVLT(II),RADX(II)
17 CONTINUE
C
1403 FORMAT (1X,F10.0,F16.4,7X,E12.5) 003530
DO 2997 II=2,NWAV 003560
TRADX=TRADX+(WAVLT(II-1)-WAVLT(II))*
S (RADX(II-1)+RADX(II))*0.5
2997 CONTINUE
IF (I.EQ.1) GO TO 2999 003710
CENTER = CENTER+(X(I)-X(I-1))*0.5*(TRADX+TRADXL)*X(I)
TOTRAD=TOTRAD+(TRADX+TRADXL)*0.5*(X(I)-X(I-1)) 003730
STRAD(I)=(TRADX+TRADXL)*0.5*(X(I)-X(I-1))
DO 2998 II=1,NWAV 003740
RADXT(II)=RADXT(II)+((RADX(II)+RADXL(II))/2.
S )*(X(I)-X(I-1))*SC
2998 CONTINUE
WRITE(6,5000) TOTRAD
5000 FORMAT(1H0,24HTOTAL RADIATION EMITTED=,E10.4,11H WATTS/STER)
WRITE(6,1402)
DO 5001 II=1,NWAV,2
IF((II+1).GT.NWAV) GO TO 5002
WRITE(6,16) WAVNO(II),WAVLT(II),RADXT(II),WAVNO(II+1),WAVLT(II+1)
+,RADXT(II+1)
5001 CONTINUE
5002 CONTINUE
WRITE(6,16) WAVNO(II),WAVLT(II),RADXT(II)
2999 TRADXL=TRADX 003760
DO 2995 II=1,NWAV 003770
2995 RADXL(II)=RADX(II) 003780
RRN(I)=TRADX 003790
WRITE (6,1492) I,TRADX 003800
1492 FORMAT(/5X,'TOTAL RADIATION FOR STATION',I3,1X,1H=,1X,E12.5,
1 'WATTS/SR/CM')
C
C CHECK FOR COMBINE RUN. IF SO OUTPUT DATA TO DISC
C
IF(LSWIT .EQ.2) GO TO 22
WRITE(IDISC) NWAV,I,TRADX
WRITE(IDISC) (WAVNO(JJ),WAVLT(JJ),RADX(JJ),JJ=1,NWAV)
IDISC=IDISC+1
22 CONTINUE
3000 CONTINUE 003820
WRITE (6,1401) 003830
CENTR=0.0 003880
3001 CONTINUE
3003 CENTR=CENTER *SC/TOTRAD
TOTRAD=TOTRAD*SC 003920
IF (TOTRAD .EQ.0.) TOTRAD=1.*10.**(-30)
IF(LSWIT .EQ.1) WRITE(IDISC) TOTRAD,(RADXT(J),J=1,NWAV)
3004 WRITE (6,1404) (TITLE(J),J=1,8) 003910
WRITE (6,1405)
1405 FORMAT (70X,'STATION RADIATION (W/SR)')
WRITE (6,1407)(X(II),PRE(II),STRAD(II),RCALC(II),II=1,NXS)
1407 FORMAT (F15.2,2HX,E13.7,20X,E13.7,1X,A1)
C
WRITE (6,8421) 004030
C CALL FICAL PLOT ROUTINE
C
C
8421 FORMAT(/10X,80H* RADIATION VALUE WAS CALCULATED BASED ON THE SLOPE 004040
* BETWEEN THE PREVIOUS TWO POINTS) 004050
WRITE(6,1211) TOTRAD
1211 FORMAT(1H0,24HTOTAL RADIATION EMITTED=,E10.4,11H WATTS/STER) 004080

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      IF(LSWIT .EQ.1) WRITE(IDISC) TOTRAD
      WRITE(6,1254) CENTR
1254 FORMAT(1H0,9HCENTRID=,E11.4,'CM')
C
C CHECK FOR COMBINED RUN. SENSE SWITCH 4 WILL OVER RIDE THE ISTOP OPTION
C
      IF(LSWIT .EQ.1) GO TO 24
      GO TO (107,103,172),ISTOP
C
C DETERMINE IF ANOTHER PASS IS TO BE PERFORMED
C
24 CONTINUE
      JPASS=JPASS+1
      IDISC=IDISC+1
      KSWIT=1
      IF(JPASS.LT.2) GO TO 47
      KSWIT=2
      JPASS=0
      IDISC=8
C
C READ H2O AND CO2 RADIATION BACK FROM DISC.
C SUM AND OUTPUT
C
      END FILE 9
      REWIND 9
      END FILE 8
      REWIND 8
      WRITE(6,27)
27 FORMAT(1H1,30X,' C O M B I N E D H 2 O C O 2 R A D I A T I O N
S S U M S ')
      WRITE(6,28) NWAV1,NWAV2
28 FORMAT(1H //,43X,'SUMMED OVER WAVE NUMBERS ',2I6)
C
      DO 25 JK=1,NXS
      READ(8) NWAV,I,CTRADX
      READ(9) NWAV,I,HTRADX
      CALL ERRCHK(2)
      READ(8) (WAVNO(JJ),WAVLT(JJ),CRADX(JJ),JJ=1,NWAV)
      READ(9) (WAVNO(JJ),WAVLT(JJ),HTRADX(JJ),JJ=1,NWAV)
      CALL ERRCHK(3)
      DO 26 JJ=1,NWAV
      SRADX(JJ)=HTRADX(JJ)+CRADX(JJ)
26 CONTINUE
      HCTOTX(JK)=HTRADX+CTRADX
C
      WRITE(6,9911)
      WRITE(6,1402)
      DO 29 JJ=1,NWAV,2
      IF((JJ+1).GT.NWAV) GO TO 30
      WRITE(6,16) WAVNO(JJ),WAVLT(JJ),SRADX(JJ),
S WAVNO(JJ+1),WAVLT(JJ+1),SRADX(JJ+1)
29 CONTINUE
      GO TO 31
30 CONTINUE
      WRITE(6,16) WAVNO(JJ),WAVLT(JJ),SRADX(JJ)
31 CONTINUE
C
      WRITE(6,32) I,HCTOTX(JK)
32 FORMAT(//10X,' TOTAL COMBINED RADIATION FOR STATION ',I3,
S ' IS ',E12.4)
      IF(HCTOTX(18).EQ.0.) GO TO 25
25 CONTINUE
C
      READ(8) CTRAD , (CRADXT(J),J=1,NWAV)
      READ(9) HTRAD , (HTRADT(J),J=1,NWAV)
      CALL ERRCHK(4)

```

