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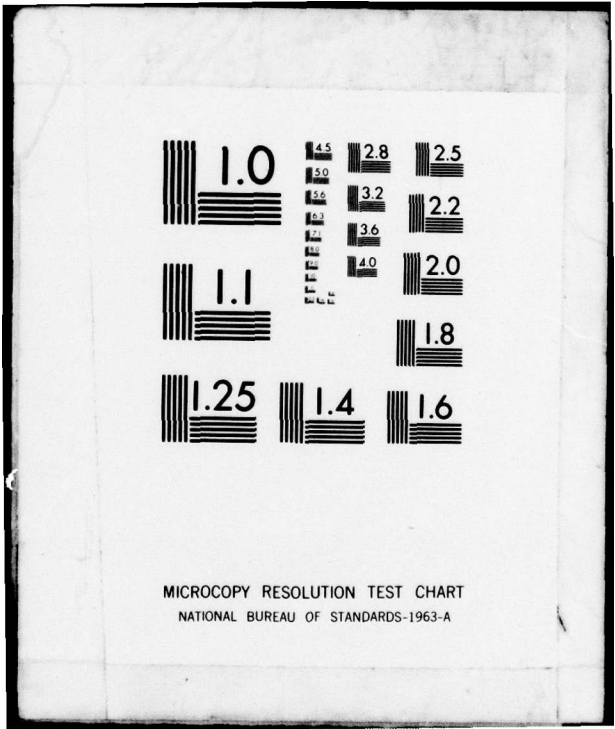
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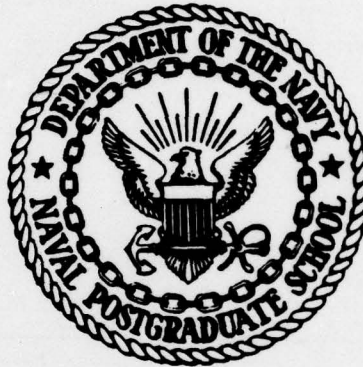
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A METHOD OF CHARACTERISTICS APPROACH TO THE PROBLEM  
OF SUPERSONIC FLOW PAST OSCILLATING CASCADES  
WITH FINITE BLADE THICKNESS

K. Vogeler

October 1978

Approved for public release; distribution unlimited

Prepared for:  
Chief of Naval Research  
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EDG REPORT DOCUMENTATION PAGE		READ INSTRUCTIONS BEFORE COMPLETING FORM
1. REPORT NUMBER <b>14</b> NPS67-78-112	2. GOVT ACCESSION NO.	3. RECIPIENT'S CATALOG NUMBER
4. TITLE (and Subtitle) <b>6</b> A METHOD OF CHARACTERISTICS APPROACH TO THE PROBLEM OF SUPERSONIC FLOW PAST OSCILLATING CASCADES WITH FINITE BLADE THICKNESS		5. TYPE OF REPORT & PERIOD COVERED
7. AUTHOR(s) <b>10</b> → R. Vogeler K. R. R. D.		8. CONTRACT OR GRANT NUMBER(s) <b>12</b> 463
9. PERFORMING ORGANIZATION NAME AND ADDRESS Naval Postgraduate School Monterey, California <sup>251 410</sup>		10. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS 61152M;RR 000 01 01 000-1478-WR-80023
11. CONTROLLING OFFICE NAME AND ADDRESS		12. REPORT DATE
14. MONITORING AGENCY NAME & ADDRESS (if different from Controlling Office) <b>16</b> R. R. 00001		13. NUMBER OF PAGES
16. DISTRIBUTION STATEMENT (of this Report) Approved for public release; distribution unlimited <b>17</b> R. R. 0000101		15. SECURITY CLASS. (of this report) UNCLASSIFIED
		18a. DECLASSIFICATION/DOWNGRADING SCHEDULE
17. DISTRIBUTION STATEMENT (of the abstract entered in Block 20, if different from Report)		<b>11</b> OCT 78
18. SUPPLEMENTARY NOTES		
19. KEY WORDS (Continue on reverse side if necessary and identify by block number) Unsteady Aerodynamics Cascade Aerodynamics Supersonic Blade Flutter Aeroelasticity of Turbomachines Method of Characteristics		
20. ABSTRACT (Continue on reverse side if necessary and identify by block number) A method of characteristics approach is described for the computation of supersonic flow past oscillating cascades having blades of finite thickness distribution. The transonic small disturbance equation is solved for cascades with subsonic leading-edge locus. Sample calculations are presented for the inlet and passage flow and comparisons are given with calculations for flat-plate cascades and available measured pressure distributions.		

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#### ACKNOWLEDGEMENT

This study was performed under the sponsorship of the Office of Naval Research and the Naval Postgraduate School Foundation Program. I want to thank both institutions for their financial support and hence for the possibility to work for one year at the Naval Postgraduate School in Monterey, CA.

The project would have been far less successful without the help of Prof. M. F. Platzler, Chairman of the Department of Aeronautics in the NPS. The contents of this report is a reflection of the many discussions with him over his recent work and his ideas for the future.

Finally, I want to thank Prof. H. E. Gallus from the Technical University of Aachen, (W. Germany), who made this opportunity available to me.

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## 1. INTRODUCTION

The demand for increasing performance and efficiency of turbines and compressors in jet engines forces the industry to advanced designs. This means that it is no longer acceptable to use blades in turbomachines which are loaded considerably below their mechanical limits. There are two major ways to improve the engine thrust-to-weight and thrust-to-volume characteristics:

- Reduced weight and size
- Increased massflow, temperature and pressure.

It is obvious that only a compromise can be successful, as advances in the aerothermodynamics oppose those of the structure.

These trends require the use of slender and thin blades which are increasingly susceptible to flutter and vibration problems. The most important ones are:

1. Supersonic unstalled flutter
2. Forced response
3. Subsonic stall flutter
4. Choke flutter
5. Supersonic stall flutter

While the latter three are quite difficult to describe in a fluid mechanical model, the unstalled supersonic blade flutter is amenable to analysis with reasonable effort. What makes it even more interesting and at the same time highly important is the possibility of its occurrence at the design condition of the engine. Especially, modern fans with large diameters operate with the outer part of their blades in the transonic flow region ( $1. < M < 1.5$ ). Their flutter susceptibility therefore makes the analysis of supersonic unstalled flutter increasingly important. The problem has been attacked not only in the United States but in all major industrialized countries, which shows that it is

a general one to be considered in the design of modern high-performance turbo-machines.

During the past decade, several methods have been developed to predict supersonic blade flutter. All were based on the idealization of the actual flow by the planar flow through a staggered cascade of oscillating blades. However, the two-dimensional flow and the flat plate assumptions impose severe simplifications whose range of validity needs to be better understood.

Since the incorporation of three-dimensional flow effects is rather difficult, it seems logical to first explore the effect of blade thickness and shape on supersonic blade flutter while retaining the cascade concept. To this end, the nonlinear transonic small perturbation equation is adopted in this report as the governing equation. The use of this equation rather than the full potential equation or the Euler equations is suggested by Teipel's success to analyze the thickness effect of a single oscillating airfoil in low supersonic flow. Hence, the present work is an extension of Teipel's method to oscillating supersonic cascades for the purpose of determining the influence of steady non-uniform flow effects due to blade thickness, shape, camber or angle of attack on the oscillatory pressure distributions, forces and moments. The ultimate goal of this study is to replace the transonic small disturbance equation by the Euler equations so that the range of applicability of this simpler equation can be ascertained.

## 2. THE BASIC EQUATIONS

### 2.1 The Nonlinear Transonic Equation

Landahl presents in (1) the differential equation which is valid for the transonic flow region

$$\left[ M^2 - 1 + \frac{U(\kappa+1)}{a^2} \frac{\partial \phi}{\partial X} \right] \frac{\partial^2 \phi}{\partial X^2} - \frac{\partial^2 \phi}{\partial Y^2} + \frac{1}{a^2} + \frac{\partial^2 \phi}{\partial T^2} + \frac{2M}{a} \frac{\partial^2 \phi}{\partial X \partial T} = 0 \quad (1)$$

with

- U - Free stream velocity
- a - Local velocity of sound
- M - Free stream Mach number
- $\phi$  - Velocity potential (X,Y,T)
- $\kappa$  - Ratio of specific heats

Eq. (1) can be written non-dimensionally by using the terms

$$x = \frac{X}{c}, \quad y = \frac{Y}{c}, \quad t = \frac{TU}{c}, \quad \phi = \frac{\phi}{cU},$$

where c is the chord.

Thus we get the new form

$$\left[ M^2 - 1 + M^2 (\kappa+1) \frac{\partial \phi}{\partial x} \right] \frac{\partial^2 \phi}{\partial x^2} - M^2 \frac{\partial^2 \phi}{\partial t^2} + 2M^2 \frac{\partial^2 \phi}{\partial x \partial t} = 0 \quad (2)$$

Following Teipel, who developed in (2) for single airfoils a method of characteristics using Eq. (1), in this work Eq. (2) is to be used. This has already successfully been done by Platzler, Chadwick and Strada (3,4,5,6,7) for a single oscillating airfoil and for oscillating cascades of wedges and thick blades with flat upper surfaces. Nevertheless, the basic steps, which lead to the solution of this problem shall be repeated in this report, to make it at the same time a summary of the work already done by the authors above.

It is an extension in so far as it introduces the thickness effect of the upper blade surface in an oscillating, staggered cascade and shows a method-of-characteristics-approach to the unsteady supersonic wake of not only a flat plate but also airfoils with small but finite thickness.

As we consider only small perturbations of the freestream flow values, the potential function  $\phi$  of Eq. (2) can be split into a steady and an unsteady one. Furthermore, we assume only harmonic oscillations so that we can write

$$\phi(x,y,t) = \varphi(x,y) + \Psi(x,y) \cdot e^{ikt} \quad (3)$$

where  $k = \frac{\omega \cdot c}{U}$  is the reduced frequency.

Introducing (3) in (2), we can separate the unsteady from the steady problem and we obtain a set of two differential equations

$$[M^2 - 1 + (\kappa+1) M^2 \varphi_x] \varphi_{xx} - \varphi_{yy} = 0 \quad (4)$$

and

$$[M^2 - 1 + (\kappa+1) M^2 \varphi_x] \Psi_{xx} - \Psi_{yy} + [M^2(\kappa+1)\varphi_{xx} + 2ik M^2] \Psi_x - M^2 k^2 \Psi = 0 \quad (5)$$

The boundary conditions for the flow over an oscillating airfoil can also be written as the sum of steady and unsteady influences:

$$h(x,t) = h_0(x) + h_1(x) \cdot e^{ikt} \quad (6)$$

$(h(x,t); x)$  is the true location of a surface point. Thus the boundary conditions for the steady and unsteady problem can be expressed

$$\varphi_y = \frac{\partial h_0}{\partial x} \equiv \text{slope of the surface} \quad (7a)$$

$$\psi_y = \frac{\partial h_1}{\partial x} + \frac{\partial h_1}{\partial t} \equiv \text{the unsteady movement of a flat plate} \quad (7b)$$

Eq. (4) together with (7a) describes the transonic flow field over a fixed airfoil.

Following Sauer (8), we can attack the problem with the method of characteristics. The left- and right- running characteristics shall be indicated by  $\alpha$  and  $\beta$ . We find for their slopes

$$\left(\frac{\partial y}{\partial x}\right)_{\alpha, \beta} = \pm \frac{1}{\sqrt{M^2 - 1 + (\kappa + 1) M^2 \varphi_x}} \quad (8)$$

or with

$$\lambda = M^2 - 1 + (\kappa + 1) M^2 \varphi_x \quad (9)$$

$$\left(\frac{\partial y}{\partial x}\right)_{\alpha, \beta} = \pm \frac{1}{\sqrt{\lambda}} \quad (10)$$

The upper sign indicates the  $\alpha$  - direction. Introducing a second substitution

$$\mu = \frac{3}{2} (\kappa + 1) M^2 \varphi_y \quad (11)$$

The compatibility relation  $\pm \sqrt{\lambda} \varphi_{xx} + \varphi_{yx} = 0$

can now be written as  $\left(\frac{\partial \lambda^{3/2}}{\partial x}\right)_{\alpha, \beta} \mp \left(\frac{\partial \mu}{\partial x}\right)_{\alpha, \beta} = 0 \quad (12)$

That Eq. (12) holds, can easily be verified by resubstituting (9) and (11) in (12) and executing the differentiation. The result will be Eq. (4).

In (12) we find only derivatives in the x-direction along the characteristics. Therefore we can integrate (12) easily and obtain

$$\lambda^{3/2} \mp \mu = \text{const} = C_{\alpha, \beta} \quad (13)$$

We changed our variables from  $\varphi_x$  and  $\varphi_y$  to  $\lambda$  and  $\mu$ . Consequently, we have to convert our boundary conditions for the steady problem:

$y = 0$ :

$$\mu = \frac{3}{2} (\kappa + 1) M^2 \frac{\partial h_0}{\partial x} \quad (14)$$

Eq. (13) says: as long as we move along one characteristic,

$$\left( \lambda^{3/2} \mp \mu \right)_{\alpha, \beta}$$

will not change. This makes (13) a tool to evaluate the original desired unknowns  $\varphi_x$  and  $\varphi_y$  in the field.

In the free-stream field  $\varphi_x$  and  $\varphi_y$  are zero. Hence, there we have

$$\lambda_{\infty} = M^2 - 1$$

and

$$\mu_{\infty} = 0$$

Therefore all the characteristics have here the slope

$$\left( \frac{\partial y}{\partial x} \right)_{\alpha, \beta} = \pm \frac{1}{\sqrt{M^2 - 1}}$$

and from (12) we obtain

$$C_{\infty} = \lambda_{\infty}^{3/2} = (M^2 - 1)^{3/2} \quad (15)$$

Fig. 1 shows a  $\beta$ -characteristic of the free stream hitting the surface of the airfoil.

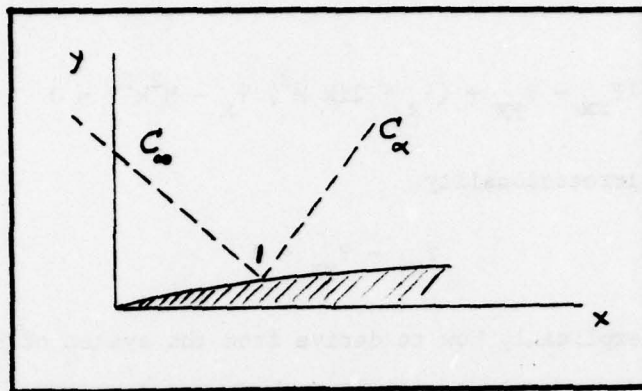


Fig. 1. Characteristics on the Airfoil Surface

In point 1 there has to be  $C_\infty = C_\alpha$  with Eq. (9), (13) and (14) we now can find

$$\lambda_1 = (M^2 - 1)^{3/2} \pm \frac{3}{2} (\kappa + 1) M^2 \frac{\partial h_0}{\partial x}^{2/3} \quad (16)$$

In the two-dimensional case there can be shown: for the characteristic leading away from the boundary there is not only

$$C_\alpha = \text{const}$$

but also

$$\mu_\alpha = \text{const} \quad \text{and}$$

$$\lambda_\alpha = \text{const}$$

With Eq. (5) through (16) the steady problem (4) can be solved. We get a net of characteristics. On the gridpoints the values of  $\varphi_x$  and  $\varphi_y$  are known. The computational procedure is shown later.

With the knowledge of the characteristic net and the values for  $\varphi_x$ ,  $\varphi$ ,  $\lambda$  and  $\mu$  in each grid point we can now solve the unsteady part of the problem.

We put Eq. (9) and (11) into (5). The result is a new differential equation for the unsteady flow.

$$\lambda \Psi_{xx} - \Psi_{yy} + (\lambda_x + 2ik M^2) \Psi_x - M^2 k^2 \Psi = 0 \quad (17)$$

Further we assume irrotationality

$$\Psi_{xy} - \Psi_{yx} = 0 \quad (18)$$

Teipel shows explicitly how to derive from the system of Eq. (17) and (18) the compatibility relations for the unsteady characteristics.

$$\Psi_{xx} \mp \frac{1}{\sqrt{\lambda}} \Psi_{yx} + \frac{1}{\lambda} (\lambda_x + 2ik M^2) \Psi_x - \frac{1}{\lambda} k^2 M^2 \Psi = 0 \quad (19)$$

The geometry of the characteristic net is determined by the coefficients connected with the highest order terms. Thus we can see from Eq. (4) and (5) that the net remains the same for the unsteady flow problem.

To solve Eq. (17) via Eq. (19) by moving along the already known characteristics of the steady field, we need the unsteady boundary values along the airfoil and along the shock.

The first one is given with Eq. (7b).

$$\Psi_y = \frac{\partial h_1}{\partial x} + ikh_1$$

It reads for pitch - movement

$$\bar{\Psi}_y = - [1 + ik(x - b)]$$

and for plunge - movement

$$\Psi_y = - ik$$

where  $b$  is the normalized location of the pitching axis. These expressions are derived in section 2.2. The boundary conditions along the oscillating shock are much more difficult to obtain. This is done in section 2.3.

After obtaining the solution for  $u_1$ ,  $v_1$  and  $\psi$  the unsteady pressure coefficients can be computed from

$$c_{pl} = 2 \frac{p_1}{\rho_\infty U} \quad (20)$$

$$c_{pl} = -2(u_1 + ik\psi)$$

## 2.2 Boundary Conditions Along the Airfoil

The general expression of the location of an oscillating airfoil is given by Eq. (6)

$$h(x,t) = h_0(x) + h_1(x) \cdot e^{ikt}$$

$h_0(x)$  represents the surface and we assume for now that it is described by an analytical function for which the second derivative exists. With Eq. (7a) we can calculate  $\varphi_y$  on the airfoil:

$$\varphi_y = \frac{\partial h_0}{\partial x}$$

However, as we only consider slender bodies, we project the point down to the  $x$ -axis  $y = 0$ , so that the boundary coordinates of the characteristic net are always  $(x,0)$  instead of  $(x,y)$ . This has two reasons

1.  $y = 0$  makes the steady boundary step much less complicated, without introducing a considerable mistake;
2. The oscillating movement of the airfoil can be reduced to the movement of a flat plate.

We examine two moving modes. The pitch- and plunge- mode.

PITCH:

Fig. 2 shows the deflected airfoil (flat plate) in a system of coordinates

$$h_1 = (b - x) \cdot \text{tg}\alpha$$

As  $\alpha$  is small and a harmonic motion, we can say

$$\text{tg}\alpha = \alpha_0 \cdot e^{ikt}$$

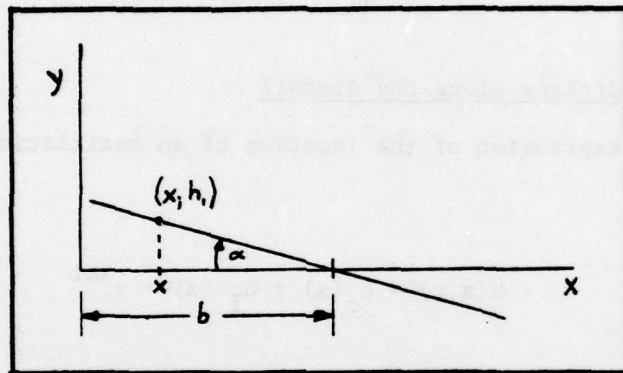


Fig. 2. Pitching Flat Plate

Thus we get

$$h_1 = (b - x) \alpha_0 \cdot e^{ikt}$$

The unsteady boundary condition is with Eq. (7b)

$$\psi_y = [-\alpha_0 + \alpha_0 ik(b - x)] e^{ikt}$$

or

$$\psi_y = -\alpha_0 [1 + ik(x - b)]$$

where the  $\exp(ikt)$  - term is omitted. This can be normalized by  $|\alpha_0|$ . So the final expression is the well known form

$$\Psi_y = - [1 + ik(x - b)] \quad (21)$$

PLUNGE:

For the plunge mode the deflection  $h$  is no function of  $x$ . Again a harmonic motion is assumed.

$$h = - h_0 e^{ikt}$$

Eq. (7b) gives us then the boundary value for

$$\Psi_y = - h_0 \cdot ik e^{ikt}$$

To be consistent with the previous work (3 to 7), the downward deflection is defined as positive. Again the expression is normalized by  $|h_0|$  and the

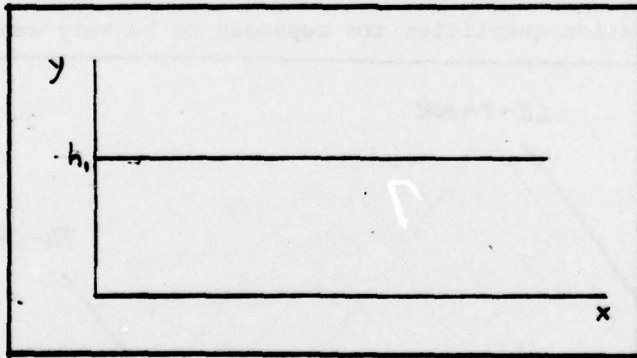


Fig. 3. Plunge Movement

$\exp(ikt)$  - term is omitted

$$\Psi_y = - ik \quad (22)$$

This approach to the boundary conditions is a rather physical one. A more rigorous derivation is shown by Bell in (9), including the derivation of Eq. (7b).

### 2.3 The Oscillating Shock Wave in an Oscillating Flow Field

The basic problem in using the method of characteristics is finding and introducing the proper boundary conditions. Section 2.2 gives us the influence of the moving airfoil into the field. The second boundary in the field between surface and shock are the unsteady flow properties immediately downstream of the shock. See Fig. 4.

It is

$w$  - Velocity

$u, v$  -  $x, y$  components of  $w$

$w_n, w_t$  - Normal and tangential components of

$\wedge$  - Indicates properties behind the shock

$\gamma_0$  - Slope of the shock in the steady problem

$\gamma'$  - Deflection of the shock due to oscillation of the airfoil

All the perturbation quantities are supposed to be very small.

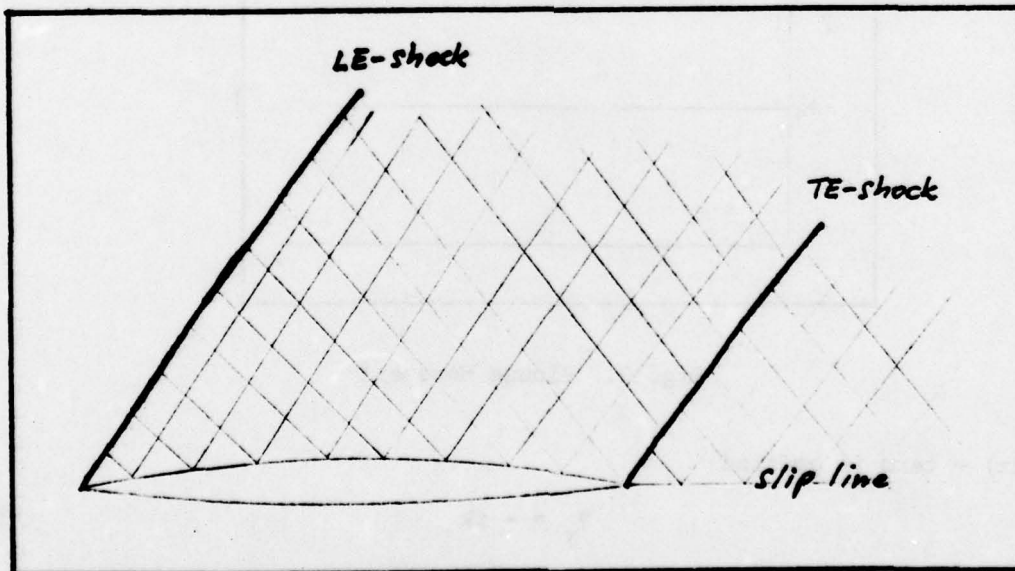


Fig. 4. Field of Characteristics Past an Airfoil

Teipel (10) found a way to obtain those properties for an airfoil in undisturbed supersonic flow. As the final object is to compute the flow in a staggered cascade, his work had to be extended. This was first done by Chadwick (4) who applied it to a cascade of wedges. Strada used Chadwick's equations in (5) to compute the inlet flow of a cascade with airfoils which have flat upper and curved lower surfaces.

Fig. 5 shows the velocities upstream and downstream of an arbitrary shock:

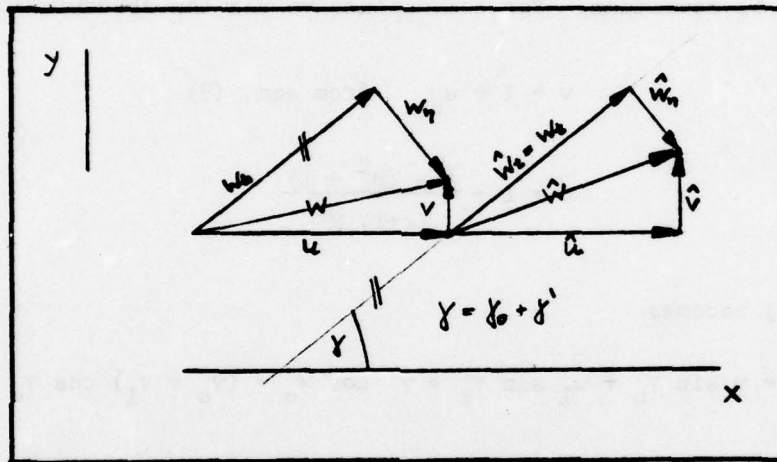


Fig. 5. Velocities at the Shock

The general case in the cascade is  $w \neq u$ . Thus we write

$$u = 1 + u_0 + u_1$$

$$v = v_0 + v_1$$

$$\hat{u} = 1 + \hat{u}_0 + \hat{u}_1$$

$$\hat{v} = \hat{v}_0 + \hat{v}_1$$

(23)

The only unknowns are  $\hat{u}_1$  and  $\hat{v}_1$ , as the steady problem is already expected to be solved and the field in front of the shock is presumed to be well known. Then we can take from the geometry (Fig. 5).

$$\begin{aligned} w_n &= u \cdot \sin(\gamma_0 + \gamma') - v \cos(\gamma_0 + \gamma') \\ w_t &= u \cos(\gamma_0 + \gamma') + v \sin(\gamma_0 + \gamma') \end{aligned} \quad (24)$$

If we apply Eq. (23) to (24) and neglect higher order terms like  $u_0 \gamma^1$ , Eq. (24) can be rewritten. For convenience we use the abbreviation

$$\begin{aligned} v &= 1 + u_0 \quad \text{from eqn. (9)} \\ v &= 1 + \frac{\lambda - (M^2 + 1)}{(\kappa + 1) M^2} \end{aligned} \quad (25)$$

Eq. (24) becomes

$$\begin{aligned} w_n &= v \sin \gamma_0 + u_1 \sin \gamma_0 + \gamma' \cos \gamma_0 - (v_0 + v_1) \cos \gamma_0 \\ w_t &= v \cos \gamma_0 + u_1 \cos \gamma_0 - \gamma' \sin \gamma_0 + (v_0 + v_1) \sin \gamma_0 \end{aligned} \quad (26)$$

Now we take a look at the equation for normal moving shocks (11).

$$\hat{w}_n - w_n = \frac{2}{\kappa + 1} \cdot W \left(1 - \frac{a^2}{W^2}\right), \quad (27)$$

where  $W$  is the relative velocity of the shock with respect to the fluid. To transform from the system moving with the fluid into an airfoil - fixed system of coordinates, we find the velocity of a point oscillating with the shock as

$$W^* = W + w_n \quad (28)$$

See Fig. 6.

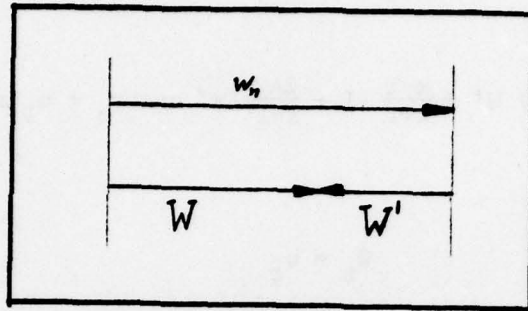


Fig. 6. Definition of Relative Velocities in a Shock-Point

Connecting Eq. (26) with Eq. (28) we obtain

$$W = W' - v \sin \gamma_0 - u_1 \sin \gamma_0 - \gamma' \cos \gamma_0 + (v_0 + v_1) \cos \gamma_0 \quad (29)$$

Forming  $W^2$  and again neglecting all terms of higher order, recasting Eq. (29) leads to

$$v^2 \sin^2 \gamma_0 = W(W + 2v \sin \gamma_0) \quad (30)$$

Reintroducing this into the shockpolar Eq. (27), we get

$$w_n - w_n' = \frac{2}{\kappa+1} W - \frac{2}{\kappa+1} \cdot \frac{a^2}{v^2 \sin^2 \gamma_0} (W + 2v \sin \gamma_0) \quad (31)$$

By making use of Eq. (26), (29), (31) and the substitution

$$A = \frac{a_\infty^2}{v^2 \sin^2 \gamma_0} = \frac{1}{v^2 M^2 \sin^2 \gamma_0} = \frac{1}{v^2 M_n^2} \quad (32)$$

We obtain, again neglecting all higher order terms

$$\hat{w}_n = \frac{\kappa-1}{\kappa+1} v \sin \gamma_0 \left(1 + \frac{2A}{\kappa-1}\right) - \frac{\kappa-1}{\kappa+1} v_0 \cos \gamma_0 \left(1 - \frac{2A}{\kappa-1}\right) + \quad (33)$$

$$+ \frac{2}{\kappa+1} (1+A) W^* + \frac{\kappa-1}{\kappa+1} \left(1 - \frac{2A}{\kappa-1}\right) (\gamma^* \cos \gamma_0 + u_1 \sin \gamma_0 - v_1 \cos \gamma_0)$$

together with

$$\hat{w}_t = w_t \quad (34)$$

Out of Eq. (26) we have now expressions for the velocities behind the moving shock in terms of the known values and  $\gamma^*$  and  $W^*$ . In the next step we look at a point of the oscillating shock, Fig. 7.

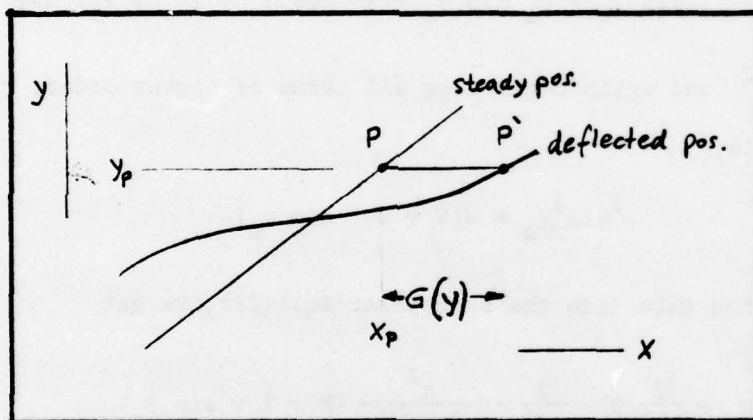


Fig. 7. Geometry of an Oscillating Shock

If we assume  $G$  to be a harmonic oscillation around the steady middle-position  $P$ , we can say

$$x = x_p + G(y) \cdot e^{ikt}, \quad (35)$$

where  $G(y)$  is the amplitude of the shock vibration and  $k$  the same reduced frequency as given for the airfoil.

The normal motion of P is

$$W' = \left( \frac{\partial x}{\partial t} \right)_y \cdot \sin (\gamma_0 + \gamma)$$

$$W' = ik \sin \gamma_0 G(y)$$
(36)

again the exponential term is omitted.

As we know  $\gamma_0(y)$  from the solution of the steady problem, we can find the x-coordinates of P :

$$x_p = \int_0^y \text{ctg} (\gamma_0 (y)) dy$$

This completes Eq. (35) to

$$x = \int_0^y \text{ctg} (\gamma_0 (y)) \cdot dy + G(y) \cdot e^{ikt}$$
(37)

The total differential gives us

$$dx = \left( \frac{\partial x}{\partial t} \right)_y \cdot dt + \left( \frac{\partial x}{\partial y} \right)_t \cdot dy$$

$$dx = ik G(y) \cdot e^{ikt} \cdot dt + \left[ \text{ctg} (\gamma_0 (y)) + \left( \frac{\partial G}{\partial y} \right) e^{ikt} \right] \cdot dy$$

with

$$\left( \frac{\partial x}{\partial y} \right)_t = \text{ctg} (\gamma_0 (y) + \gamma' (y))$$

and now considering always a fixed y we can formulate

$$\text{ctg} (\gamma_0 + \gamma') = \text{ctg} \gamma_0 + G_y \cdot e^{ikt}$$

Trigonometric relations and higher order terms going to zero give us

$$\gamma' = -\sin^2 \gamma_0 G_y \quad (38)$$

The transformation back to our Cartesian system of coordinates (Fig. 1) is done by

$$\begin{aligned} \hat{u} &= \hat{w}_t \cos(\gamma_0 + \gamma') + \hat{w}_n \sin(\gamma_0 + \gamma') \\ \hat{v} &= \hat{w}_t \sin(\gamma_0 + \gamma') - \hat{w}_n \cos(\gamma_0 + \gamma') \end{aligned} \quad (39)$$

The following steps have to be executed on Eq. (39):

- Introducing Eq. (33), (34)
- Applying trigonometric relations
- Neglecting all higher order products
- Introducing Eq. (36) and (38)
- Separating the purely steady expressions from the rest of the equations

The result is a substitute for the shock polar (27) which gives us the unsteady velocities  $u_1$  and  $v_1$  behind the shock.

$$\begin{aligned} \hat{u}_1 - m_1 G_y - i m_2 G - m_3 u_1 - m_4 v_1 &= 0 \\ \hat{v}_1 - n_1 G_y - i n_2 G - n_3 u_1 - n_4 v_1 &= 0 \end{aligned} \quad (40)$$

with A defined in Eq. (32) it is

$$\begin{aligned} m_1 &= \frac{2}{\kappa+1} v \sin 2\gamma_0 \sin^2 \gamma_0 \\ m_2 &= \frac{2k}{\kappa+1} (1+A) \sin^2 \gamma_0 \\ m_3 &= \sin^2 \gamma_0 \frac{\kappa-1}{\kappa+1} \left(1 - \frac{2A}{\kappa-1}\right) + \cos^2 \gamma_0 \\ m_4 &= \frac{\sin^2 \gamma_0}{\kappa+1} (1+A) \end{aligned} \quad (41)$$

$$\begin{aligned}
n_1 &= -\frac{2v}{\kappa+1} (\cos \gamma_0 + A) \sin^2 \gamma_0 \\
n_2 &= -\frac{2k}{\kappa+1} (1 + A) \cos \gamma_0 \sin \gamma_0 \\
n_3 &= m_4 \\
n_4 &= \sin^2 \gamma_0 + \cos^2 \gamma_0 \frac{\kappa-1}{\kappa+1} \left(1 - \frac{2A}{\kappa-1}\right)
\end{aligned}
\tag{42}$$

These are exactly the coefficients Chadwick presented in (4). For the airfoil in undisturbed supersonic flow the perturbation quantities  $u_1$  and  $v_1$  in front of the shock are zero. With this in mind, Eq. (40) reduces to the shockpolar derived by Teipel in (10). It should be stated explicitly at this point that we followed Teipel very closely in this extension of his work and that Chadwick indicates the way in (4).

We have to do a last step, to make Eq. (40) a tool for computing the unsteady boundary values along the shock:

On the leading edge we know  $G = 0$  because the shock is always attached.  $\hat{v}_1|_{y=0}$  is known from the boundary conditions on the airfoil, Eq. (21). Now we can isolate  $\hat{u}_1|_{y=0}$  in Eq. (40):

$$G_y|_{y=0} = \frac{1}{n_1} \left[ \hat{v}_1|_{y=0} - (n_3 u_1 + n_4 v_1) \right] \tag{43}$$

and

$$\hat{u}_1|_{y=0} = \frac{m_1}{n_1} \hat{v}_1|_{y=0} + n_1 \left( m_3 - \frac{m_1}{n_1} n_3 \right) + v_1 \left( m_4 - \frac{m_1}{n_1} n_4 \right) \tag{44}$$

After this initial step we develop finite differences along the steady shock to solve gradually for the unsteady values of  $\hat{u}_1$ ,  $\hat{v}_1$  and  $G$  as shown later.

## 2.4 Connections Between the Linear and the Nonlinear System of Equations

In an earlier work Teipel developed in (12) a method of characteristics for an oscillating single flat plate. He derived an analytical solution for the unsteady boundary values along the shock. Using the perturbation velocity of sound rather than the perturbation potential, his concept was taken by Bell in (9) and by Platzler and associates in (3,13,14) to obtain results for a cascade of flat plates. They started from the Euler and continuity equations with the substitution

$$U(x,y) \cdot e^{ikt} = u_1 \quad (45a)$$

$$V(x,y) \cdot e^{ikt} = \frac{1}{\sqrt{M^2 - 1}} \cdot v_1 \quad (45b)$$

$$C(x,y) \cdot e^{ikt} = \frac{2}{\kappa - 1} \frac{1}{M^2} \frac{a - a_\infty}{a_\infty} \quad (45c)$$

where  $U$ ,  $V$  and  $C$  are complex nondimensional amplitudes. The result is a set of differential equations which reads

$$\text{Continuity: } \frac{\partial V}{\partial x} + \sqrt{M^2 - 1} \cdot \frac{\partial U}{\partial y} + M^2 \frac{\partial C}{\partial x} + ik M^2 C = 0 \quad (46)$$

$$\text{Euler: } \frac{\partial U}{\partial x} + \frac{\partial C}{\partial x} + ikU = 0 \quad (47)$$

$$\text{Irrotationality: } \frac{\partial U}{\partial x} - \sqrt{M^2 - 1} \frac{\partial V}{\partial x} = 0 \quad (48)$$

Furthermore it is shown by Bell in (9) that the pressure in terms of Eq. (45) can be expressed as

$$p - p_\infty = \frac{2}{\kappa - 1} \rho_\infty a_\infty (a - a_\infty)$$

or

$$c_p = 2C$$

As  $(p - p_\infty)$  is our unsteady pressure disturbance (the steady pressure disturbance is zero for a flat plate) we can say with Eq. (20)

$$2C = -2(u_1 + ik\Psi) \quad (50)$$

or

$$C = -(u_1 + ik\Psi)$$

Now we take Eq. (45a), (45b) and (50) and substitute into Eq. (46) and we get

$$(M^2 - 1) \frac{\partial u_1}{\partial x} - \frac{\partial v_1}{\partial y} + 2 ik M^2 u_1 - k^2 M^2 \Psi = 0 \quad (51)$$

With  $u_1 = \Psi_x$  and  $v_1 = \Psi_y$  this is our Eq. (17), when we consider that for a flat plate the values for  $\varphi_x$  and  $\lambda_x$  are always zero.

$$\lambda = \sqrt{M^2 - 1}$$

Applying Eq. (50) on (47) gives us

$$u_1 = \frac{\partial \Psi}{\partial x}$$

which is one of our basic equations. Obviously, Eq. (48) can be transformed into Eq. (18). Hence it is shown that the basic equations previously used by Teipel and Platzler for the flat plate are equivalent to the perturbation potential equations used in this work, when these are reduced to the linear case.

The next step shows how to obtain an analytical expression for the oscillating shock generated by a single flat plate from the rather complicated differential equation (40).

As there are no perturbations in front of the shock, Eq. (40) reduces to the Teipel-form

30

$$\hat{u}_1 = m_1 G_y + i m_2 G$$

$$\hat{v}_1 = n_1 G_y + i n_2 G$$

In addition the steady shock angle  $\gamma_0$  can be expressed as

$$\sin \gamma_0 = \frac{1}{M} \quad (52)$$

and with this  $A$  from Eq. (41), defined in Eq. (32), becomes

$$A = \frac{1}{\sqrt{2} M^2 \sin^2 \gamma_0} = 1 \quad (53)$$

( $v = 1 + u_0 = 1$ )

The shockpolar written down explicitly reads

$$\hat{u}_1 = \frac{2}{\kappa+1} \sin \gamma_0 \sin^2 \gamma_0 G_y + i \frac{2\kappa}{\kappa+1} 2 \sin^2 \gamma_0 G$$

$$\hat{v}_1 = -\frac{2}{\kappa+1} (\cos 2\gamma_0 + 1) \sin^2 \gamma_0 G_y - i \frac{2\kappa}{\kappa+1} 2 \cos \gamma_0 \sin \gamma_0 G$$

Using the trigonometrical relations and Eq. (52) we obtain after some adding and subtracting

$$\hat{u}_1 = -\frac{1}{\sqrt{M^2 - 1}} \hat{v}_1 \quad (54)$$

for the velocities immediately downstream of the shock. We now go into the compatibility relation Eq. (19). Linearized with  $\lambda_x = 0$  it is for the upper side

$$\frac{\partial u_1}{\partial x} - \frac{1}{\sqrt{M^2 - 1}} \frac{\partial v_1}{\partial x} + \frac{2\kappa M^2}{M^2 - 1} u_1 - \frac{1}{M^2 - 1} \kappa^2 M^2 \psi = 0$$

Eq. (54) shows that

$$\frac{\partial u_1}{\partial x} - \frac{1}{\sqrt{M^2 - 1}} \frac{\partial v_1}{\partial x} = 2 \frac{\partial u_1}{\partial x}$$

Landahl (1) shows that  $(\Psi + \phi)$  should be continuous across a shock. As all the steady components are zero in this case his result applies here to the unsteady potential alone. On the other hand we consider the shock thickness as infinitely small. Assuming  $\hat{d}\Psi$  to be nonzero,

$$\frac{\partial \hat{\Psi}}{\partial x}$$

would become infinitely large across the shock. Thus  $\hat{\Psi}$  can only be zero along the unsteady shock of a single flat plate, or

$$\hat{\Psi} = \Psi$$

for a cascade.

The remaining steps are fairly straight forward. With the conditions shown above, Eq. (19) becomes

$$2 \frac{\partial u_1}{\partial x} + \frac{2ik M^2}{M^2 - 1} u_1 = 0$$

or

$$\frac{\partial u_1}{\partial x} = - ik \frac{M^2}{M^2 - 1} u_1$$

this can be integrated

$$u_1 = \text{const} \cdot e^{-i \frac{k M^2}{M^2 - 1} x}$$

for  $x = 0$  we get the perturbation for the leading edge

$$\hat{u}_{1LE} = - \frac{1}{\sqrt{M^2 - 1}} \cdot \hat{\phi}_{1LE}$$

$\phi_{1LE}$  can be obtained from the unsteady boundary conditions. The final expression for the unsteady velocities behind the shock attached to an oscillating flat plate is now

$$\hat{u}_1(x) = \hat{u}_{1LE} \cdot e^{-i \frac{k M^2}{M^2 - 1} x} \quad (55)$$

This result was also obtained by Bell (9) and considering the different systems of coordinates, it is equal to the solution found by Teipel in (12).

### 2.5 The Wake

Fig. 8 shows the characteristic net for the airfoil and wake regions. The fluid in the fields 2 and 11 has to adjust via the two trailing edge shocks in such a way, that the wake condition on the slip-line in the fields 4 and 13 is not violated. This condition requires flow tangency, continuity of pressure and normal velocity across the slip-line, i.e.,

$$\theta_4 = \theta_{13}$$

and

$$cp_4 = cp_{13} \quad (56)$$

$$w_{n4} = w_{n13}$$

$w_n$  is the velocity component normal to the slip-line. Again this problem is split into steady and unsteady parts.

