

AD-A110 299 NEW YORK UNIV NY COURANT INST OF MATHEMATICAL SCIENCES F/G 12/1
COMPUTING SOLUTIONS OF THE REDUCED WAVE EQUATION. (U)
SEP 81 C S MORAWETZ F49620-79-C-0193

UNCLASSIFIED

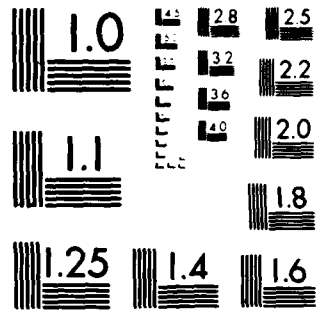
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FINAL REPORT

AIR FORCE GRANT NO. F-49620-79-C-0193

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The final report for Air Force Grant No. F-49620-79-C-0193 consists of the annual progress reports dated November 21, 1979 and August 25, 1980, as well as the following report for 1980-81:

During 1980-81 various numerical experiments were made on one-dimensional model inverse problems modelled by a method that is extendable to higher dimensions. The underlying problem is to recover the speed of propagation or the shape of an object from a scattered field.

The one dimensional problem that was investigated was based on the wave equation with potential

$$u_{tt} - u_{xx} + q(x)u = 0$$

where the function to be recovered is the potential $q(x)$. The model initial conditions were $u = 0$, $u_t = -2\delta'(x)$. The model boundary condition was $u = 0$ on $x = 0$. And the "extra" condition which determines $q(x)$ is either (a) $\partial u / \partial x$ on $x = 0$ for all time or (b) the scattered field $u \sim \delta(t-x) + R(t-x)$ at $x = +\infty$. The free space ($q=0$) solution is $u = \delta(x-t) - \delta(x+t)$. The distorted plane wave $P(x,t)$ is defined by the condition that it is a solution of the equation and at $x = +\infty$, $t \rightarrow +\infty$,

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it becomes $\delta(t-x)$ (exactly if q has compact support) and P is written as $\delta(t-x) + P_{\text{scat}}$.

The linearized equation based on the assumption that $q \ll 1$ and that terms of order q^2 may be neglected, yields for (a)

$$q = + 2 \frac{d}{dx} \left[\mathcal{A}(2x) - P_{\text{scat}}(0, -2x) \right]$$

with $\int_{-\infty}^t \frac{\partial u}{\partial x} dt \Big|_{x=0_{\infty}} = \mathcal{A}(t)$ for $t > 0$;

$$q = -2 \frac{d}{dx} R(2x)$$

for (b).

The numerical method is based on the well-known formula that follows from the propagation of singularities

$$q = 2 \frac{d}{dx} \lim_{t \rightarrow x+} u(x, x)$$

and then representing $u(x, x)$ by means of distorted plane waves. (All these have analogous formulas in higher dimensions.)

The analogues of the above equation are for (b)

$$q = -2 \frac{d}{dx} R(2x) + \frac{d}{dx} \int_0^{\infty} (P_{\text{scat}}(x, x-s) - P_{\text{scat}}(-x, x-s)) R(s) ds$$

and for (a) we insert in this formula



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$$R(s) = -\mathcal{H}(s) + P_{\text{scat}}(0, -s) + \mathcal{H}(0) \int_{-\infty}^s P_{\text{scat}}(0, -s') ds' \\ + \int_0^{\infty} \mathcal{H}(t) P_{\text{scat}}(0, t-s) dt$$

The idea is to use one of these formulas as part of an iteration. We guess q and P and recompute q using the above formulas for cases (a) or (b).

Note that if the given data vanishes in case (b) we recover $q \equiv 0$ but in case (a) we donot.

We have been unable to implement (a) as planned. The reason is that we must operate in a finite region and use a radiation condition at a finite distance. In our case for example with q a simple quadratic with support in $|x| < 2$, over $-8 \leq x \leq 8$ and $0 \leq t \leq \dots$. The errors produced are compounded in case (a). The potential could only be recovered to 40% of its original value with an error of about 5% of the maximum potential. However, the potentials were well beyond the range of the linear theory.

In case (b) even with a coarse mesh the results are surprisingly worse.

The results are being written up as a technical report to be issued at the Courant Institute.

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REPORT DOCUMENTATION PAGE		READ INSTRUCTIONS BEFORE COMPLETING FORM	
1. REPORT NUMBER AFOSR-TR- 31 -0879	2. GOVT ACCESSION NO AD A110 299	3. RECIPIENT'S CATALOG NUMBER	
4. TITLE (and Subtitle) COMPUTING SOLUTIONS OF THE REDUCED WAVE EQUATION,		5. TYPE OF REPORT & PERIOD COVERED FINAL, 21 NOV 79 - 30 SEP 81	6. PERFORMING ORG. REPORT NUMBER
7. AUTHOR(s) Cathleen S. Morawetz		8. CONTRACT OR GRANT NUMBER(s) F49620-79-C-0193	
9. PERFORMING ORGANIZATION NAME AND ADDRESS Courant Institute of Mathematical Sciences, New York University, 251 Mercer Street, New York NY 10012		10. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS 61102F; 2304/A3	
11. CONTROLLING OFFICE NAME AND ADDRESS Mathematical & Information Sciences Directorate Air Force Office of Scientific Research Bolling AFB DC 20332		12. REPORT DATE SEP 1981	
14. MONITORING AGENCY NAME & ADDRESS (if different from Controlling Office)		13. NUMBER OF PAGES 4	
		15. SECURITY CLASS. (of this report) UNCLASSIFIED	
		15a. DECLASSIFICATION DOWNGRADING SCHEDULE	
16. DISTRIBUTION STATEMENT (of this Report) Approved for public release; distribution unlimited.			
17. DISTRIBUTION STATEMENT (of the abstract entered in Block 20, if different from Report)			
18. SUPPLEMENTARY NOTES			
19. KEY WORDS (Continue on reverse side if necessary and identify by block number)			
20. ABSTRACT (Continue on reverse side if necessary and identify by block number) This report summarizes progress made in various numerical experiments made on one-dimensional model inverse problems.			

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