```

WRITE(6,9911)
WRITE(6,34)
34 FORMAT(20X,'WAVNO',10X,'WAVELENGTH',10X,'HCRADXT'//)
C PUNCH CARDS FOR ATMOSPHERIC TRANSMISSION PROGRAM
NN=NWAV
DO 33 JJ=1,NWAV
HCRDXT(JJ)=HRADXT(JJ)+CRADXT(JJ)
WRITE(6,35) WAVNO(JJ),WAVLT(JJ),HCRDXT(JJ)
35 FORMAT(17X,E12.4,6X,E12.4,6X,E12.4)
33 CONTINUE
STORAD=HTORAD+CTORAD
WRITE(6,43) (TITLE(J),J=1,8)
43 FORMAT(1H ///,' SUMMARIES FOR ',7A10,A2//)
WRITE(6,36) HTORAD
WRITE(6,37) CTORAD
36 FORMAT(10X,'TOTAL H2O RADIATION',E14.4)
37 FORMAT(10X,'TOTAL CO2 RADIATION',E14.4)
WRITE(6,41)
41 FORMAT(29X,'-----')
WRITE(6,42) STORAD
42 FORMAT(' TOTAL COMBINED RADIATION',4X,E14.4)
C
GO TO (107,107,172),ISTOP
107 CONTINUE
STOP
END
SUBROUTINE READER

```

004420

```

C THIS ROUTINE WILL READ ONE COMPLETE BLOCK OF
C MOTOR DATA PER CALL
C

```

```

COMMON /INDATA/ X(40),R(40,30),T(40,30),PCO2(40,30),
S PH2O(40,30),PM2(40,30),PO2(40,30),PCO(40,30),PH2(40,30)
COMMON /CTROL/ TITLE(12),KOUNT,NXS,NORUN
COMMON /OTH/ PP,RJ,NRP(40)
COMMON /RASP/RASP
DATA JJ/1/

```

```

C
READ(5,100) (TITLE(J),J=1,12)
WRITE(6,100) (TITLE(J),J=1,12)

```

```

100 FORMAT(12A4)

```

```

READ(5,101)RASP

```

```

READ(5,300) PP,RJ,NRP(1)

```

```

RJ=RJ*100.

```

```

WRITE(6,300) PP,RJ,NRP(1)

```

```

300 FORMAT(2E10.4,I4)

```

```

KOUNT=0

```

```

C
C NOW READ IN THE ENTIRE BLOCK OF MOTOR DATA
C

```

```

30 CONTINUE
READ(5,400) X(JJ),ILAST
X(JJ)=X(JJ)*RJ
WRITE(6,700)X(JJ)

```

```

400 FORMAT (E10.3,69X,I1)

```

```

700 FORMAT (1X,E10.3)

```

```

IF(ILAST.EQ.1) GO TO 14

```

```

103 FORMAT(I5)

```

```

READ(5,103)NRP(JJ)

```

```

IX=NRP(JJ)

```

```

DO 12 KK=1,IX

```

```

C THE PUNCH DATA INPUT FOR RADIATION PROGRAM IN
C CHANGE READ (5,500) R(JJ,KK), T(JJ,KK), CARD 10
C THE FOLLOWING ORDER PCO2, PH2O, PM2, PO2
C +PH2(JJ,KK),PO2(JJ,KK)
C READ(5,500) R(JJ,KK),T(JJ,KK),PCO2(JJ,KK),PH2O(JJ,KK),

```

```

      READ(5,500)R(JJ, KK), T(JJ, KK), PCO2(JJ, KK), PH2O(JJ, KK)
500  FORMAT(E10.3, 20X, 3E10.3)
      R(JJ, KK)=R(JJ, KK)*RJ
      PCO(JJ, KK)=1.0E-20
      PH2(JJ, KK)=1.0E-20
      PO2(JJ, KK)=1.0E-20
      PN2(JJ, KK)=1.0E-20
      WRITE(6,600) R(JJ, KK), T(JJ, KK), PN2(JJ, KK), PCO(JJ, KK), PCO2(JJ, KK),
+PH2O(JJ, KK), PO2(JJ, KK)
600  FORMAT (1X, 8E10.4)
      12 CONTINUE
      JJ=JJ+1
101  FORMAT(E10.4)
      KOUNT=KOUNT+1
      GO TO 30
      14 CONTINUE
      CALL LIMIT(KOUNT, 40, 1)
      CALL LIMIT(NRP(1), 30, 2)
C
      RETURN
      END
      SUBROUTINE LIMIT(IARG, LIM, NPS)
C
      IF(IARG.GE.0.AND.IARG.LE.LIM) RETURN
C
      WRITE(6,10) IARG, LIM, NPS
10  FORMAT(' ILLEGAL INPUT PARAMETER ',
$      2I10/' CALLED FROM LOCATION ', I8)
C
      STOP
      END
      SUBROUTINE SWAP
C
      COMMON /BANDP/ XXX(254, 7), DDINV(254, 7), UK(254), AKU(254)
      COMMON /RPN1/ NWAV, WAVVJ(254), WAVLT(254)
C
      DIMENSION XSAVE(254, 7), DSAVE(254, 7)
C
      DO 10 JJ=1, NWAV
      DO 10 KK=1, 7
      XHOLD=XXX(JJ, KK)
      DHOLD=DDINV(JJ, KK)
      XXX(JJ, KK)=XSAVE(JJ, KK)
      DDINV(JJ, KK)=DSAVE(JJ, KK)
      XSAVE(JJ, KK)=XHOLD
      DSAVE(JJ, KK)=DHOLD
10  CONTINUE
C
      NWVOLD=NWAV
      NWAV=NWLST
      NWLST=NWVOLD
C
      RETURN
C
      ENTRY SAVE1
C
      DO 11 JJ=1, NWAV
      DO 11 KK=1, 7
      XSAVE(JJ, KK)=XXX(JJ, KK)
      DSAVE(JJ, KK)=DDINV(JJ, KK)
11  CONTINUE
C
      NWLST=NWAV
C
      RETURN
      END

```

SUBROUTINE PRMD

C
COMMON /INDATA/ X(40),R(40,30),T(40,30),PCO2(40,30),
S PH2O(40,30),PH2(40,30),PO2(40,30),PCO(40,30),PH2(40,30)
COMMON /CTRL/ TITLE(12),KOUNT,NXS,NORUN
COMMON /OTH/ PP,RJ,NRP(40)

C
WRITE(6,10)
10 FORMAT(1H1)
WRITE(6,11) NORUN
11 FORMAT(1H ,25X,' M O T O R ',I3//)
WRITE(6,12) KOUNT
12 FORMAT(1H ,9X,'NUMBER OF X STATIONS IS',I4)

C
DO 14 JJ=1,KOUNT
NRP1=NRP(JJ)
WRITE(6,15) JJ,X(JJ),NRP1
15 FORMAT(//,' X(',I2,') = ',E12.4,
S 8X,' NUMBER OF POINTS AT THIS X STATION IS',I4)
WRITE(6,16)
16 FORMAT(1H ,4X,'R',13X,'T',11X,'CO2',10X,
S 'H2O',11X,'N2',11X,'O2'//)
DO 14 KK=1,NRP1
WRITE(6,18) R(JJ,KK),T(JJ,KK),PCO2(JJ,KK),
S PH2O(JJ,KK),PN2(JJ,KK),PO2(JJ,KK)
18 FORMAT(1H ,6(E12.4,1X))
14 CONTINUE

C
RETURN
END
SUBROUTINE BPPRT

C
C THIS ROUTINE PRINTS THE BAND PARAMETERS TOGETHER
C WITH THE ASSOCIATED WAVENUMBERS AND WAVELENGTHS

C
COMMON /BANDP/ XKK(254,7),DOINV(254,7),UK(254),AKU(254)
COMMON /RPN1/ NWAV,WAVNO(254),WAVLT(254)

C
WRITE(6,11)
11 FORMAT(1H1,51X,'B A N D P A R M E T E R S'//)
LINCT=0
WRITE(6,12)
12 FORMAT(10X,'WAVENUMBER',1X,'WAVELENGTH',20X,
S 'XKK AND DOINV PARAMETERS'//)

C
DO 10 JJ=1,NWAV
WRITE(6,14) WAVNO(JJ),WAVLT(JJ),(XKK(JJ,K),K=1,7)
WRITE(6,14) WAVNO(JJ),WAVLT(JJ),(DOINV(JJ,K),K=1,7)
14 FORMAT(14X,F5.0,4X,F8.1, 7X,7E12.4)
WRITE(6,15)
15 FORMAT(1H)
LINCT=LINCT+3
IF(LINCT.LT.56) GO TO 10
LINCT=0
WRITE(6,16)
16 FORMAT(1H1)
WRITE(6,12)
10 CONTINUE

C
RETURN
END
SUBROUTINE LIMITS(ARRAY,ND,RLBW,KUPPER)

C
REAL ARRAY(254)
C
WRITE(6,10)