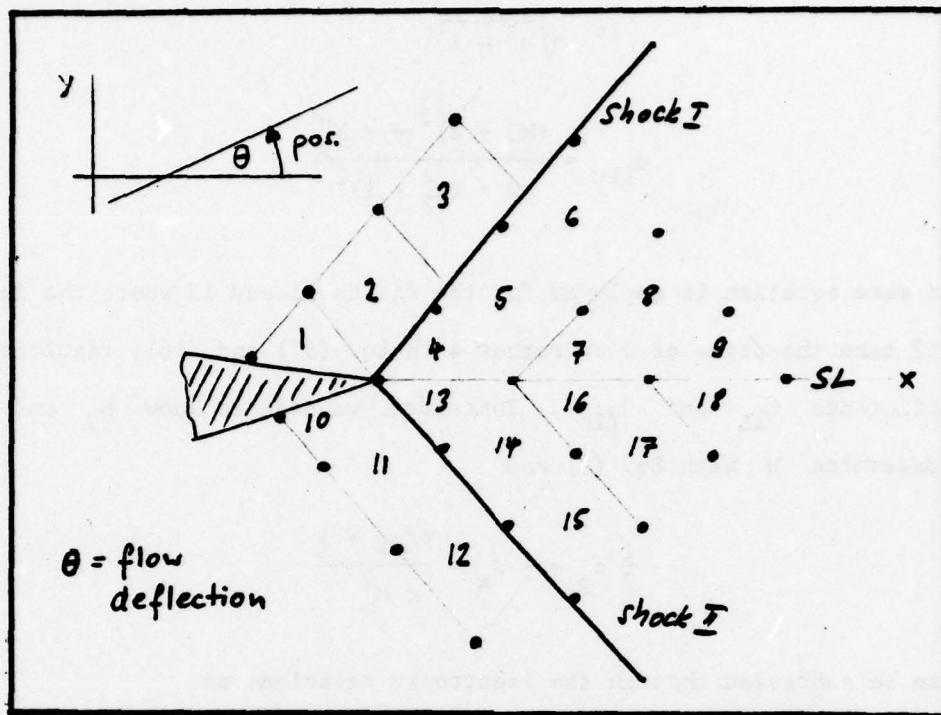


Fig. 8. The Net of Characteristics for the Steady Wake

Shapiro (11) gives a series solution for the pressure difference connected to a direction-change  $\Delta\theta = \theta_4 - \theta_2$  for supersonic flow.

$$2 \frac{P_4 - P_2}{\kappa P_4 M_2^2} = \pm C_{Iu} \cdot (\theta_4 - \theta_2) + C_{IIu} \cdot (\theta_4 - \theta_2)^2 \quad (57)$$

+ left running Mach lines

- right running Mach lines

with

$$C_{Iu} = \frac{2}{\sqrt{M^2 - 1}} \quad (58)$$

$$C_{IIu} = \frac{(M_2^2 - 2)^2 + \kappa M_2^4}{2 \cdot (M_2^2 - 1)^2}$$

The same equation is employed for the fields 11 and 13 where the indices 11 and 13 take the place of 2 or rather 4 in Eq. (57) and (58), resulting in the coefficients  $C_{IL}$  and  $C_{IIL}$ . Therefore, we need to know  $M_2$  and  $M_{11}$ . We can determine  $M$  with Eq. (9) and

$$\frac{1}{2} c_p = -\varphi_x = \frac{p/p_\infty - 1}{\kappa M_\infty^2}$$

$p/p_\infty$  can be expressed through the isentropic relations as

$$p/p_\infty = f(\kappa, M_\infty, M)$$

with this we obtain

$$M^2 = \frac{2}{\kappa - 1} \left[ \frac{1 + \frac{\kappa - 1}{2} M_\infty^2}{\left[ 1 - \frac{\kappa}{\kappa + 1} (\lambda - M_\infty^2 + 1) \right]^{\frac{\kappa - 1}{\kappa}}} - 1 \right]$$

For a given point in which  $\lambda$  is known we can find the flow direction for  $\theta_\infty = 0$  as parallel to the x-axis.

$$\operatorname{tg} \theta = \frac{v_o}{1 + u_o}$$

from Eq. (11)

$$v_o = \frac{2}{3} \frac{\mu}{(\kappa + 1) M_\infty^2}$$

and (13)

$$\mu = \mp \left[ (M_\infty^2 - 1)^{3/2} - \lambda^{3/2} \right]$$

we arrive through Eq. (25) at

$$\operatorname{tg} \theta = \mp \frac{2}{3} \cdot \frac{(M_\infty^2 - 1)^{3/2} - \lambda^{3/2}}{\lambda + M_\infty^2 \kappa + 1} \quad (60)$$

The upper sign indicates the left-running (- lower surface) and the lower sign the right running characteristic (= upper surface).

Now we have the coefficients  $C_{Iu}$ ,  $C_{IIu}$ ,  $C_{IL}$  and  $C_{IIL}$ . In addition it must be

$$p_4 = p_{13}$$

and

$$\theta_4 = \theta_{13}$$

Thus we can say

$$\begin{aligned} \frac{p_4}{p_\infty} &= \frac{1}{2} \kappa M_\infty^2 \left[ C_{Iu} \cdot (\theta_4 - \theta_2) + C_{IIu} \cdot (\theta_4 - \theta_2)^2 \right] + 1 \\ \frac{p_{13}}{p_\infty} &= \frac{1}{2} \kappa M_\infty^2 \left[ -C_{IL} \cdot (\theta_4 - \theta_{11}) + C_{IIL} \cdot (\theta_4 - \theta_{11})^2 \right] + 1 \end{aligned} \quad (61)$$

The equality of these two equations gives us a way to solve for  $\theta_4$ . As we took only second order terms of  $\Delta\theta_1$  the result is  $\theta_4 = \theta_{13} = 0$  for zero angle of attack. This is correct, because we assume isentropic flow across the shocks. The reason why it is done in this rather difficult way is that the upwash of the nonisentropic flow can be easily added to the program by simply adding the third coefficient  $C_{III}$  given by Shapiro. Actually, this is already done in the program. It has been reversed through setting  $C_{III} = 0$  in order to be consistent with the assumption of isentropic flow across the weak shocks.

After we have determined the flow in field 4 and 13 we compute the trailing edge shocks and the rest of the wake field. The step from the slip-line is repeated from the fields 5 and 14 to 7 and 16 as described, etc. All other steps in the wake field are general characteristic steps like those over the airfoil.

It is now possible to solve for the unsteady flow field of the wake. This is again done on the same grid locations which are known from the steady solution.

The conditions for this part of the solution are

- No Pressure-Jump Across the Slip-Line

$$u_{12u} + ik \Psi_{2u} = u_{12L} + ik \Psi_{2L} \quad (62)$$

(see Fig. 10).

- Tangential Velocity on the Slip-Line

$$\frac{v_0 + v_1}{1 + u_0 + u_1} = \text{tg} (\epsilon_0 + \epsilon') \quad (63)$$

(see Fig. 9)

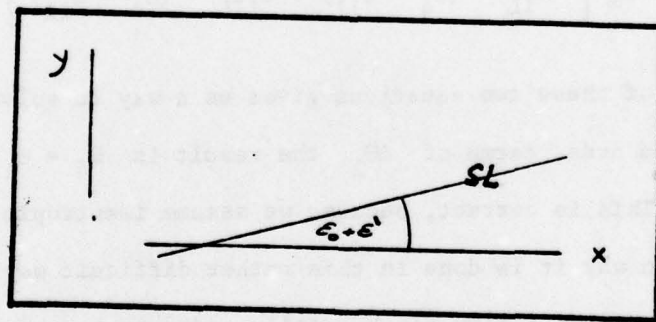


Fig. 9 Slip-Line Geometry

With trigonometric relations Eq. (63) becomes

$$\varepsilon' = v_0 - (1 + u_0) \operatorname{tg} \varepsilon_0 + v_1 - u_1 \operatorname{tg} \varepsilon_0$$

from the steady solution we know

$$v_0 - (1 + u_0) \operatorname{tg} \varepsilon_0 = 0$$

hence we obtain for the unsteady deflection of the slip-line

$$\varepsilon' = v_1 - u_1 \operatorname{tg} \varepsilon_0$$

as  $\varepsilon'$  has to be equal for the upper and the lower side of the flow field,

$$(v_1 - u_1 \operatorname{tg} \varepsilon_0)_{2u} = (v_1 - u_1 \operatorname{tg} \varepsilon_0)_{2L} \quad (64)$$

is the first equation in the system, which we want to solve for  $(u_1, v_1)_{2u}$  and  $(u_1, v_1)_{2L}$ . Indices refer to Fig. 10.

The next step uses Eq. (64). Fig. 10 shows the simultaneous step from the upper and lower wake field to the slip-line

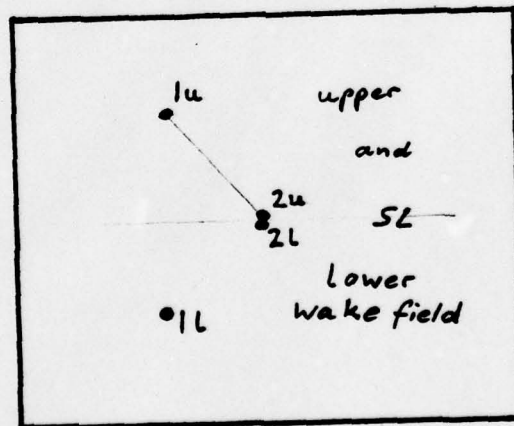


Fig. 10. Slip-Line Step

with

$$\psi_2 = \psi_1 + \frac{x_2 - x_1}{2} \left( u_{12} + u_{11} \pm 2 \frac{v_{12} + v_{11}}{\sqrt{\lambda_1} + \sqrt{\lambda_2}} \right) \quad (65)$$

We can substitute  $\psi_2$  by the known  $\psi_1$  and the desired unknowns. This gives us the second equation of the system when properly recast in left- and right- hand sides.

The two missing statements which make the system solvable comes from the compatibility relation (19). Applied to the upper and lower slip-line steps, respectively, they complete the system.

The computational procedure is explained in more detail in chapter 3.

### 3. THE SINGLE OSCILLATING AIRFOIL: WING

#### 3.1 The Program Organization Over the Airfoil

The name of the program for the single oscillating airfoil is W I N G. It was designed to use the subroutines as often as possible, that is over the surface as well as in the wake field. For this purpose a special notation was introduced which is shown in Fig. 11 and Fig. 12.

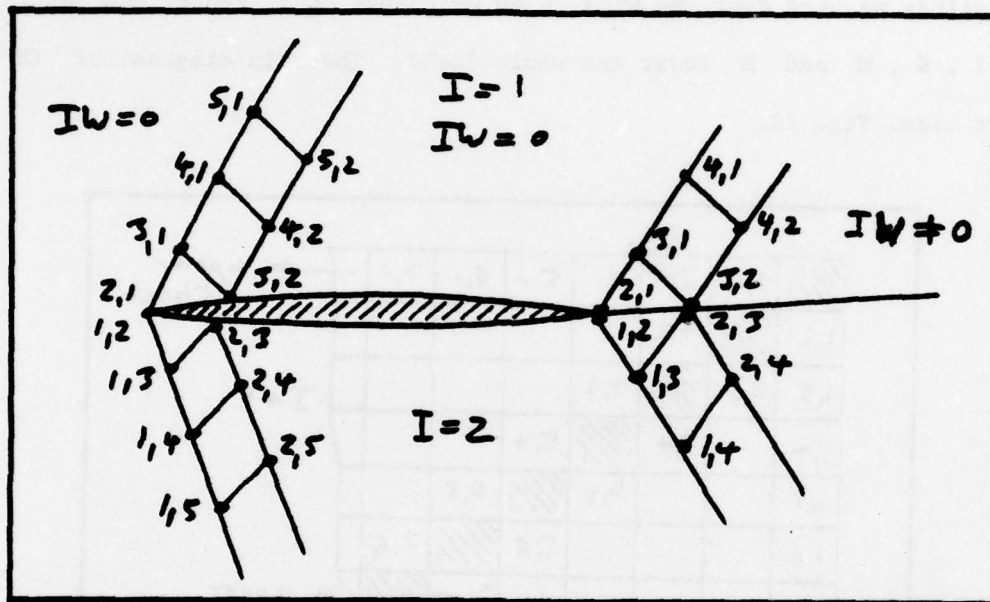


Fig. 11. Notation of the Fields in Wing

The arrays which contain properties of the flow field have the dimension  $(K, M, N)$ .  $M$  and  $N$  give the location of the gridpoint: Point  $M$  on the  $N$ -th characteristic for the upper field and vice versa for the lower field.

$K$  and  $IW$  indicate where we are:

- $IW = 0$  Not yet in the wake
- $IW = 0$  Computation of the wake
- $K = 1$  Field already known
- $K = 2$  Field computed now

The upper and lower side is indicated by  $I = 1$  and  $I = 2$  respectively. Thus we can switch the changing sign of the characteristics with

$$(-1)^{(I+1)}$$

This notation has the advantage that we can compute the field on both sides of the wing ( $IW = 0$ ) by switching  $I$ . For the wake we still can use the subroutines we used over the wing, when we change  $K$ . Proper combinations from  $IW$ ,  $I$ ,  $K$ ,  $M$  and  $N$  cover the whole field. The main diagonal of  $(M, N)$  is not used, Fig. 12.

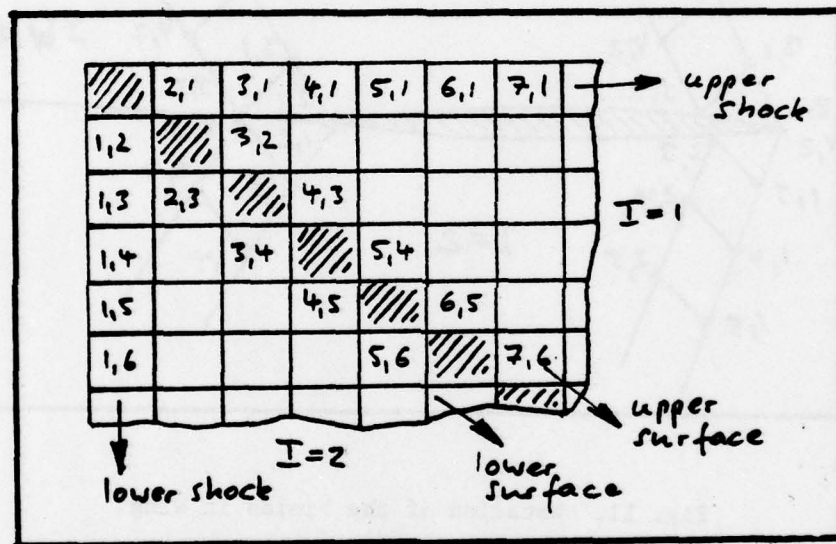


Fig. 12. Array Organization in Wing

Originally for the first shock we need only the Teipel - form of Eq. (2.40). But we have to use the whole expression in our step from the airfoil into the wake because the unsteady velocities over the airfoil are nonzero. So the subroutine R A N D S includes the extended shockpolar. To be consistent in our treatment of the shock at the leading- and trailing edge, an initial field in

front of the airfoil is generated by the characteristics with the slope

$$\frac{dy}{dx} = \pm \frac{1}{\sqrt{M^2 - 1}}$$

all perturbation properties are set to zero. During the computation of the flow field over the airfoil (IW = 0) the index K has the meaning

K = 1 Initial field

K = 2 Field over the airfoil

After this calculation step we do not need the initial field any longer and we copy (2, M, N) to (1, M, N). Now the K=2 - array is free for the wake field results and K indicates

K = 1 Field over the airfoil

K = 2 Wake field

In the shown listing of W I N G the arrays have the size (2, 25, 25) to keep the storage area reasonable. However, with 16 points on the airfoil the wake is only computed to a distance of 1.7 times the chordlength from the leading edge. If this is not enough, the program can be blown up by increasing the arrays to (2, KV, KV).

The value of KV has to correspond with the first statement of the program. All the DO-loops are then dimensioned correctly. Another change has to be done in subroutine L I F T. Here all the arrays stay the same besides A. This is only a dummy-array to define the space between X, U, Q, P and PX, PS, PU. A must have the size

$$A(22 \cdot KV^2 + 4 \cdot KV - 1000)$$

The next page shows an approximate block diagram of the program. On the sides are the names of the subroutines involved in the step of the respective block. This diagram is only a summary of what is really done. A lot of

details have been skipped in order to keep it clear. One of the most important routines does not even appear. `FIND` is an orientation subprogram that works only on the field  $(1, M, N)$ . If the arbitrary coordinates  $(x, y)$  are the input of `FIND`, it comes out with the indices  $(M2, N2)$  which give us the mesh of the field in which  $(x, y)$  is positioned; see Fig. 13.

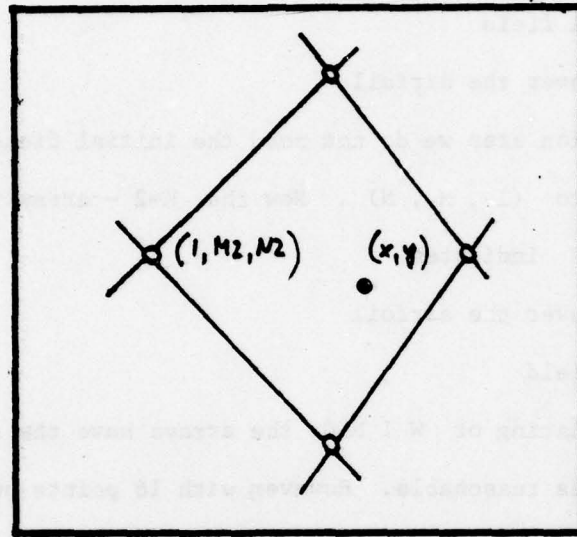
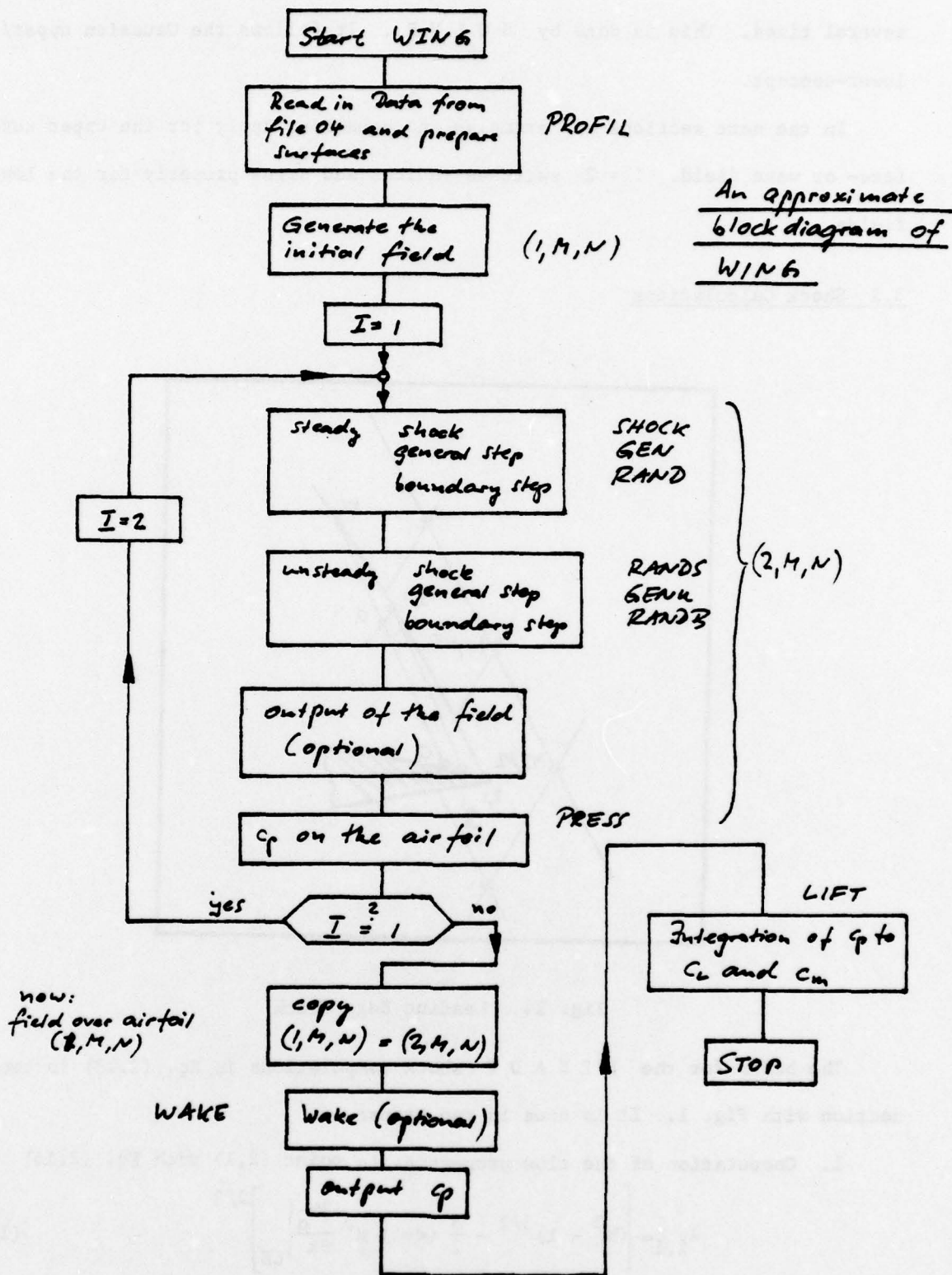


Fig. 13. A Characteristic Mesh

Again `I` is the switch for upper and lower sides. If  $(1, M, N)$  does not contain  $(x, y)$ , `FIND` indicates this by setting `IE = 1`, which is otherwise zero. Besides this there is a message in the output. `IE = 1` causes for the LE-shock to assume free stream values in front of it. For the TE-shock it terminates further computations, because there would be no sense in it. If  $(1, M, N)$  contains  $(x, y)$  and  $(1, M2, N2)$  can not be found, this will be in the output also. Then we know, something went wrong. But this should not happen. It is only a precautionary feature.

Another important background - subroutine is `SOLVE`. During the run through `WING`, complex systems of linear equations have to be solved

TABLE I



AN APPROXIMATE BLOCK DIAGRAM OF WING

several times. This is done by S O L V E . It follows the Gaussian upper/lower-concept.

In the next sections all examples and equations apply for the upper surface- or wake field. I = 2 switches indices and signs properly for the lower fields.

### 3.2 Shock Calculations

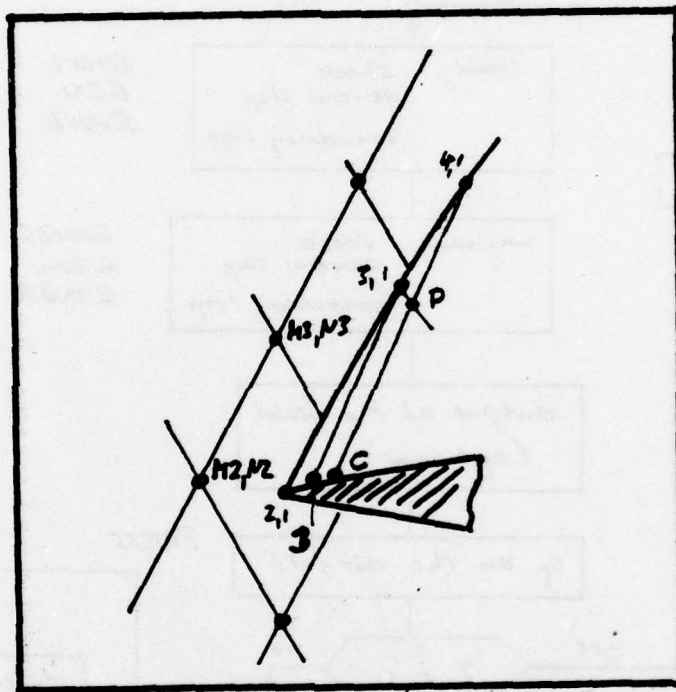


Fig. 14. Leading Edge Shock

The basis for the S T E A D Y shock computations is Eq. (2.13) in connection with Fig. 1. It is done in two steps:

1. Computation of the flow properties in point (2,1) with Eq. (2.16)

$$\lambda_{2,1} = \left[ (M^2 - 1)^{3/2} - \frac{3}{2} (\kappa + 1) M^2 \frac{\partial h_0}{\partial x} \Big|_{LE} \right]^{2/3} \quad (1)$$

which is all we want to know for a point in the steady field. The slope of the shock at the LE is taken as the average of the characteristic slopes upstream and downstream of the shock:

$$\operatorname{tg} \gamma_{o_{2,1}} = \frac{2}{\sqrt{\lambda_{2,1}} + \sqrt{\lambda_{M3, N2}}} \quad (2)$$

2. All the other points of the shock are found by the same step done repeatedly:

From the LE we make a step  $DX(I)$  to find B (Fig. 14). Employing  $\left. \frac{\partial h_o}{\partial x} \right|_B$  in Eq. (2.16) gives us

$$\lambda_{3,1} = \lambda_B$$

and so the slope of the characteristics starting in B

$$\operatorname{tg} \alpha_B = \frac{1}{\sqrt{\lambda_B}}$$

The intersection of the shock with the slope  $\operatorname{tg} \gamma_{o_{2,1}}$  and the B - characteristic is the point (3,1). Here the slope of the shock is changed to

$$\operatorname{tg} \gamma_{o_{3,1}} = \frac{2}{\sqrt{\lambda_{3,1}} + \sqrt{\lambda_{M3, N2}}}$$

we apply the same procedure for point C respectively (4,1) ; etc. for the whole shock.

The points (M2,N2) or rather (M3,N3) are part of the initial field for the LE-shock. In case of the TE-shock they are part of the field over the airfoil. They are spotted in F I N D .

All this is done in S H O C K .

The unsteady shock can be determined also with two basic steps:

1. The initial step was already described with Eq. (2.43) and (2.44).

After this we have  $\hat{u}_1$  and  $\hat{v}_1$  at the origin of the shock

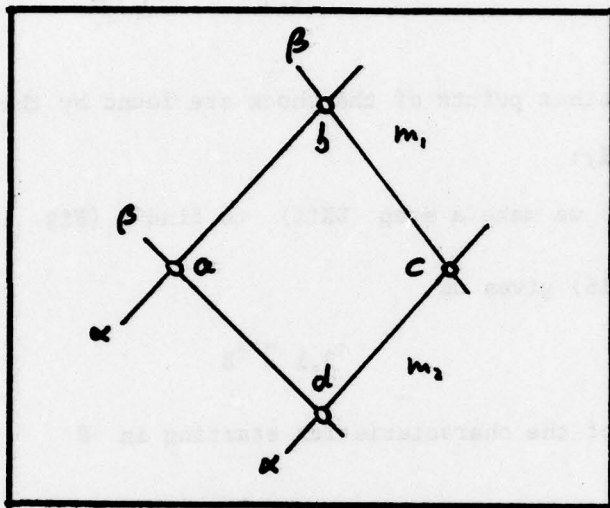


Fig. 15. Computation of  $\lambda_x$

2. For the remaining points we take the two Eq. (2.40) (Shockpolar) and Eq. (2.19) (Compatibility Relations on the Characteristics Between (3,1) and (4,1) ; see fig. 14).

This gives us a complex system of three linear equations with the solution  $\hat{u}_1$ ,  $\hat{v}_1$  and  $G$  in each point. The system is solved by subroutine `SOLVE`. The coefficients of Eq. (2.40) are computed in `COEFF1`.

However, to express Eq. (2.19) in finite difference form, we have some difficulties with  $\lambda_x$ . Consider Fig. 15.

If we are moving from  $a$  to  $b$ ,  $\Delta\lambda$  is zero because we are on a left-running characteristic.

$\lambda_x$  is in point a

$$\lambda_x|_\alpha = \frac{\lambda_c - \lambda_a}{x_c - x_a}$$

through  $\lambda_c = \lambda_\alpha$  and approximately  $x_c - x_a = 2(x_c - x_d)$ , we can say

$$\lambda_x|_\alpha = \frac{1}{2} \frac{\lambda_d - \lambda_a}{x_d - x_a}$$

or

$$\lambda_x|_\alpha = \frac{1}{2} \lambda_x|_\beta$$

We go now back to Fig. 14 and Eq. (2.19). Assuming that setting  $\lambda_x(3,1) = \lambda_x(P)$  causes only a small error because both points are very close together, we can express  $\lambda_x(P)$  as

$$\lambda_x(P) = \frac{1}{2} \frac{\lambda_{4,1} - \lambda_{3,1}}{x_p - x_{3,1}}$$

With this the finite difference form of Eq. (2.19) along the  $\alpha$ -characteristic from P to (4,1) is no problem. It is written without the second index.

Thus 3,1 becomes 3:

$$\frac{\hat{u}_{14} - \hat{u}_{13}}{x_4 - x_p} - \frac{\hat{v}_{14} - \hat{v}_{13}}{\sqrt{\lambda_4}} + \frac{\hat{u}_{14} + \hat{u}_{13}}{\sqrt{\lambda_4} + \sqrt{\lambda_3}} \cdot \left[ \frac{1}{2} \frac{\lambda_4 - \lambda_3}{x_p - x_3} + 21k M^2 \right] - k^2 M^2 \frac{\hat{\psi}_4 + \hat{\psi}_3}{\sqrt{\lambda_4} + \sqrt{\lambda_3}} = 0 \quad (4)$$

With Eq. (2.64) substituted for  $\hat{\psi}_4$  and the  $\hat{u}_3$ ,  $\hat{v}_3$  and  $\hat{\psi}_3$  already known, this is the third equation of the system after recasting into right and left hand sides.

After obtaining the results with `S O L V E`,  $\hat{\Psi}_4$  is calculated separately. For the initial step at the shock origin we set

$$\hat{\Psi}_2 = \Psi_{M2, N2}$$

The procedure for the trailing edge shock is exactly the same. In the boundary values  $\text{tg } \epsilon_0$  is substituted for

$$\frac{\partial h_0}{\partial x}$$

The unsteady boundary values of the field behind the shock are computed in subroutine `R A N D S`.

### 3.3 The Boundary Step

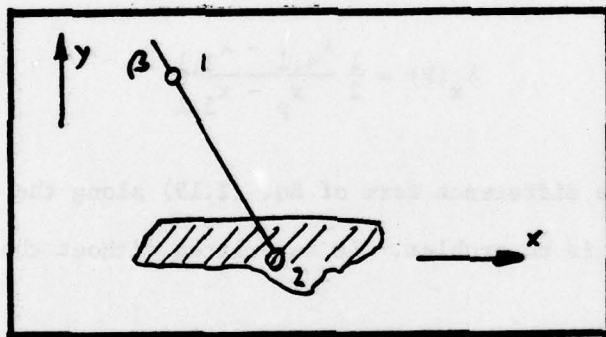


Fig. 16. Boundary Step

Fig. 16 shows the step from the field to the surface. It is done in `R A N D`. On the  $\beta$ -characteristic from 1 to 2 we have to satisfy two equations: Eq. (2.16)

$$\lambda_2 = \left[ \lambda_\infty^{3/2} - \frac{3}{2} (\kappa+1) M^2 \frac{\partial h_0}{\partial x} \Big|_2 \right]^{2/3} \quad (5)$$

and

$$\frac{y_2 - y_1}{x_2 - x_1} = - \frac{2}{\sqrt{\lambda_2 + \lambda_1}}$$

To make it easier, we set  $y_2$  to zero and get

$$\lambda_2 = \left( \frac{2(x_2 - x_1)}{y_1} - \lambda_1 \right)^2 \quad (6)$$

Setting Eq. (6) equal to (5), we can find by iteration the  $x_2$  which matches both of them. With  $x_2$  we then find  $\lambda_2$ . Here the steady problem is solved. The unsteady part is very straightforward. The known  $k$  gives us  $\Psi_1$  this point via the unsteady boundary conditions Eq. (2.21) or (2.22). Now the only unknown left is  $\Psi_x$ . It can be separated from the compatibility form. Again  $\Psi_2$  is obtained separately afterwards from Eq. (2.64). The unsteady step is computed in R A N D B .

#### 3.4 The General Step

$\lambda_b$  and  $\lambda_d$  are known (Fig. 15). From the  $\alpha$  - characteristic we know:

$$\lambda_c = \lambda_d$$

Finding the point  $c$  for the steady case is only a geometrical problem. It is the intersection of the two lines with the slopes

$$m_1 = - \frac{2}{\sqrt{\lambda_b} + \sqrt{\lambda_c}} \quad \text{and}$$

$$m_2 = \frac{1}{\sqrt{\lambda_c}}$$

they run through the points  $b$  and  $d$ . This is done in G E N .

The unsteady part is done in G E N U . Here the compatibility relation (2.19) is applied for both characteristic directions from point b and d to the point c . The unknowns are  $u_{1c}$  and  $v_{1c}$  . We obtain them as results from a complex system of two linear equations which we form out of the two compatibility relations. The system is solved through S O L V E . As always  $\Psi_c$  is evaluated separately from Eq. (2.64).

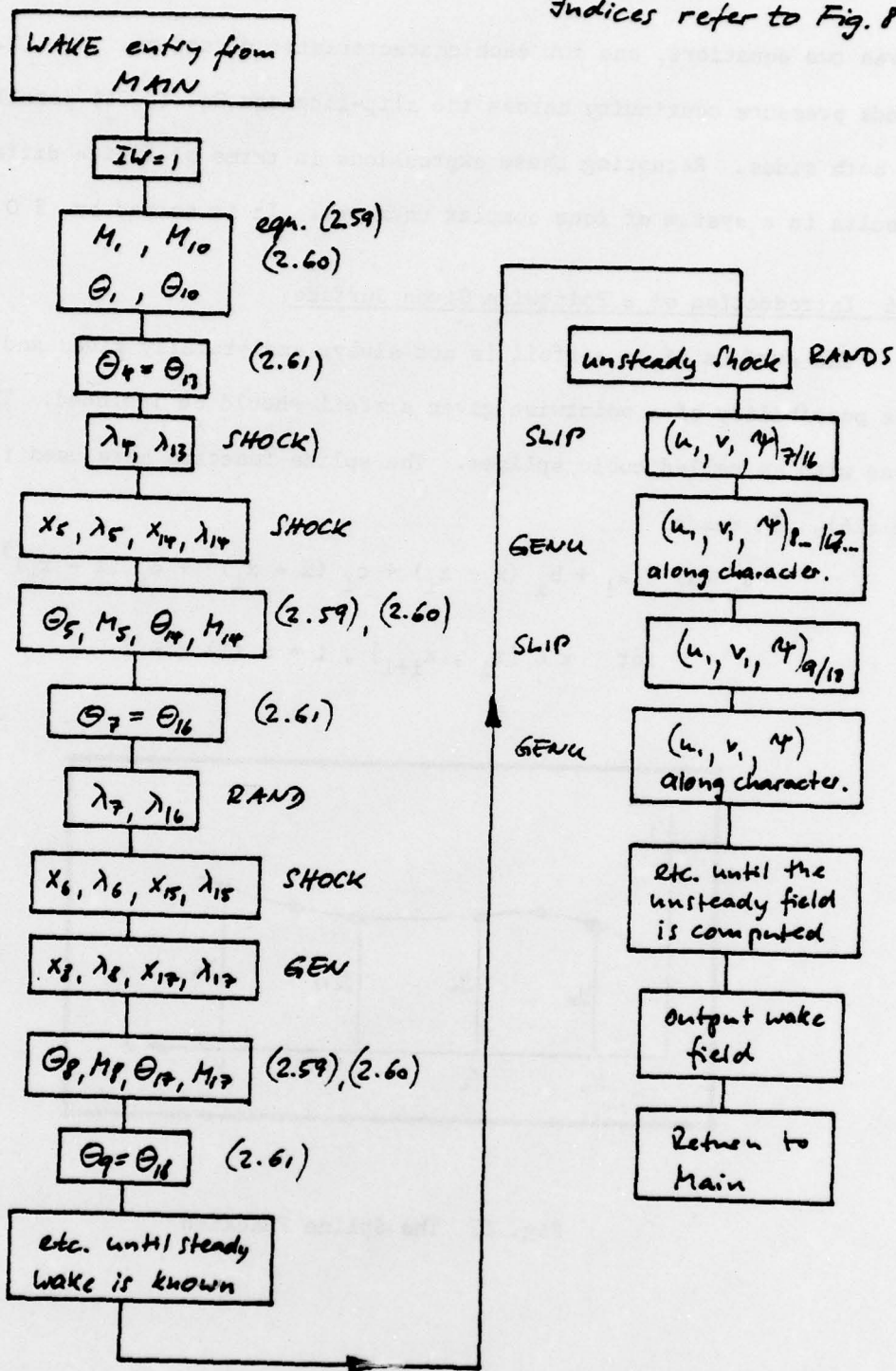
### 3.5 Computation of the Wake

For the wake computation W I N G leaves the main program completely and works in a new organization program, called W A K E . This was mainly done to separate the entirely different computation sequence from the organization from the field above the surface. In section 2.5 it was already indicated how to solve for the slip-line and then for the wake field. The wake field computation is shown in the following flow diagram. The indices refer to Fig. 8. Attached to the blocks are the names of the subroutines or the equations involved in the step.

Once the steady field is known, the calculation becomes straightforward again. The conditions directly downstream of the TE are defined by the unsteady shockpolar (2.40). Applied simultaneously for the upper and lower airfoil-side, they give the conditions for pressure continuity (2.62) and parallel flow on the slip-line (2.64). With these four equations we can solve for  $(\hat{u}, \hat{v})_{up}$  and  $(\hat{u}, \hat{v})_{lo}$  . This is the slightly more difficult initial step for the TE-shock . After it is done, the computation runs along the shock or rather along the characteristics with the same subroutines as over the airfoil. The basic difference lies in the step to the slip-line. Here R A N D B can not be used any longer. S L I P does the simultaneous step from the upper and the lower wake field to the slip-line. The compatibility relation (2.19)

TABLE II

Indices refer to Fig. 8



An Approximate Block Diagram for the Wake Computation

gives two equations, one for each characteristic direction. Eq. (2.62) demands pressure continuity across the slip-line and Eq. (2.64) parallel flow on both sides. Recasting these expressions in terms of finite differences results in a system of four complex unknowns. It is solved by S O L V E .

### 3.6 Introduction of a Pointwise Given Surface

The surface of an airfoil is not always analytically given and therefore the possibility of a pointwise given airfoil should be included. This can be done with so called cubic splines. The spline function here used is described in (16). It reads

$$S_i(x) \equiv a_i + b_i(x - x_i) + c_i(x - x_i)^2 + d_i(x - x_i)^3$$

$$\text{for } x \in [x_i, x_{i+1}], i = 0(1)n - 1$$

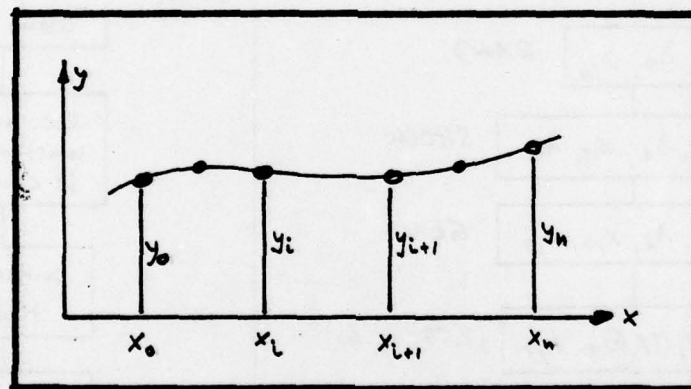


Fig. 17 The Spline Function

The demand for

$$\begin{aligned}
 S_i(x_i) &= y_i & i &= 0 \text{ (1) } n \\
 S_i(x_i) &= S_{i-1}(x_i) & i &= 1 \text{ (1) } n \\
 \left. \begin{aligned}
 S_i'(x_i) &= S_{i-1}'(x_i) \\
 S_i''(x_i) &= S_{i-1}''(x_i)
 \end{aligned} \right\} & i &= 1 \text{ (1) } n-1
 \end{aligned}$$

leads to a system of equations which give for each interval  $(x_i, x_{i+1})$  the coefficients of the spline  $S_i(x)$ . With

$$h_i = x_{i+1} - x_i$$

The coefficients can be expressed for known  $c_i$  as

$$\begin{aligned}
 a_i &= y_i \\
 b_i &= \frac{1}{h_i} (a_{i+1} - a_i) - \frac{h_i}{3} (c_{i+1} + 2c_i) \\
 d_i &= \frac{1}{3h_i} (c_{i+1} - c_i)
 \end{aligned} \tag{7}$$

The  $c_i$  are the solution of the linear algebraic system

$$\begin{aligned}
 h_{i-1} \cdot c_{i-1} + 2c_i (h_{i-1} + h_i) + h_i c_{i+1} &= \\
 \frac{3}{h_i} (a_{i+1} - a_i) - \frac{3}{h_{i-1}} (a_i - a_{i-1}) & \quad i = (1) n-1
 \end{aligned} \tag{8}$$

where  $c_0$  and  $c_n$  are presumed to be zero:

$$c_0 = c_n = 0$$

Once the coefficients  $a_i$  through  $d_i$  are known, we have a set of polynomials which are excellent for interpolation between  $x_i$  and  $x_{i+1}$ . As in each  $x_i$  not only the value but also the slope and the curve of the two neighboring polynomials are identical,  $S_i(x)$  is also applicable to determine  $S_i'(x)$  and  $S_i''(x)$  of a pointwise given function between those points:

$$S_i'(x) \equiv b_i + 2 c_i (x - x_i) + 3 d_i (x - x_i)^2$$

$$S_i''(x) \equiv 2 c_i + 6 d_i (x - x_i)$$

In W I N G the subroutine P R O F I L follows two options:

L04 = 1      Airfoil Analytical Given

L04 = 2      Pointwise Given

For L04 = 2 system (8) is solved. With (7) the coefficients can be re-created at any time. Thus it is possible to read in the geometry of any reasonable airfoil and obtain good approximations for position, slope and curve at any station of it. These values are needed for the steady boundary conditions in B O U N D called from R A N D and S H O C K .

### 3.7 RESULTS

In this chapter W I N G - results are compared with (2), Teipel's linear and nonlinear airfoil results, using the method of characteristic and with Verdon's (18) analytical results for the velocity and pressure distribution over a flat plate.

Figs. 18 through 23 provide a comparison of the Teipel - with the W I N G results. Although it can be said that the linear solutions ( $\tau = 0$ ) generally agree quite well, this is not true for higher Mach numbers, as can be seen in

Fig. 20 for  $M = 1.4$ . For this case the real part of  $C_p$  shows considerable differences distributed toward the TE. For  $\tau \neq 0$  there are differences distributed over all cases. This was already noted by Chadwick (4) and Strada (5) who argued that the deviations could be caused by different difference equations, averaging procedures and number of grid points. However, this could not be proved because Teipel does not expose his complete set of finite difference equations.

This work does not show a comparison with Chadwick's results (4) obtained for airfoils  $\tau \neq 0$  and wedges. Nevertheless this has been done. It shows excellent agreement in all cases no matter whether airfoil or wedge. Thus it can be said that W I N G is at least equivalent to Chadwick's solution over the airfoil.

Further W I N G - results are compared with Verdon's work (18), which provides analytically derived results over the flat plate and downstream of it into the wake.

First it should be noted that the  $C_p$  - distribution for  $0 \leq x \leq 1$  is very good. This can be seen in Fig. 24 and 25. The disagreement for the plunge motion in Fig. 25 looks appreciable only because of the extended scale. The original Verdon plots in this area are rather hard to read because of his desire to show the wake pressure distributions rather than those over the airfoil for the zero upwash. One of his results is that the assumption of  $v_1 = 0$  on the slip-line is wrong.

