

AD-A098 461

PRINCETON UNIV NJ JOSEPH HENRY LABS OF PHYSICS
FURTHER REMARKS ON METASTABILITY IN AMORPHOUS SEMICONDUCTORS.(U)
APR 81 D C LICCIARDELLO

F/G 20/12

N00014-77-C-0711

NL

UNCLASSIFIED

1 1 1
00000000



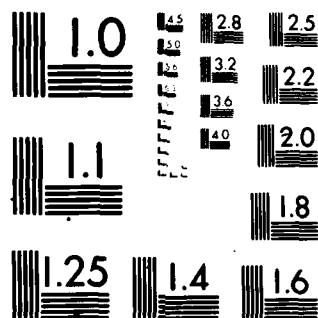
END

DATE

FILED

8-81

DTIC



MICROCOPY RESOLUTION TEST CHART
NATIONAL BUREAU OF STANDARDS 1963-A

AD A098461

LEVEL II

81 12

6 Further Remarks on Metastability
in Amorphous Semiconductors

10 D.C./Licciardello

Joseph Henry Laboratories of Physics

Princeton University
Princeton, New Jersey
08544

9 Technical rept. and

Bell Laboratories
Murray Hill
New Jersey 07974

11 Apr 81

12 12

Accession For	
NTIS GRA&I	☒
DTIC TAB	☐
Unannounced	☐
Justification	
By	
Distribution/	
Availability Codes	
Dist	Avail and/or Special
A	

ABSTRACT

The bandwidths for metastable states in a simple model of amorphous semiconductors is evaluated exactly. The region of metastability in the space of the model parameters (phase diagram) is deduced and the subregion corresponding to an infinite range model is indicated. The model predicts both paramagnetic and diamagnetic metastable states and the relationship between them is discussed.

*Research supported in part by
ONR/N00014-77-C-0711 NSF-DMR 802063

15

DISTRIBUTION STATEMENT A
Approved for public release;
Distribution Unlimited

81 5

01 021

DTIC
ELECTE
S MAY 5 1981 D

D
405629

UNC FILE COPY

In this paper we further examine the properties of metastable electronic states predicted for the negative U model (Anderson, 1975) of amorphous semiconductors by Licciardello, Stein and Haldane (1980) (see also, Licciardello 1980a). The presence of such states have been linked to photoinduced paramagnetism (Mott, Davis and Street, 1975; Licciardello and Stein, 1980) and conductivity changes (Licciardello, 1980b) observed in a wide variety of glassy semiconductors. Interest in the properties of metastable states stems from their defect-like behavior, although the model itself contains no defects per se. Thus the theory provides an alternative point of view to the generally accepted defect picture of glassy materials (Kastner, Adler, and Fritzsche, 1976).

We consider a spatially random distribution of localized states in which the Hamiltonian H_i corresponding to the i^{th} site is written

$$H_i = \epsilon_i \hat{n}_i + \omega b_i^\dagger b_i + U \hat{n}_{i\uparrow} \hat{n}_{i\downarrow} - g(b_i + b_i^\dagger) \hat{n}_i \quad (1)$$

where the quantity $\hat{n}_i$ is the number operator for an electron at the i^{th} site. Similarly $\hat{n}_{i\uparrow}$ ($\hat{n}_{i\downarrow}$) counts electrons at the i^{th} site with spin up ($\downarrow$) . I have also introduced the operator $b_i^\dagger$ which creates a phonon at the i^{th} site with frequency ω . The parameters ϵ_i, U and g are, respectively, the electron energy at the i^{th} site, the Hubbard repulsion energy and the electron phonon coupling. We have ignored the zero point lattice vibrations and employ, for simplicity, an Einstein model for the phonons. We also neglect the variation of the coupling g and the Hubbard repulsion U from site to site.

The Hamiltonian has the usual form including an electronic part H_e , a pure phonon part H_{ph} and an interaction term $H_{e.ph}$:

$$H_i = H_e + H_{ph} + H_{e \cdot ph} \quad (2)$$

It is useful, following Licciardello, Stein and Haldane (1980), to renormalize the pure terms by introducing displaced phonon operator d_i through the relation

$$d_i = b_i - g \hat{n}_i / \omega \quad (3)$$

This immediately gives the form

$$H_i = \tilde{H}_{el} + \tilde{H}_{ph} \quad (4)$$

where

$$\tilde{H}_{el} \equiv \epsilon_i \hat{n}_i - \frac{1}{2} C \hat{n}_i \hat{n}_i + U \hat{n}_{i\uparrow} \hat{n}_{i\downarrow} \quad (5)$$

and

$$\tilde{H}_{ph} \equiv \omega d_i^\dagger d_i \quad (6)$$

and where we have defined the coupling energy $C \equiv 2 g^2 / \omega$. The stationary solutions are described by quantum numbers n_i, v_i which give the number of electrons and phonons, respectively, localized at the site i .