```

      DO 10 JJ=1,N0
      IF (ARRAY(JJ).GT.RHIGH) RHIGH=ARRAY(JJ)
10  CONTINUE
C
      IUP=ALOG10(RHIGH)+1
      ILP=IUP-3
      RUPPER=10.**IUP
      RLOW=10.**ILP
C
      RETURN
      END
      SUBROUTINE ERRCHK(ILOC)
      RETURN
      END
      SUBROUTINE PARAM(TAV,XL,PH20,PN2,P02,PC02,PC0,PH2)
      COMMON/CHECK/IREAD
      COMMON /BANDP/ XKK(254,7),DDINV(254,7),UK(254),AKU(254)
      COMMON /RPN1/ NWAV,NAVNO(254),WAVLT(254)
      COMMON /SWITCH/ KSWIT ,IXXJG
      REAL TT(7)
      REAL LOG1,LOG2
C
      DATA TT(1),TT(2),TT(3),TT(4),TT(5),TT(6),TT(7)/300.,600.,1000.,
      $1500.,2000.,2500.,3000./
      DATA IDBG/-1/
      T=TAV
      IXXJG=IXXJG+ 1
      IF (KSWIT.NE.2) GO TO 23
      TT(3)=1200.
      TT(4)=1500.
      TT(5)=1800.
      TT(6)=2400.
23  CONTINUE
      IDBG=IDBG+1
      CALCULATE OPTICAL DEPTH, U, AND WIDTH, GAMMA
      T273=273./T
      SQT273=SQRT(T273)
      IF (KSWIT .EQ.2) GO TO 18
      U=T273*PH20*XL
      GAMMA=120.1*PH20/T+SQT273*(.09*(PH20+PN2)+.04*P02+.12*PC02
      $+.05*PH2+0.1*PC0)
      GO TO 19
18  CONTINUE
      U=T273*PC02*XL
      GAMMA=0.01*PC02*T273+(.07*PN2+0.055*P02+0.09*PC02+PH2*.08
      $+PC0*0.06)*SQT273
19  CONTINUE
C
C  SELECTION OF TEMPERATURES TO BE USED IN INTERPOLATIONS
      I=1
      DO 30 KK=1,7
      K=KK
      TT1=ABS(T-TT(K))
      IF (TT1.LE.2-20) GO TO 31
30  CONTINUE
      GO TO 32
31  KK=XKK(I,K)
      DDIV=DDINV(I,K)
      GO TO 16
32  CONTINUE
      DO 8 J=3,7
      K=J
      IF (TT(J)-T) 8,8,9
8  CONTINUE
9  CONTINUE

```

004470

004480

004490

004530