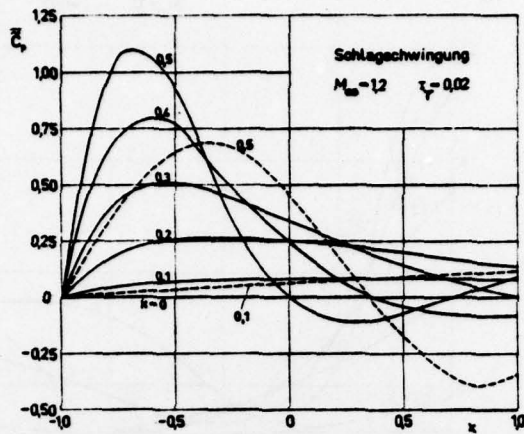
It has to be  $C_p = 0$ . With this assumption he gets a  $v_1$  - distribution on the slip-line shown in Fig. 26 and 27.

The system of equations used in W I N G does not presume  $cp_u = 0$  and  $cp_L = 0$  but only  $cp_u = cp_L$ . It turns out that  $|c_p|$  on the slip-line is of the order 0.000001. The upwash behind the airfoil obtained by W I N G is shown in Fig. 26 and 27 for  $1 \leq x \leq 4.6$ . For this range the agreement is

considered to be excellent for pitch and plunge. The reason for not showing results in the area  $x > 4.6$  is storage difficulties in the time sharing system of the used IBM 360/67.

If we proceed from the slip-line into the wake field, Verdon results versus W I N G - results are shown in Fig. 28 and 29. Again the agreement is very good. The total unsteady pressure along characteristics with origin  $x = 1.08$  and  $x = 2.0$  from the LE is shown. In this way the pressure field over the wake can be represented. Fig. 28 and 29 are very important, as the single airfoil is only the first step to the cascade.

The good agreement with Verdon's results justifies the method of characteristic approach for the problem of the oscillating cascade. The advantage of this approach is the flexibility with which characteristic computational procedures can be designed.



$$\tau_T = 2\tau$$

Abb. 8 a. — Druckverteilung  $\bar{C}_p$  (Realteil) für ein Parabelbogenprofil mit dem Dickenverhältnis  $\tau = 0,02$  (Parameter : reduzierte Frequenz  $k$ ).

Fig. 18A. From (2)

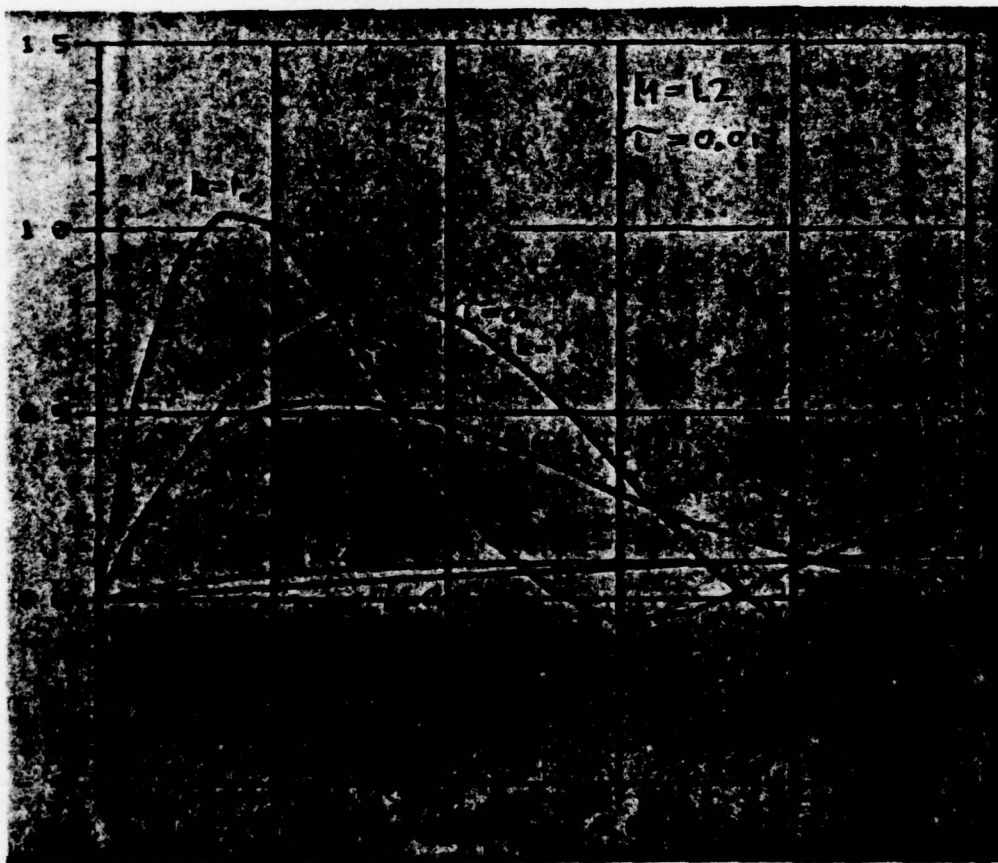


Fig. 18B. Pressure Distribution (Real Part) from W I N G , Plunge Motion

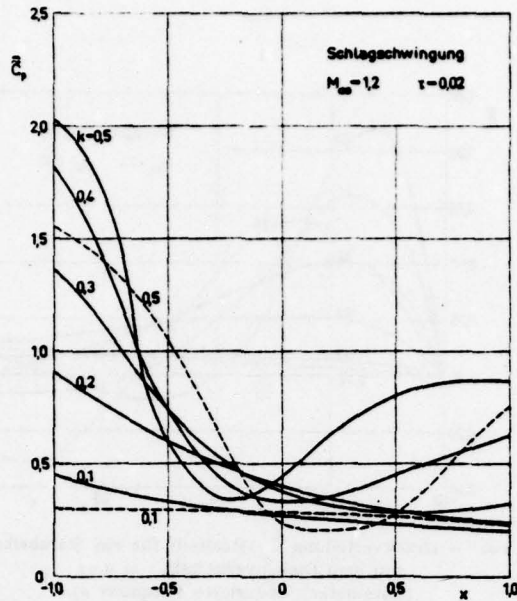


Abb. 8 b. — Druckverteilung  $\bar{C}_p$  (Imaginärteil) für ein Parabelbogenprofil mit dem Dickenverhältnis  $\tau = 0,02$  (Parameter : reduzierte Frequenz  $k$ ).

Fig. 19A. From (2)

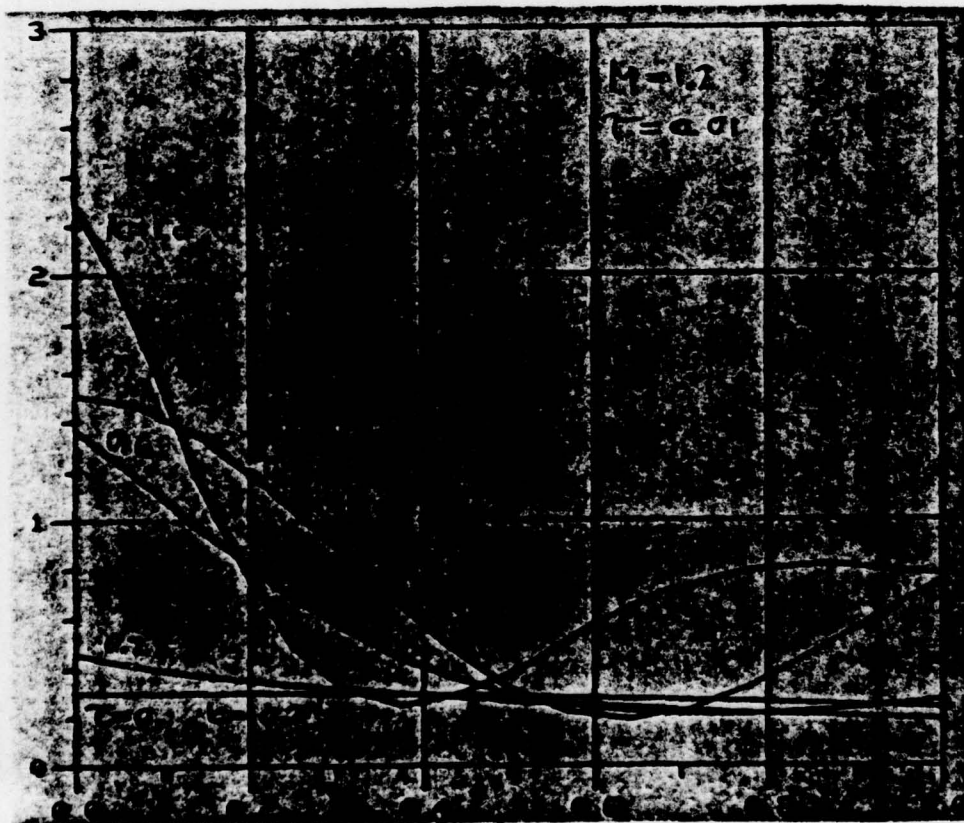


Fig. 19B. Pressure Distribution (Imaginary Part) from W I N G , Plunge Motion

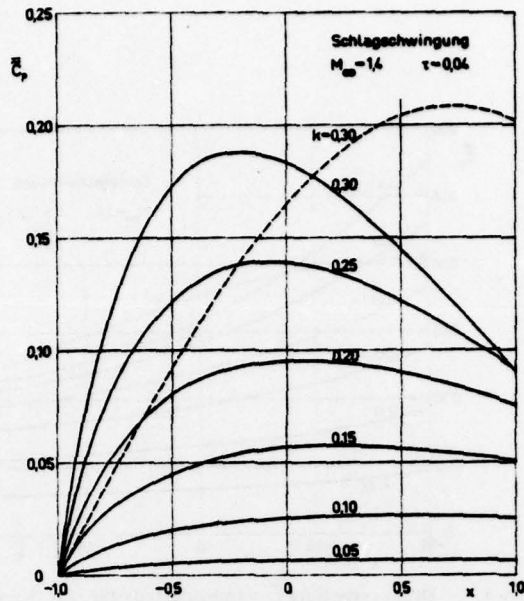


Abb. 9 a. — Druckverteilung  $\bar{C}_p$  (Realteil) für ein Parabelbogenprofil mit dem Dickenverhältnis  $\tau = 0,04$  (Parameter : reduzierte Frequenz  $k$ ).

Fig. 20A. From (2)

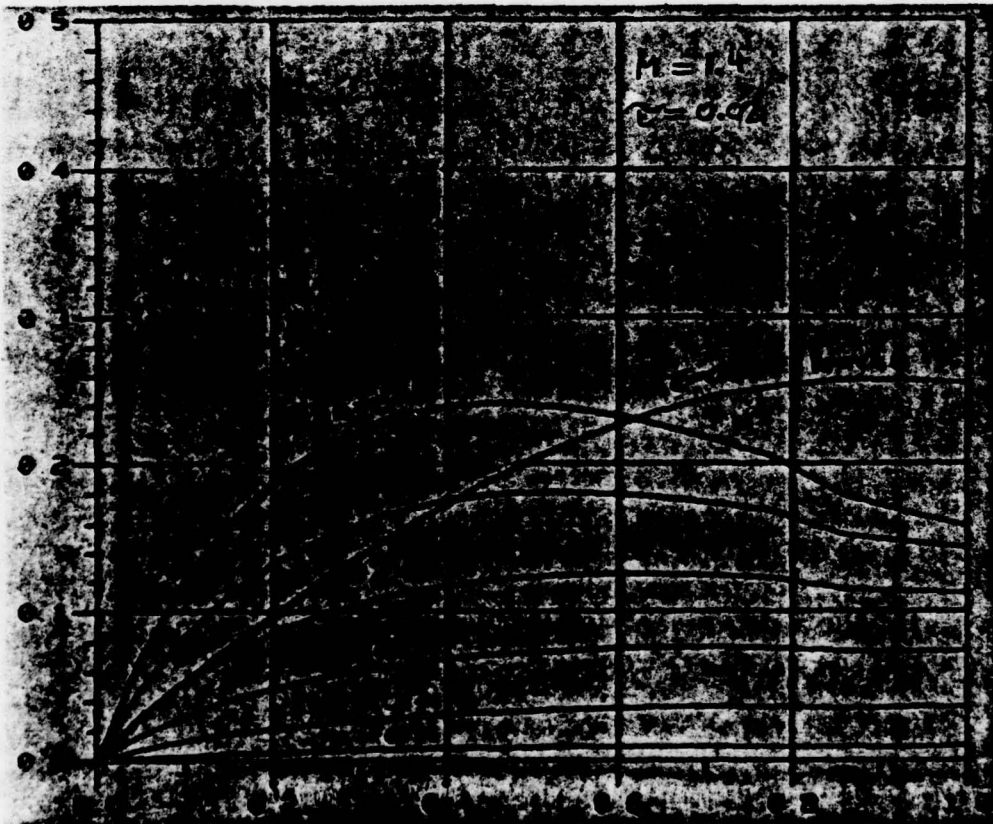


Fig. 20B. Pressure Distribution (Real Part) from W I N G , Plunge Motion

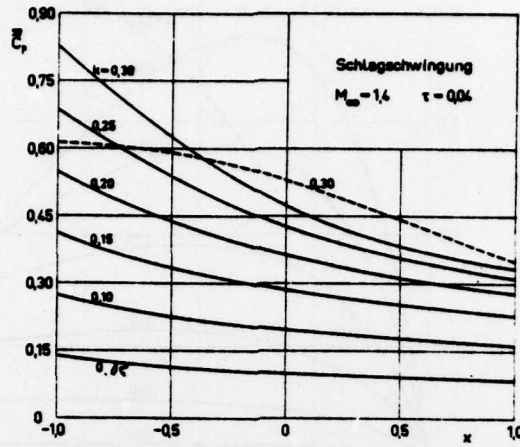


Abb. 9 b. -- Druckverteilung  $\bar{C}_p$  (Imaginärteil) für ein Parabelbogenprofil mit dem Dickenverhältnis  $\tau = 0,04$  (Parameter : reduzierte Frequenz  $k$ ).

Fig. 21A. From (2)

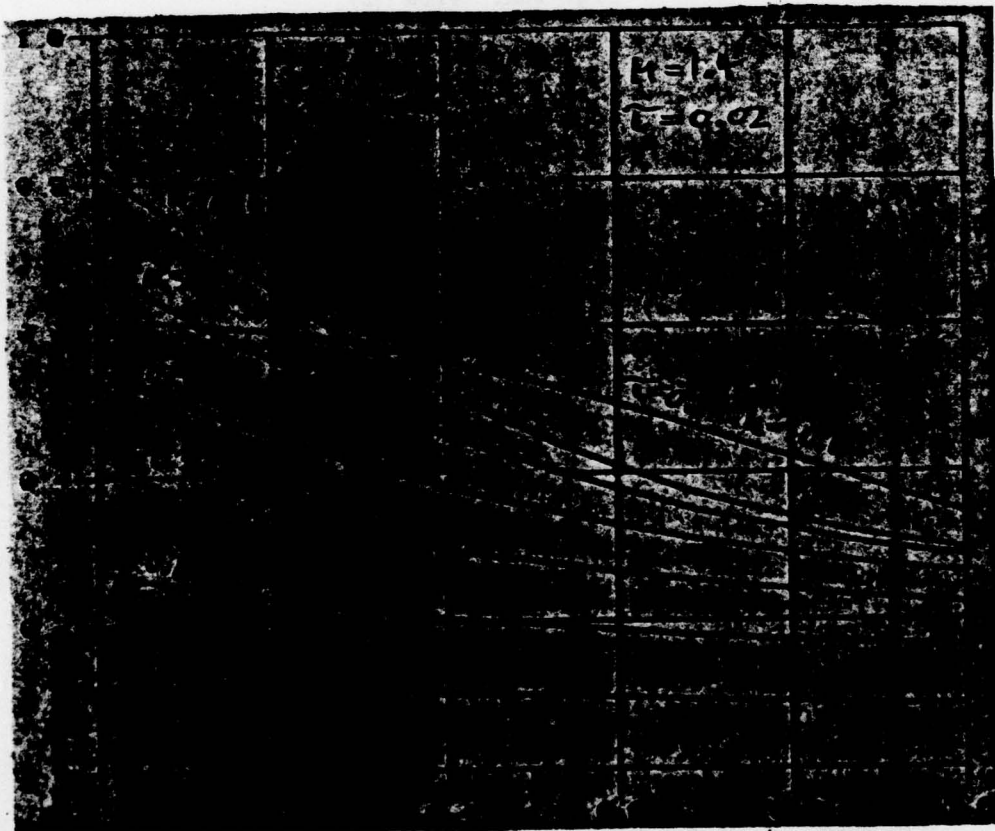


Fig. 21B. Pressure Distribution (Imaginary Part) from W I N G , Plunge Motion

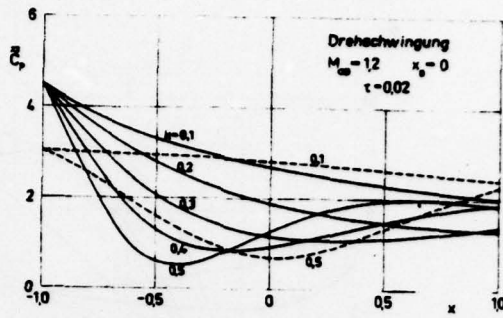


Abb. 13 a. — Druckverteilung  $\bar{C}_p$  (Realteil) für ein Parabelbogenprofil mit dem Dickenverhältnis  $\tau = 0,02$  (Parameter : reduzierte Frequenz  $k$ ).

Fig. 22A. From (2)

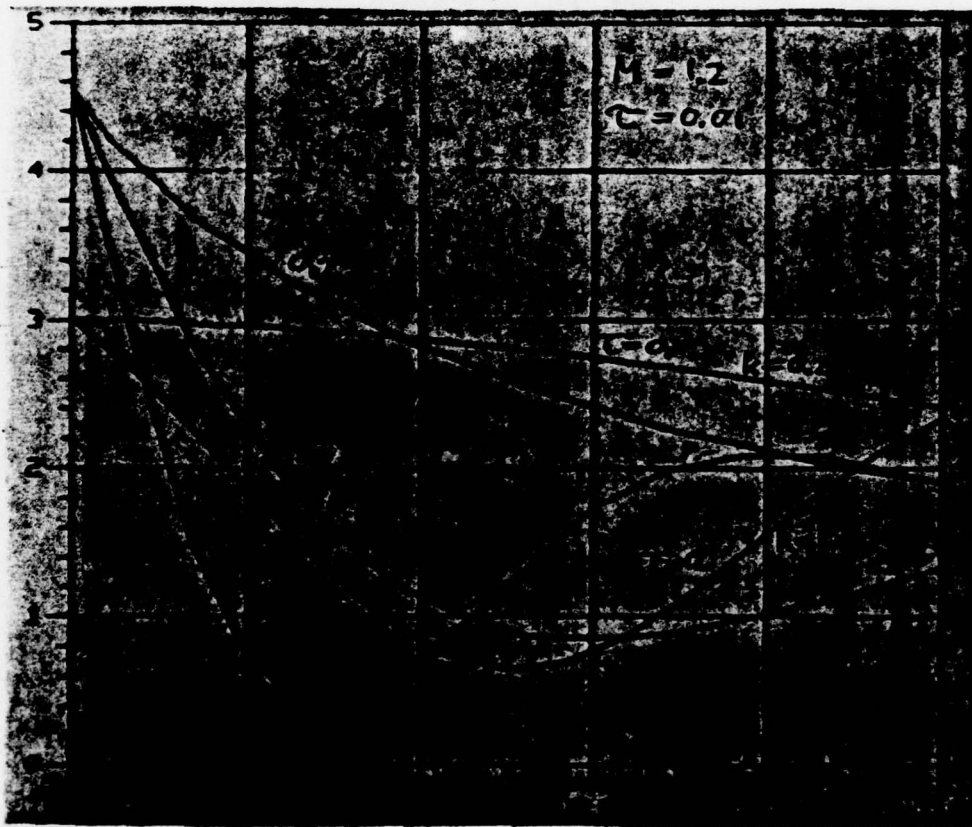


Fig. 22B. Pressure Distribution (Real Part) from W I N G , Pitching Motion

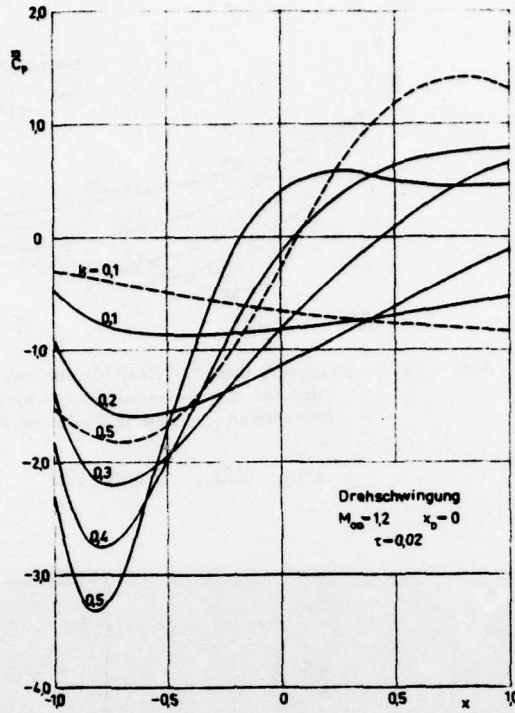


Fig. 23A. From (2)

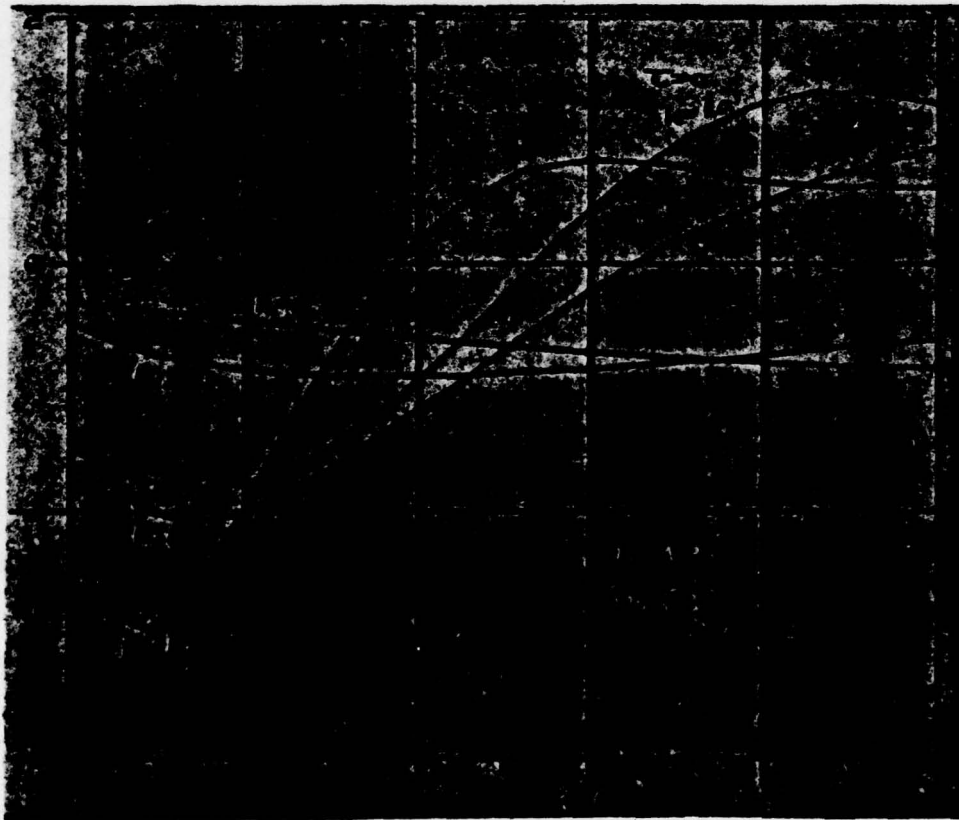


Fig. 23B. Pressure Distribution (Imaginary Part) from W I N G , Pitching Motion

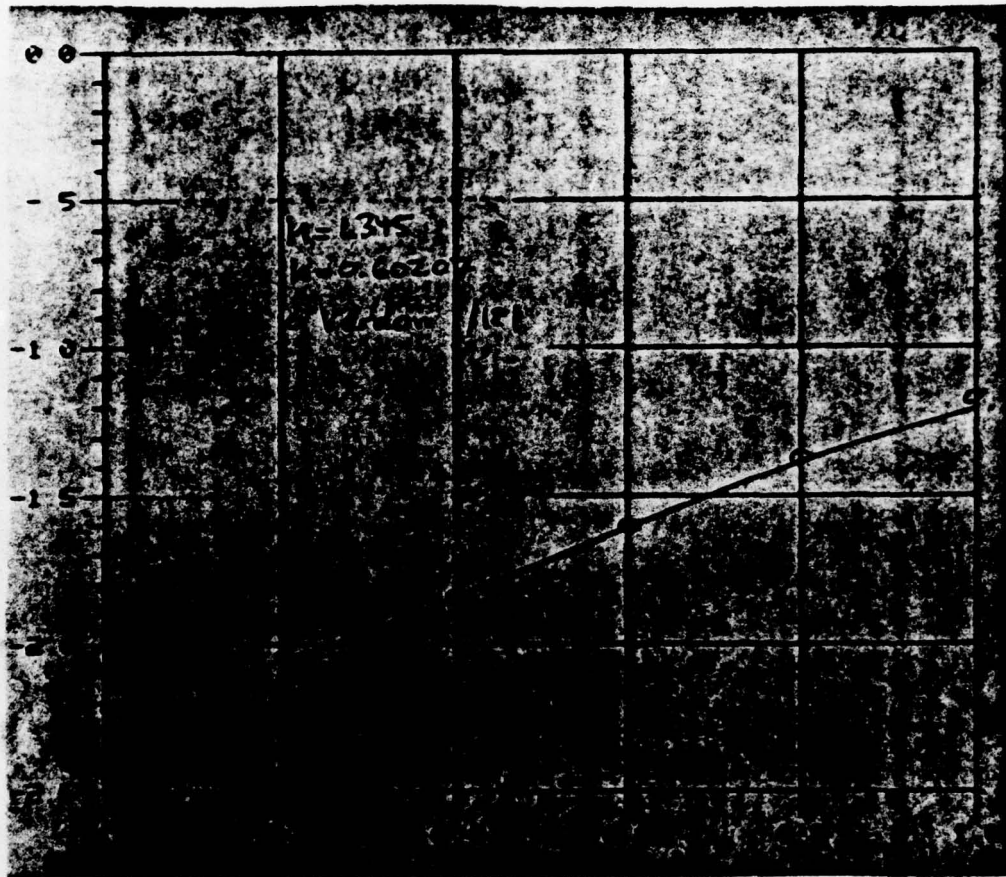


Fig. 24A. Pressure Distribution on Airfoil (Real Part) for Pitching Motion

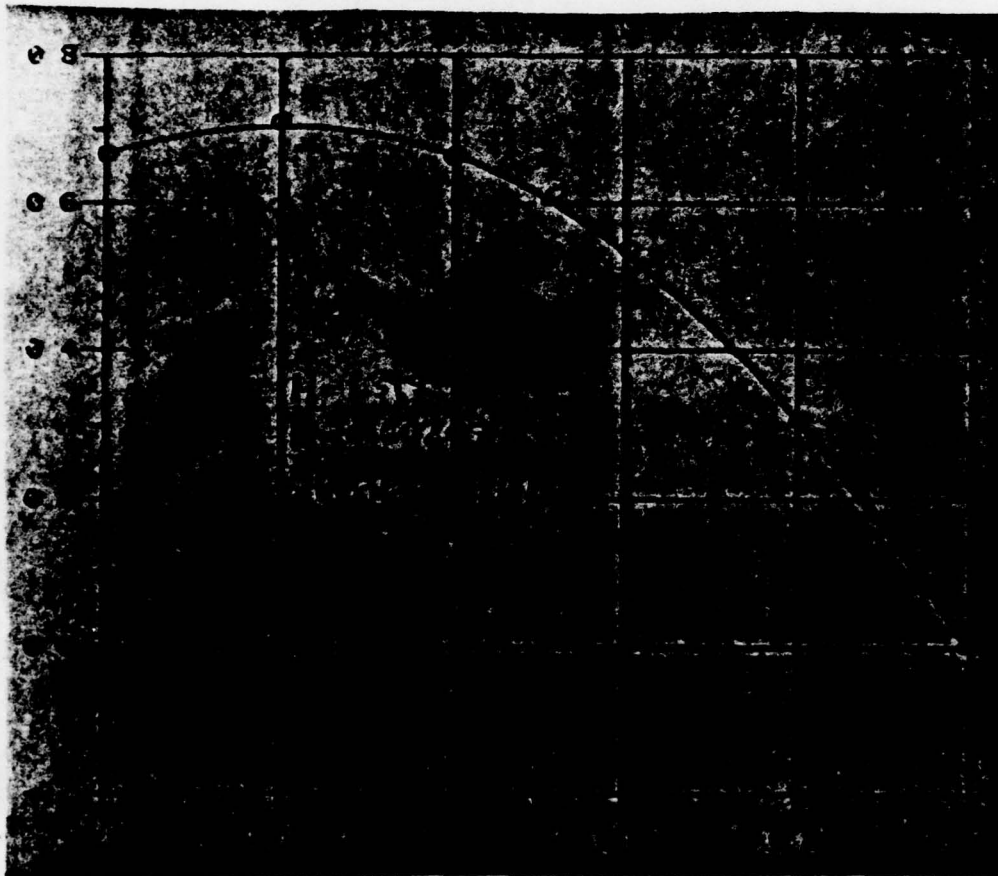


Fig. 24B. Pressure Distribution on Airfoil (Imaginary Part) for Pitching Motion

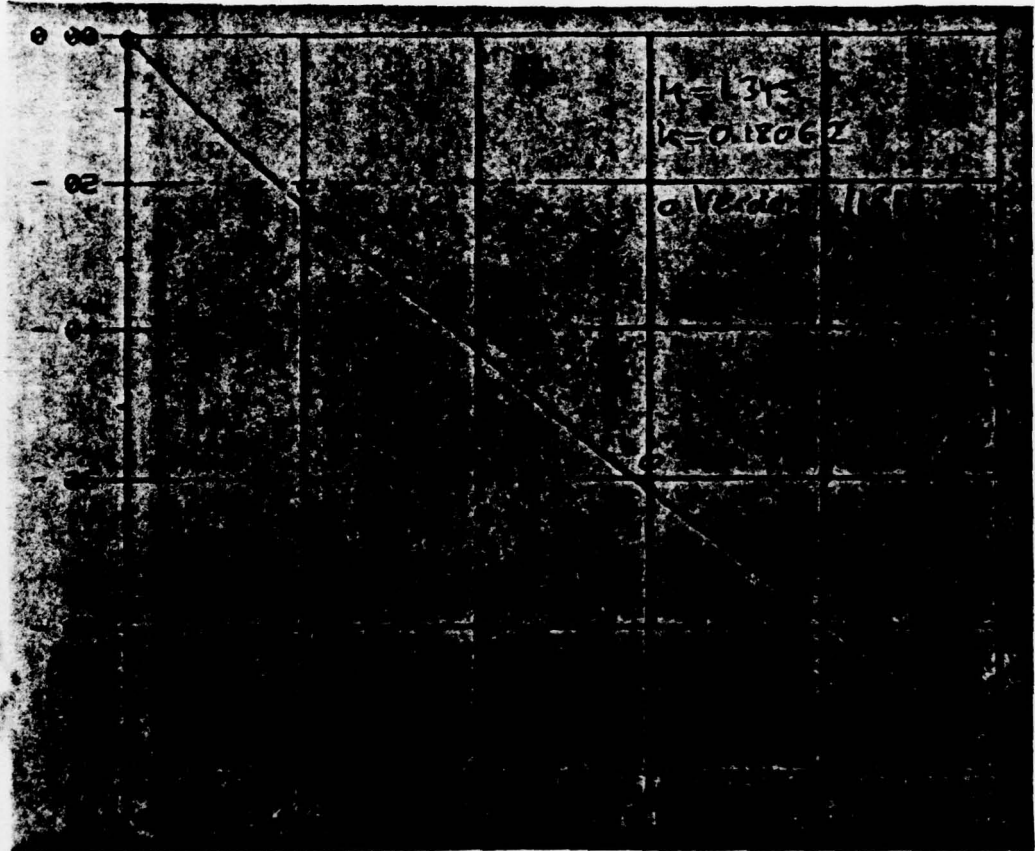


Fig. 25A. Pressure Distribution on Airfoil (Real Part) for Plunge Motion

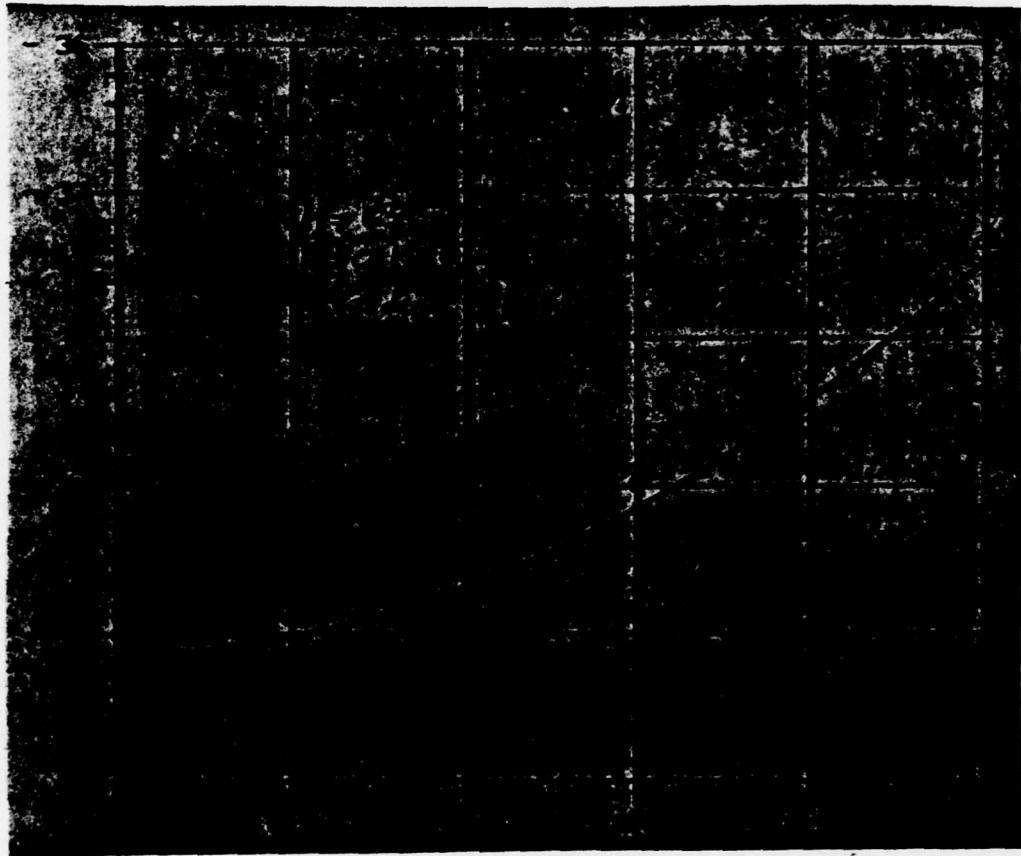


Fig. 25B. Pressure Distribution on Airfoil (Imaginary Part) for Plunge Motion

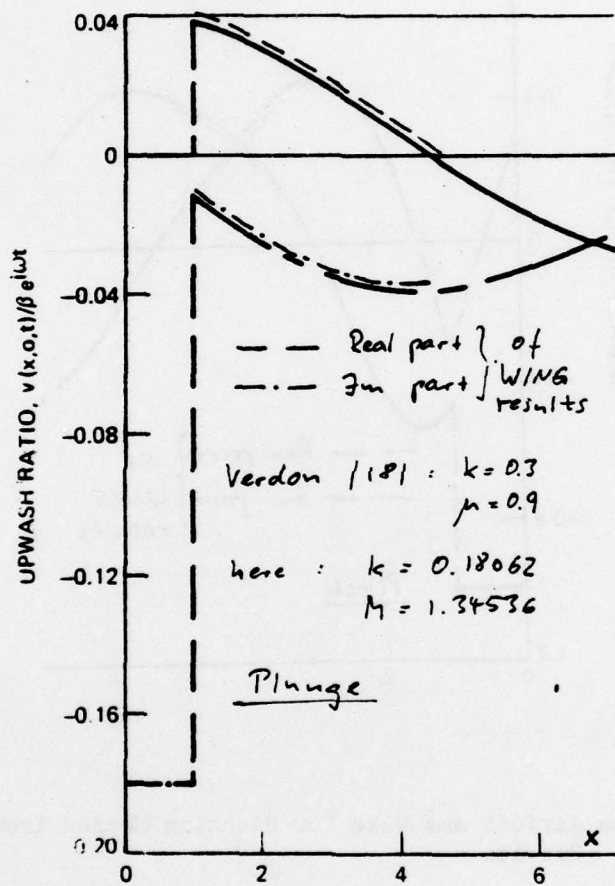


Fig. 26. Upwash on Airfoil and Wake for Plunge Motion from (18) Compared with W I N G -Results

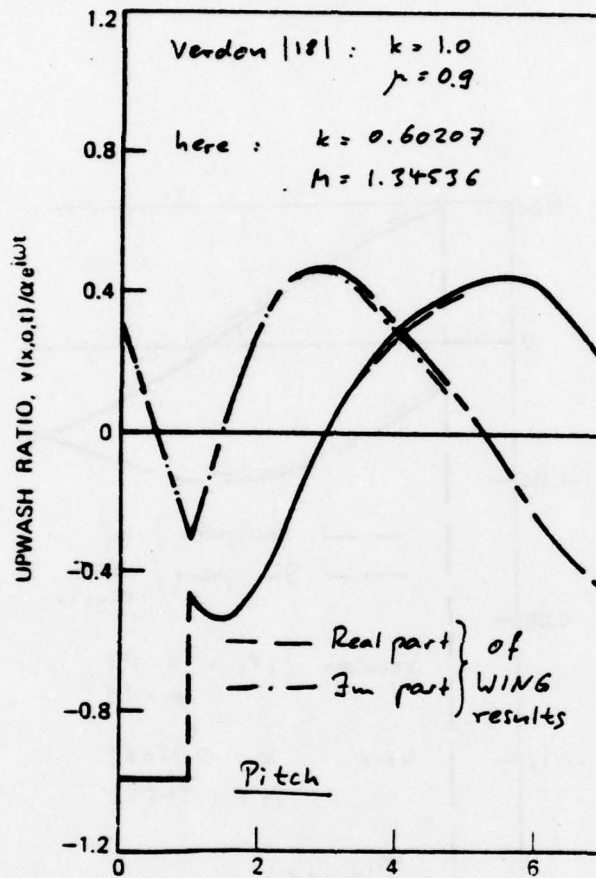


Fig. 27. Upwash on Airfoil and Wake for Pitching Motion from (18) Compared with W I N G -Results

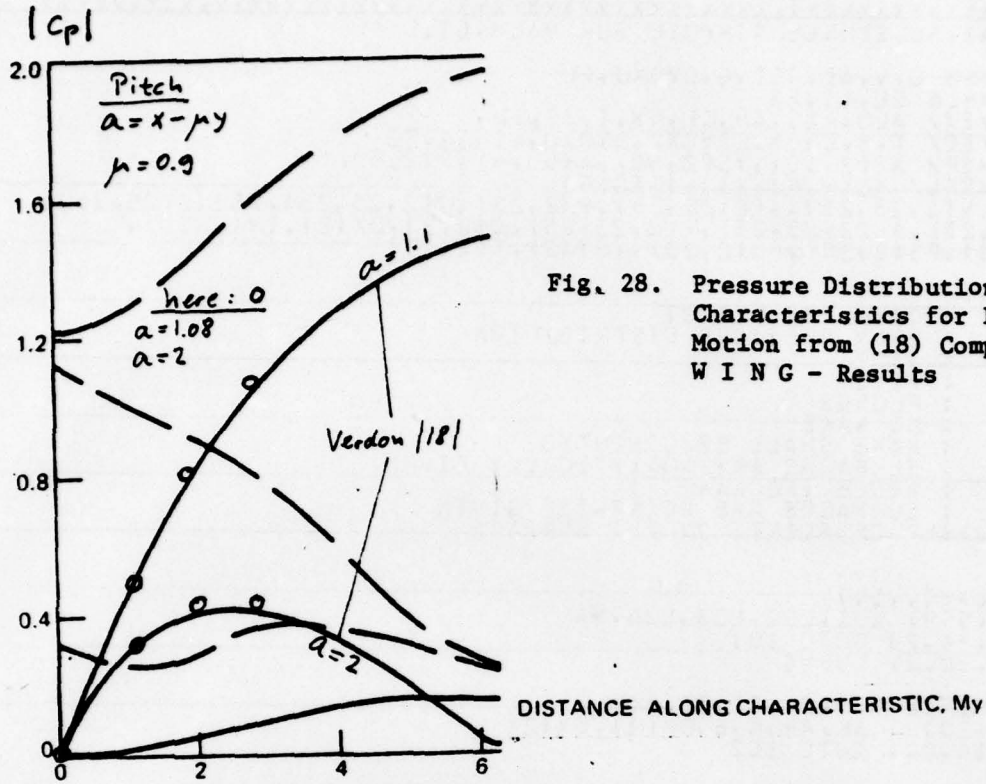


Fig. 28. Pressure Distribution Along Wake Characteristics for Pitching Motion from (18) Compared With W I N G - Results

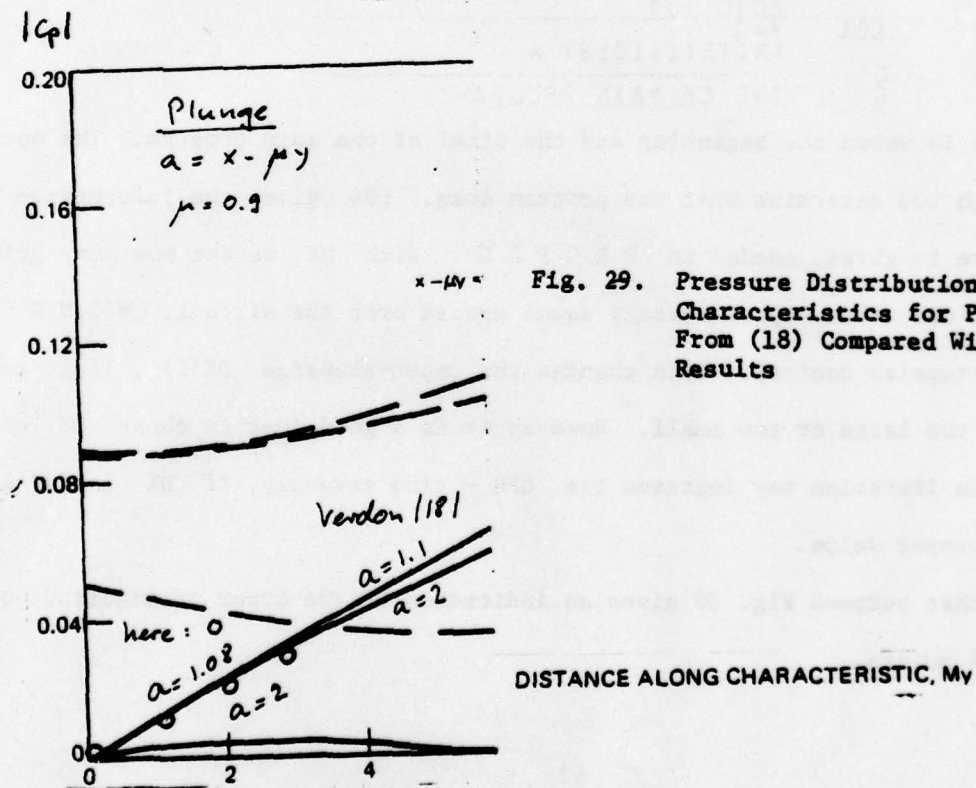


Fig. 29. Pressure Distribution Along Wake Characteristics for Plunge Motion From (18) Compared With W I N G - Results

### 3.8 Manual, Sample Data, Listing and Output of W I N G

```

CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
C
C OSCILLATING SINGLE AIRFOIL FOR MACH.GT.1
C
C GENFLX=8 L,V,AI,PSI,G,DYDXU,PU
C COMPLEX#16 EL,RI,ES
C COMMON/EA/ ALO,ALO,AM,C1,AK,I,IE,IW
C COMMON/EC/ T,B,CYCX,C2YDX2,DYDXU,AI,I1,T3
C COMMON/SP/ XS(2,50),YS(2,50),A(50,4),R(2,50)
C COMMON/SCL/ EL(4,4),FI(4),ES(4)
C COMMON V(2,25,25),X(2,25,25),P(2,25),U(2,25,25),PSI(2,25,25),
FG(2,25,25),AL(2,25,25),Y(2,25,25),C(2,3),DX(2),IN(2,25),
FPX(2,25),PS(2,20),PG(2,20),TE(2),TEE(20)
C
C OPTICS:
C LC1 =C : COMPLETE OUTPUT
C      =1 : ONLY PRESSURE DISTRIBUTION
C      =2 : STOP
C LC2 =C : PITCH
C      =1 : PLUNGE
C LC3 =C : NO WAKE
C      =1 : WAKE SHALL BE COMPUTED
C LC4 =1 : SURFACES ARE ANALYTICALLY GIVEN
C      =2 : WEDGE (NO WAKE)
C      =0 : SURFACES ARE POINTWISE GIVEN
C MA = NUMBER OF POINTS ON THE SURFACE
C
C KV=25
C AI=CMPLX(3.,1.)
102 READ(4,999) LO1,LO2,LO3,LO4,MA
C IF(LC1.EC.2) GOTC 101
C IF(LC4.EC.2) LO3=C
C MAX=MA
C CALL FRCFIL(WE,LC4,T1,T2,AX)
100 READ(4,1000) AK,AM,C,B,DX(1),DX(2)
C IF(AM.EC.C.) GOTC 102

```

```

C MAIN PROGRAM
C GOTC 100
C K=0
C WRITE(1,1016) K
C
C END OF MAIN PROGRAM

```

Above is shown the beginning and the final of the main program. The options LO1 through LO3 determine what the program does. LO4 gives the information how the surface is given, needed in P R O F I L . With MA we set how many gridpoints we want to distribute approximately equal spaced over the airfoil. W I N G has a built-in stepsize control, which changes the input-stepsize DX(I) , if it turns out to be too large or too small. However it is a good idea to chose DX correctly because the iteration may increase the CPU - time severely, if DX is far away from the proper value.