We define a local equilibrium state as one which satisfies

$$\tilde{H}_{ph} |n_i, v_i\rangle_{eq} = 0 \quad (7)$$

or

$$d_i \equiv 0$$

It is the purpose of this note to examine conditions on the Hamiltonian parameters under which the state $|n_i, v_i\rangle_{eq}$ is metastable.

We consider two sites i and j whose corresponding states $|n_i, v_i\rangle_{eq}$ and $|n_j, v_j\rangle_{eq}$ each satisfy equation (7). At low temperatures, we may roughly define stability of the equilibrium states i and j with respect to spontaneous charge transfer between them by

requiring the energy for the process to be endothermic, i.e.

$$\begin{aligned} & \text{eq} \langle n_i v_i | H_i | n_i v_i \rangle_{\text{eq}} + \text{eq} \langle n_j v_j | H_j | n_j v_j \rangle_{\text{eq}} \\ & < \langle n_i + \Delta n_i v_i | H_i | n_i + \Delta n_i v_i \rangle \\ & + \langle n_j - \Delta n_j v_j | H_j | n_j - \Delta n_j v_j \rangle \end{aligned} \quad (8)$$

where $\Delta n = \pm 1, \pm 2 \Rightarrow 0 \leq n_{i,j} \pm \Delta n \leq 2$. It is instructive to consider the range of electronic energies ϵ_i under which eq. (8) is satisfied for diamagnetic $n_i = 2$ and paramagnetic states $n_i = 1$. For the system in equilibrium defined by eq. (7) for each site, the energy eigenvalue is directly given through equation (5) for the site i :

$$E_i = \epsilon n_i + U \delta_{2,n_i} - \frac{1}{2} C n_i^2 \quad (9)$$

where n_i is the number of electrons at the site i and δ is the Kronicker delta. As has been noted (Anderson, 1975), the system ground state is determined by the occupation set $\{n_i\} = \{0, 2\}$ and for convenience we define the Fermi energy ϵ_F by the relation

$$2\epsilon_F + U - 2C = 0 \quad (10)$$

Thus, in the ground state, sites for which $\epsilon_i > \epsilon_F$ have occupation $n_i = 0$ and sites for which $\epsilon_i < \epsilon_F$ are occupied by $n_i = 2$.

Infinite Range Model:

We define a metastable state i as an equilibrium state i (cf eq. 7) for which $\epsilon_i > \epsilon_f$ with $n_i \neq 0$ or $\epsilon_i < \epsilon_f$ with $n_i \neq 2$ and which is

stable according to eq. (8) for all j . The latter condition defines an infinite range model which we consider here. This assumption has been relaxed in other works (Licciardello, Haldane, Stein, 1980; Licciardello 1980b).

Metastable Diamagnetism:

We consider a metastable state i which contains $n_i = 2$. The bandwidth for stability is given through eq. (8) and we evaluate the case $n_j = 0$ and $\Delta n_i = -1$. It has been shown that the two-electron emission process $\Delta n_i = \pm 2$ is always more severe (i.e. states are more stable energetically with respect to this process) so that it may be ignored. We obtain, using eqs. (8), (4), (5) and (6)

$$\epsilon_i - \epsilon_j < 2C - U \quad (11)$$

as the condition for metastable diamagnetism. The region is sketched in Figure 1 where the stable regions $\epsilon_i < \epsilon_F$ and $\epsilon_j > \epsilon_F$, and the Fermi point ϵ_F , defined through eq. (10), are also indicated. The metastable states occur generally in the shaded triangle. In the infinite range model which assumes non-zero electronic matrix elements between all i, j , it is easy to show that the allowed solution corresponds to the largest rectangle, with sides parallel to the axes, which may be inscribed in the metastable region. This is also indicated yielding the bandwidths

$$C - U/2 < \epsilon_i < 2C - U \quad (12)$$

for $n_i = 2$ and

$$0 < \epsilon_j < C - U/2 \quad (13)$$

for $n_j = 0$, in agreement with the results of Ref. 2. It is interesting

to note that any relaxation of the range of the interaction increases the range of metastability. In particular, if only decays to the ground state are permitted (as may be arranged experimentally for a system with few metastable states) the bandwidths are increased by a factor of 2 (indicated by the lengths of the sides of the isosceles triangle).