```

      IF(T.GE.TT(K-1)) K=K+1
      TTK=TT(K)-TT(K-2)
      TTK1=TT(K)-TT(K-1)
      TTK2=TT(K-1)-TT(K-2)
      DIV=TTK*TTK1*TTK2
C
1000 CONTINUE
      LOG1=ALOG(XKK(I,K-1)/XKK(I,K-2))
      LOG2=ALOG(XKK(I,K)/XKK(I,K-2))
      B=((TTK**2)*LOG1-(TTK2**2)*LOG2)/DIV
      C=(TTK2*LOG2-TTK*LOG1)/DIV
      XK=EXP(ALOG(XKK(I,K-2))+B*(T-TT(K-2))+C*(T-TT(K-2))**2)
C
      IF(T.LT.TT(2)) GO TO 14
      DLOG=ALOG(DDINV(I,K-2))
      DINV=EXP((ALOG(DDINV(I,K-1))-DLOG)*(T-TT(K-2))/(TT(K-1)-TT(K-2))
C      +DLOG)
      GO TO 16
C
14 CONTINUE
      DLOG=ALOG(DDINV(I,1))
      DINV=EXP((ALOG(DDINV(I,2))-DLOG)*(T-TT(1))/(TT(2)-TT(1))+DLOG)
C
16 CONTINUE
C
C CALCULATE PARAMETERS
C
      UK(I)=XK*U
      AKU(I)=GAMMA*DINV*UK(I)
      IF (TT1.LE.E-20) GO TO 17
C
      KDEX=K-2
      IDEX=K-1
      FLOW=.9*AMIN1(XKK(I,KDEX),XKK(I,IDEX))
      FHIGH=1.1*AMAX1(XKK(I,KDEX),XKK(I,IDEX))
      IF(XK.GE.FLOW.AND.XK.LE.FHIGH) GO TO 17
      B=ALOG(XKK(I,K-1)/XKK(I,K-2))/TTK2
      XK=EXP(ALOG(XKK(I,K-2))+B*(T-TT(K-2)))
      IF(XK.LT.FLOW.OR.XK.GT.FHIGH) WRITE(6,20)
      S WAVNO(I),I,XK,T,H,GAMMA,TTK,B,C,DINV,UK(I),AKU(I)
      20 FORMAT(1X,/, ' XK STILL OUT OF RANGE AT WAVNO = ',F10.0,' AND I = '
      S ,I5,/,1X,5E14.4/,1X,5E14.4)
      17 CONTINUE
      I=I+1
      IF (I-NWAV)29,29,15
      29 IF (TT1.LE.E-20)GO TO 31
      IF(I-NWAV) 1000,1000,15
C
15 CONTINUE
      IREAD=IREAD+1
      RETURN
      END
//GO.FT08F001 DD DSN=HARWE.PLOT8,UNIT=SYSDA,DISP=(NEW,DELETE),
// DCB=(RECFM=VSB,LRECL=100,BLKSIZE=1000,BUFNO=1),
// SPACE=(CYL,(5,5),RLSE)
//GO.FT09F001 DD DSN=HARWF.PLOT9,UNIT=SYSDA,DISP=(NEW,DELETE),
// DCB=(RECFM=VSB,LRECL=100,BLKSIZE=1000,BUFNO=1),
// SPACE=(Cyl,(5,5),RLSE)
//GO.SYST1 DD *

```

APPENDIX II

INPUT INSTRUCTIONS AND SAMPLE INPUT DATA

APPENDIX II

The computer program input is composed of two parts: the flow field section and the band parameter section. The flow field of a plume can be generated by several computer codes like the NASA MOC code or the AeroChem code [3 and 5]. The flow field program provides the temperature and partial pressure at different locations along the plume. This information is necessary to calculate the optical depth and the radiance of the plume. The band parameter input is composed of the coefficient of absorption part and the line density part. It is tabulated according to the wave number and temperature. Subroutine READER reads in the flow field data and subroutine PARAM calculates the transmissivity τ .

Table A.1 lists all input parameters, format of the read statements, and their meaning, in the same order as encountered in the program for the flow field. Table A.2 lists the input parameters required for the radiation calculations. Tables A.1 and A.2 indicate all input data needed for operation of the code. A listing of a sample input data case is included as Table A.3.

TABLE A.1

<u>Card #</u>	<u>Parameter</u>	<u>Format</u>	<u>Meaning</u>
1	Title	12A4	Title of the Flowfield
2	PP	E10.4	Total Pressure
	RJ	E10.4	Radius of the Nozzle
	NRP(1)	I4	# of Radial Points at the Beginning of the Plume
3	X(JJ)	E10.3	Axial Location
	ILAST	69X,I1	Control value. ILAST = 1 for the end of the flow field
4	NRP(JJ)	I5	# of Radial Points at Axial Location X(JJ)
5	R(JJ,KK)	E10.3	Radial Location at Point (JJ,KK)
	T(JJ,KK)	E10.3	Temperature at Point (JJ,KK)
	PCO ₂ (JJ,KK)	E10.3	Partial Pressure of CO ₂ at Point (JJ,KK)
	PH ₂ O(JJ,KK)	E10.3	Partial Pressure of H ₂ O at Point (JJ,KK)

Repeat card 5 for NRP(JJ) times.

Go back to card 3 for another axial point.

At the end of the flow field, set ILAST = 1. Such value of ILAST will terminate the reader subroutine and return to the MAIN ROUTINE.

TABLE A.2

<u>Card #</u>	<u>Parameter</u>	<u>Format</u>	<u>Meaning</u>
1	NWAV 1	I10	Wave number, lower limit
	NWAV 2	I10	Wave number, upper limit
2	WAV NO(I)	F9.1	Wave number
	XKK(I,J)	7E9.2	Coefficient of absorption K @ 300,600,1200,1500, 1800,2400,3000 °K of wave number WAVNO(I).
3	WAVNO(I)	F9.1	Wave number (must be the same as WAVNO(I) in card 2).
	DDINV(I,J)	7E9.2	Line density at temperature of 300,600,1200,1500,1800, 2400,3000 °K of wave number WAVNO(I).

Repeat card 2 and 3 for $N = (NWAV\ 2 - NWAV\ 1)/25 + 1$

4	ALPHA	F10.5	Aspect angle of the line of sight ISTOP = 1, will terminate the program after single alpha calculation. ISTOP = 2 will return to this point for another ALPHA calculation. ISTOP = 3, will read in a new set of flow field.