For that purpose Fig. 30 gives an indication of the order of magnitude of DX for 14 points.

The next read - statements are in P R O F I L :

```

SUBROUTINE PRCFIL(WE,LC4,T1,T2,N)
CCNCA/SP, XS(2,50),YS(2,50),A(50,4),R(2,50)
C
C PREPARATION OF THE PROFIL - SURFACES
C
READ(4,1000) T1,T2,NT,IKP
IF(LC4.EC.C) READ(4,1001) SP
IF(LC4.EC.2) READ(4,1001) WE
DO 6 J=1,2
12=0
N=NT
T=T1
IF(J.EC.2) T=T2
IF(LC4.NE.C) GOTC 12
DO 5 K=1,4
N=(K-1)*5+1
READ(4,1001) XS(J,M),XS(J,M+1),XS(J,M+2),XS(J,M+3),XS(J,M+4)
READ(4,1001) YS(J,M),YS(J,M+1),YS(J,M+2),YS(J,M+3),YS(J,M+4)
IF(YS(J,M+4).EC.100.) GOTC 7
5
CONTINUE
7
CONTINUE
ETC

```

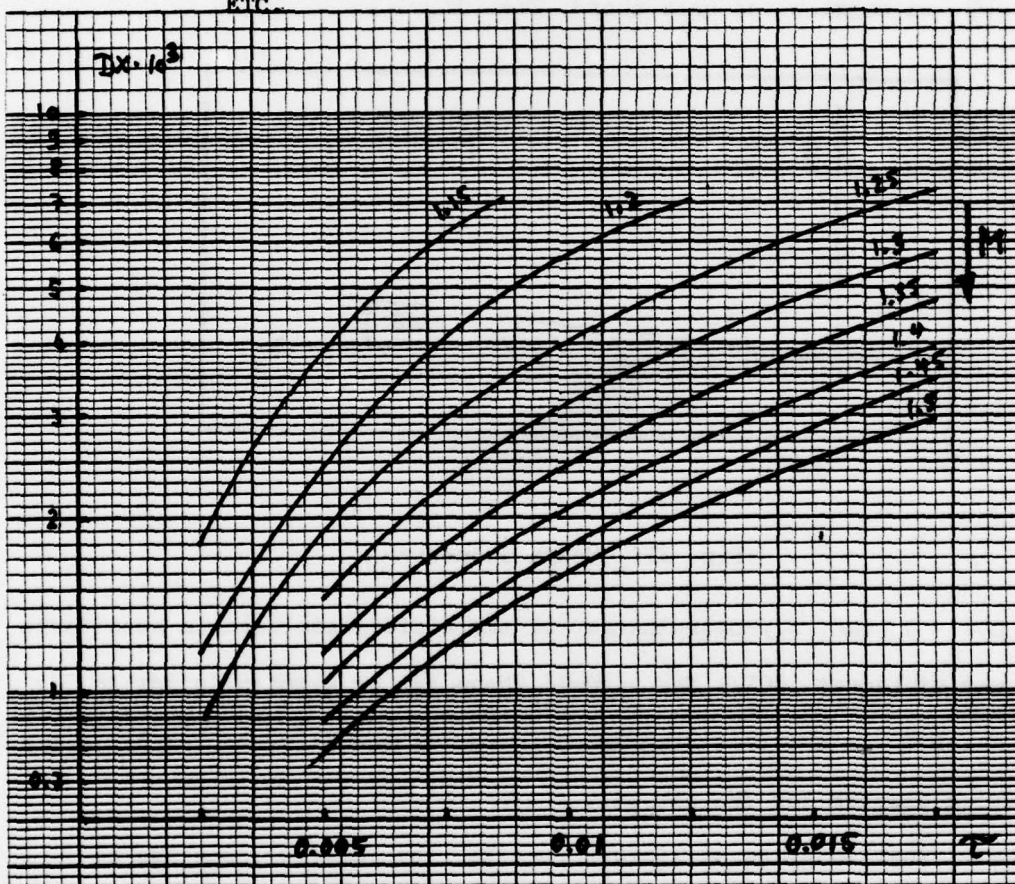


Fig. 30. Stepsize for N = 14 Points on the Airfoil

T1 and T2 represent the thickness of the airfoil. As W I N G is presented here, each surface is given by the parabola

$$y = 4tx(1 - x)$$

with

T = T1 for the upper

and

T = T2 for the lower side

For  $L04 = 0$  NT determines the number of points on the blade and IKP sets a flag to print out the input data.

IKP = 0 NO OUTPUT OF (X,Y)

= 1 OUTPUT

SP is a scaling factor to make the input data non-dimensional. If the chord is already 1, then  $SP = 1$ . Otherwise SP is the characteristic length, normally the chord.  $L04 = 2$  changes the airfoil into a symmetric wedge with the slope WE in degrees and unit chord, Fig. 31.

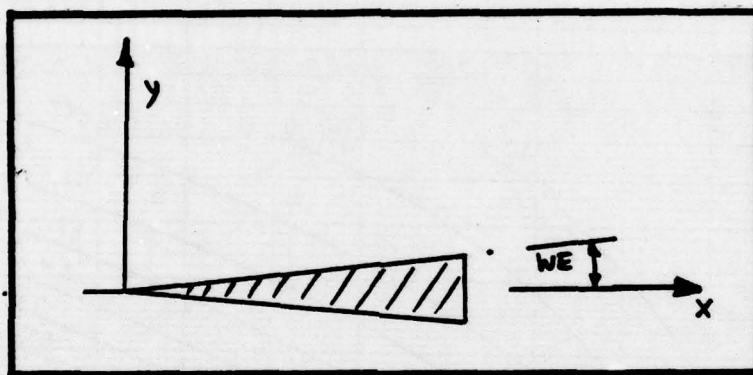


Fig. 31. Wedge Geometry

For  $L04 = 0$  the  $x,y$ -coordinates of the airfoil are read in accordance to the format statement 1001. The airfoil shape can be examined prior to entering WING in the test program TEST as shown in chapter 5. TEST is a preparing program for the airfoil.

If the airfoil is given by another function than Eq. (9), one has to change the statements in BOUND which compute  $y' = DYDX$  and  $y'' = D2YDX2$ .

After leaving P R O F I L there is another read for

AK = k

AM = M

C = ae

B =  $x_0$  = b

DX(1) = STEPSIZE ON THE UPPER SIDE

DX(2) = STEPSIZE ON THE LOWER SIDE

If the run is finished, W I N G jumps back to this line and expects changed data for this airfoil. For M = 0 it jumps to the start of the main program for mode and/or airfoil changes. However for L01 = 2 W I N G is finally terminated. This way a whole series of airfoil shapes and aerodynamic conditions can be examined in a single run. It should be noted that L01 = 0 produces a complete field - output which is normally not necessary and rather long.

For L03 = 1 the wake field output can not be suppressed as it is the only information which is printed about the wake. The whole wake may be skipped with L03 = 0 .

Besides the normal output file 06 W I N G has a second one: File 01 . Only the unsteady pressure distribution over the surfaces is written here. This file is used after the run by the plot program P L O T 1 to make plots of the pressure distribution on the Tektronix terminal in the computer center of the NPS. (S. Ch. 5).

W I N G contains still about 150 troubleshooting statements. All important subroutines start with

IKK = 0

The next statement is usually an if-statement converted into a comment card. If this is activated, one can choose here the conditions for a complete tracing of this particular subprogram. Thus

If (IW.NE.0) IKK = 1

would trace this program during the whole wake - computation. One can imagine that a general IKK = 1 will result in irresponsible amounts of paper.

SAMPLE DATA SETS :

AIRFOIL ANALYTICALLY GIVEN:

FILE: FILE	FTC4F001 P1			NAVAL POSTGRADUATE		
1101 16	0.01	20 0				
1.00	1.20	1.4	0.5	0.467E-02	0.467E-02	
0.60	1.20	1.4	0.5	0.467E-02	0.467E-02	
0.20	1.20	1.4	0.5	0.467E-02	0.467E-02	
0.30	0.00	1.4	0.5	0.467E-02	0.467E-02	
1101 16						
0.00	0.00	20 0				
1.00	1.20	1.4	0.5	0.467E-02	0.467E-02	
0.20	1.20	1.4	0.5	0.467E-02	0.467E-02	
0.0	0.0	0.0	0.0	0.0	0.0	
2222 22						

OPTIONS:

(Refer to W I N G)

- Only Pressure Distribution (1)
- Plunge (1)
- No Wake (0)
- Analytical Surface (1)
- 16 Points on each side

AIRFOIL DATA

For

T1 = T2 = 0.01  
M = 1.2  
= 1.4  
b = 0.5

DX (1) = DX (2) = 0.00467

The program runs through k = 1.0 , 0.6 and 0.2 . After M = 0 we want with the same options the linear case T1 = T2 = 0 for k = 1.0 and 0.2 . Finally we stop W I N G .

For the option-set 1001 the same data would be computed for the pitch mode. The sample output is the result for the first run of this data. Figs. 18b and 19b show the plots of the pressure distribution which we obtain with the data above. Figs. 22B and 23B show the same for the pitching mode.

AIRFOIL POINTWISE GIVEN :

FILE: FILE	FIT4FOJ1 P1				N	P
1000 14						
0.01	0.03	17 0				
100.						
0.0	5.21	11.53	17.84	24.14		
-0.28	-0.185	-0.110	-0.050	-0.035		
30.45	36.75	43.05	49.35	55.65		
0.029	0.056	0.074	0.087	0.097		
61.95	68.25	74.55	80.88	87.21		
0.008	0.095	0.08	0.0516	0.016		
0.035	0.10	100.	100.	100.		
0.0	6.29	12.59	18.90	25.21		
-0.26	-1.08	-1.85	-2.12	-2.49		
31.53	37.85	44.18	50.51	56.83		
-2.76	-2.93	-3.00	-2.97	-2.84		
63.15	69.47	75.78	82.06	88.36		
-2.60	-2.29	-1.94	-1.57	-1.17		
94.66	100.00	100.	100.	100.		
0.70	0.10	100.	100.	100.		
1.5	1.55	1.4	0.5	0.544E-03	0.454E-02	
0.0	0.0	0.0	0.0	0.0	0.0	
2222 22	C.C	C.C	C.C	C.C	C.C	

OPTIONS:

- ONLY PRESSURE DISTRIBUTION
- NO PITCH
- NO WAKE
- POINTWISE GIVEN SURFACES
- 16 POINTS ON EACH SIDE

IN PROFIL :

- T1 = 0.01 NOT IMPORTANT
- T2 = 0.03
- N = 17 POINTS GIVEN FOR EACH SIDE
- IKP = 0 NO READ-BACK OF DATA
- SP = CHORD = 100

This example shows the different input-type for  $L04 = 0$ . The profil thickness  $T1$  and  $T2$  is not important any longer. But the variables have to be defined because they appear in the output. They can be used for identifying purposes.

Here  $x$ -location and  $y$ -value are given in percent of chord, therefore  $SP = 100$ .

PLUNGE - MCEE

W\*C/U= 1.000, M= 1.20, K= 1.40, B/C= 0.50, DX= 0.467E-02, T/C= 0.0100

PRESSURE-DISTRIBUTION UPPER SURFACE:

PCINT	X	CPS	RCPU	ICPU
2, 1	0.000	0.143E 00	0.0	-0.455E 01
3, 2	0.043	0.128E 00	-0.111E 01	-0.402E 01
4, 3	0.090	0.113E 00	0.181E 01	-0.326E 01
5, 4	0.140	0.968E -01	0.209E 01	-0.235E 01
6, 5	0.193	0.807E -01	0.203E 01	-0.159E 01
7, 6	0.251	0.643E -01	0.169E 01	-0.957E 00
8, 7	0.312	0.475E -01	0.123E 01	-0.621E 00
9, 8	0.378	0.304E -01	0.725E 00	-0.492E 00
10, 9	0.448	0.128E -01	0.332E 00	-0.604E 00
11, 10	0.522	-0.531E -02	-0.381E -01	-0.821E 00
12, 11	0.601	-0.239E -01	-0.133E -01	-0.116E 01
13, 12	0.685	-0.429E -01	-0.787E -01	-0.141E 01
14, 13	0.774	-0.625E -01	-0.112E 00	-0.158E 01
15, 14	0.869	-0.826E -01	0.297E 00	-0.163E 01
16, 15	0.969	-0.103E 00	0.501E 00	-0.157E 01
17, 16	1.000	-0.110E 00	0.564E 00	-0.155E 01

W\*C/U= 1.000, M= 1.20, K= 1.40, B/C= 0.50, DX= 0.467E-02, T/C= 0.0100

PRESSURE-DISTRIBUTION LOWER SURFACE:

PCINT	X	CPS	RCPU	ICPU
1, 2	0.000	0.143E 00	0.0	0.455E 01
2, 3	0.043	0.128E 00	0.111E 01	0.402E 01
3, 4	0.090	0.113E 00	0.181E 01	0.326E 01
4, 5	0.140	0.968E -01	0.209E 01	0.235E 01
5, 6	0.193	0.807E -01	0.203E 01	0.159E 01
6, 7	0.251	0.643E -01	0.169E 01	0.957E 00
7, 8	0.312	0.475E -01	0.123E 01	0.621E 00
8, 9	0.378	0.304E -01	0.725E 00	0.492E 00
9, 10	0.448	0.128E -01	0.332E 00	0.604E 00
10, 11	0.522	-0.531E -02	-0.381E -01	0.821E 00
11, 12	0.601	-0.239E -01	-0.133E -01	0.116E 01
12, 13	0.685	-0.429E -01	-0.787E -01	0.141E 01
13, 14	0.774	-0.625E -01	-0.112E 00	0.158E 01
14, 15	0.869	-0.826E -01	0.297E 00	0.163E 01
15, 16	0.969	-0.103E 00	0.501E 00	0.157E 01
16, 17	1.000	-0.110E 00	0.564E 00	0.155E 01

MOMENTUM- AND LIFT - COEFFICIENTS FOR A SINGLE AIRFOIL WITH A SURFACE DESCRIBED ABOVE:

	RCL	ICL	RCM	ICM
UNSTEADY	-2.0585	1.2519	0.2039	-0.2964
STEADY	0.0	0.0	0.0	0.0

#### 4. The Oscillating Finite Cascade: C A S C A D E

##### 4.1 The Physical Difference of the Two Problems

The name of the program, developed to calculate the inlet flow of a staggered finite cascade with thick airfoils is C A S C A D E .

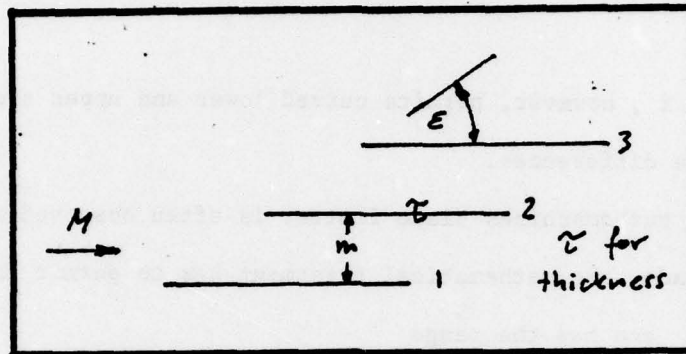


Fig. 32. Staggered Cascade of Airfoils

Nearly all basic steps to simulate this problem are already done in W I N G and therefore described in Chapter 3. However, there are two major differences which have to be considered in its mathematical treatment.

1. Only the first blade is exposed to the free stream flow without any disturbances. This means for the cascade of flat plates already that the perturbations of  $(n - 1)$  blades hit the  $n$ -th blade and have influence on the development of the flow and the shocks at this airfoil. For the cascade of curved blades there is an additional consequence:

The constant value

$$C_{\infty} = (M_{\infty}^2 - 1)^{3/2}$$

from Eq. (2.15) can not be used any longer, as in general the steady flow will not have the free stream velocity.

Similarly, the wake slip line will not be parallel to the x-axis but it will rather have the direction of the field in front of that particular blade. This is the reason why the wake slope in W I N G was not simply set to zero. Because of these difficulties the first approach to the cascade including a thickness effect was made with airfoils whose upper surface were flat. This way the steady field in front of each new blade was identical with the free-stream field (see /5/).

C A S C A D E , however, permits curved lower and upper sides and is able to consider these differences.

2. As in actual turbomachines blade flutter is often observed with a phase lag from blade to blade, the mathematical treatment has to permit this. The phase lag is called  $\mu$  and has the range

$$- 180^\circ \leq \mu \leq + 180^\circ$$

which covers all cases.

The introduction of the phase lag can be done relatively easily and is shown in Section 4.3.

#### 4.2 The Constant Value Along the Characteristics

We consider Eq. (2.13) :

$$\lambda^{3/2} + \mu = C_{\alpha,\beta}$$

(where in this case  $\mu$  is not the phase lag)

as long as the incoming characteristics originate in the freestream,  $\mu$  is zero and therefore

$$C_{\alpha,\beta} = (M^2 - 1)^{3/2}$$

(Fig. 1)

This still holds in a cascade of flat plates, because here the steady field is identical with the freestream field. This is also true for a cascade with

blades which have flat upper surfaces.

It is not valid any longer when we permit curved surfaces on both sides, as the steady inlet flow field for the blade ( $n \geq 2$ ) will in general be deflected. Therefore  $\mu$  on the characteristics is not zero and has influence on  $C_{\alpha, \beta}$ .

Fig. 33 shows the geometry of an incoming characteristic which is reflected into the cascade.

In A it has to be

$$C_{\infty} = \lambda_{\infty}^{3/2} = (\lambda_A^{3/2} + \mu_A) B \quad (1)$$

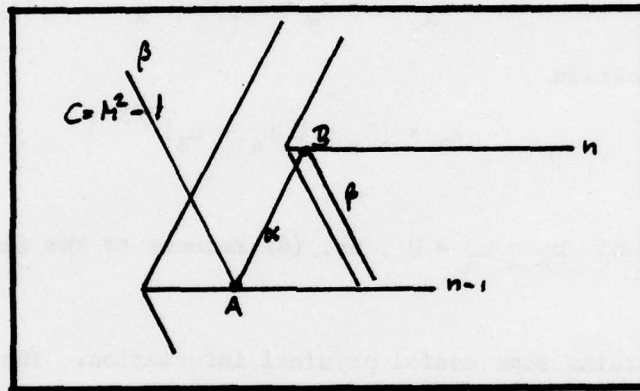


Fig. 33, Reflection of Characteristics

Thus we get  $\lambda_A$  by solving for it, as  $\mu_A$  is known from the steady boundary conditions.

The next step provides the new constant along the  $\alpha$ -characteristic.

$$C_{\alpha\alpha} = \lambda_A^{3/2} - \mu_A \quad (2)$$

This is now the value which has to be applied in B in order to solve for  $\lambda_B$ .

$$C_{A\alpha} = C_{B\beta}$$

$$\lambda_A^{3/2} - \mu_A = \lambda_B^{3/2} - \mu_B$$

$$\lambda_B = [\lambda_A^{3/2} - \mu_A + \mu_B]^{2/3} \quad (3)$$

with

$$\lambda_A^{3/2} - \mu_A = C_\infty - 2\mu_A$$

From Eq. (1) we obtain

$$\lambda_B = [C_\infty - 2\mu_A + \mu_B]^{2/3} \quad (4)$$

In the case of  $\mu_A = \mu_B = 0$ , Eq. (4) reduces to the flat-plate solution  $\lambda_B = \lambda_\infty$ .

Eq. (4) contains some useful physical information. The maximum deflection  $\delta_{\max}$  which is connected to each Mach number M cannot be used for the leading-edge-slope of an airfoil in a cascade. Instead it has to be considerably smaller due to the double deflection in A and B. Only in case of a flat upper surface (5), the slope in B could be  $\delta_{\max}$ . The general case,  $\mu_A$  and  $\mu_B$  not zero, demands smaller deflections if detached leading edge shocks shall be avoided. This is a necessity according to our earlier assumption of weak shocks.

The calculation of the particular constant value, valid along the respective considered characteristic is done in C O N S T 1 .

### 4.3 The Phase Lag

Each time the unsteady properties  $u_1$ ,  $v_1$  or  $\Psi$  are computed, the result is always the amplitude of the oscillating function

$$F(x,y,t) = A(x,y) \cdot e^{ikt}$$

Therefore, if we consider the two neighbouring blades (n-1) and n, where n leads the oscillation with the phase lag  $\mu$ , we can express this for the time t as follows

$$F_n = A_n \cdot e^{ikt}$$

$$F_{n-1} = A_{n-1} \cdot e^{i(kt-\mu)}$$

For the single airfoil we did not need to write the exponent expression because it was always a common factor.

We reconsider this. For an example we take Eq. (2.40) which is the used unsteady shock polar. The first of those two formulas reads now, if written explicitly

$$\hat{u}_1 \cdot e^{ikt} = (m_1 G_y + i m_2 G) \cdot e^{ikt} + (m_3 u_1 + m_4 v_1) e^{i(kt-\mu)} \quad (5)$$

$$\hat{u}_1 = m_1 G_y + i m_2 G + (m_3 u_1 + m_4 v_1) e^{-i\mu}$$

We see that the phase lag does not cancel out. If we do this for all the equations used in W I N G, we find as a general rule: the term  $\exp(ikt)$  always can be dropped. Each time when properties of the previous blade appear in the equation, they have to be multiplied by  $\exp(-i\mu)$ . Thus we reduce the magnitude of the (n-1)st -blade-values with the phase lag to the actual size they have

when they influence the  $n$ th -blade-properties. These are the unknown amplitudes of the oscillating functions connected to the  $n$ th -blade.

Eq. (5) gives an example how the equations change from W I N G to C A S C A D E . It is not necessary to show the whole set of finite difference formulas because they can be generated simply by using the concept outlined above.

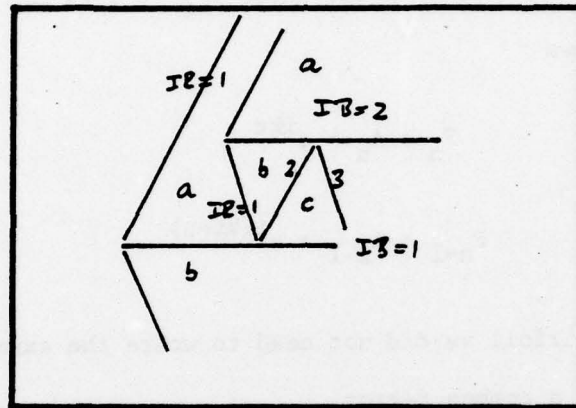


Fig. 34. Reflected Shocks in the Passage

The subroutines for the general and the boundary step can be taken as they were used in W I N G for the unsteady field. R A N D S is changed in so far, as the phase lag has to be added to the system of equations. The steady field is not affected.

For the field behind the reflected shocks in the passage one has to be careful, see Fig. 34.

It has not only to be considered that the oscillation of field  $b$  is ahead of that in field  $a$ , according to the phase lag  $\mu$ . But also it is ahead of field  $c$ , which is again connected to the movement of blade 1. This is done by reversing the sign of  $\mu$  in R A N D S for this case.

#### 4.4 The Organization of the Program

As the computation of the reflected shocks and the fields behind them in the passages of a cascade is considered a rather complex procedure, the goal for the design of the program was to keep it as straightforward as possible. Therefore the convenient array organization from W I N G was adopted and extended. The main variable fields have the form

$$X( IB,IR,M,N )$$

The part (... ,M,N) corresponds to Fig. 11 and 12. It allows the separation of fields on the upper and lower surfaces. IB indicates the blade, which is connected to the respective field and IR counts the shock in the passage, Fig. 34. This notation allows again repeated use of the same routines when only IB and IR are set correctly. Again throughout the whole program upper and lower sides are indicated by I=1 and I=2. The index I is used to control the M,N-notation and the sign like it was done in W I N G.

In C A S C A D E there appear three different kinds of flow fields which are shown in Fig. 34.

Type a is the field over the whole upper surface. The procedure is here the same as in W I N G. Computation of the b-field follows also W I N G, but it is stopped when the shock crosses blade 1. At this point the shock is reflected and the main difference in the subroutines is a transformed system of coordinates, because the origins of the shock and the system of coordinates have to be identical. This transformation causes considerable changes of the statements in S H O C K and F I N D, compared to their counterparts in W I N G. However, the computational sequence is not really changed, but the extension of those subroutines enables them to identify the type of the field and to make all necessary changes for signs, coordinates and termination points.

C A S C A D E starts with blade 2 (IB=2) exposed to an initial field where all perturbations are zero. After evaluating the complete upper field, the b -field is computed. This is terminated when the shock hits blade 1, which is assumed to be a fixed flat plate for this first step. So here blade 1 could be considered a wind tunnel wall. When all the reflections in this first channel are done, the IB=2-field is copied to the IB=1-field. Now the former can be used again. From this point on blade 1 is an oscillating airfoil with the given shape.

C A S C A D E has three output files:

- File 07 for complete field output and/or the final pressure distribution
- File 06 is a documentation which allows to follow the iterations and the main steps from blade to blade. Like W I N G, C A S C A D E contains trouble shooting statements which can be activated by IKK=1 for the desired subprogram. For a general IKK=0 File 06 will be limited to a few pages. IKK=1 will print a complete tracing of the particular routine, which increases File 06 significantly.
- File 01 contains geometric data of the cascade and shock configuration.

It is used as an input file for P L O T 2 which produces plots like Fig. 35 on the Tektronix - terminal. P L O T 2 is described in Section 5.3. The input data have to be in File 03. They are shown in Section 4.6.

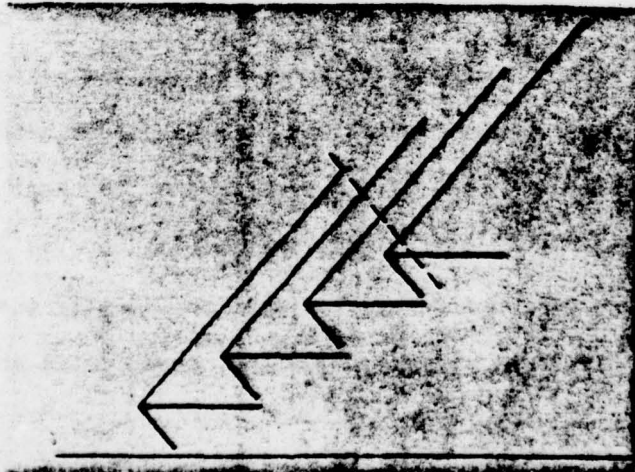


Fig. 35. Cascade A and B Shock Geometry [15]

Finally, it should be said that the organization of the storage arrays is actually a waste of memory space because fields of the b - and c - types (Fig. 34) need only small areas whereas a -fields demand large ones. This expensive way was chosen in order to make a general program clear. If one considers that the next step is the addition of the wake fields for each blade and that these have to be much larger than those in the passages, this solution gives an approach which can be extended analogous to W I N G . The desire to save storage place would certainly complicate the main organization and increase the number of subroutines significantly.

The main disadvantage, besides a slow-down of the computer is the limitation of the fields. The dashed line in Fig. 35 was not produced by P L O T 2 but it was added to show the limit of C A S C A D E . Beyond the end of the headshock for the first blade its influence has vanished. This results in completely wrong fields behind the dashed line. Hence the fields in Fig. 35 are usable only up to blade 3. The arrays in C A S C A D E used in this work have the size

$$X(2,3,50,20)$$

They can be extended to

$$X(2,3,KV,20)$$

where KV has to correspond with the first FORTRAN-STATEMENT of C A S C A D E . Then all the counters and Do-loops are dimensioned properly. The number of possible blade-computation is thus given by the limitation of the computer.

#### 4.5 Results

As a test for C A S C A D E three different supersonic cascades are checked for which results already exist. Cascade A and B are those introduced by Verdon (15;17) and cascade T is a model for an experimental cascade used by Fleeter (19) for which Strada (5) gives computer results for the inlet flow.

##### Cascade Data:

	A	B	T
M	1.345	1.281	1.550
m	0.4	0.301	0.335
$\tau_{up}$	0.0	0.0	0.0
$\tau_{Lo}$	0.0	0.0	0.03
$\epsilon$	30.5	26.6	23.9
k	0.90226	0.75054	0.28

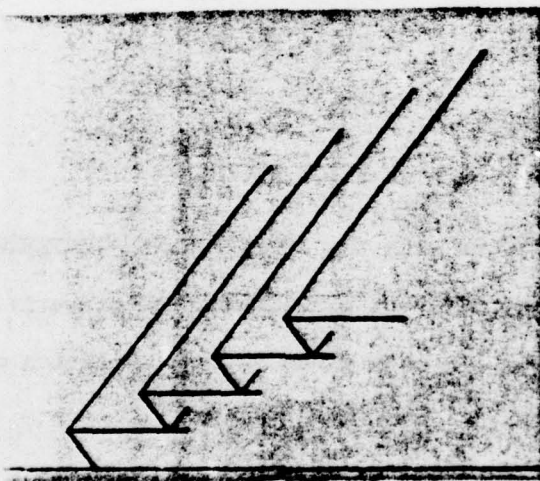


Fig. 36. Cascade B (15)

Fig. 35 shows the geometry of cascade A . Cascade B and T are shown in Fig. 36 and 37. As C A S C A D E is not completely programmed, it is only possible to show results of the inlet flow. We compare with theoretical aerodynamic data obtained by Verdon (15) for the linear cases, which are identical with those computed by Bell (9) and with results for the cascade T , given by Strada in (5).

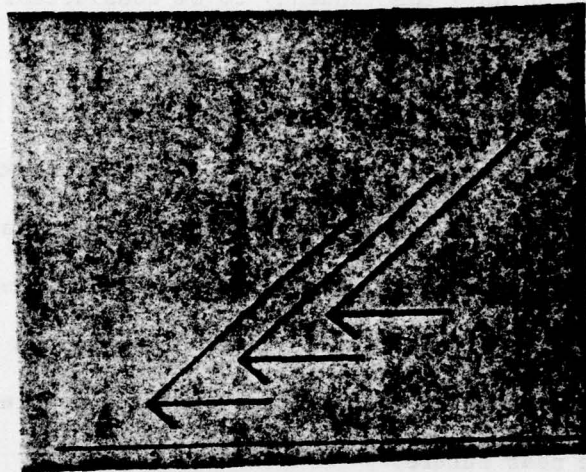
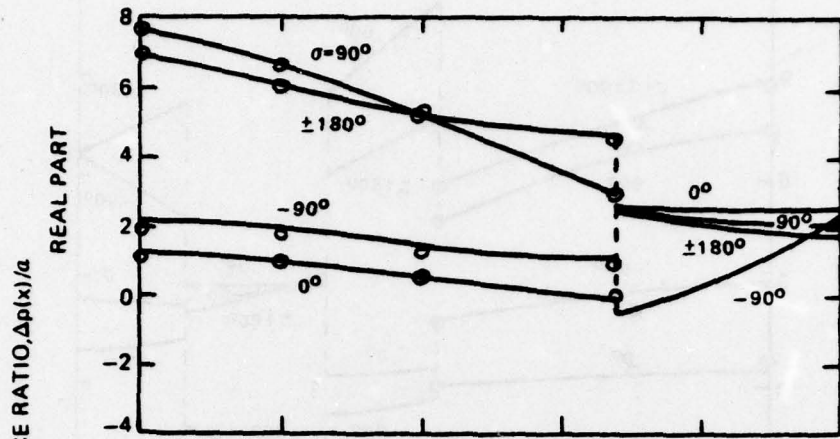


Fig. 37. Cascade T (5)

Fig. 38 and 39 show pressure difference distributions calculated by C A S C A D E for the third blade. Considering that Verdon introduced an infinite cascade and that Bell gave results for the 14th blade, the inlet flow computed with C A S C A D E looks very promising. The differences are not too significant and can be explained by the low number of blades which were examined.

As C A S C A D E and Verdon-results agree sufficiently, it could be expected that the linear data given by Strada and Platzer ( P-) results would coincide very well with C A S C A D E computations. Actually, as both programs examine the second blade, the distribution of the total pressure coefficient agrees exactly for both sides of the A blade.

This is not true for the included slope effect. Fig. 40 through 43 show computations of this study in direct comparison with Strada (5). The upper surface represents the linear case and shall not be discussed because of good agreement. However, the lower surface is shaped and the agreement is good only at the leading edge. In all cases results of C A S C A D E for the total pressure distribution have the tendency to be considerably smaller than those presented by Strada. There is no real progress compared to the Fleeter Experiments (19) but with some imagination one could say that C A S C A D E meets the tendency of those data slightly better. It is remarkable that in spite of significant differences in the magnitude, the phase angle between imaginary and real part of the pressure coefficients agrees very well with Strada. The difference between the Strada results and those obtained here may be caused by the different treatment of the shock and remains a problem to be investigated more thoroughly in the near future.



o CASCADE - Results

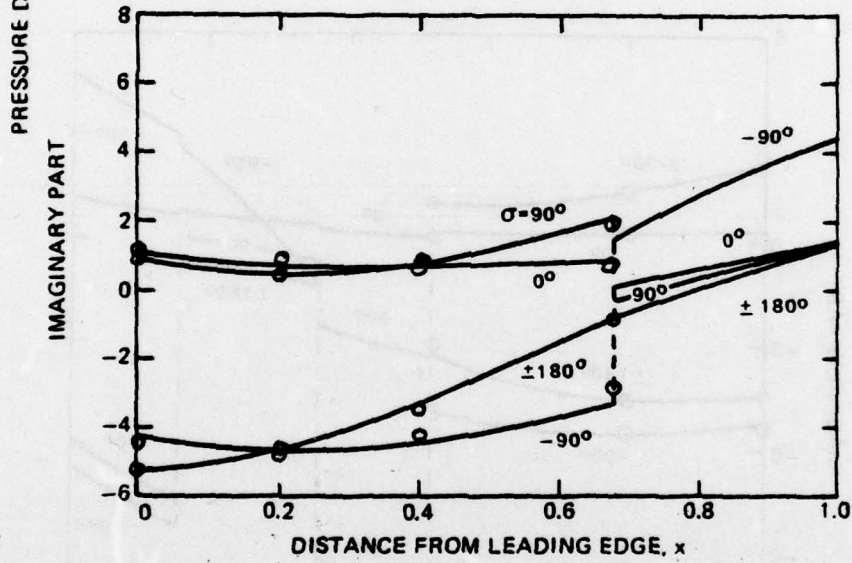


Fig. 38. From (15) Comparison of Linear Results Casc. A

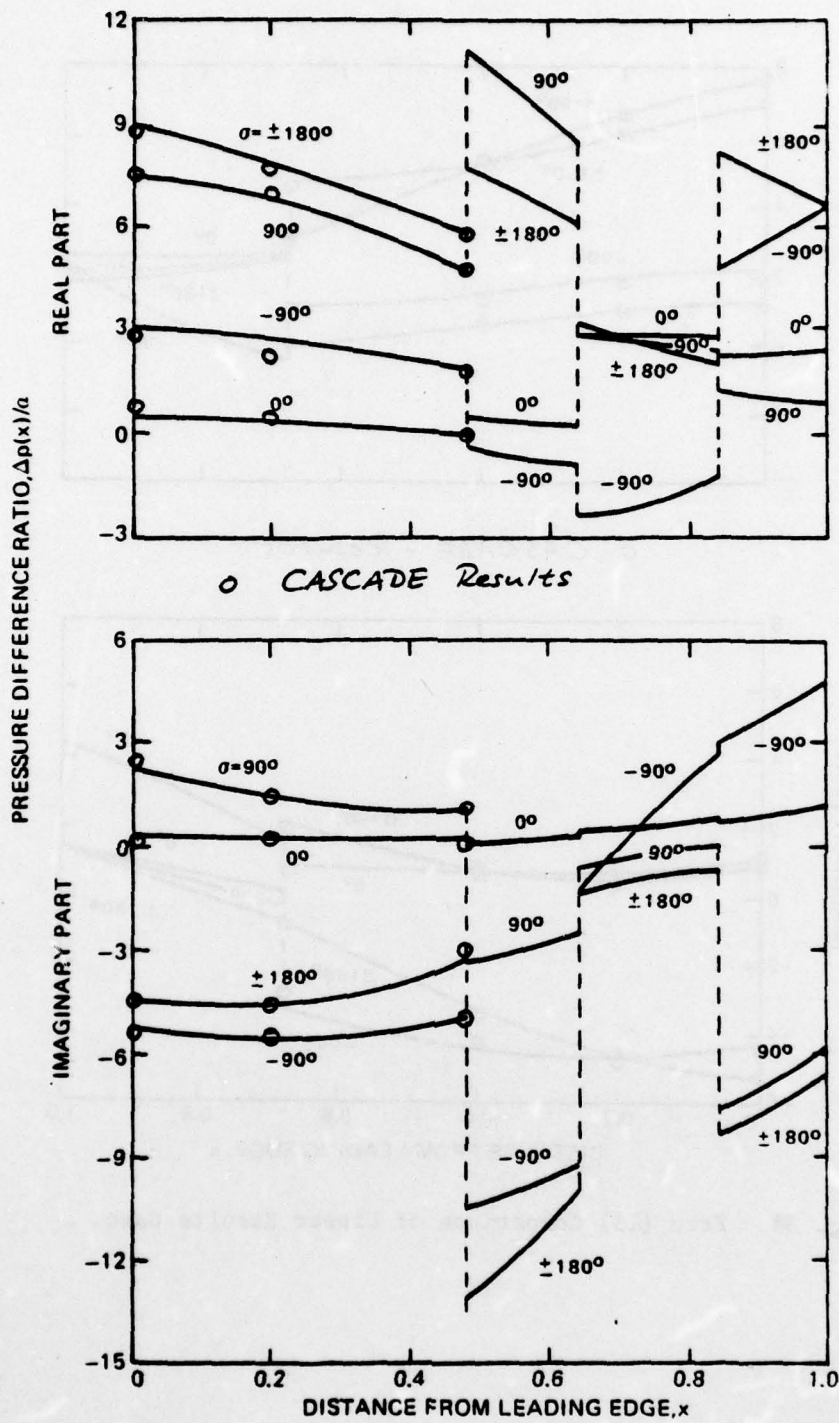


Fig. 39. From (15) Comparison of Linear Results Casc. B

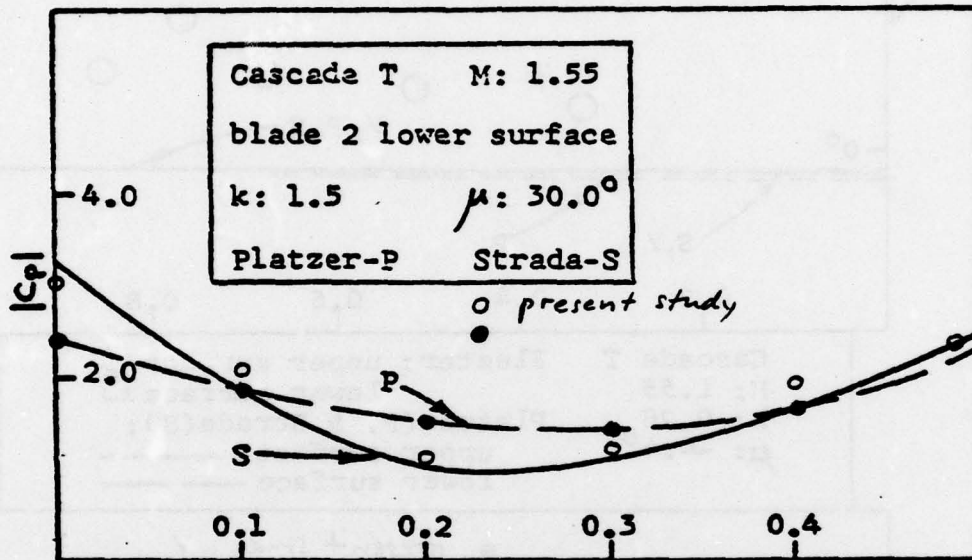


Figure 3.10.2 High-k comparison of the linear and non-linear solutions.