Metastable Paramagnetism:

We also consider metastable states occupied by a single electron which, as described above, have no stable regions. Thus we consider a site i where $n_i = 1$ and consider the absorption $\Delta n = 1$ and emission $\Delta n = -1$ of a single particle through a process involving the site j . Again using the criterion (8) to define the region of metastability, the range of energy ϵ_i is given through the relations

$$n_j = 0 \quad \epsilon_i - \epsilon_j \leq C \quad (14)$$

$$n_j = 1 \quad \epsilon_i - \epsilon_j \leq U \quad (15)$$

$$n_j = 2 \quad \epsilon_j - \epsilon_i \leq C \quad (16)$$

where we have employed eqs. (4), (5), and (6). The paramagnetic case is sketched in Figure 2, where the lines defining the regions satisfying the inequalities (14) and (16) are drawn. The vertical boundaries of the shaded region where paramagnetic metastable states may occur are defined through the limits of $n_j = 2$ states (eq.(12), where we use $i \rightarrow j$) within the infinite range approximation and correspondingly for $n_j = 0$ (eq. (13)). These are the vertical lines at $\epsilon_j = 2C - U$ and at $\epsilon_j = 0$. The same algorithm applies for determining the allowed paramagnetic

states which is the largest rectangle with sides parallel to the axes. This region is indicated as (a) for the infinite range model. The bandwidth satisfies $\delta\epsilon_1 < 2U$, as is seen by inspection, first predicted by Licciardello et al., and is consistent with the inequality (15) without further restriction.

It is interesting to examine the consequences of ^{ignoring} fast particle exchange between paramagnetic states, i.e. process (15). This is entirely realizable at densities in which paramagnetism has been induced in amorphous materials up to now. If the range of metastable diamagnetic states is suppressed, (indicated in rectangle b, in Figure 2) the energy range of allowed paramagnetic states increases as shown. This is an interesting prediction and may be tested experimentally. In general, metastable diamagnetic states are produced thermally (no gap) whereas metastable paramagnetic states must be optically induced. Thus materials annealed at low temperature for long times should, according to these considerations, be susceptible to a larger density of induced paramagnetism.

Acknowledgement

I have benefitted from conversations with P. Littlewood.

References

- Anderson, P.W., 1975, Phys. Rev. Lett. 34, 953
- Kastner, M., Adler, D., and Fritzsche, H., 1976, Phys. Rev. Lett. 37, 1504
- Licciardello, D.C., 1980a, Comments in Solid State Physics 9, 217
- Licciardello, D.C., 1980b, to be published
- Licciardello, D.C. and Stein, D.L., 1980, J. Phys. C. 13 L391-396 (1980)
- Licciardello, D.C., Stein, D.L., and Haldane, F.D.M., 1980, Phil. Mag. (in press)
- Mott, N.F., Davies, E.A., and Street, R.A., 1975, Phil. Mag. 32, 961.

Figure Captions

Figure 1. Phase diagram for diamagnetic metastable states with energy ϵ_i and occupation $n_i = 2$ to fast emission of an electron to a site g with $n_j = 0$ and energy ϵ_j . The rectangle defines the metastable region when the process $i \rightarrow j$ is permitted for all i, j .

Figure 2. Phase diagram for paramagnetic metastable states with $n_i = 1$ and energy ϵ_i . The shaded region defines the region of stability for such states to fast emission or absorption of a single electron from metastable states with energy ϵ_j . The rectangular area (a) defines the metastable region when the process $i \rightarrow j$ is allowed for all i, j . The region (b) ignores processes involving only paramagnetic states and considers processes involving a narrower region of metastable states with $n_j = 0, 2$.