	ISTOP	I5	

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TABLE A.3

PR1.OX

.1000E 010.3930E 00 12

0.0
0.0 0.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19
0.3930E-010.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19
0.7860E-010.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19
0.1179E 000.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19
0.1572E 000.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19
0.1965E 000.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19
0.2358E 000.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19
0.2751E 000.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19
0.3144E 000.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19
0.3537E 000.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19
0.3930E 000.2400E 040.1000E-190.1000E-190.7546E-010.2563E 000.1000E-19
0.4606E 000.2880E 030.1000E-190.1000E-190.2885E-060.2885E-060.1000E-19
0.410E 00
0.0 0.2477E 040.1000E-190.1000E-190.7954E-010.2625E 000.1000E-19
0.3993E-010.2477E 040.1000E-190.1000E-190.7954E-010.2625E 000.1000E-19
0.7982E-010.2477E 040.1000E-190.1000E-190.7954E-010.2625E 000.1000E-19
0.1197E 000.2477E 040.1000E-190.1000E-190.7954E-010.2625E 000.1000E-19
0.1596E 000.2477E 040.1000E-190.1000E-190.7954E-010.2625E 000.1000E-19
0.1996E 000.2477E 040.1000E-190.1000E-190.7954E-010.2625E 000.1000E-19
0.2395E 000.2477E 040.1000E-190.1000E-190.7954E-010.2625E 000.1000E-19
0.2794E 000.2477E 040.1000E-190.1000E-190.7953E-010.2625E 000.1000E-19
0.3193E 000.2475E 040.1000E-190.1000E-190.7947E-010.2623E 000.1000E-19
0.3591E 000.2448E 040.1000E-190.1000E-190.7834E-010.2592E 000.1000E-19
0.3997E 000.2228E 040.1000E-190.1000E-190.6844E-010.2279E 000.1000E-19
0.4531E 000.1212E 040.1000E-190.1000E-190.2709E-010.8465E-010.1000E-19
0.5439E 000.5520E 030.1000E-190.1000E-190.8866E-020.2061E-010.1000E-19
0.6449E 000.3539E 030.1000E-190.1000E-190.1660E-020.4986E-020.1000E-19
0.7459E 000.3020E 030.1000E-190.1000E-190.3467E-030.1067E-020.1000E-19
0.8363E 000.2880E 030.1000E-190.1000E-190.2885E-060.2885E-060.1000E-19
0.814E 00
0.0 0.2487E 040.1000E-190.1000E-190.8209E-010.2609E 000.1000E-19
0.3997E-010.2487E 040.1000E-190.1000E-190.8209E-010.2609E 000.1000E-19
0.7994E-010.2487E 040.1000E-190.1000E-190.8209E-010.2609E 000.1000E-19
0.1199E 000.2487E 040.1000E-190.1000E-190.8209E-010.2609E 000.1000E-19
0.1599E 000.2487E 040.1000E-190.1000E-190.8209E-010.2609E 000.1000E-19
0.1999E 000.2487E 040.1000E-190.1000E-190.8209E-010.2609E 000.1000E-19
0.2398E 000.2487E 040.1000E-190.1000E-190.8209E-010.2609E 000.1000E-19
0.2798E 000.2486E 040.1000E-190.1000E-190.8207E-010.2607E 000.1000E-19
0.3197E 000.2473E 040.1000E-190.1000E-190.8185E-010.2594E 000.1000E-19
0.3598E 000.2412E 040.1000E-190.1000E-190.8035E-010.2512E 000.1000E-19
0.4016E 000.2173E 040.1000E-190.1000E-190.7212E-010.2137E 000.1000E-19
0.4559E 000.1479E 040.1000E-190.1000E-190.4261E-010.1220E 000.1000E-19
0.5357E 000.8166E 030.1000E-190.1000E-190.1581E-010.4558E-010.1000E-19
0.6308E 000.5012E 030.1000E-190.1000E-190.5747E-020.1687E-010.1000E-19
0.7318E 000.3739E 030.1000E-190.1000E-190.2240E-020.6671E-020.1000E-19
0.8269E 000.3231E 030.1000E-190.1000E-190.9029E-030.2708E-020.1000E-19
0.9188E 000.3017E 030.1000E-190.1000E-190.3501E-030.1053E-020.1000E-19
0.1005E 010.2925E 030.1000E-190.1000E-190.1148E-030.3454E-030.1000E-19
0.122E 01
0.0 0.2497E 040.1000E-190.1000E-190.8402E-010.2598E 000.1000E-19
0.4005E-010.2497E 040.1000E-190.1000E-190.8402E-010.2598E 000.1000E-19
0.8009E-010.2497E 040.1000E-190.1000E-190.8402E-010.2598E 000.1000E-19
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 0.0 0.2521E 040.1000E-190.1000E-190.1131E 000.1942E 000.1000E-19
 0.9511E-010.2490E 040.1000E-190.1000E-190.1127E 000.1890E 000.1000E-19
 0.1891E 000.2460E 040.1000E-190.1000E-190.1122E 000.1828E 000.1000E-19
 0.2861E 000.2350E 040.1000E-190.1000E-190.1090E 000.1660E 000.1000E-19
 0.3901E 000.2239E 040.1000E-190.1000E-190.1057E 000.1502E 000.1000E-19
 0.6300E 000.1816E 040.1000E-190.1000E-190.8450E-010.1061E 000.1000E-19
 0.9129E 000.1323E 040.1000E-190.1000E-190.5351E-010.6543E-010.1000E-19
 0.1233E 010.9345E 030.1000E-190.1000E-190.3035E-010.3857E-010.1000E-19