Fig. 40. From (15) Comparison with Strada

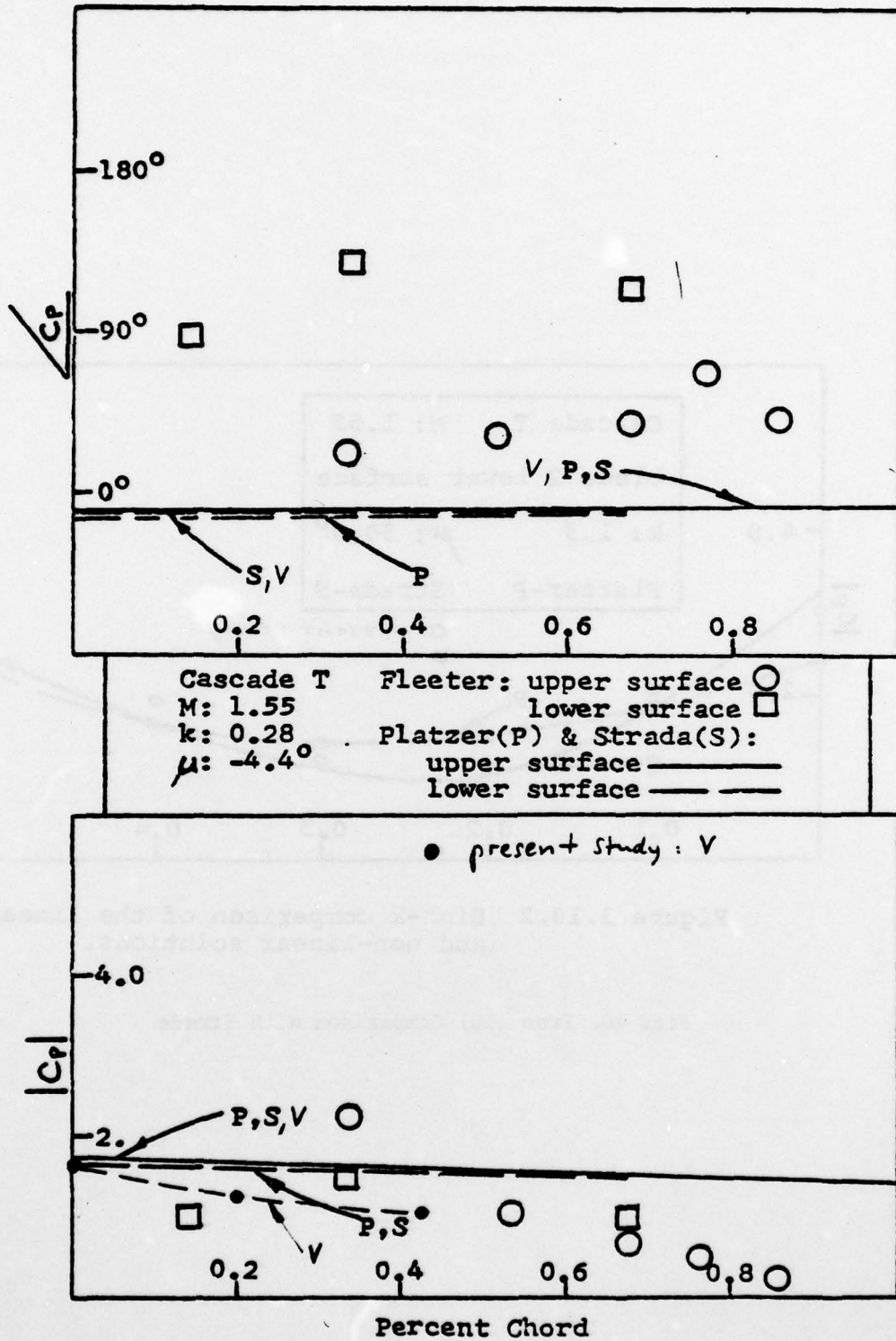


Fig. 41. From (5) Comparison with Strada

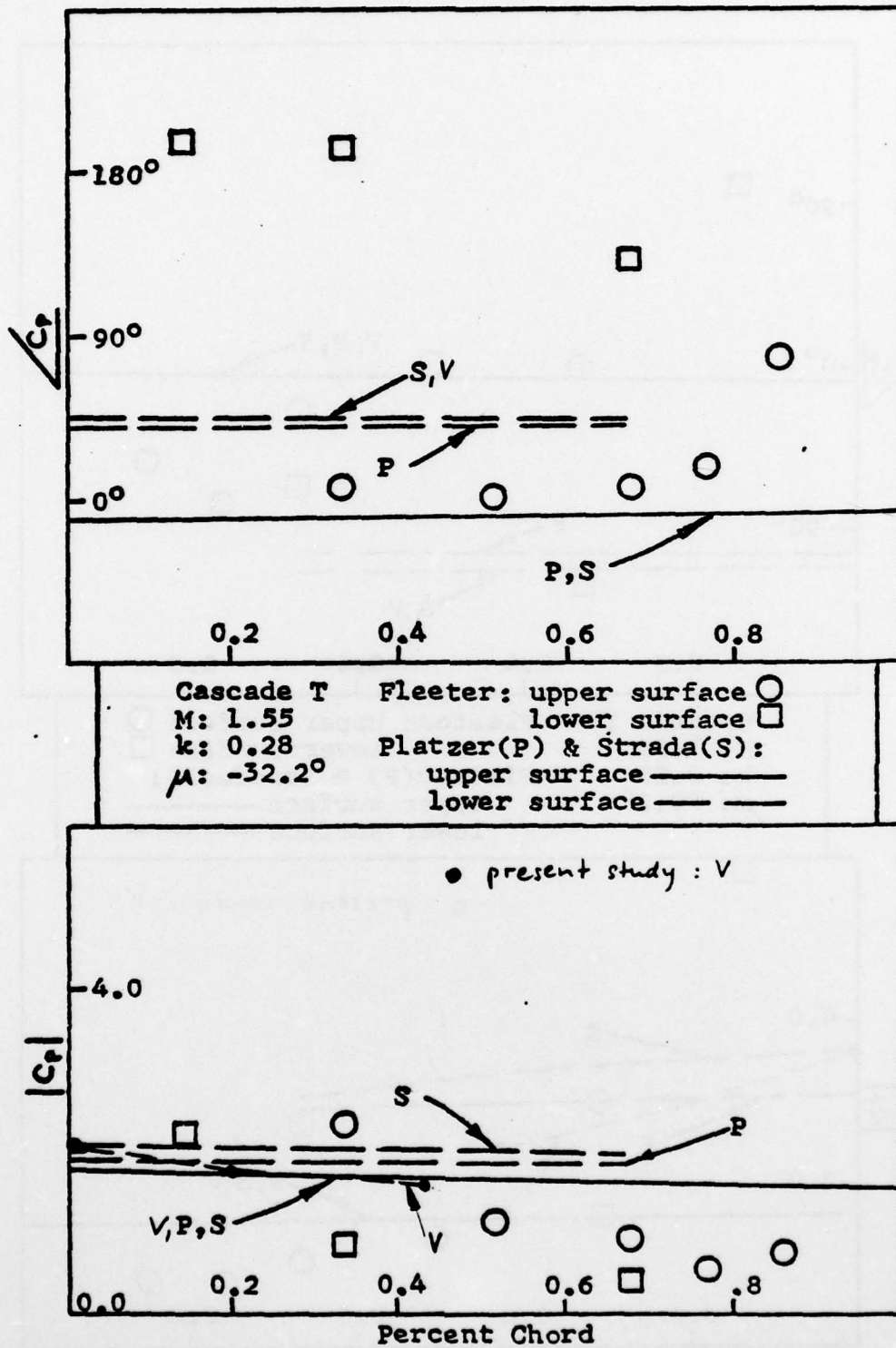


Fig. 42. From (5) Comparison with Strada



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A METHOD OF CHARACTERISTICS APPROACH TO THE PROBLEM OF SUPERSON--ETC(U)

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2 OF 2

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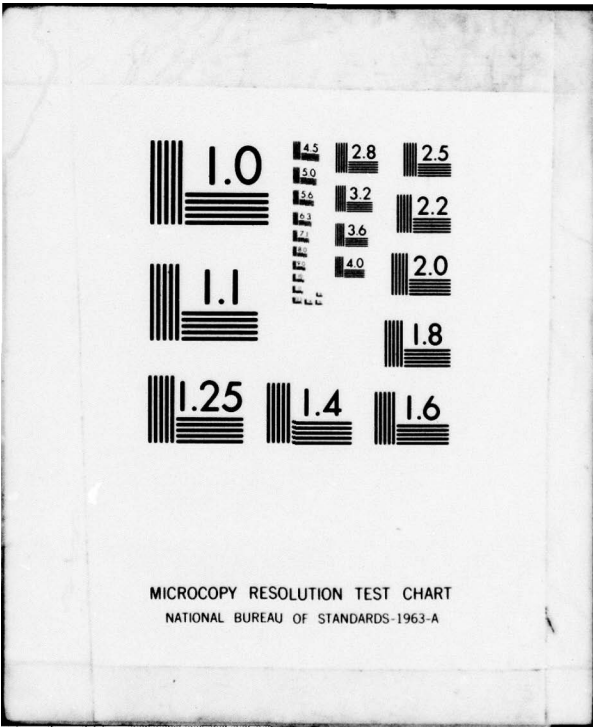


END

DATE  
FILMED

12-79

DDC



MICROCOPY RESOLUTION TEST CHART  
NATIONAL BUREAU OF STANDARDS-1963-A



The other three variables define the Cascade:

ET =  $\epsilon$  Stagger Angle  
EM = m Distance Between the Blades  
AMY =  $\mu$  Phase Lag

See Fig. 32.

The rest of the input data have the same definition as those in W I N G .

Two sample data sets shall be shown:

1. Analytically Given Blades

```
FILE: FILE      FTC3EQ01 P1      NAVAL P.  
10C1160301  
0.000 0.00 2) 0  
0.75054 1.281 1.4 0.392E-02 0.392E-02  
0.0 0.5 26.6 0.351  
0.28000 0.000 1.4 0.392E-02 0.392E-02  
222222222
```

This would cause a run as follows.

- 1st Row : No field output, pitch motion, no computation between the first blade and the imaginary wall (first channel), 16 grid points on each surface, they are analytically given, the third blade shall be examined.
- 2nd Row : The blades have flat upper and thick lower sides
- 3rd Row : k, M, x , Dx
- 4th Row :  $\mu$ , b,  $\epsilon$  , m
- 5th and 6th Row : STOP

2. Pointwise Given Blades

FILE: FILE	ETC3E001 P1	NAVAL		
1000160201				
0.01	0.03	17	1	
100.				
0.0	5.21	11.53	17.84	24.14
-0.26	-0.185	-0.110	-0.050	-0.005
30.45	36.75	43.03	49.35	55.65
0.029	0.056	0.084	0.117	0.157
61.95	68.25	74.55	80.88	87.21
0.098	0.095	0.088	0.0816	0.076
93.53	100.00	100.00	100.00	100.00
-0.035	-0.10	100.00	100.00	100.00
0.0	6.29	12.58	18.90	25.21
-0.26	-1.08	-1.65	-2.14	-2.45
31.53	37.35	44.18	50.51	56.83
-2.76	-2.25	-3.00	-2.97	-2.84
63.15	69.47	75.78	82.06	88.36
-2.60	-2.29	-1.94	-1.51	-1.17
94.66	100.00	100.00	100.00	100.00
-0.70	-0.10	100.00	100.00	100.00
0.28	1.55	1.4	0.5448-03	0.4548-C2
-4.4	0.5	2.9	0.335	
0.28	1.55	1.4	0.5448-03	0.4548-02
-22.2	0.5	2.3	0.335	
0.28	1.55	1.4	0.5448-03	0.4548-C2
53.2	0.5	2.3	0.335	
0.28	0.00	1.4	0.5448-03	0.4548-C2
222222222				

1st Row : No field output, pitch motion, no first passage, pointwise given blades, 16 gridpoints on the surface, 2nd blade shall be examined;

2nd Row : Identification for upper and lower side (no real importance here), 17 points given as input for each side, the geometry of the blades shall be printed;

3rd Row : Chord is 100

46h to 19th Row : Input geometry of blades, according to subprogram P R O F I L ;

20th Row : k , m , , Dx

21st Row :  $\mu$  , b , m , e

Two other cases and STOP

The output of the pressure distribution for the first case is shown next, followed by a listing of C A S C A D E in Appendix B.

FITCH - MOCF

PHASE= -4.40 / 2. ELACE

W\*C/U= 0.28000, M= 1.550, K= 1.40, B/C= 0.50, T/C= 0.0100

PRESSURE-COEFFICIENTS UPPER SURFACE:

PCINT	X	CPS	RCPU	ICPU
2, 1	C.C	0.2555	-0.1695	0.3648
3, 2	0.043	0.2111	-0.1671	0.3511
4, 3	0.167	0.1491	-0.1641	0.3341
5, 4	0.251	0.9955	-0.1621	0.3171
6, 5	0.334	0.7381	-0.1611	0.3111
7, 6	0.416	0.4181	-0.1581	0.2621
8, 7	0.503	0.3255	-0.1571	0.2631
9, 8	0.568	0.8701	-0.1555	0.2481
10, 9	0.629	-0.2822	-0.1388	0.6891
11, 10	0.691	-0.2522	-0.2322	0.6841
12, 11	0.752	-0.6461	-0.1341	0.6501
13, 12	0.814	-0.6511	-0.2022	0.6691
14, 13	0.876	-0.1131	-0.1341	0.6111
15, 14	0.938	-0.1571	-0.2011	0.7081
16, 15	1.000	-0.1741	-0.1261	0.5681

W\*C/U= 0.28000, M= 1.550, K= 1.40, B/C= 0.50, T/C= 0.0100

PRESSURE-COEFFICIENTS LOWER SURFACE:

PCINT	X	CPS	RCPU	ICPU
1, 2	C.C	0.3066	-0.1481	0.4401
2, 3	0.043	0.3955	-0.1411	0.4681
3, 4	0.058	0.2355	-0.1231	0.3491
4, 5	0.089	0.1801	-0.1121	0.3341
5, 6	0.124	0.1541	-0.1111	0.2851
6, 7	0.160	0.1441	-0.1101	0.2111
7, 8	0.200	0.1201	-0.1061	0.2571
8, 9	0.242	0.9791	-0.1041	0.2111
9, 10	0.263	0.6391	-0.1021	0.2201
10, 11	0.266	0.7221	-0.1011	0.2331
11, 12	0.341	0.4421	-0.9981	0.1921
12, 13	0.396	0.1921	-0.9871	0.1111
13, 14	0.462	-0.1681	-0.9801	0.3371
14, 15	0.491	0.4861	0.1571	-0.2171

## 5. The Programs as a Working System

### 5.1 TEST

TEST was written to examine the blade geometry if it is pointwise given. CASCADE and WING are not prepared to handle shocks which are caused by a changing slope over the airfoil. Therefore it has to be checked that the surface has no turning point.

For a given set of points one has to deform the PROFIL in such a way that the curvature of it never changes the sign. This would be indicated by a changing sign for the second derivative of the surface function.

An example for TEST - input data is shown below:

0.01	0.03	17	Y		
100.					
0.0	5.21	11.53	17.84	24.14	
-0.26	-0.185	-0.110	-0.057	-0.005	
30.45	36.73	43.05	49.35	55.65	
0.029	0.056	0.074	0.087	0.097	
61.55	68.25	74.55	80.88	87.21	
0.059	0.095	0.128	0.1516	0.116	
93.52	100.00	0.0	0.0	0.0	
100.00	0.10	100.	100.	100.	
0.0	5.29	12.59	19.92	25.21	
-0.26	-1.08	-1.63	-2.12	-2.45	
31.53	37.33	44.18	50.91	56.83	
-2.72	-2.53	-3.00	-2.97	-2.84	
63.15	69.47	75.78	82.06	88.36	
-2.60	-2.25	-1.94	-1.57	-1.17	
94.66	100.00	0.0	0.0	0.0	
0.10	0.10	100.	100.	100.	
-0.70	-0.10	100.	100.	100.	

TEST reads from File 05 and writes the result  $x$ ,  $y'_L$ ,  $y'_u$ ,  $y''_L$  and  $y''_u$  on File 06 and 01. File 06 can be printed and shows the original input data, the interpolated surfaces and besides the derivatives mentioned above, the coefficients of the cubic splines for each side. File 06 contains only  $x$ ,  $y$ ,  $y'$  and  $y''$ . It is used as an input file for PLOT1 which produces plots like Fig. 44 and Fig. 45 with it. Fig. 44 shows  $y'$  of the two PROFIL sides and Fig. 45 shows  $y''$ . PLOT1 is explained in Section 5.2.

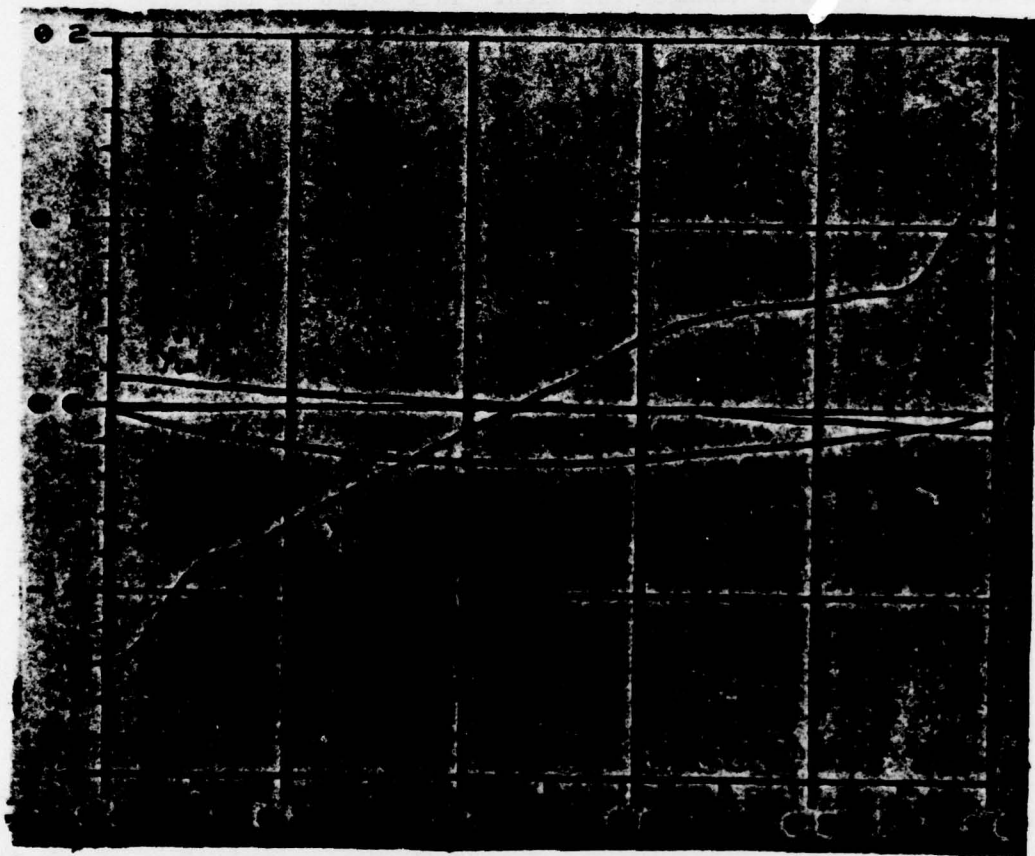


Fig. 44. Blade Geometry and  $y'$

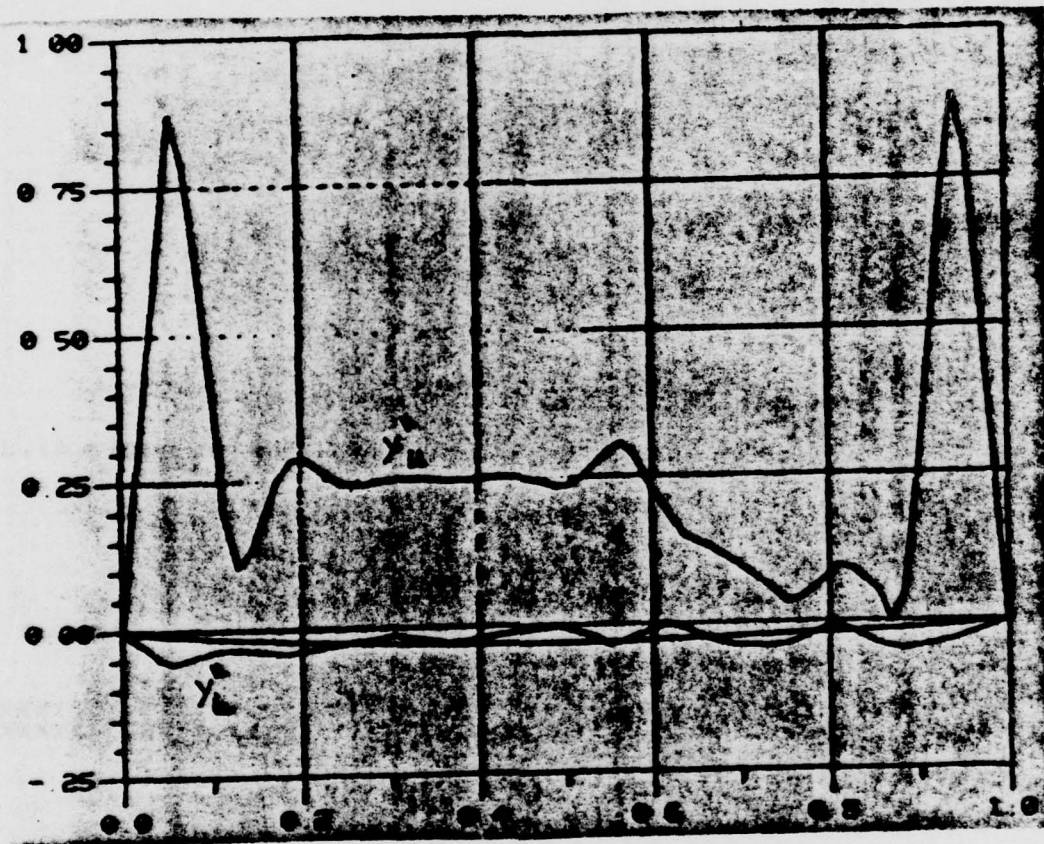


Fig. 45. Blade Geometry and  $y''$





```

DC 49 K=1,KV
A(K,1)=YS(J,K)
A(K,3)=R(J,K)
H=XS(J,K+1)-XS(J,K)
49 A(K,2)=(YS(J,K+1)-YS(J,K))/H-H*(R(J,K+1)+2.*R(J,K))/3.
A(K,4)=(R(J,K+1)-R(J,K))/(3.*H)
WRITE(6,1001)
DX=1./49.
DC 4 M=1,50
X=DX*(M-1)
DC 5 K=2,50
I=K-1
IF(XS(J,K).GE.X) GOTO 6
5 CONTINUE
6 H=X-XS(J,I)
Y=A(I,1)+A(I,2)*H+A(I,3)*H*H+A(I,4)*H*H*H
R(J,M)=X
YS(J,M)=Y
4 CCNTINUE
I2=1
N=50
DC 13 I=1,50
13 XS(J,I)=R(J,I)
GOTO 3
8 CCNTINUE
WRITE(6,1003) N
DC 15 M=1,N
15 WRITE(6,1002) XS(1,M),YS(1,M),XS(2,M),YS(2,M)
N=N-1
1002 FORMAT(5X,4(E12.5,2X))
1003 FORMAT(1H1,5X,'THE COMPLETE EXTRAPOLATED SURFACES N=',I3,/)
1005 FORMAT(3F10.5)
RETURN
END

```

## 5.2 P L O T 1

P L O T 1 was written to produce plots from two pointwise given functions

$$Y1 = F1(x)$$

$$Y2 = F2(x)$$

The  $x$ -stations are for both functions identical. This is useful to make diagrams for real and imaginary part of pressure distributions (see Fig. 18 through 25) and to visualize the  $y'$  and  $y''$  values of the blade geometry (Fig. 44 and 45). The input is read from File 01. This was produced from the respective program whose results shall be plotted. That is either W I N G or T E S T in this work.

The next step is to look in File 01 for the highest and lowest values of the functions and to decide the scale of the diagram. The last line of File 01 is a zero. Behind this zero there has to be inserted now the coordinates of Points 1 to 4 from Fig. 46.

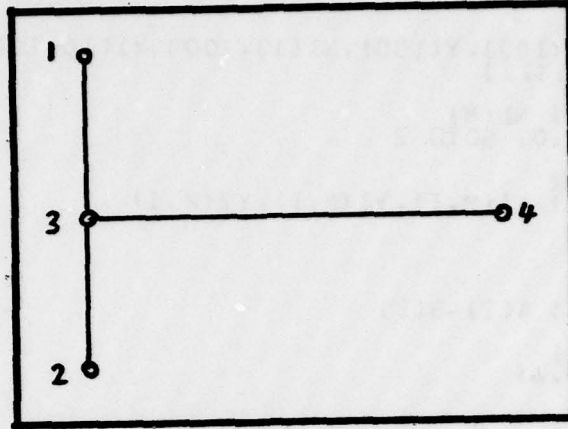


Fig. 46. Dimensioning of the Plot

This must be done for both expected diagrams. An example for the final lines of File 01 is shown below.

This file produced Fig. 44 and 45.

x	y1	y2
0.89855	C.06626	0.21406
0.91374	C.07111	0.44111
0.92764	C.07919	0.65786
0.94203	C.08115	0.80999
0.95652	C.10344	0.91529
0.97101	C.11344	0.94123
0.98551	C.11513	0.94650
1.00000	C.12099	0.00000
0.	C.15	
0.	-0.15	
0.	C.	
1.	C.	
0.	-0.1	
0.	C.	
1.	C.	

One can see the inserted eight additional lines behind the zero.

P L O T 1 works only on the Tektronix-terminal. After unpacking and compiling, it can be called on this device with

\$\$ PLOT 1

After the first plot is done, the terminal makes a tone. Then P L O T 1 hits a pause-statement. By striking an arbitrary key on the keyboard and CTRLS , it produces the second plot.

```

C      P L O T 1  L I S T I N G  :
C
C      DIMENSION X(100),Y(100),X1(10,100),Y1(10,100),Y2(10,100),
1      FA(4),B(4),N1(11)
      M=1
      READ(1,1000) N1(M)
      IF(N1(M).EQ.0) GOTO 2
      N2=N1(M)
      DC 10 I=1,N2
10     READ(1,1001) X1(M,I),Y1(M,I),Y2(M,I)
      M=M+1
      GOTO 1
      M=M-1
      DC 14 I=1,4
14     READ(1,1001) A(I),B(I)
      CALL INIT
      CALL NPTS(I)
      CALL PLCT(B,A)
      DC 3 I=1,M
      N=N1(I)
      DC 4 J=1,N
      X(J)=X1(I,J)
4      Y(J)=Y1(I,J)
      CALL NPTS(N)
      CALL CPLOT(X,Y)
3      CCNTINUE
      CALL PAUSE
      DC 15 I=1,4
15     READ(1,1001) A(I),B(I)
      CALL INIT
      CALL NPTS(I)
      CALL PLOT(B,A)
      DC 5 I=1,M
      N=N1(I)
      DC 6 J=1,N
      X(J)=X1(I,J)
6      Y(J)=Y2(I,J)
      CALL NPTS(N)
      CALL CPLOT(X,Y)
5      CONTINUE
      CALL FIN
1000  FORMAT(I3)
1001  FORMAT(3F10.5)
      STOP
      END
-5.3 P L O T 2

```

P L O T 2 is a special program which graphs the shock and cascade geometry on the Tektronix terminal. It is called with

\$\$ PLOT 2.

Input is read from File 01 which is prepared by C A S C A D E in each run.

The resulting plots are for example Fig. 35, 36 and 37. P L O T 2 asks for the size of the diagram by SCALE:

By enterine 1.5, 1., 0.5 or any other digital number from the keyboard, one may change the graph to the desired size. Each time the program has to be started again with the \$\$ - command.



This means for TEST, WING or CASCADE to be copied to the T-Disk, unpacked, altered to a FORTRAN file and compiled. After this, the input/output files are defined correctly. Finally the execution can be initialized by

\$ FILE NAME

If a compiled version of the desired program is already on the T-Disk,

DOTEST

DOWING

or

DOCAS

will also define the input/output files. Again \$ starts the execution.

See the three sample sessions on the next pages.

CONTENTS OF THE DISK:

1

FILENAME	FILETYPE	MODE	NO.REC.	DATE
TEST	DATA	P1	19	9/30
FILE4	DATA	P1	3	9/29
PLOT1	PACKED	P5	2	9/19
FILE	FT03F001	P1	1	9/30
FILE	FT04F001	P1	2	9/29
DOTEST	EXEC	P1	1	9/29
START	EXEC	P1	1	9/29
DOWING	EXEC	P1	1	9/29
CLOSE	EXEC	P1	1	10/01
TEST	PACKED	P5	8	9/29
PLOT2	PACKED	P5	2	9/28
FILE3	DATA	P1	3	9/29
.WING	PACKED	P5	77	9/26
FILE	FT05F001	P1	2	9/29
CASCADE	PACKED	P5	81	9/28
GETWING	EXEC	P1	1	9/29
GETCAS	EXEC	P1	1	9/29
DOCAS	EXEC	P1	1	9/29
GETTEST	EXEC	P1	1	9/29

this is a sampel session for t e s t .  
condition : you have Just lossed in.

```
start
17.30.06 VSET RDYMSG OFF
17.30.07 VSET BLIP /
17.30.08 CP DEFINE T2314 192 10
17.30.09 FORMAT T
** "FORMAT T" WILL ERASE ALL YOUR T-DISK (192) FILES **
**DO YOU WISH TO CONTINUE? ENTER "YES" OR "NO":
>yes
FORMATTING T-DISK (2314)...
T (192): 010 CYL
17.30.30 OFFLINE READ *
READER EMPTY OR NOT READY.
!!! E(00009) !!!
17.30.33 OFFLINE READ *
READER EMPTY OR NOT READY.
!!! E(00009) !!!
R;
>settest
17.30.42 VSET BLIP /
17.30.43 COMBINE TES PACKED T5 TEST PACKED P5
17.30.45 UNPACK TES
17.30.46 ALTER TES UNPACKED T5 TEST FORTRAN T1
17.30.47 ERASE TES PACKED T5
17.30.49 F TEST
////17.31.02 FILEDEF FT05F001 DSK
17.31.03 FILEDEF FT01F001 DSK-T1
17.31.05 FILEDEF FT06F001 DSK-T1
R;
>$ test
/EXECUTION BEGINS...
<R;
```

this is a sampel session for w i n s .

condition : t - disk is already lossed in

```
setwins
17.34.46 VSET BLIP /
17.34.47 COMBINE WIN PACKED T5 WING PACKED P5
17.34.53 UNPACK WIN
/17.35.01 ALTER WIN UNPACKED T5 WING FORTRAN T1
17.35.03 ERASE WIN PACKED T5
17.35.04 F WING
////////////////////17.37.09 FILEDEF FT01F001 DSK-T1
17.37.10 FILEDEF FT04F001 DSK
17.37.12 FILEDEF FT06F001 DSK-T1
R;
>$ wins
/EXECUTION BEGINS...
//////////////////R;
>
```

this is a sampel session for cascade .

condition : you have already logged in the t - disk

```
setcas
14.12.23 VSET BLIP /
14.12.24 COMBINE CAS PACKED T5 CASCADE PACKED P5
14.12.35 UNPACK CAS
/14.12.46 ALTER CAS UNPACKED T5 CASCADE FORTRAN T1
14.12.48 ERASE CAS PACKED T5
14.12.49 F CASCADE
////////////////////////////////////14.16.35 FILEDEF FT01F001 DSK-T1
14.16.37 FILEDEF FT03F001 DSK
14.16.38 FILEDEF FT06F001 DSK-T1
14.16.39 FILEDEF FT07F001 DSK-T1 RECFM F BLKSIZE 133
R/ T=122.98/127.06 14.16.39
```

```
>$ cascade
/EXECUTION BEGINS...
////////////////////////////////////R/ T=70.81/73.66 14.18.45
```

## 6. Critique, Conclusion and Outlook

A systems of programs was written, which allows a systematic approach to the problem of oscillating shaped airfoils in a staggered cascade. T E S T is only a geometrical examination of the PROFIL. W I N G is a test for the theoretical behaviour of the blade, if exposed oscillating to a supersonic stream. Finally, C A S C A D E is a computer model for a finite cascade of airfoils which may oscillate with a phase lag from blade to blade in a supersonic flow. C A S C A D E is not yet finished. Results are available up to now only for the inlet flow, as the wake could not be added because of lack of time.

All three programs do not only give numerical but also graphical results. The latter can be obtained by the two plot codes P L O T 1 and P L O T 2 .

C A S C A D E - and W I N G -results are compared in this report with solutions of Teipel (2), Verdon (15,17,18) and Platzer and collaborators (3,4,5, 6,7,9,13,14). For the linear cases the agreement is considered to be very good. The included thickness effect gives different results to those from Strada (5). This has to be investigated more closely in the near future.

It is clearly understood that C A S C A D E is a rather expensive approach to the problem. This is considered worth while due to the complexity of the problem. After the model of shaped, oscillating airfoils is better understood, it would be desirable to write a more efficient computer code.

To finish this work, there are two more steps to do:

1. Adding the procedure of wake computation to C A S C A D E and then obtaining results over the whole blades.
2. Optimizing the program in this form.

This will be done in the near future.

When this is accomplished, a rather flexible logic structure for the problem of an oscillating supersonic cascade is available. This could be used to examine different systems of basic equations and their influence on the results with the final purpose to substitute the potential equations by those of Euler. Thus we could get rid of the assumption of constant entropy in the field which would be a new step forward.



```

C      CCMPUTATION OF THE SPLINE - COEFFICIENTS
C
Ih=0
IE=0
T3=0.
I=0
30  I=I+1
    T=T1
    IF(I.EQ.2) T=T2
    DC 46 K=1,NX
    A(K,1)=YS(I,K)
    A(K,3)=R(I,K)
    H=XS(I,K+1)-XS(I,K)
    A(K,2)=(YS(I,K+1)-YS(I,K))/H-H*(R(I,K+1)+2.*R(I,K))/3.
46  A(K,4)=(R(I,K+1)-R(I,K))/(3.*H)
    I1=(-1)**(I+1)
16  MA=MAX
    M2=1
C
C      THE STEADY BOUNDARY-PROPERTIES BEHIND THE SHOCK
C
CX1=1.07/(MA-1)
DC 9 J=2,KV
KT=J
XF=DX(I)*(J-2)
IF(T.EQ.0.) XP=0.
IM=1
9  CALL SWITCH(J,IM,M,N,I)
CC  CALL SHOCK(KV,WE,LC4,M,N,XP,C3,MA,CX1)
C
C      ALL OTHER STEPS OF THE STEADY FLOW FIELD
C
KA=KT
DC 1 J=2,KV
J1=J
IM=J+1
CALL SWITCH(IM,J,M,N,I)
CALL RAND(WE,LC4,M,N,C3)
IF(X(2,M,N).GT.1.) M2=0
KB=KA+1
DC 3 J2=J,KB
IF(J2.EC.KB) GOTO 3
IM=J-1
CALL SWITCH(J2,IM,K1,K2,I)
3  IF(AL(2,K1,K2).GT.AL(2,M,N)) GOTO 6
6  CONTINUE
   J2=J2-1
   IN(I,J-1)=J2
   L=J+2
   DC 7 K=L,J2
7  CALL SWITCH(K,J,M,N,I)
   CALL GEN(M,N,I,IW)
   IF(M2.EQ.0) GOTO 15
   IF(J2.EC.KA) GOTO 1
   J3=M+1
   J4=N-1
   J5=M
   J6=N-1
   IF(I.NE.2) GOTO 12
   J3=M-1
   J4=N+1
   J5=M-1
   J6=N
12  D1=(Y(2,J3,J4)-Y(2,J5,J6))/(X(2,J3,J4)-X(2,J5,J6))*I1
   C2=1./SQRT(AL(2,M,N))
   K1=M+1
   K2=J
   IF(I.NE.2) GOTO 14
   K1=J
   K2=N+1
14  AL(2,K1,K2)=AL(2,M,N)
   X(2,K1,K2)=(Y(2,J5,J6)-Y(2,M,N))*I1
   X(2,K1,K2)=-D2*X(2,M,N)+C1*X(2,J5,J6)-X(2,K1,K2)
   X(2,K1,K2)=X(2,K1,K2)/(C1-D2)

```

```

Y(2,K1,K2)=C2*(X(2,K1,K2)-X(2,M,N))*I1+Y(2,M,N)
IF(1.EC.2) J3=J4
DC 11 K=J3,KA
IF(K.EC.KV) GOTO 11
IM=K+1
CALL SWITCH(IM,J,M,N,I)
IM=J-1
CALL SWITCH(K,IM,K1,K2,I)
X(2,M,N)=X(2,K1,K2)
Y(2,M,N)=Y(2,K1,K2)
AL(2,M,N)=AL(2,K1,K2)
11 CONTINUE
IF(KA.LT.KV) KA=KA+1
IF(MA.GT.IN(I,J-1)) MA=MA+1
1 CCNTINUE
15 CCNTINUE
IF(T.EQ.0.) GOTO 47
C
C
C AUTOMATIC STEP - SIZE CCNTRCL
IF(J1.GT.MA) DX(I)=DX(I)*1.05
IF(J1.LT.MA) DX(I)=DX(I)*0.96
IF(J1.NE.MA) GOTO 16
C
C
C THE UNSTEADY FLOW FIELD
C THE UNSTEADY BOUNDARY PROPERTIES BEHIND SHOCK
47 KA=KT
IF(L01.EQ.0) WRITE(6,1007)
DC 18 J=2,KT
IM=1
18 CALL SWITCH(J,IM,M,N,I)
CALL RANCS(KV,M,N,DX(I),LC2,M2,N2)
X(2,1,2)=0.
Y(2,1,2)=0.
X(2,2,1)=0.
Y(2,2,1)=0.
C
C
C ALL OTHER STEPS OF UNSTEADY FLOW FIELD
J5=J1-1
DO 20 J=2,J5
IP=J+1
CALL SWITCH(IP,J,M,N,I)
CALL RANDB(M,N,AM,AK,I,L02)
L=J+2
J2=IN(I,J-1)
DO 22 K=L,J2
22 CALL SWITCH(K,J,M,N,I)
CALL GENU(M,N,I,AK,AM)
IF(J2.EC.KA) GOTO 20
K1=M+1
K2=J
K3=K1
IF(I.NE.2) GOTO 24
K1=J
K2=N+1
K3=K2
24 P(I,K3)=X(2,M,N)
CALL RANCS(KV,K1,K2,DX(I),L02,M2,N2)
J3=J2+1
DC 29 K=J3,KA
IF(K.EC.KV) GOTO 20
IM=K+1
CALL SWITCH(IM,J,M,N,I)
IM=J-1
CALL SWITCH(K,IM,K1,K2,I)
U(2,M,N)=U(2,K1,K2)
V(2,M,N)=V(2,K1,K2)
G(2,M,N)=G(2,K1,K2)
29 PSI(2,M,N)=PSI(2,K1,K2)
IF(J2.LT.KA .AND. KA.LT.KV) KA=KA+1
20 CCNTINUE
C

```

```

C      OUTPUT STEADY AND UNSTEADY FIELD
C
IF(LO1.NE.0) GOTO 38
WRITE(6,1004)
IF(I.EQ.1) WRITE(6,1002)
IF(I.EC.2) WRITE(6,1003)
J5=J1+1
DC 25 J=2,J5
K=J-1
LL=IN(I,K)
WRITE(6,1007)
WRITE(6,1008)
N=K
IF(I.NE.2) GOTO 26
M=K
DC 31 N=J,LL
31 WRITE(6,1009) M,N,X(2,M,N),Y(2,M,N),AL(2,M,N),U(2,M,N),
FV(2,M,N),PSI(2,M,N),G(2,M,N)
GOTO 17
26 CONTINUE
DC 32 M=J,LL
32 WRITE(6,1009) M,N,X(2,M,N),Y(2,M,N),AL(2,M,N),U(2,M,N),
FV(2,M,N),PSI(2,M,N),G(2,M,N)
17 WRITE(6,1005)
25 CCNTINUE
C      COMPUTATION OF THE PRESSURE - COEFFICIENTS
C
38 K=J1+1
DC 27 J=2,K
IM=J-1
27 CALL SWITCH(J,IM,M,N,I)
CALL PRESS(M,N,J1)
IF(I.EC.1) IUS=K-1
IF(I.EC.2) ILS=K-1
IF(I.EQ.1) GOTO 30
C
C      CHANGING THE FIELDS AND WAKE - COMPUTATION
C
DC 48 M=1,KV
DC 48 N=1,KV
X(1,M,N)=X(2,M,N)
Y(1,M,N)=Y(2,M,N)
U(1,M,N)=U(2,M,N)
V(1,M,N)=V(2,M,N)
G(1,M,N)=G(2,M,N)
AL(1,M,N)=AL(2,M,N)
48 PSI(1,M,N)=PSI(2,M,N)
CALL WAKE(KV,WE,T1,T2,LC4,IUS,ILS,LO3,NX,MA)
C
C      CORRECTION OF THE INDEX NUMBERS
C
DC 35 I=1,2
K=IUS
IF(I.EQ.2) K=ILS
KA=K
DC 43 J=1,K
IF(IN(I,J).GT.K) GOTO 35
KA=KA-1
J1=IN(I,J)
CC 44 L=J1,KA
M=L+1
PX(I,L)=PX(I,M)
PL(I,L)=PL(I,M)
PS(I,L)=PS(I,M)
44 IF(I.EQ.1) IUS=KA
IF(I.EC.2) ILS=KA
435 CCNTINUE
C
C      OUTPUT PRESSURE DISTRIBUTION
C
WRITE(6,1004)
IF(LO2.EC.0) WRITE(6,556)
IF(LO2.EQ.1) WRITE(6,595)

```



```

2      DC 2 LL=1,N
      XS(J,LL)=XS(J,LL)/SP
      YS(J,LL)=YS(J,LL)/SP
12     IF(L04.EC.0) GOTO 3
      DX=1.0/(N-1)
      I1=(-1)**(J+1)
      DC 20 K=1,N
      XS(J,K)=(K-1)*CX-0.25
      YS(J,K)=I1*4.*T*XS(J,K)*(1.-XS(J,K))
20     C
      TC ENTER THIS PART,THE SURFACES SFCULD ALREADY BE
      GIVEN POINTWISE
      EXTRAPOLATING POINTS
      C
      CCNTINUE
      C
      INTERPOLATION THROUGH CUBIC SPLINES
      C
      IF(T.NE.0. .OR. L04.EC.C) GCTC 51
      DC 52 I=1,N
      A(I,1)=0.
      GCTC 53
      CCNTINUE
      DO 10 I=1,N
      A(I,3)=XS(J,I)
      A(I,4)=YS(J,I)
10     C
      C
      MATRIX OF COEFFICIENTS AND RIGHT SIDES
      C
      K=N-2
      DC 25 I=1,K
      A(I,1)=A(I+1,3)-A(I,3)
      A(I,3)=A(I+2,3)-A(I+1,3)
      A(I,2)=2.*(A(I,1)+A(I,3))
      A(I,4)=3.*(A(I+2,4)-A(I+1,4))/A(I,3)-
      F3.*(A(I+1,4)-A(I,4))/A(I,1)
25     CCNTINUE
      A(1,1)=0.0
      A(N-2,3)=0.0
      C
      C
      THE STEP OF GAUSS
      C
      K=N-3
      DC 30 I=1,K
      CC 35 M=3,4
      A(I,M)=A(I,M)*A(I+1,1)/A(I,2)
      A(I,1)=0.0
      A(I,2)=A(I+1,1)
      A(I+1,2)=A(I+1,2)-A(I,3)
      A(I+1,4)=A(I+1,4)-A(I,4)
30     C
      C
      SOLUTION
      C
      A(1,1)=0.
      A(N,1)=0.
      A(N-2,1)=0.0
      L=N-1
      DC 40 I=2,L
      K=N-I
      M=K+1
      A(M,1)=(A(K,4)-A(K,3)*A(M+1,1))/A(K,2)
40     CCNTINUE
      DO 11 M=1,N
      R(J,M)=A(M,1)
      IF(IZ.EQ.1) GCTC E
      KV=N-1
      DC 49 K=1,KV
      A(K,1)=YS(J,K)
      A(K,3)=R(J,K)
      H=XS(J,K+1)-XS(J,K)
      A(K,2)=(YS(J,K+1)-YS(J,K))/H-H*(R(J,K+1)+2.*R(J,K))/3.
49     A(K,4)=(R(J,K+1)-R(J,K))/(3.*H)
      DX=1./49.
      DC 4 M=1,50