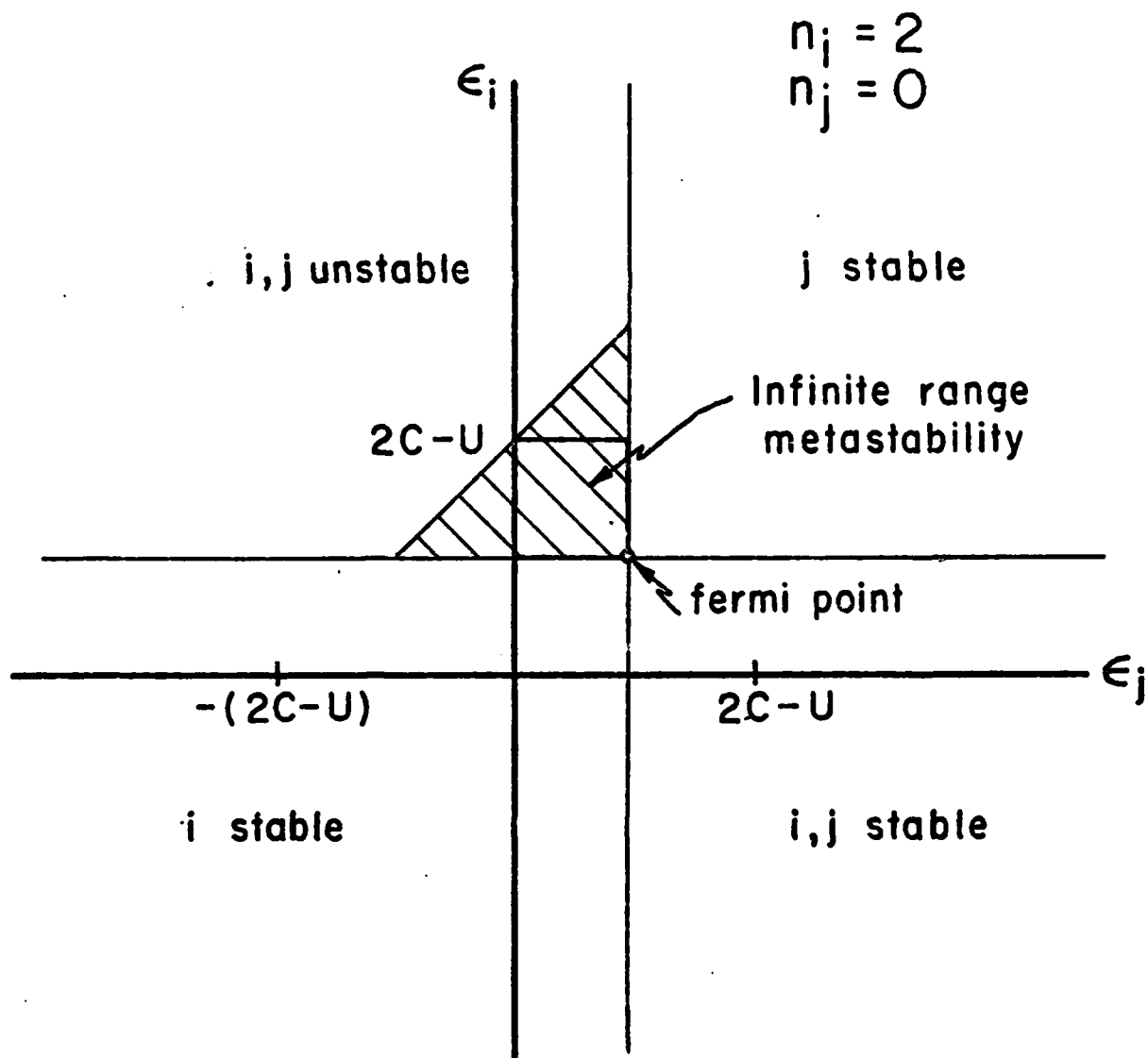


Fig. 1

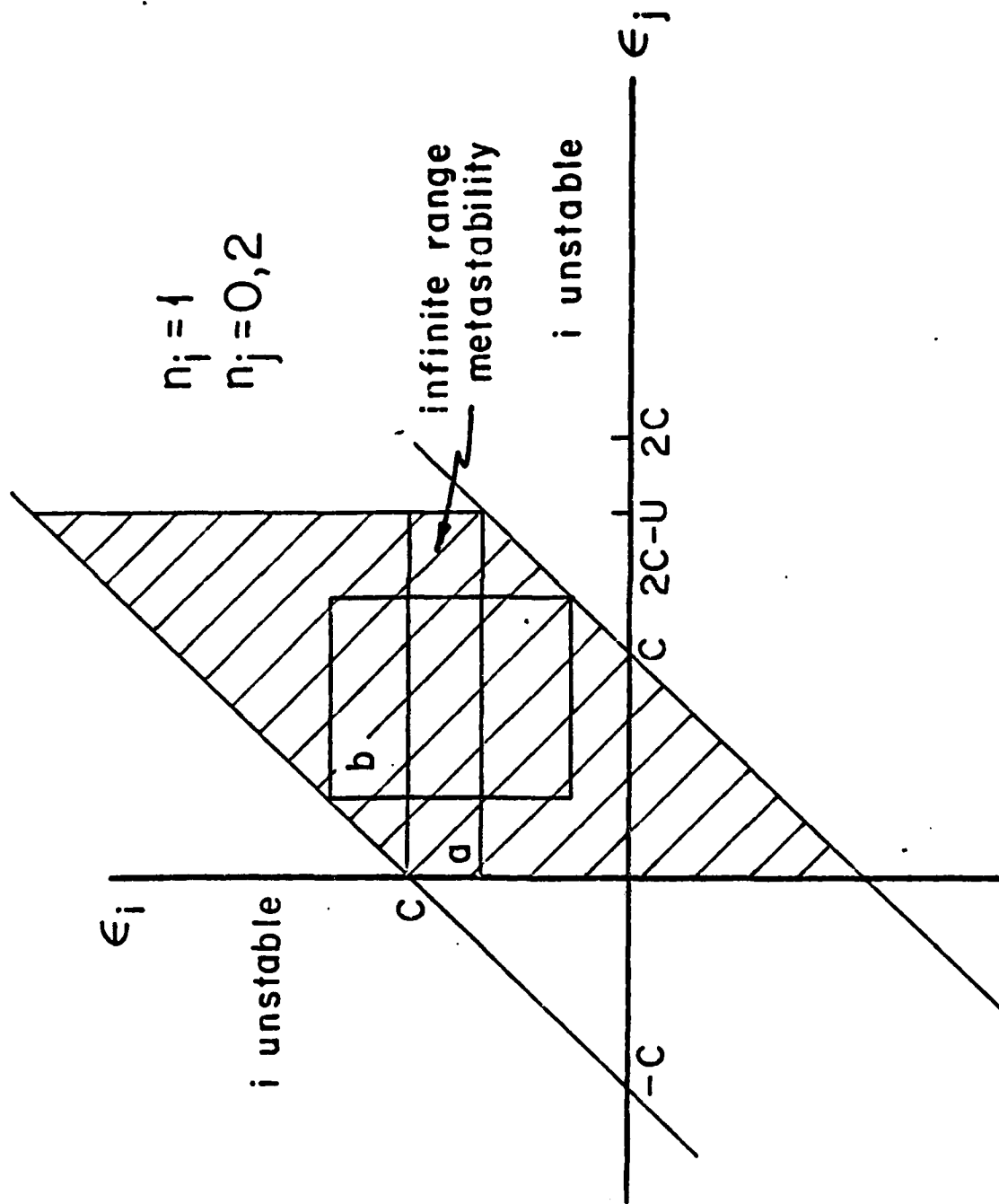


Fig. 2

UNCLASSIFIED

SECURITY CLASSIFICATION OF THIS PAGE (When Data Entered)

REPORT DOCUMENTATION PAGE		READ INSTRUCTIONS BEFORE COMPLETING FORM
1. REPORT NUMBER	2. GOVT ACCESSION	3. RECIPIENT'S CATALOG NUMBER
	AD A098461	
4. TITLE (and Subtitle)		5. TYPE OF REPORT & PERIOD COVERED
Further Remarks on Metastability in Amorphous Semiconductors		Technical Report
7. AUTHOR(s)		6. PERFORMING ORG. REPORT NUMBER
D.C. Licciardello		
9. PERFORMING ORGANIZATION NAME AND ADDRESS		8. CONTRACT OR GRANT NUMBER(s)
Department of Physics, Princeton University Princeton, N.J. 08544		N00014-77-C-0711
11. CONTROLLING OFFICE NAME AND ADDRESS		10. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS
Office of Naval Research (Code 427) Arlington, Va. 22217		NR 318-058
14. MONITORING AGENCY NAME & ADDRESS (if different from Controlling Office)		12. REPORT DATE
		April 1981
		13. NUMBER OF PAGES
		7 + 2 figures
		15. SECURITY CLASS. (of this report)
		unclassified
		15a. DECLASSIFICATION/DOWNGRADING SCHEDULE
16. DISTRIBUTION STATEMENT (of this Report)		
Approved for public release; distribution unlimited.		
17. DISTRIBUTION STATEMENT (of the abstract entered in Block 20, if different from Report)		
18. SUPPLEMENTARY NOTES		
Submitted to Journal of Physics C		
19. KEY WORDS (Continue on reverse side if necessary and identify by block number)		
glass, amorphous semiconductors, electronic states		
20. ABSTRACT (Continue on reverse side if necessary and identify by block number)		
The bandwidths for metastable states in a simple model of amorphous semiconductors is evaluated exactly. The region of metastability in the space of the model parameters (phase diagram) is deduced and the subregion corresponding to an infinite range model is indicated. The model predicts both paramagnetic and diamagnetic metastable states and the relationship between them is discussed.		

DD FORM 1473
1 JAN 73EDITION OF 1 NOV 65 IS OBSOLETE
S/N 0702-LF-014-6601

UNCLASSIFIED

SECURITY CLASSIFICATION OF THIS PAGE (When Data Entered)