 0.1569E 010.6676E 030.1000E-190.1000E-190.1750E-010.2278E-010.1000E-19
 0.1906E 010.5409E 030.1000E-190.1000E-190.1062E-010.1414E-010.1000E-19
 0.2235E 010.4523E 030.1000E-190.1000E-190.6730E-020.9144E-020.1000E-19
 0.2554E 010.3954E 030.1000E-190.1000E-190.4387E-020.6063E-020.1000E-19
 0.2861E 010.3602E 030.1000E-190.1000E-190.2897E-020.4060E-020.1000E-19
 0.3161E 010.3357E 030.1000E-190.1000E-190.1904E-020.2699E-020.1000E-19
 0.3453E 010.3188E 030.1000E-190.1000E-190.1223E-020.1749E-020.1000E-19
 0.3741E 010.3069E 030.1000E-190.1000E-190.7486E-030.1078E-020.1000E-19
 0.4032E 010.2946E 030.1000E-190.1000E-190.4176E-030.6042E-030.1000E-19
 0.4311E 010.2926E 030.1000E-190.1000E-190.1823E-030.2645E-030.1000E-19
 0.105E 02

0.0 0.2504E 040.1000E-190.1000E-190.1139E 000.1880E 000.1000E-19
 0.9432E-010.2470E 040.1000E-190.1000E-190.1132E 000.1825E 000.1000E-19
 0.1914E 000.2441E 040.1000E-190.1000E-190.1126E 000.1770E 000.1000E-19
 0.2928E 000.2330E 040.1000E-190.1000E-190.1090E 000.1615E 000.1000E-19
 0.3946E 000.2220E 040.1000E-190.1000E-190.1053E 000.1460E 000.1000E-19
 0.6398E 000.1801E 040.1000E-190.1000E-190.8391E-010.1025E 000.1000E-19
 0.9208E 000.1341E 040.1000E-190.1000E-190.5492E-010.6639E-010.1000E-19
 0.1237E 010.9680E 030.1000E-190.1000E-190.3242E-010.4012E-010.1000E-19
 0.1570E 010.7205E 030.1000E-190.1000E-190.1925E-010.2440E-010.1000E-19
 0.1905E 010.5682E 030.1000E-190.1000E-190.1195E-010.1550E-010.1000E-19
 0.2235E 010.4741E 030.1000E-190.1000E-190.7724E-020.1023E-010.1000E-19
 0.2555E 010.4139E 030.1000E-190.1000E-190.5140E-020.6924E-020.1000E-19
 0.2867E 010.3737E 030.1000E-190.1000E-190.3473E-020.4749E-020.1000E-19
 0.3170E 010.3464E 030.1000E-190.1000E-190.2354E-020.3257E-020.1000E-19
 0.3467E 010.3272E 030.1000E-190.1000E-190.1574E-020.2200E-020.1000E-19
 0.3757E 010.3134E 030.1000E-190.1000E-190.1015E-020.1428E-020.1000E-19
 0.4044E 010.3030E 030.1000E-190.1000E-190.5980E-030.8466E-030.1000E-19
 0.4327E 010.2948E 030.1000E-190.1000E-190.2727E-030.3872E-030.1000E-19
 0.4606E 010.2901E 030.1000E-190.1000E-190.8526E-040.1212E-030.1000E-19
 0.110E 02
 0.0 0.2485E 040.1000E-190.1000E-190.1145E 000.1820E 000.1000E-19
 0.9628E-010.2453E 040.1000E-190.1000E-190.1136E 000.1767E 000.1000E-19
 0.1938E 000.2421E 040.1000E-190.1000E-190.1128E 000.1714E 000.1000E-19
 0.2959E 000.2310E 040.1000E-190.1000E-190.1090E 000.1570E 000.1000E-19
 0.3989E 000.2200E 040.1000E-190.1000E-190.1049E 000.1419E 000.1000E-19
 0.6351E 000.1809E 040.1000E-190.1000E-190.8426E-010.1043E 000.1000E-19
 0.9169E 000.1333E 040.1000E-190.1000E-190.5430E-010.6648E-010.1000E-19
 0.1235E 010.9522E 030.1000E-190.1000E-190.3143E-010.3941E-010.1000E-19
 0.1569E 010.7046E 030.1000E-190.1000E-190.1840E-010.2363E-010.1000E-19
 0.1906E 010.5549E 030.1000E-190.1000E-190.1129E-010.1484E-010.1000E-19
 0.2235E 010.4634E 030.1000E-190.1000E-190.7232E-020.9697E-020.1000E-19
 0.2555E 010.4053E 030.1000E-190.1000E-190.4766E-020.6502E-020.1000E-19
 0.2865E 010.3669E 030.1000E-190.1000E-190.3186E-020.4409E-020.1000E-19
 0.3188E 010.3411E 030.1000E-190.1000E-190.2130E-020.2983E-020.1000E-19
 0.3462E 010.3231E 030.1000E-190.1000E-190.1402E-020.1982E-020.1000E-19
 0.3752E 010.3103E 030.1000E-190.1000E-190.8881E-030.1264E-020.1000E-19
 0.4036E 010.3010E 030.1000E-190.1000E-190.5152E-030.7373E-030.1000E-19
 0.4319E 010.2939E 030.1000E-190.1000E-190.2323E-030.3335E-030.1000E-19
 0.0

NWAV1	2000	2500							
WAV	2000.0	0.0	0.0	0.0	0.13E-03	0.0	0.13E	00	0.66E 00
WAV	2000.0	0.25E 01	0.15E 02	0.18E 03	0.40E 03	0.87E 03	0.28E	04	0.70E 04
WAV	2025.0	0.0	0.0	0.0	0.78E-03	0.0	0.26E	00	0.11E 01
WAV	2025.0	0.25E 01	0.15E 02	0.18E 03	0.40E 03	0.87E 03	0.28E	04	0.63E 04
WAV	2050.0	0.0	0.0	0.0	0.37E-02	0.0	0.50E	00	0.17E 01
WAV	2050.0	0.25E 01	0.15E 02	0.18E 03	0.40E 03	0.87E 03	0.27E	04	0.51E 04
WAV	2075.0	0.0	0.0	0.0	0.15E-01	0.0	0.95E	00	0.26E 01
WAV	2075.0	0.25E 01	0.15E 02	0.18E 03	0.40E 03	0.87E 03	0.27E	04	0.44E 04
WAV	2100.0	0.0	0.0	0.0	0.52E-01	0.29E 00	0.17E	01	0.37E 01
WAV	2100.0	0.25E 01	0.15E 02	0.18E 03	0.40E 03	0.86E 03	0.25E	04	0.36E 04
WAV	2125.0	0.0	0.0	0.0	0.18E	00	0.69E	00	0.28E 01
WAV	2125.0	0.25E 01	0.15E 02	0.18E 03	0.40E 03	0.82E 03	0.21E	04	0.29E 04
WAV	2150.0	0.0	0.0	0.0	0.52E	00	0.15E	01	0.45E 01
WAV	2150.0	0.25E 01	0.15E 02	0.18E 03	0.40E 03	0.75E 03	0.17E	04	0.23E 04