```

```

X=DX*(M-1)
DC 5 K=2,50
I=K-1
IF(XS(J,K).GE.X) GOTO 6
5 CCNTINUE
6 H=X-XS(J,I)
Y=A(I,1)+A(I,2)*H+A(I,3)*H*H+A(I,4)*H*H*H
R(J,M)=X
YS(J,M)=Y
4 CONTINUE
IZ=1
N=50
DC 13 I=1,50
13 XS(J,I)=R(J,I)
GOTO 3
8 CCNTINUE
IF(IKP.EQ.0) GOTO 17
DC 15 M=1,N
15 WRITE(6,1002) XS(1,M),YS(1,M),XS(2,M),YS(2,M)
17 N=N-1
1000 FORMAT(2F10.5,2I2)
1001 FCRMAT(5F10.5)
1002 FORMAT(5X,4(E12.5,2X))
RETURN
END
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
SUBROUTINE FIND(KV, IW, XS, YS, M, N, IE, I)
C COMPLEX*8 U, V, AI, PSI, G, CYOXU
C COMMON /BC/ T, B, DYCX, DZCYX2, DYDXU, AI, I1
C COMMON V(2,25,25), X(2,25,25), P(2,25), U(2,25,25),
F PSI(2,25,25), G(2,25,25), AL(2,25,25), Y(2,25,25),
FC(2,3), CX(2), IN(2,25)
C
C FIND LOOKS FOR THE MESHINDEX, RESPONSIBEL FGR
C THE POINT XS;YS
C
IF(IE.EQ.1) GOTO 19
IKK=0
C IF(IW.GT.5 .AND. XS.GE.1 .AND. XS.LE.1.10) IKK=1
IF(IKK.EQ.1) WRITE(6,2011) I, IE, I1
2011 FCRMAT(IX, 'FIND ENTRY: ',3I4)
IZ=0
KA=0
L1=1
L2=1
21 I9=21
IF(IKK.EQ.1) WRITE(6,2003) I9, I, IE, XS, YS
DO 22 I2=L1,20
IM=I2+L2
CALL SWITCH(IM, I2, M, N, I)
L3=I2+L2
I9=26
IF(IKK.EQ.1) WRITE(6,2000) I9, M, N, X(1, M, N), Y(1, M, N)
IF(L3.GE.KV) GOTO 13
J1=M+1
J2=N+1
24 KP=J1-J2
IF(I.EQ.2) KM=J2-J1
IC=24
IF(IKK.EQ.1) WRITE(6,2000) I9, M, N, X(1, M, N), Y(1, M, N)
IF(XS.GE.X(1, M, N) .AND. XS.LT.X(1, J1, J2)) GOTO 23
IF(X(1, J1, J2).GT.1. .AND. KM.EQ.1) KA=J2
IF(I.EQ.2 .AND. KA.EQ.J2) KA=J1
IF(I.EQ.1 .AND. KA.NE.J2) GOTO 22
IF(I.EQ.2 .AND. KA.NE.J1) GOTO 22
IF(I.EQ.1) J1=J1+1
IF(I.EQ.2) J2=J2+1
IF(J1.EQ.KV .OR. J2.EQ.KV) GOTO 13
GOTO 24
IC=22
22 IF(IKK.EQ.1) WRITE(6,2000) I9, M, N, X(1, M, N), Y(1, M, N)
23 IC=23
IF(IKK.EQ.1) WRITE(6,2000) I9, M, N, X(1, M, N), Y(1, M, N)

```

```

IF(I.EC.2) GOTO 27
IF(YS.LT.Y(1,J1,N)) GCTC 1
IF(N.EQ.1) GOTO 13
L2=L2+1
L1=N-1
GCTO 21
27 IS=27
IF(IKK.EQ.1) WRITE(6,2000) I9,M,N,X(1,M,N),Y(1,M,N)
IF(YS.GT.Y(1,M,J2)) GCTC 1
IF(M.EQ.1) GOTO 13
L2=L2+1
L1=M-1
GCTO 21
1 IS=1
IF(IKK.EQ.1) WRITE(6,2000) I9,M,N,X(1,M,N),Y(1,M,N)
M2=M
N2=N
GCTO 14
28 IS=28
IF(IKK.EC.1) WRITE(6,2000) I9,M,N,X(1,M,N),Y(1,M,N)
IZ=IZ+1
IF(IZ.EQ.5) GOTO 29
M=M2
N=N2
18 IS=18
IF(IKK.EQ.1) WRITE(6,2000) I9,M,N,X(1,M,N),Y(1,M,N)
IF(IZ.EC.0) GOTO 28
IF(IZ.EC.1 .AND. I.EC.1) N=N-1
IF(IZ.EC.1 .AND. I.EC.2) M=M-1
IF(M.EC.0 .OR. N.EC.0) GOTO 28
IF(IZ.EC.2 .AND. I.EC.1) M=M+1
IF(IZ.EC.2 .AND. I.EC.2) N=N+1
IF(M.EC.KV .OR. N.EC.KV) GCTC 28
IF(IZ.EC.3 .AND. I.EC.1) N=N+1
IF(IZ.EC.3 .AND. I.EC.2) M=M+1
IF(IZ.NE.3) GCTC 3
3 IF(N.EC.M) GOTO 28
IF(IZ.EC.4 .AND. I.EC.1) M=M-1
IF(IZ.EC.4 .AND. I.EC.2) N=N-1
IF(IZ.NE.4) GOTO 28
2 IF(M.EC.N) GOTO 28
CONTINUE
GCTO 14
29 WRITE(6,1000) KK,M,N,XS,YS,X(1,M,N),Y(1,M,N)
13 IE=1
WRITE(6,1001) M,N,XS,YS,IE
GOTO 19
C
C
C
14 CCONTRCL - STEP
J1=M+1
J2=N+1
K1=M+1
K2=N
K3=M
K4=N+1
IF(I.NE.2) GOTO 20
K1=M
K2=N+1
K3=M+1
K4=N
20 IS=20
IK1=0
IF(IKK.EC.1 .AND. M.EC.14 .AND. N.EC.13) IK1=1
IF(IK1.EC.1) WRITE(6,2000) I9,M,N,X(1,M,N),Y(1,M,N)
IF(IK1.EC.1) WRITE(6,2000) I9,J1,J2,X(1,J1,J2),X(1,J1,J2)
IF(IK1.EC.1) WRITE(6,2000) I9,K1,K2,X(1,K1,K2),Y(1,K1,K2)
IF(IK1.EC.1) WRITE(6,2000) I9,K3,K4,X(1,K3,K4),Y(1,K3,K4)
IF(XS.EC.X(1,M,N) .AND. YS.EC.Y(1,M,N)) GCTC 15
D1=I1*(Y(1,K1,K2)-Y(1,M,N))/(X(1,K1,K2)-X(1,M,N))
IF(K4.EC.K3) GOTO 11
D2=I1*(Y(1,K3,K4)-Y(1,M,N))/(X(1,K3,K4)-X(1,M,N))
11 IF(K4.EC.K3) D2=0.
D12=I1*(YS-Y(1,M,N))/(XS-X(1,M,N))

```

```

IF(K4.EQ.K3) GOTO 25
D3=I1*(Y(1,J1,J2)-Y(1,K3,K4))/(X(1,J1,J2)-X(1,K3,K4))
25 IF(K4.EQ.K3) D3=0.
D4=I1*(Y(1,J1,J2)-Y(1,K1,K2))/(X(1,J1,J2)-X(1,K1,K2))
D4=I1*(Y(1,J1,J2)-YS)/(X(1,J1,J2)-XS)
IF(D1.LT.D12 .OR. C2.GT.D12) GOTO 18
IF(D3.LT.D34 .OR. D4.GT.C34) GOTO 18
IF(XS.LT.X(1,M,N) .OR. XS.GE.X(1,J1,J2)) GOTO 18
1000 FCRMAT(IX,'NO-FIND: ',I2,' ',I2,' ',I2,4(2X,F8.3))
1001 FCRMAT(IX,'FIND: ',I2,' ',I2,2(2X,E10.3),' IE= ',I2)
2000 FCRMAT(IX,3I4,2X,F8.3,2X,F8.3)
19 RETURN
END

```

CXX  
CXX

```

SUBROUTINE BOUND(WE,LC4,I,X,IW)
COMPLEX*8 L,V,AI,PSI,G,DYDXU
COMMON/BC/ T,B,DYDX,D2YCX2,DYDXU,AI,I1,T3
COMMON/SP/ XS(2,50),YS(2,50),A(5C,4),R(2,50)
IF(IW.NE.0) GOTO 1

```

C  
C  
C

```

STEADY BOUNDARY CONDITIONS ALONG THE AIRFOIL
IF(LO4.EC.1) GOTO 3
IF(LO4.EC.2) GOTO 4
DC 5 K=2,50
J=K-1
IF(XS(I,K).GE.X) GOTO 6
5 CONTINUE
6 H=X-XS(I,J)
DYDX=A(J,2)+2.*A(J,3)*H+3.*A(J,4)*H*H
D2YDX2=2.*A(J,3)+6.*A(J,4)*H
3 GOTO 2
DYDX=I1*4.*T*(1.-2.*X)
D2YDX2=-I1*8.*T
2 GOTO 2
4 W=WE*4.*ATAN(1.)/180.
DYDX=I1*TAN(W)
D2YDX2=0.
GOTO 2

```

C  
C  
C

```

STEADY BOUNDARY CONDITIONS ALONG WAKE-SLIP-LINE
1 DYCX=TAN(T3)
2 D2YDX2=0.
RETURN
END

```

CXX  
CXX

```

SUBROUTINE SHOCK(KV,WE,LC4,M,N,XP,C3,MA,CX1)
COMPLEX*8 U,V,AI,PSI,G,DYDXU
COMMON/BA/ ALO,ALO,AM,C1,AK,I,IE,IW
COMMON/BC/ T,B,DYDX,D2YCX2,DYDXU,AI,I1,T4
COMMON V(2,25,25),X(2,25,25),P(2,25),U(2,25,25),
PSI(2,25,25),G(2,25,25),AL(2,25,25),Y(2,25,25),
FQ(2,3),DX(2),IN(2,25)

```

C  
C  
C

COMPUTATION OF THE FIELD NEAR BEHIND THE SHOCK

```

ITE=0
IKK=0
C IF(IW.NE.0) IKK=1
2000 IF(IKK.EC.1) WRITE(6,2000) M,N,I,IE
FORMAT(IX,'SHOCK ENTRY: ',4I4)
TX=0.
IF(IW.NE.0) TX=1.
J=M-1
K=N
L=M
IF(I.NE.2) GOTO 5
J=M
K=N-1
L=N
5 CALL BOUND(WE,LO4,I,XP,IW)

```

```

AL(2,M,N)=(ALD-C3*DYDX*I1)**(2./3.)
IF(L.EC.2) GOTO 17
D1=I1/SQRT(AL(2,M,N))
X5=X(2,J,K)
Y5=Y(2,J,K)
CALL FIND(KV,IW,X5,Y5,M2,N2,IE,I)
9 IK=9
IF(IKK.EQ.1) WRITE(6,1000) IK,IE,AL(2,M,N),X5,Y5
IL=9
IF(IE.EC.1 .OR. M2.GE.KV) GOTO 16
L1=M2+1
L2=N2
IF(I.NE.2) GOTO 7
L1=M2
L2=N2+1
7 X1=X(1,M2,N2)
Y1=Y(1,M2,N2)
X2=X(1,L1,L2)
Y2=Y(1,L1,L2)
X3=X(1,M2+1,N2+1)
Y3=Y(1,M2+1,N2+1)
I7=0
14 D2=SQRT(AL(1,M2,N2))+SQRT(AL(2,J,K))
D2=2.*I1/D2
IK=7
IF(IKK.EC.1) WRITE(6,1000) IK,M2,X1,X2,X3
IF(T.NE.O. .AND. IW.EC.C .AND. IE.EC.O) GOTO 2
IF(T4.NE.O. .AND. IW.NE.C .AND. IE.EQ.O) GCTC 2
X4=X(2,J,K)+DX1/2.
Y4=Y(2,J,K)+(X4-X(2,J,K))*D2
IK=14
IF(IKK.EQ.1) WRITE(6,1000) IK,IE,X4,Y4,X5,Y5
GOTO 13
2 X4=(D1*XP-D2*X5+Y5)/(D1-D2)
Y4=(X4-XP)*D1
IF(X4.LE.TX) ITE=1
IF(ITE.EC.1) GOTO 14
IK=2
13 IF(IKK.EQ.1) WRITE(6,1000) IK,IE,X4,Y4,X5,Y5
CALL FIND(KV,IW,X4,Y4,M1,N1,IE,I)
IF(M1.EC.M2 .AND. N1.EC.N2 .AND. IE.EQ.O) GOTO 8
IF(IW.NE.O .AND. IE.EQ.1) GOTO 11
X6=(Y2-Y5)/D2+X5
IF(X6.GE.X2) GOTO 10
D3=(Y2-Y1)/(X2-X1)
X5=(Y1-Y5+D2*X5-D3*X1)/(D2-D3)
Y5=D3*(X5-X1)+Y1
IL=131
IF(N2.EQ.1) GOTO 16
N2=N2-1
IF(I.EC.2) GOTO 19
IL=132
IF(AL(1,M2,N2).NE.AL(1,M2-1,N2)) GCTC 16
GOTO 20
19 IL=19
IF(AL(1,M2,N2).NE.AL(1,M2,N2-1)) GOTO 16
20 CCNTINLE
IF(I.NE.2) GOTO 9
M2=M2-1
N2=N2+1
GOTO 9
10 D3=(Y3-Y2)/(X3-X2)
X5=(Y2-Y5+D2*X5-D3*X2)/(D2-D3)
Y5=D3*(X5-X2)+Y2
M2=M2+1
IL=10
IF(M2.EQ.KV) GOTO 16
IF(I.NE.2) GOTO 9
N2=N2+1
M2=M2-1
GOTO 9
8 X(2,M,N)=X4
Y(2,M,N)=Y4
GOTO 11

```

```

16      IE=1
        D2=2.*SQRT(ALO)+SQRT(AL(2,J,K))+SQRT(AL(2,M,N))
        D2=4.*I1/D2
        IK=16
        IF(IKK.EC.1) WRITE(6,1000) IK,I1,XP,X5,Y5,X(2,M,N)
        IF(IKK.EQ.1) WRITE(6,1000) M,N,C1,C2,DX1
        IF(T.EQ.0.) GOTO 1
        X(2,M,N)=(D1*XP-D2*X5+Y5)/(D1-D2)
        Y(2,M,N)=(X(2,M,N)-XP)*C1
        GOTO 11
1       X(2,M,N)=X5+DX1/2.
        Y(2,M,N)=D2*DX1/2.+Y5
C
C
C
C
11      D3=1./SQRT(AL(2,M,N))
        D4=1./SQRT(AL(2,J,K))
        P(I,L)=(D4*X(2,J,K)+D3*XP+Y(2,J,K)*I1)/(D3+D4)
        IF(L.EQ.3) P(I,L)=XP
        IF(T.EC.0. .OR. I.NE.0) P(I,L)=X(2,J,K)
        GOTO 12
17      X(2,M,N)=0.
        Y(2,M,N)=0.
        IF(IW.EC.0) GOTO 12
        X(2,M,N)=1.
12      IF(IKK.EC.1) WRITE(6,1000) M,N,X(2,M,N),Y(2,M,N),AL(2,M,N),
1000    F(I,L)
        FFORMAT(1X,2I4,4(2X,F8.3))
        RETURN
        END
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
        SLBROUTINE RAND(WE,LC4,M,N,C1)
        CCMPLX*8 U,V,AI,PSI,G,CYDXU
        CCOMMON/PA/ ALO,ALC,AM,C3,AK,I,IE,IW
        CCOMMON/BC/ T,B,DYDX,D2YCX2,DYDXU,AI,I1
        CCOMMON V(2,25,25),X(2,25,25),P(2,25,25),U(2,25,25),PSI(2,25,25),
        FG(2,25,25),AL(2,25,25),Y(2,25,25),C(2,25)
C
C
C
C
        BCUNDERYSTEP OF THE STEADY FIELD TO THE AIRFOIL
C
        IKK=0
        IF(LO4.EC.0) IKK=1
        IF(IKK.EQ.1) WRITE(6,1002)
1002    FFORMAT(1X,'RAND-ENTRY')
        A=-C1*I1
        J1=M
        J2=N-1
        IF(I.NE.2) GOTO 5
        J1=M-1
        J2=N
5       X(2,M,N)=X(2,J1,J2)
        DC I1 KX=1,15
        CALL BOUNC(WE,LO4,I,X(2,M,N),IW)
        IF(IKK.EC.1) WRITE(6,1001) M,N,KX,X(2,M,N),DYDX,D2YCX2
1001    FFORMAT(1X,3I3,3E12.5)
        AL(2,M,N)=(ALD+A*DYOX)**(2./3.)
        IF(IW.NE.0) GOTO 2
        IF(DYOX.EQ.0. AND. D2YCX2.EQ.0.) GOTO 2
        F=2.*(X(2,M,N)-X(2,J1,J2))*I1/Y(2,J1,J2)
        F=(SQRT(AL(2,J1,J2))+SQRT(AL(2,M,N)))
        FS=2.*I1/Y(2,J1,J2)-A*C2YDX2/(3.*AL(2,M,N))
        D=F/FS
        X(2,M,N)=X(2,M,N)-C
        IF(ABS(C).LE.0.000001) GOTO 10
11      CCONTINUE
        WRITE(6,1000) M,N,D
1000    FFORMAT(1X,'BOUNDERYSTEP DID NOT CONVERGE: ',I2,', ',I2,
10      F' D= ',E10.3)
        Y(2,M,N)=0.
        CALL BOUNC(WE,LO4,I,X(2,M,N),IW)
        AL(2,M,N)=(ALD+A*DYOX)**(2./3.)
        GOTO 1

```

```

2      Y(2,M,N)=0.
      D=-2.*I1/(SQRT(AL(2,J1,J2))+SQRT(AL(2,M,N)))
      X(2,M,N)=X(2,J1,J2)-Y(2,J1,J2)/D
1      RETURN
      END
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
      SUBROUTINE GEN(M,N,I,IW)
      COMPLEX*8 U,V,AI,PSI,G,CYDXU
      COMMON/BC/ T,B,DYCX,D2YCX2,DYDXU,AI,I1
      COMMON V(2,25,25),X(2,25,25),P(2,25),U(2,25,25),PSI(2,25,25),
      FG(2,25,25),AL(2,25,25),Y(2,25,25),C(2,25)
C
C      GENERAL STEP OF STEADY FIELD
C
      IKK=0
      IF(IKK.EQ.1) WRITE(6,10CC)
1000  FORMAT(1X,'GEN - ENTRY')
      J1=M
      J2=N-1
      K1=M-1
      K2=N
      IF(I.NE.2) GOTO 5
      J1=M-1
      J2=N
      K1=M
      K2=N-1
5      D1=1./SQRT(AL(2,K1,K2))
      D2=2./SQRT(AL(2,J1,J2))+SQRT(AL(2,K1,K2))
      XY=D2*X(2,J1,J2)+D1*X(2,K1,K2)+(Y(2,J1,J2)-Y(2,K1,K2))*I1
      X(2,M,N)=XY/(D1+D2)
      Y(2,M,N)=D1*I1*(X(2,M,N)-X(2,K1,K2))+Y(2,K1,K2)
      AL(2,M,N)=AL(2,K1,K2)
      IF(IKK.EQ.1) WRITE(6,1CC1) M,N,X(2,M,N),Y(2,M,N),AL(2,M,N)
1001  FORMAT(1X,2I3,3(2X,E12.5))
      RETURN
      END
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
      SUBROUTINE RANDS(KV,M,N,X1,L,M2,N2)
      COMPLEX*8 U,V,AI,PSI,G,CYDXU,DU,DV,C2,C6
      COMPLEX*16 A,RI,ES
      COMMON/BA/ ALO,ALD,AM,C,AK,I,IE,IW
      COMMON/BC/ T,B,DYCX,C2YCX2,DYDXU,AI,I1,T4
      COMMON/SCL/ A(4,4),RI(4),ES(4)
      COMMON V(2,25,25),X(2,25,25),P(2,25),U(2,25,25),
      FPSI(2,25,25),G(2,25,25),AL(2,25,25),Y(2,25,25),
      FG(2,3),CX(2),IN(2,25)
C
C      COMPUTATION OF UNSTEADY BOUNDARY-PROPERTIES ALONG SPOCK
C
      IKK=0
      IF(IW.NE.0) IKK=1
      J=M
      IF(I.NE.2) GOTO 5
      J=N
5      IF(IKK.EQ.1) WRITE(6,2000)
2000  FORMAT(1X,'RANDS-ENTRY')
      CALL COEF1(KV,I1,M,N,M2,N2,AM1,AM2,AM3,AM4,AN1,AN2,AN3,
      FAN4,S)
      IK=5
      IF(IKK.EQ.1) WRITE(6,1001) IK,M,N,U(1,M2,N2),V(1,M2,N2),
      FAL(1,M2,N2)
      TX=1.
      DV=V(1,M2,N2)
      DU=U(1,M2,N2)
      ALT=AL(1,M2,N2)
      IF(IW.EQ.0) TX=0.
      IF(IE.EQ.1) ALT=A LO
      IF(IE.EQ.1) DV=0.
      IF(IE.EQ.1) DU=0.
      IF(J.GT.2) GOTO 6
      CALL BOLNDU(IW,TX,I,AK,L)
      V(2,M,N)=DYDXU

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D2=AM1*V(2,M,N)/AN1+DU*(AM3-AM1*AN3/AN1)
U(2,M,N)=D2+DV*(AM4-AM1*AN4/AN1)
K1=M+1
K2=N
IF(I.NE.2) GOTO 4
K1=M
K2=N+1
4 D1=(AL(2,M,N)+AL(2,K1,K2))/2.
IK=4
IF(IKK.EQ.1) WRITE(6,1001) IK,M2,N2,C1,X1,TX,S
D1=I1/SQRT(D1)
D3=I1/SQRT(AL(2,K1,K2))
XX=-Y(2,K1,K2)/D3+X(2,K1,K2)-TX
IF(T.EQ.0.) XX=0.
X(2,M,N)=D1*XX/(TAN(S)+D1)+TX
Y(2,M,N)=(X(2,M,N)-TX)*TAN(S)
PSI(2,M,N)=PSI(1,M2,N2)
P(I,3)=XX
6 GOTO 9
J1=M-1
J2=N
LP=M
IF(I.NE.2) GOTO 8
J1=M
J2=N-1
LP=N
8 CONTINUE
IF(AL(2,J1,J2).LT.AL(2,M,N).OR.T.EQ.0.) GOTO 12
IF(IW.NE.0) GOTO 12
J2=N-1
IF(I.NE.2) GOTO 12
12 J1=M-1
DO=X(2,M,N)-P(I,LP)
DAL=AL(2,M,N)
IE=0
CALL FIND(KV,IW,X(2,J1,J2),Y(2,J1,J2),M3,N3,IE,I)
DU=U(1,M2,N2)+U(1,M3,N3)
DV=V(1,M2,N2)+V(1,M3,N3)
NF=1000000*(P(I,LP)-X(2,J1,J2))
IF(MP.LT.0) WRITE(6,1002) M,N,P(I,LP),X(2,J1,J2)
1002 FCORMAT(1X,'P-X,LT,0 :',2I4,2(2X,E12.5))
IF(MP.NE.0) GOTO 3
O1=0.
GOTO 2
3 D1=(AL(2,M,N)-AL(2,J1,J2))*DO
IK=3
IF(IKK.EQ.1) WRITE(6,1001) IK,IW,IE,CU,CV,CAL
2 C1=O1/(4.*DAL*(P(I,LP)-X(2,J1,J2)))
D2=AI*AK*AM*DO/DAL
D3=-0.25*AM*(AK*DO)**2/DAL
D4=1./SQRT(AL(2,M,N))
D5=-D3*D4
D6=1.*AK*AK*AM*PSI(2,J1,J2)*DO/CAL
D7=Y(2,M,N)-Y(2,J1,J2)
A(1,1)=1.+D1+D2+D3
A(1,2)=-(C4+D5)*I1
A(1,3)=C.
RI(1)=U(2,J1,J2)*(2.-A(1,1))-V(2,J1,J2)*(D4-D5)*I1+C6
A(2,3)=-2.*AM1/C7-AI*AM2
RI(2)=-2.*AM1/D7-AI*AM2)*G(2,J1,J2)-U(2,J1,J2)
RI(2)=RI(2)+(AM3*CU+AN4*DV)
A(2,1)=1.
A(2,2)=0.
A(3,3)=-2.*AN1/D7-AI*AN2
RI(3)=-2.*AN1/D7-AI*AN2)*G(2,J1,J2)-V(2,J1,J2)
RI(3)=RI(3)+(AN3*CU+AN4*CV)
A(3,1)=0.
A(3,2)=1.
C
NC=3
CALL SOLVE(NO)
U(2,M,N)=RHS(1)
V(2,M,N)=RHS(2)
G(2,M,N)=RHS(3)

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```

7    CU=U(2,M,N)+U(2,J1,J2)+(V(2,M,N)+V(2,J1,J2))*D4*I1
    PSI(2,M,N)=CU*0.5*DO+PSI(2,J1,J2)
9    RETURN
1001 FCRMAT(1X,3I4,5(2X,E11.4))
    END
C>XXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
    SLBROUTINE BOUNDU(IW,X,I,AK,L)
    CCMPLX*8 U,V,AI,PSI,G,CYDXU
    COMMON/BC/ T,B,DYDX,D2YCX2,DYDXU,AI,I1
C
C    UNSTEADY BOUNDARY - CCNCITIONS ALONG THE AIRFCIL
C
    IF(L.EQ.C) GOTO 10
    IF(L.EQ.1) GOTO 11
10   DYDXU=-1.-(X-B)*AK*AI
    GOTO 13
11   DYDXU=-AI*AK
13   RETURN
    END
C>XXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
    SLBROUTINE RANCB(M,N,AM,AK,I,L)
    CCMPLX*8 U,V,AI,PSI,G,CYCXU,D2,A,C
    COMMON/BC/ T,B,DYDX,D2YCX2,DYDXU,AI,I1
    CCMON V(2,25,25),X(2,25,25),P(2,25),U(2,25,25),PSI(2,25,25),
    FG(2,25,25),AL(2,25,25),Y(2,25,25)
C
C    BOUNDARY STEP OF THE UNSTEADY FIELD TO THE AIRFOIL
C
    CALL BCUNDU(IW,X(2,M,N),I,AK,L)
    V(2,M,N)=CYCXU
    J1=M
    J2=N-1
    IF(I.NE.2) GOTO 5
    J1=M-1
    J2=N
5    D1=0.5*(AL(2,M,N)-AL(2,J1,J2))
    DC=X(2,M,N)-X(2,J1,J2)
    C6=V(2,M,N)+V(2,J1,J2)
    D2=AI*AK*AM*DO*2.
    D3=-0.5*(AK*DO)**2*AM
    D4=AL(2,M,N)+AL(2,J1,J2)
    D5=SQRT(AL(2,M,N))+SQRT(AL(2,J1,J2))
    A=(D1+D2+D3)/D4
    U(2,M,N)=U(2,J1,J2)*(1.-A)-2.*(V(2,M,N)-V(2,J1,J2))*I1/D5
    D=2.*PSI(2,J1,J2)-DO*D6*I1/D5
    U(2,M,N)=U(2,M,N)+AK*AK*AM*DO*D/D4
    U(2,M,N)=U(2,M,N)/(1.+A)
    PSI(2,M,N)=U(2,M,N)+U(2,J1,J2)-2.*C6*I1/D5
    PSI(2,M,N)=PSI(2,J1,J2)+C0*PSI(2,M,N)/2.
    RETURN
    END
C>XXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
    SLBROUTINE GENU(M,N,I,AK,AM)
    CCMPLX*8 U,V,AI,PSI,G,CYDXU,D2,D6,B2,B6
    CCMPLX*16 A,RI,ES
    COMMON/BC/ T,B,DYCX,C2YCX2,DYDXU,AI,I1
    COMMON/SOL/ A(4,4),RI(4),ES(4)
    CCMON V(2,25,25),X(2,25,25),P(2,25),U(2,25,25),PSI(2,25,25),
    FG(2,25,25),AL(2,25,25),Y(2,25,25)
C
C    GENERAL STEP OF UNSTEADY FIELD
C
    J1=M
    J2=N-1
    J3=M-1
    J4=N
    IF(I.NE.2) GOTO 5
    J1=M-1
    J2=N
    J3=M
    J4=N-1

```

```

5 DC=X(2,M,N)-X(2,J1,J2)
D1=0.5*(AL(2,M,N)-AL(2,J1,J2))
D2=AI*2.*AK*AM*DO
D3=-0.5*AM*(AK*DO)**2
D4=AL(2,M,N)+AL(2,J1,J2)
D5=SQRT(AL(2,M,N))+SQRT(AL(2,J1,J2))
DE=(D1+D2+D3)/D4
A(1,1)=1.+DE
A(1,2)=2.*(1.-D3/C4)*I1/D5
RI(1)=U(2,J1,J2)*(1.-DE)+V(2,J1,J2)*I1*2.*(1.+D3/D4)/D5
RI(1)=2.*AK*AK*AM*PSI(2,J1,J2)*C0/D4+RI(1)
BC=X(2,M,N)-X(2,J3,J4)
B1=(AL(2,M,N)-AL(2,J1,J2))*B0/(2.*DO)
B2=2.*AI*AK*AM*B0
B3=-0.5*AM*(AK*B0)**2
B4=1./SQRT(AL(2,M,N))
B5=-0.5*B3*B4/AL(2,M,N)
BE=0.5*(B1+B2+B3)/AL(2,M,N)
A(2,1)=1.+BE
A(2,2)=- (B4+B5)*I1
RI(2)=U(2,J3,J4)*(1.-B6)+V(2,J3,J4)*I1*(-B4+B5)
RI(2)=PSI(2,J3,J4)*AM*BC*AK*AK/AL(2,M,N)+RI(2)

```

```

C
NO=2
CALL SCLVE(NO)
U(2,M,N)=ES(1)
V(2,M,N)=ES(2)
D2=V(2,J1,-2)+V(2,M,N)
D2=(U(2,J1,J2)+U(2,M,N)-I1*2.*D2/C5)*D0/2.
DE=(V(2,J3,J4)+V(2,M,N))*B4
D6=(U(2,J3,J4)+U(2,M,N)+I1*D6)*B0/2.
D2=PSI(2,J1,J2)+PSI(2,J3,J4)+D2+DE
PSI(2,M,N)=C2/2.
RETURN
END

```

```

CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
SUBROUTINE PRESS (M,N,J1)
COMMON/8 U,V,AI,PSI,G,CYDXU,PU
COMMON/BA/ ALO,ALC,AM,C1,AK,I,IE,IA
COMMON/BC/ T,B,DYDX,C2YCX2,DYDXU,AI,I1
COMMON V(2,25,25),X(2,25,25),P(2,25),U(2,25,25),PSI(2,25,25),
FG(2,25,25),AL(2,25,25),Y(2,25,25),Q(2,3),DX(2),IN(2,25),
FPX(2,20),PS(2,20),PU(2,2C)

```

```

CCCC
COMPUTATION OF THE PRESSURE - COEFFICIENTS ALCNG
THE AIRFOIL
PU=CPU, PS=CPS

K=N
IF(I.EQ.2) K=M
IF(K.EQ.J1) GOTO 5
PS(I,K)=-2.*(AL(2,M,N)-ALC)/(C1*AM)
PU(I,K)=-2.*(U(2,M,N)+AI*AK*PSI(2,M,N))
PX(I,K)=X(2,M,N)
GOTO 6
5
K1=K-1
K2=K-2
D2=PX(I,K1)-PX(I,K2)
D1=(1.-PX(I,K1))/C2
PS(I,K)=PS(I,K1)+(PS(I,K1)-PS(I,K2))*D1
PU(I,K)=PU(I,K1)+(PU(I,K1)-PU(I,K2))*D1
PX(I,K)=1.
6
RETURN
END

```

```

CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
SUBROUTINE LIFT (IUS,ILS,LC2)
COMMON/8 U,AI,CL,CM,CCL,DCM,PU,Q
COMMON/8 P,DYDX,CLS,DCLS,CMS
COMMON/BC/ T,B,DYDX,D2YCX2,DYDXU,AI,I1
COMMON X(2,100),U(2,100),Q(2,100),P(2,2),A(12850),
FPX(2,20),PS(2,20),PU(2,20)

```

C  
C  
C

COMPUTATION OF THE LIFT- AND MOMENTUM - COEFFICIENTS  
FOR BOTH OF THE SURFACES AND FOR THE COMPLETE AIRFOIL  
LIFT: CL,CLS ; MOMENTUM: CM,CMS

```

DC 4 I=1,2
I1=(-1)**(I+1)
K=IUS
IF(I.EC.2) K=ILS
DC 5 J=1,K
L=2*J-1
X(I,L)=FX(I,J)
U(I,L)=FU(I,J)
Q(I,L)=PS(I,J)
M=K-1
DO 7 J=1,M
N=2*J
L1=N-1
L2=N+1
X(I,N)=(X(I,L1)+X(I,L2))/2.
U(I,N)=(U(I,L1)+U(I,L2))/2.
Q(I,N)=(Q(I,L1)+Q(I,L2))/2.
CL=0.
CM=0.
CLS=0.
CMS=0.
N=2*K-2
DC 8 J=1,N
D1=X(I,J+1)-X(I,J)
D2=((X(I,J+1)+X(I,J))/2.-B)*I1
DCL=0.5*D1*(U(I,J)+U(I,J+1))
DCLS=0.5*D1*(Q(I,J)+Q(I,J+1))
CM=CM-DCL*D2
CMS=CMS-DCLS*D2
CL=CL-I1*DCL
CLS=CLS-I1*DCLS
P(I,1)=CLS
Q(I,1)=CMS
P(I,2)=CL
U(I,1)=CM
CLS=(P(2,1)+P(1,1))
CL=(P(2,2)+P(1,2))
CM=(U(2,1)+U(1,1))
CMS=(Q(2,1)+Q(1,1))
IF(LO2.EC.1) CM=AI*CM
IF(LO2.EC.1) CL=AI*CL
WRITE(6,1000)
WRITE(6,1001)
WRITE(6,1002) CL,CM
WRITE(6,1003) CLS,CMS
1000 FCRMAT(2X,'MOMENTUM- AND LIFT - COEFFICIENTS FOR A SINGLE'
F,/,2X,'AIRFOIL WITH A SURFACE DESCRIBED ABOVE:',/)
1001 FCRMAT(14X,'RCL',6X,'ICL',6X,'RCM',6X,'ICM',/)
1002 FCRMAT(2X,'UNSTEADY',4(2X,F7.4))
1003 FCRMAT(2X,'STEADY',4(2X,F7.4))
RETURN
END
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
SUBROUTINE WAKE(KV,WE,TA,TB,LO4,IUS,ILS,LO3,NX,PA)
COMPLEX*8 L,V,AI,PSI,CL,G,DYDXU,PL,AUU,ALL
COMMON/BA/ ALO,ALO,AM2,E1,AK,I,IE,Ih
COMMON/BC/ T,Z,DYCX,D2YCX,DYDXL,AI,I1,T4
COMMON/SP/ XS(2,50),YS(2,50),A(50,4),R(2,50)
COMMON V(2,25,25),X(2,25,25),P(2,25),U(2,25,25),FSI(2,25,25),
FG(2,25,25),AL(2,25,25),Y(2,25,25),Q(2,3),DX(2),IN(2,25),
FPX(2,20),PS(2,20),PU(2,20),TEF(2),TED(20)

```

C  
C  
C

COMPUTATION OF THE FIELD BEHIND THE AIRFOIL

```

IKK=0
IF(LO3.EC.0) GOTO 1
DXS=0.5*(DX(1)+DX(2))
EX=E1*AM2*1.5
E=E1-1.

```

1311

```

E5=E-1.
Ih=0
IE=0
C
C
COMPUTATION OF THE SPLINE COEFFICIENTS
DC 2 I=1,2
T=TA
IF(I.EQ.2) T=TB
I1=(-1)*I+1
DC 46 K=1,NX
A(K,1)=YS(I,K)
A(K,3)=R(I,K)
H=XS(I,K+1)-XS(I,K)
IS=46
IF(IKK.EQ.1) WRITE(6,2000) I9,K,I,JX,H,XS(I,K)
A(K,2)=(YS(I,K+1)-YS(I,K))/H-H*(R(I,K+1)+2.*R(I,K))/3.
46 A(K,4)=(R(I,K+1)-R(I,K))/(3.*H)
C
C
THE STEADY WAKE
XXS=1.
YYS=0.
CALL FIND(KV,IW,XXS,YYS,M,N,IE,I)
I9=2
IF(IKK.EC.1) WRITE(6,2000) I9,IW,M,N,X(1,M,N)
TE=2.*I1*(ALD-AL(1,M,N))*1.5)/3.
TE=TE/(AL(1,M,N)+E*AM2+1.)
2 TET(I)=ATAN(TE)
I9=2
IF(IKK.EC.1) WRITE(6,2000) I9,Ih,M,N,DYDX
M=2
N=1
XF1=1.
XP2=1.
4 I3=0
I9=4
IF(IKK.EQ.1) WRITE(6,2000) I9,IW,M,N,TET(1),TET(2)
DX1=DXS
DX2=DXS
DXU=1./(MA-1)
CXL=DXU
DC 6 I4=1,30
XP3=XP1+DX1
XP4=XP2+CX2
IF(IW.EQ.0) GOTC 36
N=1
IF(IW.LE.2) T4=TED(1)
IF(IW.LE.2) GOTC 28
I1=IW-1
DC 29 I2=1,II
29 IF(X(2,I2+1,I2).GT.XP3) GOTO 30
I9=29
30 IF(IKK.EQ.1) WRITE(6,2000) I9,IW,M,N,XP3
T4=TED(I2-1)
I9=30
28 IF(IKK.EQ.1) WRITE(6,2000) I9,IW,M,N,T4
I=1
I9=28
IF(IKK.EQ.1) WRITE(6,2000) I9,Ih,M,N,XP3,XP4
T=TA
I1=1
CALL SHOCK(KV,WE,LO4,M,N,XP3,EX,MA,DXU)
IF(IE.EC.1) GOTO 17
I=2
I1=-1
T=TB
IF(IW.LE.2) GOTO 35
CO 33 I2=2,II
33 IF(X(2,I2,I2+1).GT.XP4) GOTO 34
I9=33
34 IF(IKK.EC.1) WRITE(6,2000) I9,IW,M,N,XP4
35 T4=TED(I2-1)
I9=35

```

```

IF(IKK.EQ.1) WRITE(6,2000) I9,IW,M,N,T4
IF(IKK.EC.1) WRITE(6,2000) I,I1,IW,M,EX
CALL SHCCK(KV,WE,LO4,N,M,XP4,EX,MA,CXL)
IF(IKK.EQ.1) WRITE(6,2000) I9,IW,M,N,X(2,1,M)
IF(IE.EQ.1) GOTO 17
IF(IW.GT.1) GOTC 12
Ih=Iw+1
M=M+1
GOTO 4
12 IF(IKK.EQ.1) WRITE(6,2000) I9,Ih,M,N,X(2,M,1)
IF(IW.EC.2) GOTO 7
IX=Iw-1
DC 5 N=2,IX
I=1
I1=1
CALL GEN(M,N,I,IW)
I=2
I1=-1
5 CALL GEN(N,M,I,IW)
I9=5
IF(IKK.EQ.1) WRITE(6,2000) I9,IW,M,N,X(2,M,N)
DC 8 I2=1,2
IM=N+1
CALL SWITCH(IM,N,L1,L2,I2)
IS=11
IF(IKK.EC.1) WRITE(6,2000) I9,IW,L1,L2,AL(2,L1,L2)
I1=(-1)**(I2+1)
TE=2.*I1*(ALD-AL(2,L1,L2)**1.5)/3.
TE=TE/(AL(2,L1,L2)+E*AM2+1.)
8 TET(I2)=ATAN(TE)
36 I9=36
IF(IKK.EQ.1) WRITE(6,2000) I9,IW,M,N,TET(1),TET(2)
T1=TET(1)
T2=TET(2)
T3=0.5*(T1+T2)
CC 47 IA=1,2
IM=N+1
CALL SWITCH(IM,N,L1,L2,IA)
IF(N.GT.1) GOTO 53
I1=(-1)**(IA+1)
XXS=1.
YYS=0.
IE=0
53 CALL FIND(KV,IW,XXS,YYS,M2,N2,IE,IA)
IF(N.EQ.1) ALT=AL(1,M2,N2)
IF(N.GT.1) ALT=AL(2,L1,L2)
XM=(1.-(ALT-ALD)*E/E1)**(E5/E)
XM=2.*((1.+E5*AM2/2.)/XM-1.)/E5
IS=47
47 IF(IKK.EQ.1) WRITE(6,2000) I9,M,N,IA,XM,ALT
IF(IA.EQ.1) XM1=XM
IF(IA.EQ.2) XM2=XM
PII=((1.+E5*XM1/2.)/(1.+E5*XM2/2.))**((E/E5)
CALL KONST(Q(1,1),C(1,2),Q(1,3),XM1,E1)
CALL KONST(Q(2,1),Q(2,2),Q(2,3),XM2,E1)
DC 9 K=1,20
T4=T3-T1
T5=T3-T2
F=E*XM1*(Q(1,1)*T4+Q(1,2)*T4*T4+C(1,3)*T4**3)/2.+1.
FS=E*XM2*(-Q(2,1)*T5+Q(2,2)*T5*T5-Q(2,3)*T5**3)/2.+1.
F=F-FS*PII
FS=E*XM2*(Q(2,1)-2.*C(2,2)*T5+3.*C(2,3)*T5*T5)/2.*FII
FS=FS+E*XM1*(Q(1,1)+2.*C(1,2)*T4+3.*C(1,3)*T4*T4)/2.
D=F/FS
T3=T3-D
9 IF(ABS(D).LE.0.000001) GOTO 10
CCONTINUE
10 WRITE(6,1000) M,N,T3,C
I9=10
IF(IKK.EQ.1) WRITE(6,2000) I9,IW,M,N,T3
IF(IW.EQ.0) Iw=1
TED(IW)=T3
IF(IW.EQ.1) GOTC 4