WAV	2175.0	0.0	0.0	0.31E 00	0.14E	01	0.31E	01	0.67E 01
WAV	2175.0	0.25E 01	0.15E 02	0.18E 03	0.39E 03	0.65E 03	0.13E	04	0.17E 04
WAV	2200.0	0.0	0.75E-03	0.11E 01	0.32E	01	0.56E	01	0.93E 01
WAV	2200.0	0.25E 01	0.15E 02	0.16E 03	0.35E 03	0.55E 03	0.96E	03	0.13E 04
WAV	2225.0	0.0	0.18E-01	0.32E 01	0.66E	01	0.93E	01	0.12E 02
WAV	2225.0	0.25E 01	0.14E 02	0.14E 03	0.25E	03	0.40E	03	0.73E 03
WAV	2250.0	0.78E-04	0.28E 00	0.78E 01	0.12E	02	0.14E	02	0.12E 02
WAV	2250.0	0.25E 01	0.14E 02	0.11E 03	0.17E	03	0.28E	03	0.63E 03
WAV	2275.0	0.21E-01	0.26E 01	0.15E 02	0.17E	02	0.17E	02	0.15E 02
WAV	2275.0	0.25E 01	0.12E 02	0.75E 02	0.14E	03	0.20E	03	0.37E 03
WAV	2300.0	0.19E 01	0.16E 02	0.22E 02	0.20E	02	0.19E	02	0.16E 02
WAV	2300.0	0.23E 01	0.78E 01	0.85E 02	0.11E	03	0.22E	03	0.34E 03
WAV	2325.0	0.37E 02	0.36E 02	0.23E 02	0.20E	02	0.18E	02	0.13E 02
WAV	2325.0	0.14E 01	0.57E 01	0.42E 02	0.73E	02	0.79E	02	0.11E 03
WAV	2350.0	0.11E 02	0.21E 02	0.23E 02	0.20E	02	0.17E	02	0.11E 02
WAV	2350.0	0.28E 01	0.80E 01	0.24E 02	0.34E	02	0.37E	02	0.49E 02
WAV	2375.0	0.24E 02	0.34E 02	0.20E 02	0.14E	02	0.11E	02	0.60E 01
WAV	2375.0	0.11E 01	0.21E 01	0.55E 01	0.75E	01	0.70E	01	0.68E 01
WAV	2400.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
WAV	2400.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
WAV	2425.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
WAV	2425.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
WAV	2450.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
WAV	2450.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
WAV	2475.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
WAV	2475.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
WAV	2500.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
WAV	2500.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0

APPENDIX III

SAMPLE CASE OUTPUT

WAVE NUMBER	WAVE LENGTH	RADIATION	WAVE NUMBER	WAVE LENGTH	RADIATION
		W/SR/CM/MICRON			
2000.	5.0000	0.98137E-03	2025.	4.9383	0.20115E-02
2050.	4.8780	0.39446E-02	2075.	4.8193	0.76348E-02
2100.	4.7519	0.17107E-01	2125.	4.7059	0.29450E-01
2150.	4.6512	0.49696E-01	2175.	4.5977	0.77929E-01
2200.	4.5455	0.11446E 00	2225.	4.4944	0.15660E 00
2250.	4.4444	0.19604E 00	2275.	4.3956	0.22414E 00
2300.	4.3478	0.25091E 00	2325.	4.3011	0.22054E 00
2350.	4.2553	0.17916E 00	2375.	4.2105	0.78166E-01
2400.	4.1567	0.0	2425.	4.1237	0.0
2450.	4.0816	0.0	2475.	4.0404	0.0
2500.	4.0000	0.0			

TOTAL RADIATION EMITTED=0.8259E-01 WATTS/STER

WAVE NUMBER	WAVE LENGTH	RADIATION	WAVE NUMBER	WAVE LENGTH	RADIATION
2000.	5.0000	0.11605E-02	2025.	4.9383	0.23865E-02
2050.	4.8780	0.45994E-02	2075.	4.8193	0.91333E-02
2100.	4.7619	0.19295E-01	2125.	4.7059	0.33014E-01
2150.	4.6512	0.55326E-01	2175.	4.5977	0.85780E-01
2200.	4.5455	0.12457E 00	2225.	4.4944	0.16815E 00
2250.	4.4444	0.20694E 00	2275.	4.3956	0.23378E 00
2300.	4.3478	0.25036E 00	2325.	4.3011	0.22397E 00
2350.	4.2553	0.18512E 00	2375.	4.2105	0.82563E-01
2400.	4.1567	0.0	2425.	4.1237	0.0
2450.	4.0816	0.0	2475.	4.0404	0.0
2500.	4.0000	0.0			

TOTAL RADIATION FOR STATION 4 = 0.78207E-01 WATTS/SR/CM

PRI.OX
STATION NO.= 28 AXIAL DIST. (CM)= 11.035 ASPECT ANGLE (DEG)= 90.0

SUM UP STADSR AT X STATION 0.0

WAVE NUMBER	WAVE LENGTH	RADIATION	WAVE NUMBER	WAVE LENGTH	RADIATION
		W/SR/CM/MICRON			
2000.	5.0000	0.0	2025.	4.9383	0.0
2050.	4.8780	0.0	2075.	4.8193	0.0
2100.	4.7619	0.0	2125.	4.7059	0.0
2150.	4.6512	0.0	2175.	4.5977	0.0
2200.	4.5455	0.0	2225.	4.4944	0.0
2250.	4.4444	0.0	2275.	4.3956	0.0
2300.	4.3478	0.0	2325.	4.3011	0.0
2350.	4.2553	0.0	2375.	4.2105	0.0
2400.	4.1667	0.0	2425.	4.1237	0.0
2450.	4.0816	0.0	2475.	4.0404	0.0
2500.	4.0000	0.0			

TOTAL RADIATION EMITTED=0.1663E 01 WATTS/STER

WAVE NUMBER	WAVE LENGTH	RADIATION	WAVE NUMBER	WAVE LENGTH	RADIATION
2000.	5.0000	0.10455E-01	2025.	4.9383	0.21407E-01
2050.	4.8780	0.41942E-01	2075.	4.8193	0.51130E-01
2100.	4.7619	0.26070E 00	2125.	4.7059	0.47098E 00
2150.	4.6512	0.82891E 00	2175.	4.5977	0.13972E 01
2200.	4.5455	0.21785E 01	2225.	4.4944	0.31900E 01
2250.	4.4444	0.42830E 01	2275.	4.3956	0.50712E 01
2300.	4.3478	0.56955E 01	2325.	4.3011	0.51826E 01
2350.	4.2553	0.39996E 01	2375.	4.2105	0.16854E 01
2400.	4.1667	0.0	2425.	4.1237	0.0
2450.	4.0816	0.0	2475.	4.0404	0.0
2500.	4.0000	0.0			

TOTAL RADIATION FOR STATION 28 = 0.0 WATTS/SR/CM

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LAPP GAS DYNAMICS - ARMY RADIATION CODE

PRI.DX	W/SR/cm	STATION RADIATION (W/SR)
0.0	0.5837254E-01	~~0.4022077E-01~~
0.41	0.6468284E-01	0.2522010E-01
0.81	0.7061887E-01	0.2730462E-01
1.22	0.7820714E-01	0.3006312E-01
1.61	0.8649528E-01	0.3262293E-01
1.97	0.9388489E-01	0.3232552E-01
2.37	0.1016850E 00	0.3915963E-01
2.78	0.1099497E 00	0.4252582E-01
3.18	0.1185645E 00	0.4656448E-01
3.58	0.1266638E 00	0.4842820E-01
3.98	0.1350878E 00	0.5225714E-01
4.39	0.1435786E 00	0.5804345E-01
4.82	0.1524286E 00	0.6281853E-01
5.22	0.1602380E 00	0.6266773E-01
5.63	0.1665467E 00	0.6678170E-01
6.03	0.1720112E 00	0.6852239E-01
6.44	0.1782773E 00	0.7089573E-01
6.85	0.1849118E 00	0.7493514E-01
6.94	0.1851284E 00	0.1672365E-01
7.41	0.1922510E 00	0.8824474E-01
7.85	0.1983629E 00	0.8673418E-01
8.30	0.2038116E 00	0.9088081E-01
8.76	0.2088537E 00	0.9325266E-01
9.21	0.2134077E 00	0.9542012E-01
9.66	0.2167625E 00	0.9805322E-01
10.12	0.2199754E 00	0.9955025E-01
10.58	0.2249740E 00	0.1014207E 00
11.04	0.0	0.5172247E-01

* RADIATION VALUE WAS CALCULATED BASED ON THE SLOPE BETWEEN THE PREVIOUS TWO POINTS

TOTAL RADIATION EMITTED=0.1663E 01 WATTS/STER

CENTROID= 0.6655E 01CM

NWAV1	2000	2500							
WAV	2000.0	0.78E-02	0.40E-01	0.77E-01	0.11E 00	0.13E 00	0.14E 00	0.15E 00	
WAV	2000.0	0.95E-01	0.17E 00	0.45E 00	0.14E 01	0.39E 01	0.98E 01	0.24E 02	
WAV	2025.0	0.54E-02	0.35E-01	0.65E-01	0.87E-01	0.11E 00	0.12E 00	0.13E 00	
WAV	2025.0	0.12E 00	0.21E 00	0.50E 00	0.15E 01	0.36E 01	0.91E 01	0.23E 02	
WAV	2050.0	0.38E-02	0.25E-01	0.51E-01	0.71E-01	0.89E-01	0.10E 00	0.11E 00	
WAV	2050.0	0.11E 00	0.19E 00	0.48E 00	0.15E 01	0.40E 01	0.10E 02	0.26E 02	
WAV	2075.0	0.26E-02	0.18E-01	0.51E-01	0.55E-01	0.72E-01	0.83E-01	0.96E-01	
WAV	2075.0	0.12E 00	0.22E 00	0.55E 00	0.18E 01	0.45E 01	0.12E 02	0.28E 02	
WAV	2100.0	0.18E-02	0.12E-01	0.29E-01	0.44E-01	0.61E-01	0.69E-01	0.84E-01	
WAV	2100.0	0.14E 00	0.25E 00	0.57E 00	0.17E 00	0.41E 01	0.12E 02	0.35E 02	
WAV	2125.0	0.13E-02	0.85E-02	0.21E-01	0.38E-01	0.58E-01	0.64E-01	0.73E-01	
WAV	2125.0	0.20E 00	0.39E 00	0.68E 00	0.17E 01	0.36E 01	0.11E 02	0.36E 02	
WAV	2150.0	0.88E-03	0.68E-02	0.15E-01	0.28E-01	0.45E-01	0.52E-01	0.63E-01	
WAV	2150.0	0.22E 00	0.42E 00	0.13E 01	0.27E 01	0.50E 01	0.17E 02	0.65E 02	
WAV	2175.0	0.62E-02	0.40E-02	0.11E-01	0.21E-01	0.37E-01	0.45E-01	0.53E-01	
WAV	2175.0	0.23E 00	0.43E 00	0.10E 01	0.21E 01	0.40E 01	0.14E 02	0.89E 02	
WAV	2200.0	0.48E-03	0.30E-02	0.93E-02	0.19E-01	0.33E-01	0.40E-01	0.46E-01	
WAV	2200.0	0.23E 00	0.45E 00	0.11E 01	0.25E 01	0.49E 01	0.19E 02	0.12E 03	
WAV	2225.0	0.41E-03	0.18E-02	0.70E-02	0.15E-01	0.30E-01	0.37E-01	0.40E-01	
WAV	2225.0	0.24E 00	0.45E 00	0.11E 01	0.27E 01	0.53E 01	0.21E 02	0.13E 03	
WAV	2250.0	0.32E-03	0.12E-02	0.45E-02	0.10E-01	0.25E-01	0.33E-01	0.35E-01	
WAV	2250.0	0.25E 00	0.48E 00	0.12E 01	0.29E 01	0.58E 01	0.23E 02	0.13E 03	
WAV	2275.0	0.23E-03	0.72E-03	0.36E-02	0.93E-02	0.23E-01	0.31E-01	0.31E-01	
WAV	2275.0	0.25E 00	0.50E 00	0.12E 01	0.31E 01	0.63E 01	0.25E 02	0.14E 03	
WAV	2300.0	0.20E-03	0.54E-03	0.32E-02	0.75E-02	0.20E-01	0.28E-01	0.29E-01	
WAV	2300.0	0.27E 00	0.52E 00	0.12E 01	0.34E 01	0.69E 01	0.28E 02	0.15E 03	
WAV	2325.0	0.15E-03	0.38E-03	0.19E-02	0.60E-02	0.18E-01	0.27E-01	0.28E-01	
WAV	2325.0	0.27E 00	0.53E 00	0.13E 01	0.37E 01	0.75E 01	0.31E 02	0.16E 03	
WAV	2350.0	0.10E-03	0.26E-03	0.12E-02	0.48E-02	0.16E-01	0.25E-01	0.28E-01	
WAV	2350.0	0.28E 00	0.55E 00	0.13E 01	0.40E 01	0.82E 01	0.34E 02	0.17E 03	
WAV	2375.0	0.85E-04	0.19E-03	0.91E-03	0.36E-02	0.13E-01	0.24E-01	0.26E-01	
WAV	2375.0	0.29E 00	0.56E 00	0.14E 01	0.43E 01	0.90E 01	0.37E 02	0.18E 03	
WAV	2400.0	0.76E-04	0.14E-03	0.71E-03	0.32E-02	0.12E-01	0.24E-01	0.25E-01	
WAV	2400.0	0.30E 00	0.58E 00	0.14E 01	0.46E 01	0.98E 01	0.41E 02	0.19E 03	
WAV	2425.0	0.62E-04	0.13E-03	0.61E-03	0.23E-02	0.10E-01	0.22E-01	0.25E-01	
WAV	2425.0	0.31E 00	0.60E 00	0.15E 01	0.50E 01	0.11E 02	0.45E 02	0.21E 03	
WAV	2450.0	0.48E-04	0.11E-03	0.52E-03	0.20E-02	0.92E-02	0.20E-01	0.25E-01	
WAV	2450.0	0.32E 00	0.63E 00	0.15E 01	0.54E 01	0.12E 02	0.50E 02	0.22E 03	
WAV	2475.0	0.37E-04	0.11E-03	0.44E-03	0.17E-02	0.79E-02	0.18E-01	0.25E-01	
WAV	2475.0	0.33E 00	0.65E 00	0.16E 01	0.58E 01	0.13E 02	0.55E 02	0.24E 03	
WAV	2500.0	0.36E-04	0.10E-03	0.38E-03	0.17E-02	0.67E-02	0.17E-01	0.25E-01	
WAV	2500.0	0.34E 00	0.67E 00	0.16E 01	0.64E 01	0.14E 02	0.61E 02	0.25E 03	