```

```

7      N=N+1
      IC=7
      IF(IKK.EQ.1) WRITE(6,2000) I9,IW,M,N,X(2,M,N)
      T4=TED(IW)
      I=1
      T=TA
      I1=1
      CALL RAND(WE,LO4,M,N,E1)
      I1=-1
      I=2
      T=TB
      CALL RAND(WE,LO4,N,M,E1)
      D=X(2,M,N)-X(2,N,M)
      IF(ABS(D).LE.0.0001) GCTC 14
      IF(X(2,M,N).LT.X(2,N,M) .OR. I3.EC.1) GCTC 15
16     I3=2
      I9=16
      IF(IKK.EQ.1) WRITE(6,2000) I9,IW,M,N,X(2,M,N),D
      DX1=DX1*((X(2,N,M)-1.)/(X(2,M,N)-1.))**(1.+0.08*(N-2))
      GCTC 6
15     IC=15
      IF(IKK.EC.1) WRITE(6,2000) I9,IW,M,N,X(2,N,M),D
      IF(I3.EQ.2) GOTO 16
      I3=1
      DX2=DX2*((X(2,M,N)-1.)/(X(2,N,M)-1.))**(1.+0.08*(N-2))
6      CCNTINUE
      WRITE(6,1000) M,N,X(2,M,N),X(2,N,M),C
14     IS=14
      IF(IKK.EQ.1) WRITE(6,2000) I9,IW,M,N,X(2,M,N)
      X(2,M,N)=(X(2,M,N)+X(2,N,M))/2.
      X(2,N,M)=X(2,M,N)
      XP1=XP3
      XP2=XP4
      IW=IW+1
      M=M+1
      GCTC 4

C      THE UNSTEADY WAKE
C      THE INITIAL STEP
17     IW=IW-1
      I=1
      M=2
      N=1
      I1=1
      CALL CCEF1(KV,I1,M,N,M2,N2,AM1,AMZ,AM3,AM4,AN1,AN2,AN3,
      FAN4,S1)
      IF(IKK.EC.1) WRITE(6,1000) M2,N2,AM1,AMZ,AM3,AM4,AN1,AN2,
      FAN3,AN4,S1
      I=2
      I1=-1
      CALL CCEF1(KV,I1,N,M,M3,N3,AM5,AM6,AM7,AM8,AN5,AN6,AN7,
      FAN8,S2)
      IF(IKK.EC.1) WRITE(6,1000) M3,N3,AM5,AM6,AM7,AM8,AN5,AN6,
      FAN7,AN8,S2
      ALU=U(1,M2,N2)*(AM3-AM1*AN3/AN1)+V(1,M2,N2)*(AM4-AM1*AN4/AN1)
      ALL=U(1,M3,N3)*(AM7-AM5*AN7/AN5)+V(1,M3,N3)*(AM8-AM5*AN8/AN5)
      PSI(2,M,N)=PSI(1,M2,N2)
      PSI(2,N,M)=PSI(1,M3,N3)
      T4=TED(I)
      CL=(PSI(2,N,M)-PSI(2,M,N))*(1.-AM1*TAN(T4)/AN1)*AI*AK
      CL=ALL-AUU+CL
      V(2,N,M)=CL/(AM1/AN1-AM5/AN5)
      U(2,N,M)=ALL+AM5*V(2,N,M)/AN5
      U(2,M,N)=U(2,N,M)+AI*AK*(PSI(2,N,M)-PSI(2,M,N))
      V(2,M,N)=V(2,N,M)+TAN(T4)*(U(2,M,N)-U(2,N,M))
      G(2,M,N)=0.
      G(2,N,M)=0.
      DC 24 I=1,2
      I1=(-1)**(I+1)
      IM=M+1
      CALL SWITCH(M,N,L1,L2,I)
      CALL SWITCH(IM,N,K1,K2,I)
      S=S1

```



```

2000 FCRMAT(1X,4I4,2(2X,F8.3))
1004 FCRMAT(1H1)
1005 FCRMAT(1X,/)
1007 FCRMAT(1X,/)
1008 FCRMAT(2X,'POINT',7X,'X',10X,'Y',7X,'LAMBDA',7X,'RL',9X,'IU',
F9X,'RV',9X,'IV',8X,'RPSI',7X,'IPSI',8X,'RG',9X,'IG',/)
1009 FCRMAT(2X,I2,'',I2,11(3X,F8.5))
1010 FCRMAT(1X,'UPPER WAKE FIELD :')
1011 FCRMAT(1X,'LOWER WAKE FIELD :')
1012 FCRMAT(1X,F10.5)
1 RETURN
END

```

```

CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX

```

```

SUBROUTINE KONST(C1,C2,C3,AM2,E1)
ALC=AM2-1.
E=E1-1.
M5=E-1.
M2=E#E
M3=E2#E
M4=E2#E2
AM4=AM2#AM2
C1=2./SQRT(ALO)
C2=((AM2-2.)**2+E*AM4)/(2.*ALO**2)
C3=E1*(AM2-(5.+7.*E-2.*E2)/(2.*E1))**2/6.
C3=C3+(-4.*E4+28.*E3+11.*E2-8.*E-3.)/(24.*E1)
C3=C3-E1*((5.-3.*E)*AM4-(12.-4.*E)*AM2+8.)/48
C3=(C3*AM4+3.*(AM2-4./3.))**2/4.)/(ALO**3.5)
C3=0.
RETURN
END

```

```

CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX

```

```

SUBROUTINE SLIP(M3,N3)
CCOMPLEX*8 U,V,G,PSI,PU,AI,DYDXU,DA
CCOMPLEX*16 A,RI,ES
CCOMMON/BA/ ALO,ALC,AM,C1,AK,I,IE,IW
CCOMMON/BC/ T,B,DYCX,D2YCX2,DYDXU,AI,I1,T3
CCOMMON/SOL/ A(4,4),RI(4),ES(4)
CCOMMON V(2,25,25),X(2,25,25),P(2,25),U(2,25,25),PSI(2,25,25),
FG(2,25,25),AL(2,25,25),Y(2,25,25),Q(2,25),DX(2),IN(2,25),
FINZ(2,25),PX(2,20),PS(2,20),PU(2,20),TET(2),TEC(20)

```

C  
C  
C

SLIP DOES THE SIMULTANECUS STEPS FROM BOTH SIDES  
OF THE WAKE TO THE SLIP - LINE

```

M2=M3-1
N2=N3-1
M4=N3
N4=M3
M13=M3
N13=N3-1
M14=M4-1
N14=N4
DC3=X(2,M3,N3)-X(2,M13,N13)
DAL3=AL(2,M3,N3)+AL(2,M13,N13)
D13=SQRT(AL(2,M3,N3))+SQRT(AL(2,M13,N13))
D23=AL(2,M3,N3)-AL(2,M13,N13)
A(2,1)=1.+0.5*D23/DAL3+2.*AI*AK*AM*CO3/DAL3
A(2,1)=A(2,1)-AM*(AK*DC3)**2/DAL3*0.5
A(2,2)=2./D13+AM*(AK*DC3)**2/(CAL3#C13)
RI(2)=(2.-A(2,1))*U(2,M13,N13)+(4./D13-A(2,2))*V(2,M13,N13)
RI(2)=RI(2)+2.*AK*AK*AM*PSI(2,M13,N13)*DC3/DAL3
A(2,3)=0.
A(2,4)=C.
DC2=X(2,M3,N3)-X(2,M2,N2)
A(4,1)=-TAN(TED(N3))
A(4,2)=1.
A(4,3)=-A(4,1)
A(4,4)=-1.
RI(4)=0.
DC4=X(2,M4,N4)-X(2,M14,N14)
D14=SQRT(AL(2,M4,N4))+SQRT(AL(2,M14,N14))
A(1,1)=1.+AI*AK*DC3/2.

```







```

48      Y(1,1,M,N)=Y(1,1,K,N)+XM*D/2.
        DC 44 M=1,KV
        DC 44 N=1,20
        AL(1,1,M,N)=ALO
        PSI(1,1,M,N)=0.
        G(1,1,M,N)=0.
        U(1,1,M,N)=0.
44      V(1,1,M,N)=0.
C
C      CCMPUTATION OF THE SPLINE - COEFFICIENTS
C
        DC 49 I=1,2
        DC 49 K=1,NX
        A(I,K,1)=YS(I,K)
        A(I,K,3)=R(I,K)
        H=XS(I,K+1)-XS(I,K)
49      A(I,K,2)=(YS(I,K+1)-YS(I,K))/H-F*(R(I,K+1)+2.*R(I,K))/3.
        A(I,K,4)=(R(I,K+1)-R(I,K))/(3.*F)
C
        EI=cos(AMY)+AI*SIN(AMY)
55      TX=EM/TAN(ET)
        TY=EM
        IR=1
        IB=2
        M1=0
        N1=0
        WRITE(6,1021) IM
        XPL=IM*TX
        YPL=IM*TY
        WRITE(1,1026)XPL,YPL,IM
        I=0
        IE=0
30      CALL FIND(KV,IM,IB,IR,IX,TY,M1,N1)
        I=I+1
        IF(I.EQ.1) GOTO 67
        T2Y=2.*EM
        DC 56 J=2,KV
        IF(Y(2,1,J,1).GT.T2Y) GOTO 57
56      CCNTINUE
57      MAX2=J
C
C      THE STEADY FLOW FIELD
C
67      I2=0
16      I2=I2+1
        IF(I2.GT.30) GOTO 101
        WRITE(6,1024) I2,MA,J1,CX(I)
1024     FCRMAT(IX,I3,' ITERATION ALONG THE SURFACE',2I4,2X,ES.3)
        MA=MAX
        XX=EM*(IB-1)/TAN(ET)
        YY=EM*(IB-1)
        IF(I.EQ.2 .OR. IR.GT.1) MA=MAX2
        IF(IR.GT.1) I=IB
        I1=(-1)**(I+1)
        MZ=1
        IBACK=0
        IC=1
        KC=0
        IF((I*IM).EQ.1 .AND. IR.GT.1) IQ=C
        T=T1
        IF(I.EQ.2) T=T2
        MM=KV
        IF(I.EQ.2 .OR. IR.GT.1) MM=20
        DC 5 K=1,20
        DC 5 J=K,MM
        CALL SWITCH(J,K,M,N,I)
5        AL(IB,IR,M,N)=0.
        M2=M1
        N2=N1
C
C      THE STEADY BOUNDARY-PROPERTIES BEHIND THE SHOCK
C
        IE=0
        KW=MM+1

```

```

DC 9 J=2,KW
KT=J-1
XP=DX(I)*(J-2)+TX-XX
IF(T.EQ.0 .OR. IQ.EQ.0) XP=TX-XX
IF(IQ.NE.0 .AND. XP.GT.1 .AND. I.EC.1) GOTO 8
IF(J.GT.MM) GOTO 8
LI=1
CALL SWITCH(J,LI,M,N,I)
CALL SHOCK(MM,LO4,EM,IR,IB,M,N,XP,M2,N2,C3)
IF(IR.GT.1 .AND. IB.EC.1 .AND. Y(IE,IR,M,N).EQ.EM) GOTO 63
IF(I.EQ.1 .AND. IR.EQ.1) GOTO 9
IF(IE.EQ.2 .AND. Y(IE,IR,M,N).EQ.C.) GOTO 63
IF(X(IE,IR,M,N).LE.1.) GOTO 9
ALO=-Y(IE,IR,M,N)/(1.-X(IE,IR,M,N))
ALO=1./(ALO*ALO)
IF(ALO.GT.ALO .AND. IM.NE.1) GOTO 8
IF(X(IE,IR,M,N).GT.1.) GOTO 8
GOTO 9
63 KT=KT+1
GOTO 8
9 CCNTINUE
8 CCNTINUE
IF(X(IE,IR,2,1).GE.1.) GOTO 38
IF(I.EQ.2 .AND. X(IE,IR,M,N).GE.1.) KC=2
C
C
C ALL OTHER STEPS OF THE STEADY FLOW FIELD
KA=KT
DC 1 J=2,KA
J1=J
IF(J.EQ.KA .AND. I.EQ.2 .AND. IR.EC.1) GOTO 15
IF(J.EQ.KA .AND. IR.GT.1) GOTO 15
LI=J+1
CALL SWITCH(LI,J,M,N,I)
CALL RANC(LO4,IE,IR,M,N,C3)
IF(X(IE,IR,M,N).GT.TX1) MZ=0
IF(IR.GT.1 .AND. I.EQ.1 .AND. X(IE,IR,M,N).GT.1 .AND. IC.NE.C) MZ=0
IF(MZ.EC.0 .AND. IR.GT.1) KC=1
KE=KA+1
DC 3 J2=J,K8
IF(J2.EC.K8) GOTO 3
LI=J-1
CALL SWITCH(J2,LI,K1,K2,I)
IF(AL(IE,IR,K1,K2).GT.AL(IE,IR,M,N)) GOTO 6
3 CCNTINUE
6 J2=J2-1
IR(I,J-1)=J2
L=J+2
DC 7 K=L,J2
7 CALL SWITCH(K,J,M,N,I)
CALL GEN(IE,IR,M,N,I)
IF(MZ.EC.0) GOTO 15
IF(J2.EQ.KA) GOTO 1
C
C
C CHARACTERISTICS OVERTAKE THE SHOCK
J3=M+1
J4=N-1
J5=M
J6=N-1
IF(I.NE.2) GOTO 12
J3=M-1
J4=N+1
J5=M-1
J6=N
12 D1=Y(IE,IR,J3,J4)-Y(IE,IR,J5,J6)
D1=D1/(X(IE,IR,J3,J4)-X(IE,IR,J5,J6))
D2=I1/SQRT(AL(IE,IR,M,N))
K1=M+1
K2=J
IF(I.NE.2) GOTO 14
K1=J
K2=N+1
14 AL(IE,IR,K1,K2)=AL(IE,IR,M,N)

```

```

D3=-D2*X(IB,IR,M,N)+D1*X(IB,IR,J5,J6)
D3=D3-Y(IB,IR,J5,J6)+Y(IB,IR,M,N)
X(IB,IR,K1,K2)=D3/(D1-C2)
Y(IB,IR,K1,K2)=D2*(X(IB,IR,K1,K2)-X(IB,IR,M,N))+Y(IB,IR,M,N)
IF(I.EQ.2) J3=J4
DC 11 K=J3,KA
IF(K.EQ.MM) GOTO 11
LI=K+1
CALL SWITCH(LI,J,M,N,I)
LI=J-1
CALL SWITCH(K,LI,K1,K2,I)
X(IB,IR,M,N)=X(IB,IR,K1,K2)
Y(IB,IR,M,N)=Y(IB,IR,K1,K2)
AL(IB,IR,M,N)=AL(IB,IR,K1,K2)
11 CCNTINUE
IF(KA.LT.MM) KA=KA+1
IF(MA.GT.IN(I,J-1)) MA=MA+1
CCNTINUE

AUTOMATIC STEP-SIZE CONTROL
15 J2=J1
IF(KC.EC.1) J2=KT
IF(T.EQ.0 .OR. IC.EC.0) GOTO 28
IF(J2.EQ.MA) GOTO 28
IF(J2.GT.MA) GOTO 68
DIF=(J2*1.)/(MA*1.)
IF(DIF.GE.0.85) GOTO 69
GOTO 72
68 DIF=(J2*1.)/(MA*1.)
IF(DIF.LE.1.15) GOTO 69
72 DX(I)=DX(I)*DIF
GOTO 73
69 IF(J2.GT.MA) DX(I)=DX(I)*1.05
IF(J2.LT.MA) DX(I)=DX(I)*0.96
73 CCNTINUE
IF(J2.NE.MA) GOTO 16
C
28 CCNTINUE
IF(IR.EQ.1 .AND. I.EC.1) GOTO 59
IF(KC.NE.0) GOTO 59
C
CC
ADDITIONAL POINTS FOR THE FIELDS IN THE PASSAGE
DC 33 LY=2,20
LI=LY-1
CALL SWITCH(LY,LI,M,N,I)
IF(AL(IB,IR,M,N).EQ.0.) GOTO 35
C
33 WRITE(6,1005) M,N,X(IB,IR,M,N),Y(IB,IR,M,N),AL(IB,IR,M,N)
35 CCNTINUE
KC=LY
LI=KD-1
LZ=2
CALL SWITCH(LI,LZ,M,N,I)
LE=1
CALL SWITCH(KD,LE,L1,L2,I)
X(IB,IR,L1,L2)=X(IB,IR,M,N)
CALL SWITCH(LI,LE,M,N,I)
LZ=KD-2
CALL SWITCH(LI,LZ,L3,L4,I)
Y1=I1*(X(IB,IR,L1,L2)-X(IB,IR,M,N))/SQRT(AL(IB,IR,M,N))+EM*(IE-2)
X1=X(IB,IR,L1,L2)
Y(IB,IR,L1,L2)=Y1
D1=(Y(IB,IR,L3,L4)-Y(IB,IR,M,N))/(X(IB,IR,L3,L4)-X(IB,IR,M,N))
D1=D1/1.5
DC 43 JY=2,LZ
CALL SWITCH(KD,JY,M,N,I)
CALL SWITCH(LI,JY,L1,L2,I)
D2=I1/SQRT(AL(IB,IR,L1,L2))
X2=D1*X1-Y1+Y(IB,IR,L1,L2)-D2*X(IE,IR,L1,L2)
43 X(IB,IR,M,N)=X2/(D1-D2)
Y(IB,IR,M,N)=D1*(X(IB,IR,M,N)-X1)+Y1
CALL SWITCH(KC,LI,M,N,I)
Y(IB,IR,M,N)=EM*(IB-1)

```

```

X(IB,IR,M,N)=(Y(IB,IR,P,N)-Y1)/C1+X1
DC 65 JY=1,LZ
CALL SWITCH(LI,JY,P,N,I)
CALL SWITCH(KD,JY,L1,L2,I)
C
C
65 WRITE(6,1009) JY,KD,X(IE,IR,M,N),Y(IB,IR,P,N),X(IB,IR,L1,L2),
FY(IB,IR,L1,L2)
CCONTINUE
CALL SWITCH(KC,LI,M,N,I)
WRITE(6,1009) M,N,X(IB,IR,P,N),Y(IB,IR,P,N)
C
C
C
C
59 THE UNSTEADY FLOW FIELD
THE UNSTEADY BOUNDARY PROPERTIES BEHIND THE SHOCK

KA=KT
LZ=2
LI=1
CALL SWITCH(LZ,LI,I7,IE,I)
XFL=X(IB,IR,I7,I8)+(IM-1)*EM/TAN(ET)
YFL=Y(IB,IR,I7,I8)+(IM-1)*EM
WRITE(1,1026) XFL,YFL,IM
CALL SWITCH(KT,LI,I7,IE,I)
XFL=X(IB,IR,I7,I8)+(IM-1)*EM/TAN(ET)
YFL=Y(IB,IR,I7,I8)+(IM-1)*EM
WRITE(1,1026) XFL,YFL,IM
1025 WRITE(6,1025) IR
FORMAT(IX,I3,' UNSTEADY FIELD')
IE=0
DC 18 J=2,KA
CALL SWITCH(J,LI,M,N,I)
IF(X(IB,IR,P,N).LT.0.) GOTO 18
CALL RANDS(MM,IB,IR,M,N,LO2,AMY,M2,N2)
18 CCONTINUE
IF(I.EQ.2) X(IB,IR,1,2)=TX
IF(I.EQ.2) Y(IB,IR,1,2)=TY
IF(I.EQ.1) X(IB,IR,2,1)=TX
IF(I.EQ.1) Y(IB,IR,2,1)=TY
C
C
C
ALL OTHER STEPS OF THE UNSTEADY FIELD
IF(J1.EQ.2) GOTO 38
J5=J1-1
DC 20 J=2,J5
LI=J+1
CALL SWITCH(LI,J,M,N,I)
CALL RANCB(IB,IR,M,N,LC2)
IF(J.EQ.J5 .AND. I.EQ.2 .AND. IR.EQ.1) GOTO 20
IF(J.EQ.J5 .AND. IR.GT.1 .AND. M2.NE.0) GOTO 20
L=J+2
J2=IN(I,J-1)
DC 22 K=L,J2
CALL SWITCH(K,J,M,N,I)
22 CALL GEAU(IM,IB,IR,M,N,I,AK,AM)
IF(J2.EQ.KA) GOTO 20
IE=0
K1=M+1
K2=J
K3=K1
IF(I.NE.2) GOTO 24
K1=J
K2=N+1
K3=K2
24 P(I,K3)=X(IB,IR,M,N)
CALL RANDS(MM,IB,IR,K1,K2,LO2,AMY,M2,N2)
J3=J2+1
DC 29 K=J3,KA
IF(K.EQ.MM) GOTO 20
LI=K+1
CALL SWITCH(LI,J,M,N,I)
LI=J-1
CALL SWITCH(K,LI,K1,K2,I)
U(IB,IR,P,N)=U(IB,IR,K1,K2)
V(IB,IR,P,N)=V(IB,IR,K1,K2)
G(IB,IR,P,N)=G(IB,IR,K1,K2)
29 PSI(IB,IR,P,N)=PSI(IB,IR,K1,K2)

```

```

20 IF(J2.LT.KA .AND. KA.LT.MM) KA=KA+1
   CC
   CC
   C  OUTPUT STEADY AND UNSTEADY FIELD
     IF(IM.LT.IBLA) GOTO 38
     IF(LOI.NE.0) GOTO 38
     WRITE(7,1004)
     IF(I.EQ.1 .AND. IR.EQ.1) WRITE(7,1002) IM
     IF(I.EQ.2 .OR. IR.GT.1) WRITE(7,1003) IR,IM
     WRITE(7,1007)
     WRITE(7,1001) AK,AX,C,B,CX(I),T
     DC 25 J=2,J1
     K=J-1
   C  LL=IN(I,K)+1
     LL=IN(I,K)
     WRITE(7,1007)
     WRITE(7,1008)
     N=K
     IF(I.NE.2) GOTO 26
     M=K
   31 DC 31 N=J,LL
     WRITE(7,1005) M,N,X(IB,IR,M,N),Y(IB,IR,M,N),AL(IE,IR,M,N)
     F,U(IB,IR,M,N),V(IB,IR,M,N),PSI(IB,IR,M,N),G(IB,IR,M,N)
     GOTO 17
   26 CCNTINUE
     DC 32 M=J,LL
   32 WRITE(7,1005) M,N,X(IB,IR,M,N),Y(IE,IR,M,N),AL(IE,IR,M,N)
     F,U(IB,IR,M,N),V(IE,IR,M,N),PSI(IE,IR,M,N),G(IE,IR,M,N)
   17 WRITE(7,1005)
   25 CCNTINUE
     WRITE(7,1004)
   38 CCNTINUE
     IF(IR.GT.1) GOTO 60
     IF(I.EQ.1) IUS=K-1
     IF(I.EQ.2) ILS=K-1
     IF(I.EQ.1) GOTO 30
     IF(IM.EC.1 .AND. LC3.EC.C) GOTO 62
   C  CHANGING THE COUNTERS FOR THE COMPLETION OF
   CC THE FIELDS BEHIND THE REFLECTED SPCKCS
   C 60 IF=IR+1
     IF(IR.GE.4) GOTO 62
     IF(KC.NE.0) GOTO 62
     WRITE(6,1023) IR
     LI=1
     CALL SWITCH(KT,LI,M,N,I)
     CALL SWITCH(KT,LI,M1,N1,I)
     TX=X(IB,IR-1,M,N)
     TY=Y(IB,IR-1,M,N)
     IF(IR.EC.2 .OR. IR.EC.4) IB=1
     IF(IR.EQ.3 .OR. IR.EQ.5) IB=2
     GOTO 67
   62 CCNTINUE
     IF(KC.EC.0) GOTO 77
     CCMPUTATION OF THE BLACE - WAKE IN
     SLBROUTINE *WAKE*
   C  CALL WAKE
   CC NEXT BLADE
   C 77 IM=IM+1
     IF(IM.LE.MAXS) GOTO 58
     IR=IR-1
     CALL PRESS(AMY,IR,EM,LC2,AX,C)
     IP=0
     WRITE(1,1026) XPL,YPL,IP
     GOTO 100
   C  TRANSFORMATION TO PROFIL 1
   CC
   58 XX=EM/TAN(ET)
     DC 70 IL=1,2
     MN=20

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IF(IL.EC.1) MM=KV
OC 51 M=1,MM
OC 51 N=1,20
X(1,IL,M,N)=X(2,IL,M,N)-XX
Y(1,IL,M,N)=Y(2,IL,M,N)-EM
PSI(1,IL,M,N)=PSI(2,IL,M,N)
U(1,IL,M,N)=U(2,IL,M,N)
V(1,IL,M,N)=V(2,IL,M,N)
G(1,IL,M,N)=G(2,IL,M,N)
AL(1,IL,M,N)=AL(2,IL,M,N)
51
70 CCNTINUE
GOTO 55
101 IM=-1
WRITE(1,1026) XPL,YPL,IM
999 FCRMAT(4I1,3I2)
1000 FCRMAT(3F10.5,2E10.3)
1001 FCRMAT(1X,/,/,1X,'W*C/U= ',F5.3,', M= ',F4.2,', K= ',F4.2,
F', B/C= ',F4.2,', CX= ',E9.3,', T/C= ',F6.4,/,/)
1002 FCRMAT(1X,' THE PROPERTIES AT THE MESHPOINTS OF THE FLOW FIELD',
F/,1X,' FOR THE UPPER SURFACE OF THE ',I2,' BLADE',/,/)
1003 FCRMAT(1X,' THE PROPERTIES AT THE MESHPOINTS OF THE FLOW FIELD',
F/,1X,' BEHIND THE ',I2,' SHOCK IN THE ',I2,' PASSAGE',/,/)
1004 FCRMAT(1F1)
1005 FCRMAT(1X,/)
1007 FCRMAT(1X,/,/)
1008 FCRMAT(2X,' POINT',7X,'X',10X,'Y',7X,' LAMBDA',7X,'RU',5X,'IU',
F9X,'RV',5X,'IV',8X,'RPSI',7X,'IPSI',8X,'RG',9X,'IG',/,/)
1009 FCRMAT(2X,I2,', ',I2,11(3X,F8.5))
1020 FCRMAT(4F10.5)
1021 FCRMAT(2X,I2,' BLADE')
1023 FCRMAT(1X,I2,' REFLECTION',/)
1026 FCRMAT(2F10.3,I3)
END
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
SLBRoutine PROFIL(LG4,T1,T2,N)
COMMON/SP/ XS(2,50),YS(2,50),A(2,50,4),R(2,50)
C
C PREPARATION OF THE PROFIL - SURFACES
C
READ(3,1000) T1,T2,NT,IKP
IF(LG4.EC.0) READ(3,1001) SP
DC 8 J=1,2
I2=0
N=NT
T=T1
IF(J.EQ.2) T=T2
IF(LG4.NE.0) GOTO 12
DC 9 K=1,4
M=(K-1)*5+1
READ(3,1001) XS(J,M),XS(J,M+1),XS(J,M+2),XS(J,M+3),XS(J,M+4)
READ(3,1001) YS(J,M),YS(J,M+1),YS(J,M+2),YS(J,M+3),YS(J,M+4)
IF(YS(J,M+4).EQ.100.) GOTO 7
S
7 CCNTINUE
CCNTINUE
DC 2 LL=1,N
XS(J,LL)=XS(J,LL)/SP
YS(J,LL)=YS(J,LL)/SP
2 IF(LG4.EC.0) GOTO 3
12 DX=1.0/(N-1)
I1=(-1)**(J+1)
DC 20 K=1,N
XS(J,K)=(K-1)*DX-0.25
YS(J,K)=I1*4.*T*XS(J,K)*(1.-XS(J,K))
20
C
C TC ENTER THIS PART,THE SURFACES SHOULD ALREADY BE
C GIVEN POINTWISE
C CCNTINUE
C
C INTERPOLATION THROUGH CUBIC SPLINES
C
IF(T.NE.0 .OR. LG4.EC.0) GOTO 51
DC 52 I=1,N
A(I,I,1)=0.
52

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51 GCTO 53
CCCONTINUE
DC 10 I=1,N
A(J,I,3)=XS(J,I)
10 A(J,I,4)=YS(J,I)
CCC
C MATRIX OF COEFFICIENTS AND RIGHT-HAND SIDES
C
K=N-2
CC 25 I=1,K
A(J,I,1)=A(J,I+1,3)-A(J,I,3)
A(J,I,3)=A(J,I+2,3)-A(J,I+1,3)
A(J,I,2)=2.*(A(J,I,1)+A(J,I,3))
A(J,I,4)=3.*(A(J,I+2,4)-A(J,I+1,4))/A(J,I,3)-
F3.*(A(J,I+1,4)-A(J,I,4))/A(J,I,1)
25 CCCONTINUE
A(J,1,1)=0.0
A(J,N-2,3)=0.0
CCC
C THE STEP OF GAUSS
C
K=N-3
DC 30 I=1,K
DC 35 M=3,4
35 A(J,I,M)=A(J,I,M)*A(J,I+1,1)/A(J,I,2)
A(J,I,1)=0.0
A(J,I,2)=A(J,I+1,1)
30 A(J,I+1,2)=A(J,I+1,2)-A(J,I,3)
A(J,I+1,4)=A(J,I+1,4)-A(J,I,4)
CCC
C SOLUTION
C
A(J,1,1)=0.
A(J,N,1)=0.
A(J,N-2,1)=0.0
L=N-1
DC 40 I=2,L
K=N-1
M=K+1
40 A(J,M,1)=(A(J,K,4)-A(J,K,3)*A(J,M+1,1))/A(J,K,2)
53 CCCONTINUE
DC 11 M=1,N
11 R(J,M)=A(J,M,1)
IF(IZ.EQ.1) GOTO 8
KV=N-1
DC 49 K=1,KV
A(J,K,1)=YS(J,K)
A(J,K,3)=R(J,K)
H=XS(J,K+1)-XS(J,K)
49 A(J,K,2)=(YS(J,K+1)-YS(J,K))/H+*(R(J,K+1)+2.*R(J,K))/3.
A(J,K,4)=(R(J,K+1)-R(J,K))/(3.*H)
DX=1./49.
DC 4 M=1,50
X=DX*(M-1)
DC 5 K=2,50
I=K-1
IF(XS(J,K).GE.X) GOTO 6
5 CONTINUE
6 H=X-XS(J,I)
Y=A(J,I,1)+A(J,I,2)*H+A(J,I,3)*H*H+A(J,I,4)*H*H*H
R(J,M)=X
4 YS(J,M)=Y
CCCONTINUE
IZ=1
N=50
13 DC 13 I=1,50
XS(J,I)=R(J,I)
8 GCTO 3
CONTINUE
IF(IKP.EQ.0) GOTO 17
15 DC 15 M=1,N
WRITE(6,1002) XS(1,M),YS(1,M),XS(2,M),YS(2,M)
17 N=N-1
1000 FCRMAT(2F10.5,2I2)

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1001  FCRMAT(5F10.5)
1002  FCRMAT(5X,4(E12.5,2X))
      RETURN
      END
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
      SLBROUTINE FIND(KV,IM,IE,IR,XS,YS,M,N)
      CCMPLX*8 U,V,PSI,G,DYCX
      CCMPLX*8 AI,EI
      REAL*4 X,Y,AL,XS,YS,P,XS
      CCMON/BC/ T1,T2,T,B,DYCX,D2YDX2,CYDXU,AI,EI,IY,IBACK,IE
      CCMON V(2,3,50,20),X(2,3,50,20),P(2,50),U(2,3,50,20),
      FPSI(2,3,50,20),G(2,3,50,20),AL(2,3,50,20),Y(2,3,50,20),
      FQ(2,50),DX(2),IN(2,50),IN2(2,50)
C
C
C      FIND LOOKS FOR THE MESHINDEX,RESPONSIBEL FOR
      THE POINT XS;YS
C
      NX7=10000*XS
      NY7=10000*YS
      IC=1
      IC=1
      IF(IR.EQ.1) GOTO 26
      IC=IR-1
      IF(IB.EQ.2) IC=1
      IF(IB.EQ.1) IC=2
26      I1=(-1)**(IC+1)
      IF(IE.EQ.1) GOTO 19
      IKK=0
C      IF(I.EQ.2) IKK=1
1004      IF(IKK.EQ.1) WRITE(6,1004)
      FCRMAT(1X,'FIND-ENTRY')
      KA=0
      L1=1
      L2=1
21      IK=21
      IF(IKK.EQ.1) WRITE(6,1003) IK,M,N,NX7,NY7,NX9,NYS
      DC 22 I2=L1,20
      L3=I2+L2
      CALL SWITCH(L3,I2,M,N,IC)
      IK=22
      IF(IKK.EQ.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NYS
      IF(IC.EQ.1 .AND. L3.GE.KV) GOTO 13
      IF(IC.EQ.2 .AND. L3.GT.19) GOTO 13
      J1=M+1
      J2=N+1
24      KM=J1-J2
      IK=24
      IF(IKK.EQ.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NYS
      IF(IC.EQ.2) KM=J2-J1
      NX8=10000*X(IC,ID,M,N)
      NX9=10000*X(IC,ID,J1,J2)
      NY8=10000*Y(IC,ID,M,N)
      NY9=10000*Y(IC,ID,J1,J2)
      XS=X(IC,ID,J1,J2)
      IF(NX7.GE.NX8 .AND. NX7.LT.NX9) GOTO 23
      IF(X9.GE.1. .AND. KM.EQ.1 .AND. IR.EQ.1) KA=J2
      IF(IC.EQ.2 .AND. KA.EQ.J2) KA=J1
      IF(IC.EQ.1 .AND. KA.NE.J2) GOTO 22
      IF(IC.EQ.2 .AND. KA.NE.J1) GOTO 22
      IF(IC.EQ.1) J1=J1+1
      IF(IC.EQ.2) J2=J2+1
      IF(J1.EQ.KV .OR. J2.EQ.20) GOTO 13
      GOTO 24
22      CCNTINUE
23      CCNTINUE
      IF(IC.EQ.2) GOTO 28
      IK=23
      IF(IKK.EQ.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NYS
      NY8=10000*Y(IC,ID,J1,N)
      IF(NY7.LT.NY8) GOTO 1
      IF(N.EQ.1) GOTO 13
      L2=L2+1
      L1=N-1

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28 GCTO 21
   IK=28
   IF(IKK.EC.1) WRITE(6,1003) IK,M,N,NX8,NYE,NX9,NY9
   NY8=10000*Y(IC,ID,M,J2)
   IF(NY7.GT.NY8) GOTO 1
   IF(M.EQ.1) GOTO 13
   L2=L2+1
   L1=M-1
1  GCTO 21
   IK=1
   IF(IKK.EQ.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NY9
   IF(M.GE.KV) GOTO 13
   J=M+1
   K=N+1
   NX8=10000*X(IC,ID,M,N)
   NY8=10000*Y(IC,ID,M,N)
   IF(NX7.EQ.NX8 .AND. NY7.EQ.NY8) GCTO 14
   IF(IR.GT.1) GOTO 29
C  ORIENTATION IN THE FIELD OVER THE UPPER BLADE
C
   IF(NX7.NE.NX8) GOTO 5
   IF(NY7.GT.NY8 .AND. N.EC.1) GOTO 13
   IF(NY7.GT.NY8) N=N-1
   IF(NY7.LT.NY8) M=M-1
5  GCTO 1
   IK=5
   IF(IKK.EQ.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NY9
   I=M+2
   IF(I.GT.KV) GOTO 7
   NX8=10000*X(IC,ID,J,N)
   NX9=10000*X(IC,ID,I,K)
   NY8=10000*Y(IC,ID,J,N)
   NY9=10000*Y(IC,ID,I,K)
   IF(NX8.EQ.NX9 .AND. NY8.EQ.NY9) GCTC 6
   GOTO 7
6  IK=6
   IF(IKK.EQ.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NY9
   D1=(Y(IC,ID,J,N)-Y(IC,IC,M,N))/(X(IC,ID,J,N)-X(IC,IC,M,N))
   D2=(Y(IC,ID,M,K)-Y(IC,IC,M,N))/(X(IC,ID,M,K)-X(IC,IC,M,N))
   D12=(YS-Y(IC,ID,M,N))/(XS-X(IC,ID,M,N))
   D5=(Y(IC,ID,J,K)-YS)/(X(IC,ID,J,K)-XS)
   D3=(Y(IC,ID,J,K)-Y(IC,IC,M,K))/(X(IC,ID,J,K)-X(IC,IC,M,K))
   IF(D12.LE.D1 .AND. D12.GE.D2 .AND. D5.GE.D1 .AND. D5.LT.C3) GCTO 14
   N=N+1
   GCTO 1
7  IK=7
   IF(IKK.EQ.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NY9
   D1=(Y(IC,ID,J,N)-Y(IC,IC,M,N))/(X(IC,ID,J,N)-X(IC,IC,M,N))
   IF(K.EQ.M) GOTO 4
   D2=(Y(IC,ID,M,K)-Y(IC,IC,M,N))/(X(IC,ID,M,K)-X(IC,IC,M,N))
4  IF(K.EQ.M) D2=0.
   D5=(Y(IC,IC,J,K)-Y(IC,IC,M,N))/(X(IC,ID,J,K)-X(IC,ID,M,N))
   D12=(YS-Y(IC,ID,M,N))/(XS-X(IC,ID,M,N))
   IF(D12.LE.D1 .AND. D12.GE.D2 .AND. XS.GT.X(IC,IC,M,N)) GCTC 8
   NX8=10000*X(IC,ID,M,N)
   IF(NX7.GT.NX8) GOTO 2
   IF(N.EC.1) GOTO 13
   N=N-1
   M=M-1
   GCTO 3
2  IF(D12.GT.D1 .AND. N.EQ.1) GCTC 13
   IK=2
   IF(IKK.EQ.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NY9
   IF(D12.GT.D1) N=N-1
   IF(D12.LT.D2) M=M-1
3  IBACK=IBACK+1
   J=M+1
   IF(AL(IC,ID,M,N).NE.AL(IC,ID,J,N)) GCTO 13
   GCTO 1
8  IK=8
   NX8=10000*X(IC,ID,J,K)
   NY8=10000*Y(IC,ID,J,K)
   IF(IKK.EQ.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NY9

```

```

IF(NX7.EQ.NX8 .AND. NY7.EQ.NY8) GCTC 12
IF(NX7.NE.NX8) GOTO 15
IF(NY7.GT.NY8) M=M+1
IF(NY7.LT.NY8) N=N+1
GCTO 1
12 IK=12
IF(IKK.EQ.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NYS
M=J
N=K
GCTO 14
15 IK=15
IF(IKK.EQ.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NYS
IF(NX7.GE.NX8) GOTO 16
IF(K.EQ.M) GOTO 10
10 D3= (Y(IC, ID, J, K) - Y(IC, IC, M, K)) / (X(IC, ID, J, K) - X(IC, ID, M, K))
IF(K.EQ.M) D3=0.
D4= (Y(IC, ID, J, K) - Y(IC, IC, J, N)) / (X(IC, ID, J, K) - X(IC, ID, J, N))
D34= (Y(IC, ID, J, K) - YS) / (X(IC, IC, J, K) - XS)
IF(D3.GE.D34 .AND. C4.LE.D34) GCTC 17
GCTO 16
17 IF(D4.EQ.D34) M=M+1
IF(D3.EQ.0. .AND. D34.EC.0.) GCTC 14
IF(D3.EQ.D34) N=N+1
IK=17
IF(IKK.EQ.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NYS
GCTO 14
16 IK=16
IF(IKK.EC.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NYS
IF(D12.GT.D5) M=M+1
IF(D12.LT.D5) N=N+1
IF(D12.EQ.D5) GOTO 9
GCTO 1
9 M=M+1
IK=9
IF(IKK.EC.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NYS
N=N+1
GCTO 1
13 IE=1
GCTO 19
C
C
C
29 ORIENTATION IN THE PASSAGE
IK=29
IF(IKK.EC.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NYS
L1=M+1
L2=N
L3=M
L4=N+1
IF(IC.NE.2) GOTO 30
L1=M
L2=N+1
L3=M+1
L4=N
30 L5=M+1
L6=N+1
NX8=10000*X(IC, ID, L5, L6)
NX9=10000*X(IC, ID, L1, L2)
NY8=10000*Y(IC, ID, L5, L6)
NY9=10000*Y(IC, ID, L1, L2)
IK=30
IF(IKK.EQ.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NYS
IF(NX7.GE.NX8 .AND. NY7.EQ.NY8) GCTC 31
IF(NX7.LT.NX8 .AND. NY7.EQ.NY8) GCTC 14
IF(NX7.EQ.NX9 .AND. NY7.EQ.NY9) GCTC 36
IF(NX7.GT.NX9) GOTO 32
D1=I1*(Y(IC, ID, L1, L2) - Y(IC, ID, M, N)) / (X(IC, IC, L1, L2) - X(IC, ID, M, N))
IF(L3.EC.L4) GOTO 33
33 D2=I1*(Y(IC, ID, L3, L4) - Y(IC, ID, M, N)) / (X(IC, IC, L3, L4) - X(IC, ID, M, N))
IF(L3.EC.L4) D2=0.
D12=I1*(Y(IC, ID, M, N) - YS) / (X(IC, ID, M, N) - XS)
IF(D12.LE.D1 .AND. C12.GE.D2) GCTC 14
NY8=10000*Y(IC, ID, M, N)
IF(NY7.GT.NY8) N=N-1
IF(NY7.LT.NY8) M=M-1

```

```

32  GCTO 1
    D3=(Y(IC,IC,L5,L6)-Y(IC,IC,L1,L2))/(X(IC,IC,L5,L6)-X(IC,IC,L1,L2))
    D3=D3*I1
    IK=32
    IF(IKK.EQ.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NY9
    IF(L3.EC.L4) GOTO 34
    D4=(Y(IC,IC,L5,L6)-Y(IC,IC,L3,L4))/(X(IC,IC,L5,L6)-X(IC,IC,L3,L4))
    D4=D4*I1
34  IF(L3.EC.L4) D4=0.
    D34=I1*(YS-Y(IC,IC,L5,L6))/(XS-X(IC,IC,L5,L6))
    IF(D34.GT.C3 .AND. D34.LT.D4) GCTC 14
    NY8=10000*Y(IC,IC,L5,L6)
    IF(NY7.GT.NY8) M=M+1
    IF(NY7.LT.NY8) N=N+1
31  GCTO 1
    M=L5
    N=L6
    IK=31
    IF(IKK.EC.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NY9
36  GCTO 14
    M=L3
    N=L4
    IK=35
    IF(IKK.EQ.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NY9
14  IK=14
    IF(IKK.EQ.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NY9
C
C
C  CONTROL - STEP
    K1=M+1
    K2=N
    K3=M
    K4=N+1
    IF(IC.NE.2) GOTO 35
    K1=M
    K2=N+1
    K3=M+1
    K4=N
35  J1=M+1
    J2=N+1
    D1=99.
    D2=99.
    D3=99.
    D4=99.
    D12=99.
    D34=99.
    NX8=10000*X(IC,IC,M,N)
    NX9=10000*X(IC,IC,J1,J2)
    NY8=10000*Y(IC,IC,M,N)
    NY9=10000*Y(IC,IC,J1,J2)
    IF(NX7.EQ.NX8 .AND. NY7.NE.NY8) GCTC 18
    IF(NX7.EQ.NX8 .AND. NY7.EQ.NY8) GCTO 19
18  GCTO 20
    WRITE(6,1000) M,N,NX7,NY7,NX8,NY8,D1,D2,C3,C4,C12,C34
    GCTO 19
20  D1=I1*(Y(IC,IC,K1,K2)-Y(IC,IC,M,N))/(X(IC,IC,K1,K2)-X(IC,IC,M,N))
    IF(K4.EQ.K3) GOTO 11
    D2=I1*(Y(IC,IC,K3,K4)-Y(IC,IC,M,N))/(X(IC,IC,K3,K4)-X(IC,IC,M,N))
11  IF(K4.EQ.K3) D2=0.
    D12=I1*(YS-Y(IC,IC,M,N))/(XS-X(IC,IC,M,N))
    IF(K4.EQ.K3) GOTO 25
    D3=I1*(Y(IC,IC,J1,J2)-Y(IC,IC,K3,K4))
    D3=D3/(X(IC,IC,J1,J2)-X(IC,IC,K3,K4))
25  IF(K4.EC.K3) D3=0.
    D4=I1*(Y(IC,IC,J1,J2)-Y(IC,IC,K1,K2))
    D4=D4/(X(IC,IC,J1,J2)-X(IC,IC,K1,K2))
    D34=I1*(Y(IC,IC,J1,J2)-YS)/(X(IC,IC,J1,J2)-XS)
    IF(D1.LT.D12 .OR. D2.GT.D12) GOTO 18
    IF(D3.LT.D34 .OR. D4.GT.C34) GOTO 18
    IF(NX7.LT.NX8 .OR. NY7.GE.NX9) GCTC 18
1000 FCRMAT(1X,'FIND : ',I2,' ',I2,4(1X,I8),1C(1X,F8.3))
1003 FCRMAT(1X,3I3,4(2X,I8))
1005 FCRMAT(1X,'FIND-EXIT')
19  IF(IKK.EQ.1) WRITE(6,1005)

```

```

RETURN
END
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
SLBROUTINE BOUND(LO4,IB,IR,IM,I,X)
CCOMPLEX*8 DYDXU
CCOMPLEX*8 AI,EI
REAL*4 X
CCOMMON/BC/ T1,T2,T,B,DYCX,D2YDX2,CYCXU,AI,EI
CCOMMON/SP/ XS(2,50),YS(2,50),A(2,50,4),R(2,50)
C
C
STEADY BCUNCERY CONCITICNS ALONG BCDY
T2=T1
IF(I.EQ.2) T3=T2
IF((I*IM).EQ.1 .AND. IR.GT.1) GCTC 2
IF(LO4.EQ.1) GOTO 3
DC 5 K=2,50
J=K-1
IF(XS(I,K).GE.X) GOTO 6
5
CCONTINUE
6
H=X-XS(I,J)
DYDX=A(I,J,2)+2.*A(I,J,3)*H+3.*A(I,J,4)*H*H
D2YDX2=2.*A(I,J,3)+6.*A(I,J,4)*H
3
GCTO 1
I1=(-1)**(I+1)
DYDX=I1*4.*T3*(1.-2.*X)
D2YDX2=-I1*8.*T3
2
GCTO 1
DYDX=0.
D2YDX2=0.
1
RETURN
END
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
SUBROUTINE SHOCK(KV,LO4,EM,IR,IB,M,N,XO,M2,N2,C3)
CCOMPLEX*8 U,V,PSI,G,DYCXU
CCOMPLEX*8 AI,EI
REAL*4 X,Y,AL,D1,D3,D,X1,X2,X3,X4,X5,Y1,Y2,Y3,Y4,Y5,Y6,X6
REAL*4 TX,TY,XX,YY,XP,YF,ALD,D2,P,C3,XO,XT,ALZ,DX1
CCOMMON/BA/ ALO,AM,ALD,C1,AK,I,IM,TX,TY,IW,NA,IC,ET,XX,YY
CCOMMON/BC/ T1,T2,T,B,DYCX,D2YDX2,CYCXU,AI,EI,I1,IBACK,IE
CCOMMON V(2,3,50,20),X(2,3,50,20),F(2,50),U(2,3,50,20),
FPSI(2,3,50,20),G(2,3,50,20),AL(2,3,50,20),Y(2,3,50,20),
FQ(2,50),CX(2),IN(2,50),IN2(2,50)
C
C
CCOMPUTATION OF THE FIELD NEAR BEHIND THE SHOCK
C
IKK=0
IF(I.EQ.2) IKK=1
IF(IKK.EQ.1) WRITE(6,1003)
MAX=1
IC=1
IC=1
IF(IR.EQ.1) GOTO 19
IC=IR-1
IF(IB.EQ.1) IC=2
IF(IB.EQ.2) IC=1
19
DX1=DX(I)
IF(T.EQ.0. .OR. IQ.EQ.0) GOTO 17
GCTO 21
17
D=EM*2.*SQRT(ALO)
IF(IB.EQ.2 .AND. T.EQ.0. .AND. I.EQ.1) C=1.05
DX1=D/(MA-1)
21
IK=21
IF(IKK.EQ.1) WRITE(6,1000) M,N,IE,IK,D,DX1,XO,TX,TY
J=M-1
K=N
L=M
IF(I.NE.2) GOTO 5
J=M
K=N-1
L=N
5
CALL CCNST1(LO4,IR,IB,ALZ,C3,XO)

```

```

CALL BOUND(LO4,IB,IR,IM,I,XO)
AL(IB,IR,M,N)=(ALZ-C3*CYCX*I1)**(2./3.)
IF(L.NE.2) GOTO 3
WRITE(6,1001) M,N,AL(IC,ID,M2,N2),M2,N2
GOTO 6
3
IK=3
IF(IKK.EQ.1) WRITE(6,1000) M2,N2,IR,IK,CYCX,XO,AL(IE,IR,M,N)
D1=I1/SQRT(AL(IB,IR,M,N))
XP=XO+XX
YP=YY
X5=X(IB,IR,J,K)
Y5=Y(IB,IR,J,K)
9
IK=9
IF(IKK.EQ.1) WRITE(6,1000) J,K,I1,IK,D1,XP,YP,X5,Y5
IF(IE.EC.1 .OR. M2.GE.KV) GOTO 16
L1=M2+1
L2=N2
L3=M2+1
L4=N2+1
IF(I.NE.2) GOTO 7
L1=M2
L2=N2+1
7
CONTINUE
IF(IR.EQ.1) GOTO 22
L1=M2
L2=N2+1
IF(IC.NE.2) GOTO 22
L1=M2+1
L2=N2
22
X1=X(IC,ID,M2,N2)
Y1=Y(IC,ID,M2,N2)
X2=X(IC,ID,L1,L2)
Y2=Y(IC,ID,L1,L2)
X3=X(IC,ID,L3,L4)
Y3=Y(IC,ID,L3,L4)
IZ=0
XT=XO
IK=22
IF(IKK.EC.1) WRITE(6,1000) M2,N2,IZ,IK,X5,Y5,X1,Y1,X2,Y2
14
D2=SQRT(AL(IC,ID,M2,N2))+SQRT(AL(IB,IR,J,K))
D2=2.*I1/D2
IF(T.NE.0. .AND. IQ.NE.0) GOTO 2
X4=(L-2)*DX1/2.+TX
Y4=D2*(X4-TX)+TY
GOTO 1
2
X4=(D1*XP-D2*X5+Y5-YP)/(C1-D2)
Y4=(X4-XP)*D1+YP
IK=1
IF(IKK.EQ.1) WRITE(6,1000) IC,ID,IE,IK,X4,Y4,XP,XX,C2
IF(I.EQ.1 .AND. IR.GT.1) GOTO 20
IF(Y4.LE.0.COOL .AND. Y4.GE.0.) GCTC 8
IF(Y4.GT.0. .AND. IZ.EQ.C) GOTO 13
IF(T.NE.0.) GOTO 4
Y4=0.
X4=-EM/D2+TX
GOTO 8
20
CONTINUE
IF(Y4.LE.EM .AND. Y4.GE.(0.999*EM)) GOTO 8
IF(Y4.LT.EM .AND. IZ.EC.0) GOTO 13
IF(T.NE.0. .AND. IQ.NE.C) GOTO 4
Y4=EM
X4=TX+EM/D2
GOTO 8
4
X5=X(IB,IR,J,K)
Y5=Y(IB,IR,J,K)
X6=(EM*(2-IB)-Y5)/D2+X5
XT=(XP-X4)*EM/(I1*Y4+EM*(IB-1))+X6-XX
IK=4
IF(IKK.EQ.1) WRITE(6,1000) IC,IC,IE,IK,X6,XT,D1,X5,Y5
CALL CONST1(LO4,IR,IB,ALZ,C3,XT)
CALL BOUND(LO4,IB,IR,IM,I,XT)
AL(IB,IR,M,N)=(ALZ-C3*CYCX*I1)**(2./3.)
D1=I1/SQRT(AL(IB,IR,M,N))
IZ=IZ+1

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```

XP=XT+XX
IF(IZ.LT.30) GOTO 14
IF(IKK.EC.1) WRITE(6,1000) M,N,IZ,IE,X4,Y4,X5,Y5
X4=X6
Y4=FM*(2-IB)
AL(IB,IR,M,N)=AL(IB,IR,J,K)
GOTO 8
13 CALL FIND(KV,IM,IB,IR,X4,Y4,M1,N1)
IK=13
IF(IKK.EC.1) WRITE(6,1000) M1,N1,IE,IK,X4,Y4
F,X(IC,IO,M1,N1),Y(IC,IO,M1,N1)
IF(IE.EC.1) WRITE(6,1002) M,N,M1,N1,X4,Y4,X(IC,IO,M1,N1)
F,Y(IC,IO,M1,N1)
IF(M1.EQ.M2 .AND. N1.EC.N2 .AND. IE.EQ.O) GCTC 8
IF(T.NE.O .AND. IQ.NE.C) GOTO 15
IF(IE.EC.1) GOTO 16
M2=M1
N2=N1
GOTO 8
15 X6=(Y2-Y5)/D2+X5
IF(X6.GE.X2) GOTO 10
D3=(Y2-Y1)/(X2-X1)
X5=(Y1-Y5+D2*X5-D3*X1)/(C2-D3)
Y5=D3*(X5-X1)+Y1
IK=15
IF(IKK.EC.1) WRITE(6,1000) M1,N1,IE,IK,X5,Y5,X6,D3
IF(IC.EC.2) GOTO 23
N2=N2-1
IF(IR.GT.1) GOTO 26
IF(N2.EQ.O) GOTO 16
IF(AL(IC,IO,M2,N2).NE.AL(IC,IO,M2-1,N2) .AND. IR.EC.1) GOTO 16
IF(I.NE.2) GOTO 9
26 M2=M2-1
N2=N2+1
GOTO 9
23 N2=N2-1
GOTO 9
10 D3=(Y3-Y2)/(X3-X2)
X5=(Y2-Y5+D2*X5-D3*X2)/(C2-D3)
Y5=D3*(X5-X2)+Y2
IK=10
IF(IKK.EC.1) WRITE(6,1000) M1,N1,IE,IK,X5,Y5,X6,C3
IF(IC.EC.2) GOTO 24
M2=M2+1
IF(IR.GT.1) GOTO 25
IF(M2.EQ.KV) GOTO 16
IF(I.NE.2) GOTO 9
25 N2=N2+1
M2=M2-1
GOTO 9
24 M2=M2+1
GOTO 9
8 X(IB,IR,M,N)=X4
Y(IB,IR,M,N)=Y4
IF(Y4.LE.O.0001 .AND. Y4.GE.O .AND. IB.EC.2) Y(2,IR,M,N)=O.
IF(Y4.LE.EM .AND. Y4.GE.(O.999*EM) .AND. IB.EC.1) Y(1,IR,M,N)=EM
IK=8
IF(IKK.EC.1) WRITE(6,1000) M,N,IE,IK,X4,Y4,X1,Y1
GOTO 11
6 X(IB,IR,M,N)=TX
Y(IB,IR,M,N)=TY
GOTO 12
16 IE=1
D2=SQRT(AL(IB,IR,J,K))+SQRT(AL(IB,IR,M,N))
D2=4.*I1/(2.*SQRT(ALO)+I2)
IF(T.NE.O .AND. IQ.NE.O) GOTO 18
X(IB,IR,M,N)=(L-2)*CX1/2.+TX
Y(IB,IR,M,N)=D2*(X(IB,IR,M,N)-TX)+TY
GOTO 11
18 X(IB,IR,M,N)=(D1*XP-D2*X5+Y5-YP)/(C1-D2)
Y(IB,IR,M,N)=(X(IB,IR,M,N)-XP)*D1+YP
C
C
C COMPUTATION OF THE ADDITIONAL PCINTS FOR THE UNSTEADY
FLOW FIELD

```

```

C
11 D3=1./SQRT(AL(IB,IR,M,N))
    D4=1./SQRT(AL(IB,IR,J,K))
    P(I,L)=(D4*X(IB,IR,J,K)+C3*XP+(Y(IB,IR,J,K)-YF)*I1)/(D3+D4)
    Q(I,L)=C3*(P(I,L)-XP)*I1+YP
    IF(T.EQ.0.) P(I,L)=X(IB,IR,J,K)
    IF(L.NE.3) GOTO 12
    P(I,L)=XP
    Q(I,L)=C.
12 CCNTINUE
1000 FCRMAT(IX,4I4,2X,6(2X,F8.3))
     IN2(1,L)=M2
     IN2(2,L)=N2
1001 FCRMAT(IX,I2,'',I2,' LAMBDA= ',FS,7,' MEST: ',I2,'',I2)
1002 FCRMAT(IX,'SHOCK RUNS OUT: ',2(I2,'',I2,2X),4(F8.3,2X))
1003 FCRMAT(IX,'SHOCK - ENTRY')
    RETURN
    END
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
  SLBROUTINE RAND(LO4,IB,IR,M,N,C1)
  CCMPLX*8 L,V,PSI,G,CYCXU
  CCMPLX*8 AI,EI
  REAL*4 TX,TY,XX,YY,X,Y,AL,F,D,X2,X1,P,ALC,C1,A,Y1,FS,ALZ
  CCOMMON/BA/ALO,AM,ALD,C3,AK,I,IM,TX,TY,IW,MA,IC,ET,XX,YY
  CCOMMON/BC/ T1,T2,T,B,DYCX,D2YDX2,CYDXU,AI,EI,I1
  CCOMMON V(2,3,50,20),X(2,3,50,20),P(2,50),U(2,3,50,20),
  FPSI(2,3,50,20),G(2,3,50,20),AL(2,3,50,20),Y(2,3,50,20),
  FQ(2,50),DX(2),IN(2,50),IN2(2,50)
C
C
C BOUNDARY STEP OF THE STEADY FIELD TO THE BCDY
C
C IK=0
  IF(IR.GT.1) IK=1
  IF(IK.EC.1) WRITE(6,1003)
  A=-C1*I1
  J1=M
  J2=N-1
  IF(I.NE.2) GOTO 5
  J1=M-1
  J2=N
5 X1=X(IB,IR,J1,J2)-XX
  Y1=Y(IB,IR,J1,J2)-YY
  X2=X1
  DC 11 KX=1,30
  CALL CONST1(LO4,IB,IR,ALZ,C1,X2)
  CALL BOUND(LO4,IB,IR,IM,I,X2)
  IF(IK.EQ.1) WRITE(6,1002) IB,IR,M,N,DYDX,ALZ,X1,Y1
  AL(IB,IR,M,N)=(ALZ+A*CYCX)**(2./3.)
  IF(T.EQ.C. .OR. IQ.EQ.0) GOTO 2
  IF(IK.EC.1) WRITE(6,1002) IB,IR,M,N,DYDX,AL(IB,IR,M,N),Y1,D2YCX2
1002 FCRMAT(IX,4I4,4(F10.3))
     F=2.*(X2-X1)*I1/Y1
     F=F-(SQRT(AL(IB,IR,J1,J2))+SQRT(AL(IB,IR,M,N)))
     FS=2.*I1/Y1-A*D2YDX2/(3.*AL(IB,IR,M,N))
     D=F/FS
     X2=X2-D
  IF(ABS(D).LE.0.000001) GOTO 10
11 CCNTINUE
  WRITE(6,1000) M,N,D
10 CALL CONST1(LO4,IR,IB,ALZ,C1,X2)
  CALL BOUND(LO4,IB,IR,IM,I,X2)
  IF(IK.EC.1) WRITE(6,1002) IB,IR,M,N,CYDX,ALZ,X2,D
  IF(IK.EC.1) WRITE(6,1004)
  AL(IB,IR,M,N)=(ALZ+A*DYCX)**(2./3.)
  Y(IB,IR,M,N)=YY
  X(IB,IR,M,N)=X2+XX
  GOTO 1
2 Y(IB,IR,M,N)=YY
  D=-2.*I1/(SQRT(AL(IB,IR,J1,J2))+SQRT(AL(IB,IR,M,N)))
  X(IB,IR,M,N)=X(IB,IR,J1,J2)-(Y(IB,IR,J1,J2)-YY)/D
1 IF(IK.EQ.1) WRITE(6,1002) IB,IR,M,N,X(IB,IR,M,N),AL(IB,IR,M,N)
1000 FORMAT(IX,'RAND : ',I2,'',I2,' C= ',E10.3)
1003 FCRMAT(IX,'RAND-ENTRY :')

```

```

1004  FORMAT(1X,'RAND-ITERATION-EXIT.')
      RETURN
      END
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
      SLBROUTINE GEN(IB,IR,M,N,I)
      CCMPLEX*8 U,V,PSI,G,CYCXU
      CCMPLEX*8 AI,EI
      REAL*4 X,Y,AL,D1,C2,X1,P
      COMMON/BC/ T1,T2,T,B,DYCX,D2YDX2,CYDXU,AI,EI,I1
      COMMON V(2,3,50,20),X(2,3,50,20),F(2,50),U(2,3,50,20),
      FPSI(2,3,50,20),G(2,3,50,20),AL(2,3,50,20),Y(2,3,50,20),
      FQ(2,50),DX(2),IN(2,50),IN2(2,50)
C
C
      GENERAL STEP OF STEADY FIELD
C
      J1=M
      J2=N-1
      K1=M-1
      K2=N
      IF(I.NE.2) GOTO 5
      J1=M-1
      J2=N
      K1=M
      K2=N-1
5     D1=1./SQRT(AL(IB,IR,K1,K2))
      D2=2./((SQRT(AL(IB,IR,J1,J2))+SQRT(AL(IB,IR,K1,K2)))
      X1=D2*X(IB,IR,J1,J2)+C1*X(IB,IR,K1,K2)
      X1=X1+(Y(IB,IR,J1,J2)-Y(IB,IR,K1,K2))*I1
      X(IB,IR,M,N)=X1/(D1+D2)
      Y(IB,IR,M,N)=D1*I1*(X(IB,IR,M,N)-X(IB,IR,K1,K2))+Y(IB,IR,K1,K2)
      AL(IB,IR,M,N)=AL(IB,IR,K1,K2)
      RETURN
      END
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
      SUBROUTINE BOUNCU(IM,IR,IB,X,I,AK,L)
      CCMPLEX*8 DYDXU
      CCMPLEX*8 AI,EI
      REAL*4 X
      COMMON/BC/ T1,T2,T,B,DYCX,D2YDX2,CYCXU,AI,EI,I1
C
C
      UNSTEADY BOUNDARY - CONDITIONS ALONG BODY
C
      DYDXU=0.
      IF((I*IM).EQ.1 .AND. IR.GT.1) GOTO 13
      IF(L.EQ.1) GOTO 11
      DYDXU=-1.-(X-B)*AK*AI
      GOTO 13
11     DYDXU= AI*AK
13     RETURN
      END
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
      SUBROUTINE RANDS(KV,IB,IR,M,N,L,APY,N2,N2)
      CCMPLEX*8 U,V,PSI,G,DYCXU,D2,D6,PII,CL,DL,DV
      CCMPLEX*8 AI,EI,EI2
      CCMPLEX*16 A,RI,ES
      REAL*4 X,Y,AL,ANY,S,AMN,AM3,AM4,AN3,AN4,C0,C1,D2,D4,D5,D7,XE
      REAL*4 P,ALD,CY,CZ,Z1,Z2,W,AA1,AA2,AA3,AA4,AA5,AM1,AM2,AN1,AN2
      REAL*4 TX,TY,XX,YY,DAL,ANY2,B1
      COMMON/BA/ ALO,AM,ALC,C,AK,I,IM,TX,TY,IM,NA,IC,ET,XX,YY
      COMMON/BC/ T1,T2,T,B,DYCX,D2YDX2,CYDXU,AI,EI,I1,IBACK,IE
      COMMON/SCL/ A(4,4),RI(4),ES(4)
      COMMON V(2,3,50,20),X(2,3,50,20),P(2,50),U(2,3,50,20),
      FPSI(2,3,50,20),G(2,3,50,20),AL(2,3,50,20),Y(2,3,50,20),
      FQ(2,50),CX(2),IN(2,50),IN2(2,50)
C
C
      COMPUTATION OF UNSTEADY BOUNDARY-PROPERTIES ALONG SHOCK
C
      IKK=0
      IF(IM.EQ.4 .AND. IR.GT.1) IKK=1
      IF(IKK.EC.1) WRITE(6,1001)
1001  FORMAT(1X,'RANDS-ENTRY')

```

```

IK=0
1002 IF(IKK.EQ.1) WRITE(6,1002) M,N,IK,X(IB,IR,M,N),Y(IB,IR,M,N)
FCRMAT(3I4,5(2X,E11.4))
IC=1
IC=1
IF(IR.EQ.1) GOTO 1
IC=IR-1
IF(IB.EQ.1) IC=2
1 IF(IB.EQ.2) IC=1
EI2=1./EI
IF(IC.EQ.2) EI2=EI
J=M
IF(I.NE.2) GOTO 5
5 J=N
IE=0
CALL FIND(KV,IM,IB,IR,X(IB,IR,M,N),Y(IB,IR,M,N),M2,N2)
ALT=AL(IC, ID, M2, N2)
DV=V(IC, IC, M2, N2)
CU=U(IC, ID, M2, N2)
IF(IE.EQ.1) ALT=ALO
IF(IE.EQ.1) DV=0.
IF(IE.EQ.1) DU=0.
IK=5
IF(IKK.EQ.1) WRITE(6,1002) IC, ID, IK, X(IC, ID, M2, N2), Y(IC, ID, M2, N2)
CX=C-2.
ANY=(ALT-AM+1.)/(C*AM)+1.
ANY2=ANY*ANY
L1=M+1
L2=N
IF(I.NE.2) GOTO 22
L1=M
L2=N+1
22 IF(L1.GT.KV .OR. L2.GT.KV) IE=1
CC IF(IE.EC.1) GOTO 23
C S=Y(IB,IR,L1,L2)-Y(IB,IR,M,N)
S=S/(X(IB,IR,L1,L2)-X(IB,IR,M,N))
S=2.*I1/(SQRT(ALT)+SQRT(AL(IB,IR,M,N)))
S=ATAN(S)
CY=SIN(S)**2
CZ=COS(S)**2
AMN=1./(AM*CY*ANY2)
AM1=2.*CY*ANY*SIN(2.*S)/C
AM2=2.*AK*CY*(1.+AMN)/C
AM3=CX*CY*(1.-2.*AMN/CX)/C+CZ
AM4=SIN(2.*S)*(1.+AMN)/C
AN1=-2.*CY*ANY*(COS(2.*S)+AMN)/C
AN2=-2.*AK*SIN(S)*COS(S)*(1.+AMN)/C
AN3=AM4
AN4=CY+CZ*CX*(1.-2.*AMN/CX)/C
IF(J.GT.2) GOTO 6
XE=X(IB,IR,M,N)-XX
CALL BCUNDU(IM,IR,IB,XE,I,AK,L)
V(IB,IR,M,N)=DYDXU
CL=AM1*V(IB,IR,M,N)/AN1+CU*(AM3-AM1*AN3/AN1)*EI2
U(IB,IR,M,N)=CL+DV*(AM4-AM1*AN4/AN1)*EI2
K1=M+1
K2=N
IF(I.NE.2) GOTO 4
4 K1=M
K2=N+1
D3=(SQRT(AL(IB,IR,M,N))+SQRT(AL(IE,IR,K1,K2)))/2.
D3=-I1/C3
D1=I1/SQRT(AL(IB,IR,K1,K2))
IK=4
IF(IKK.EQ.1) WRITE(6,1002) M,N,IK,D3,D1,S
D4=(YY-Y(IB,IR,K1,K2))/C1+X(IB,IR,K1,K2)
X(IB,IR,M,N)=(D3*D4-TX*TAN(S))/(C3-TAN(S))
6 IF(T.EQ.0. .OR. IQ.EQ.C) X(IB,IR,M,N)=TX
Y(IB,IR,M,N)=(X(IB,IR,M,N)-TX)*TAN(S)+YY
J1=M-1
J2=N
LF=M
IF(I.NE.2) GOTO 8
J1=M

```

```

J2=N-1
LF=N
8 CONTINUE
IF(AL(IB,IR,J1,J2).LT.AL(IB,IR,M,N).OR.T.EQ.0.) GCTC 12
IF(IQ.EC.0) GOTO 12
J2=N-1
IF(I.NE.2) GOTO 12
J1=M-1
12 AMN=AM*CY
C AA1=-2.*(C-1.)*SIN(2.*S)*AMN*ANY2/C
C AA2=-4.*(C-1.)*AK*ANY*AMN/C
C AA5=(C-1.)*AM*(CX-2.*AMN*ANY2)/C
C AA3=4.*(C-1.)*ANY*AMN/C+AA5
C AA4=-4.*(C-1.)*ANY*AMN*CCS(S)/(C*SIN(S))
C AA5=AK*AA5
IF(J.EC.2) GOTO 7
C DO=X(IB,IR,M,N)-P(I,LP)
C4711 WRITE(6,4711) M,N,LP,X(IB,IR,M,N),X(IB,IR,J1,J2),P(I,LP)
FORMAT(1X,3I4,3(2X,E12.5))
DAL=AL(IB,IR,M,N)
IE=0
CALL FIND(KV,IM,IB,IR,X(IB,IR,J1,J2),Y(IE,IR,J1,J2),M3,N3)
CU=U(IC,ID,M3,N3)+U(IC,IC,M2,N2)
DV=V(IC,ID,M3,N3)+V(IC,IC,M2,N2)
IF(IE.EQ.1) DU=0.
IF(IE.EQ.1) DV=0.
MP=1000000.*(P(I,LP)-X(IB,IR,J1,J2))
IF(MP.LE.0.AND.T.NE.C.) WRITE(6,1003) M,N,P(I,LP),X(IB,IR,J1,J2)
1003 FORMAT(1X,'P-X.LE.0.:',2I4,2(2X,E12.5))
IF(MP.NE.C) GOTO 3
IF(T.NE.0.AND.IQ.NE.C) GOTO 3
D1=0.
GCTO 2
3 D1=(AL(IB,IR,M,N)-AL(IB,IR,J1,J2))*CG
D1=D1/(4.*DAL*(P(I,LP)-X(IB,IR,J1,J2)))
IK=3
IF(IKK.EC.1) WRITE(6,1002) M2,N2,IK,DO,D1,CU,DV,DAL
2 D2=AI*AK*AM*DO/DAL
D3=-AM*(AK*DO)**2/(4.*CAL)
D4=1./SQRT(AL(IB,IR,M,N))
D5=-D3*D4
D6=AK*AK*AM*PSI(IB,IR,J1,J2)*DO/DAL
D7=Y(IB,IR,M,N)-Y(IB,IR,J1,J2)
A(1,1)=1.+D1+D2+D3
A(1,2)=-(D4+D5)*I1
A(1,3)=0.
RI(1)=U(IB,IR,J1,J2)*(1.-D1-D2-D3)-V(IB,IR,J1,J2)*(C4-D5)*I1+C6
A(2,1)=1.
A(2,2)=C.
A(2,3)=-2.*AM1/D7-AI*AM2
RI(2)=-(2.*AM1/D7-AI*AM2)*G(IB,IR,J1,J2)-U(IB,IR,J1,J2)
RI(2)=RI(2)+(AM3*DU+AM4*CV)*EI2
A(3,1)=0.
A(3,2)=1.
A(3,3)=-2.*AM1/D7-AI*AM2
RI(3)=-(2.*AM1/D7-AI*AM2)*G(IB,IR,J1,J2)-V(IB,IR,J1,J2)
RI(3)=RI(3)+(AM3*DU+AM4*DV)*EI2
C
NC=3
CALL SOLVE(NO)
U(IB,IR,M,N)=ES(1)
V(IB,IR,M,N)=ES(2)
G(IB,IR,M,N)=ES(3)
C7 PII=V(IB,IR,M,N)-AI*AM2*G(IB,IR,M,N)
CYY PII=PII-(AN3*U(IC,ID,M2,N2)-AN4*V(IC,ID,M2,N2))*EI2
CYY PII=AA1*PII/AN1+AI*AA2*G(IB,IR,M,N)+AA3*U(IC,ID,M2,N2)*EI2
CYY PII=PII+(AA4*V(IC,ID,M2,N2)+AA5*AI*PSI(IC,IC,M2,N2))*EI2
CYY AT=AK
CYY IF(AK.EC.0.) AT=1.
CYY PSI(IB,IR,M,N)=AI*(U(IB,IR,M,N)+PII/((C-1.)*AM))/AT
7 IF(J.GT.2) GOTO 20
PSI(IB,IR,M,N)=PSI(IC,IC,M2,N2)*EI2
GCTO 21
20 PII=(V(IB,IR,M,N)+V(IB,IR,J1,J2))*C4*I1

```

```

21      PII=(PII+U(IB,IR,M,N)+L(IB,IR,J1,J2))*0.5*CO
      PSI(IB,IR,M,N)=PII+PSI(IB,IR,J1,J2)
      RETURN
      END
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
      SUBROUTINE GENU(IM,IB,IR,M,N,I,AK,AM)
      CCOMPLEX*8 U,V,PSI,G,DYDXU,D2,D6,B6,CL,B2,CM
      CCOMPLEX*8 AI,EI
      CCOMPLEX*16 A,RI,ES
      REAL*4 X,Y,AL,Z1,Z2,Z3,F,DO,D1,C3,D4,D5,B0,B1,B2,B4,B4,F1,F2
      COMMON/BC/ T1,T2,T,B,DYEX,D2YDX2,CYDXU,AI,EI,I1
      COMMON/SCL/ A(4,4),RI(4),ES(4)
      COMMON V(2,3,50,20),X(2,3,50,20),F(2,50),U(2,3,50,20),
      FPSI(2,3,50,20),G(2,3,50,20),AL(2,3,50,20),V(2,3,50,20),
      FQ(2,50),DX(2),IN(2,50),IN2(2,50)
C
C      GENERAL STEP OF UNSTEADY FIELD
C
      IKK=0
      IF(IR.EQ.2) IKK=1
      IF(IKK.EC.1) WRITE(6,1002)
      J1=M
      J2=N-1
      J3=M-1
      J4=N
      IF(I.NE.2) GOTO 5
      J1=M-1
      J2=N
      J3=M
      J4=N-1
5      DC=X(IB,IR,M,N)-X(IB,IR,J1,J2)
      D1=0.5*(AL(IB,IR,M,N)-AL(IB,IR,J1,J2))
      D2=AI*2.*AK*AM*DO
      D3=-0.5*AM*(AK*DO)**2
      D4=AL(IB,IR,M,N)+AL(IB,IR,J1,J2)
      D5=SQRT(AL(IB,IR,M,N))+SQRT(AL(IB,IR,J1,J2))
      IF(IKK.EC.1) WRITE(6,1001) M,N,AL(IB,IR,M,N),D4,CO
      D6=(D1+D2+C3)/D4
      A(1,1)=1.+D6
      A(1,2)=2.*(1.-D3/D4)*I1/D5
      RI(1)=U(IB,IR,J1,J2)*(1.-D6)+V(IB,IR,J1,J2)*I1*2.*(1.+D3/C4)/D5
      RI(1)=2.*AK*AK*AM*PSI(IB,IR,J1,J2)*DO/D4+RI(1)
      B0=X(IB,IR,M,N)-X(IB,IR,J3,J4)
      B1=(AL(IB,IR,M,N)-AL(IB,IR,J1,J2))*B0/(2.*DO)
      B2=2.*AI*AK*AM*B0
      B3=-0.5*AM*(AK*B0)**2
      B4=1./SQRT(AL(IB,IR,M,N))
      B5=-0.5*B3*B4/AL(IB,IR,M,N)
      B6=0.5*(B1+B2+B3)/AL(IB,IR,M,N)
      A(2,1)=1.+B6
      A(2,2)=- (B4+B5)*I1
      RI(2)=U(IB,IR,J3,J4)*(1.-B6)+V(IB,IR,J3,J4)*I1*(-B4+B5)
      RI(2)=PSI(IB,IR,J3,J4)*AM*B0*AK*AK/AL(IB,IR,M,N)+RI(2)
C
      NG=2
      CALL SOLVE(NO)
      U(IB,IR,M,N)=ES(1)
      V(IB,IR,M,N)=ES(2)
      CL=V(IB,IR,J1,J2)+V(IB,IR,M,N)
      CL=(U(IB,IR,J1,J2)+U(IB,IR,M,N)-I1*2.*CL/C5)*DC/2.
      CM=(V(IB,IR,J3,J4)+V(IB,IR,M,N))*B4
      CM=(U(IB,IR,J3,J4)+U(IB,IR,M,N)+I1*CM)*B0/2.
      CL=PSI(IB,IR,J1,J2)+PSI(IB,IR,J3,J4)+CL+CM
      PSI(IB,IR,M,N)=CL/2.
1001  FCRMAT(IX,2I4,3(2X,F8.3))
1002  FCRMAT(IX,'GENU-ENTRY')
      RETURN
      END
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
      SUBROUTINE RANDB(IB,IR,M,N,L)
      CCOMPLEX*8 U,V,PSI,G,DYDXU,D2,A,D,CL,D6
      CCOMPLEX*8 AI,EI

```

```

REAL*4 TX, TY, XX, YY, X, Y, AL, DO, D1, D4, D5, ALD, D3, P, XT
CCOMMON/BA/ ALO, AM, ALD, C1, AK, I, IM, TX, TY, IW, NA, IC, ET, XX, YY
CCOMMON/BC/ T1, T2, I, B, DYCX, D2YDX2, CYDXU, AI, EI, I1
CCOMMON V(2,3,50,20), X(2,3,50,20), F(2,50), U(2,3,50,20),
FPSI(2,3,50,20), G(2,3,50,20), AL(2,3,50,20), Y(2,3,50,20),
FQ(2,50), DX(2), IN(2,50), IN2(2,50)

```

C  
C  
C

COMPUTATION OF UNSTEADY BOUNDARY-PROPERTIES ON BODY

```

XT=X(IB,IR,M,N)-XX
CALL BOUNDU(IM,IR,IB,XT,I,AK,L)
V(IB,IR,M,N)=DYDXU
J1=M
J2=N-1
IF(I.NE.2) GOTO 5
J1=M-1
J2=N
5 D1=0.5*(AL(IB,IR,M,N)-AL(IB,IR,J1,J2))
DO=X(IB,IR,M,N)-X(IB,IR,J1,J2)
D6=V(IB,IR,M,N)+V(IB,IR,J1,J2)
D2=AI*AK*AM*DO*2.
D3=-0.5*(AK*DO)**2*AM
D4=AL(IB,IR,M,N)+AL(IB,IR,J1,J2)
D5=SQRT(AL(IB,IR,M,N))+SQRT(AL(IB,IR,J1,J2))
A=(D1+D2+D3)/D4
U(IB,IR,M,N)=(V(IB,IR,M,N)-V(IB,IR,J1,J2))*I1/D5
U(IB,IR,M,N)=U(IB,IR,J1,J2)*(1.-A)-2.*U(IB,IR,M,N)
D=2.*PSI(IB,IR,J1,J2)-CC*D6*I1/D5
U(IB,IR,M,N)=U(IB,IR,M,N)+AK*AK*AM*CO*D/C4
U(IB,IR,M,N)=U(IB,IR,M,N)/(1.+A)
CL=U(IB,IR,M,N)+U(IB,IR,J1,J2)-2.*C6*I1/D5
PSI(IB,IR,M,N)=PSI(IB,IR,J1,J2)+CC*CL/2.
RETURN
END

```

CXX  
CXX

SUBROUTINE CONST1(L04, IR, IB, ALZ, C3, XF)

```

COMPLEX*8 L, V, PSI, G, DYDXU
CCOMPLEX*8 AI, EI
REAL*4 X, Y, AL, D1, D2, D3, XF, ALD, P, C3, TX, TY, XX, YY, ALZ, X1, Y1, X2, X3
CCOMMON/BA/ ALO, AM, ALD, C1, AK, I, IM, TX, TY, IW, NA, IC, ET, XX, YY
CCOMMON/BC/ T1, T2, I, B, DYCX, D2YDX2, CYDXU, AI, EI, I1
CCOMMON V(2,3,50,20), X(2,3,50,20), F(2,50), U(2,3,50,20),
FPSI(2,3,50,20), G(2,3,50,20), AL(2,3,50,20), Y(2,3,50,20),
FQ(2,50), DX(2), IN(2,50), IN2(2,50)

```

C  
C  
C

DETERMINATION OF THE CONSTANT VALUE ALONG THE CHARACTERISTIC

```

IF(IR.EQ.1 .AND. IM.EC.1) GOTO 1
IF(IR.EC.1 .AND. I.EQ.1) GOTO 1
IK=0
IF(IR.GT.1) IK=1
IC=1
IF(IC.EQ.1) GOTO 2
ID=IR-1
IC=2
IF(IB.EQ.2) IC=1
2 CCNTINUE
DC 3 J=3,20
M=J
N=J-1
IF(IC.NE.2) GOTO 4
M=J-1
N=J
4 X1=X(IC, ID, M, N)
Y1=Y(IC, ID, M, N)
IF(IK.EQ.1) WRITE(7,1000) IC, ID, M, N, X1, Y1, AL(IC, ID, M, N)
1000 FCRMAT(1X, 'CONST1: ', 4I4, 3F8.3)
D1=-I1/SQRT(AL(IC, ID, M, N))
X2=(YY-Y1)/D1+X1
X3=XP+XX
IF(X2.GE.X3) GOTO 5
3 CCNTINUE

```

```

5      M=M-1
      N=N-1
      X1=X(IC, ID, P, N)-X(IC, 1, 2, 1)
      CALL BOUND(LO4, IC, ID, IM, IC, X1)
      ALZ=AL(IC, IC, M, N)**1.5+I1*C3*DYCX
1      GOTO 6
      ALZ=ALD
6      RETURN
      END
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
      SLBROUTINE PRESS(AMY, IR, EM, LO2, AX, C)
      COMPLEX*8 U, V, PSI, G, DYCXU
      COMPLEX*8 AI, EI, PU
      REAL*4 TX, TY, XX, YY, X, Y, AL, D1, D2, D3, XP, ALO, F, C3, ALZ
      COMMON/BA/ ALO, AM, ALD, C1, AK, I, IM, TX, TY, Ih, MA, IQ, ET, XX, YY
      COMMON/BC/ T1, T2, T, B, CYEX, D2YDX2, CYCXU, AI, EI, I1
      COMMON V(2, 3, 50, 20), X(2, 3, 50, 20), P(2, 50), U(2, 3, 50, 20),
      FPSI(2, 3, 50, 20), G(2, 3, 50, 20), AL(2, 3, 50, 20), Y(2, 3, 50, 20),
      FQ(2, 50), DX(2), IN(2, 50), IN2(2, 50), FX(2, 30), FS(2, 30), PU(2, 30)
C
C
C      COMPUTATION AND OUTPUT OF THE PRESSURE - CCEFFICIENTS
      ALONG THE CHOSEN BLADE
      PU=CPU , PS=CPS
C
      YN=-EM
      YM=AMY*180./(4.*ATAN(1.))
      IN2=IM-2
      DC 7 I=1, 2
      K1=0
      DC 1 K=1, IR
      K3=2
      IF(I.EQ.1 .AND. K.EQ.1) K3=1
      K3=K+K3
      DC 3 J1=1, 20
      LI=J1+1
      CALL SWITCH(LI, J1, M, N, I)
      IF(K3.GT.IR) GOTO 10
      LI=1
      LZ=2
      CALL SWITCH(LZ, LI, M1, N1, I)
      IF(X(1, K, M, N).GT.X(1, K3, M1, N1)) GOTO 6
      GOTO 3
10     D1=REAL(U(1, K, M, N))**2+AIMAG(U(1, K, M, N))**2
3      IF(D1.EQ.0.) GOTO 6
6      CCNTINUE
      IF(J1.EQ.1) GOTO 1
      J1=J1-1
      K1=K1+J1
      DO 4 J2=1, J1
      LI=J2+1
      CALL SWITCH(LI, J2, M, N, I)
      K3=J2+K1-J1
      PS(I, K3)=-2.*(AL(1, K, M, N)-ALO)/(C1*AM)
      PU(I, K3)=-2.*(U(1, K, M, N)+AI*AK*PSI(1, K, M, N))
4      PX(I, K3)=X(1, K, M, N)
17     CCNTINUE
      IN2(I, 1)=K1
C
C
C      OUTPUT
      WRITE(7, 1015)
      IF(LO2.EC.0) WRITE(7, 996)
      IF(LO2.EC.1) WRITE(7, 995)
      WRITE(7, 1016) YM, IM2
      WRITE(7, 1001) AK, AX, C, B, T1
      WRITE(7, 1010)
      WRITE(7, 1011)
      DC 8 I=1, 2
      IF(I.EQ.2) WRITE(7, 1001) AK, AX, C, B, T2
      IF(I.EQ.2) WRITE(7, 1012)
      IF(I.EQ.2) WRITE(7, 1011)
      K=IN2(I, 1)
      DC 9 J=1, K

```

```

LI=J+1
CALL SWITCH(LI,J,M,N,I)
9 WRITE(7,1013) M,N,PX(I,J),PS(I,J),PU(I,J)
8 WRITE(7,1014)
995 FCRMAT(1X,/,1X,'PLUNGE - MODE',/)
996 FCRMAT(1X,/,1X,'PITCH - MODE',/)
1001 FCRMAT(1X,/,1X,'W*C/U=',F7.5,', M=',F5.3,', K=',F4.2,
F, B/C=',F4.2,', T/C=',F6.4,///)
1010 FCRMAT(1X,'PRESSURE-COEFFICIENTS UPPER SURFACE:',/)
1011 FCRMAT(2X,'POINT',4X,'X',8X,'CPS',8X,'RCPU',8X,'ICPU',/)
1012 FCRMAT(1X,'PRESSURE-COEFFICIENTS LOWER SURFACE:',/)
1013 FCRMAT(1X,I2,',',I2,4X,F5.3,3(2X,E10.3))
1014 FORMAT(//,1X,'*****')
F,/)
1015 FCRMAT(1H1)
1016 FORMAT(1X,'PHASE= ',F7.2,' / ',I2,',. BLADE')
RETURN
END
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
SUBROUTINE SWITCH(J,K,M,N,I)
C INDEX - SWITCH
M=J
N=K
IF(I.NE.2) GOTO 1
M=K
N=J
1 RETURN
END
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
SLBROUTINE SOLVE(N)
CCOMPLEX*16 G,C,X
CCOMMON/SCL/ G(4,4),C(4),X(4)
C
C SOLVE GIVES THE SOLUTION FOR A COMPLEX SYSTEM OF
C LINEAR EQUATIONS G*X=C
C
IKK=0
IF(IKK.NE.1) GOTO 2
WRITE(6,1001)
DC 1 M=1,N
1 WRITE(6,1000) (G(M,L),L=1,N),C(M)
2 MN=N-1
DC 10 M=1,MN
K=M+1
DC 10 J=K,N
G(M,J)=C(M,J)/G(M,M)
IF(M.EQ.1) GOTO 14
MA=M-1
DC 13 L=1,MA
13 G(M,J)=G(M,J)-G(M,L)*G(L,J)/G(M,M)
14 CCNTINUE
DC 15 L=1,M
15 G(J,K)=G(J,K)-G(L,K)*G(J,L)
10 CCNTINUE
C(1)=C(1)/G(1,1)
DC 12 I=2,N
C(I)=C(I)/G(I,I)
MA=I-1
DC 12 M=1,MA
12 C(I)=C(I)-G(I,M)*C(M)/G(I,I)
X(N)=C(N)
DC 11 I=1,MN
NA=N-I
X(NA)=C(NA)
NB=NA+1
DC 11 J=NB,N
11 X(NA)=X(NA)-G(NA,J)*X(J)
1001 FCRMAT(1X,'SOLVE-ENTRY:')
1000 FORMAT(1X,1C(2X,E10.3))
RETURN
END
CXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX

```

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