

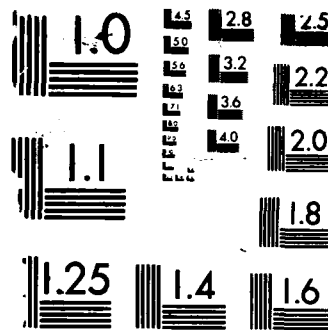
ON THE VAN DER WAALS MODES OF SOLUTE/SOLVENT CLUSTERS:
BENZENE-METHANE -D (U) COLORADO STATE UNIV FORT
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TECHNICAL REPORT #27

"ON THE VAN DER WAALS MODES OF SOLUTE/SOLVENT CLUSTERS:
BENZENE-METHANE, -DEUTEROMETHANE, AND -CARBON TETRAFLUORIDE"

by

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ABSTRACT

Clusters of benzene $(C_6H_6)_n$ and $(C_6F_6)_n$ are created in a supersonic molecular jet and are studied by two-color time-of-flight mass spectroscopy. The clusters' $S_1^n + S_0^n$ intermolecular vibronic structures are characterized by calculational modeling of the clusters' intermolecular motion. The calculations include: 1) an intermolecular normal coordinate analysis (NCA) which treats all six van der Waals modes under a harmonic oscillator assumption; and 2) a three dimensional hindered rigid rotor analysis (3D-HRRA) which treats the intermolecular torsional motion. Agreement between calculation and experiment is excellent for binding energy, symmetries, and van der Waals mode energies. The cluster spectra and calculated intermolecular modes are compared to those of benzene (C_6H_6) reported previously. A major conclusion of this work is that the clusters behave rigidly with regard to internal rotation of the cluster subunits and that the clusters possess unique equilibrium geometries. The internal torsional motion is oscillatory and that motion is constrained by an orientationally dependent intermolecular potential whose barrier height is of the order of the cluster binding energy.

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I. Introduction.

van der Waals (vdW) complexes of aromatic molecules with hydrocarbon solvents form a class of supramolecular systems whose intermolecular bonding has several interesting features. First, the intermolecular interaction is small and is dominated by long-range dispersive attractions and short-range exchange repulsions. These two features allow the interaction to be modeled using an intermolecular potential of known functional form. For example, the interaction can be modeled by additive atom-atom potentials set in a Lennard-Jones or an Exponential-Six format.¹ Second, the interaction potential surface only changes slightly, if at all, upon electronic excitation of the cluster chromophore. Thus, small spectral shifts of the chromophore electronic transition and weak intensity intermolecular vdW mode vibronic transitions are observed.^{2,3} Third, the interaction results in the formation of specific minimum energy cluster configurations. These geometries are interesting since they give insight into the nucleation processes and solvation geometry occurring in both gas and condensed phase systems.¹ Fourth, the small binding energy of the complex and the low frequency vdW vibrational modes are important since they play an essential role in intramolecular vibrational redistribution (IVR) and vibrational predissociation (VP) cluster dynamic energy transfer processes.⁴ And fifth, cluster structure, binding energy, vdW modes and dynamics (IVR and VP) are essentially dependent on the actual cluster structure and the intricacies of the intermolecular interaction.^{3,5}

In this paper, we report the spectroscopic results of the benzene deuterated methane ($\text{ben}(\text{CD}_4)_1$) and benzene carbon tetrafluoride ($\text{ben}(\text{CF}_4)_1$) clusters together with calculated modeling of selected cluster characteristics. The spectroscopic results include the $\pi^* \leftarrow \pi$ vibronic spectra of the clusters in their respective benzene constituent 6_0^1 regions.

The calculated results include the geometry, the binding energy, and the full eigenvalue/eigenvector intermolecular vibrational structure for the electronic ground state of each cluster. The intermolecular ground state vibrational structure is modeled by two methods: (1) an intermolecular normal coordinate analysis (NCA) which determines all six intermolecular vdW mode fundamentals under a harmonic oscillator assumption;³ and (2) a three dimensional hindered rigid rotor analysis (3D-HRRA) for which an anisotropic perturbation field is applied to the cluster solvents considered to be rigid internal rotors in the resulting potential well. The intermolecular torsional mode structure is thereby obtained under a hindered rigid rotor approximation.

The calculated ground state vibrational structure for each cluster is compared to its experimental vibronic spectrum and to the structure observed and calculated for the benzene methane ($\text{ben}(\text{CH}_4)_1$) cluster.^{1,3} Within this framework, vibronic selection rules which govern the clusters' $S_1 \leftarrow S_0$ transitions are discussed along with the ramifications of Herzberg-Teller (H-T) coupling for the observation of "forbidden" vibronic transitions.

The motivation for studying these systems concerns the elucidation of the vdW torsional structure. Two limiting cases can be proposed in regard to the torsional structure of $\text{ben}(\text{CH}_4)_1$, $\text{ben}(\text{CD}_4)_1$, and $\text{ben}(\text{CF}_4)_1$. In one case, the clusters possess free internal rotation between the cluster solute and solvent. The cluster solvent (CH_4 , CD_4 , or CF_4) rotates freely in three dimensions against the benzene framework and the system is considered nonrigid. In the other case, the clusters possess torsional oscillations for which the cluster solvent librates against the benzene frame with a residence time long enough to give rise to "vibration like" motion in a rigid molecule regime.

Elucidation of the vdW torsions is of particular interest in these systems since the aforementioned limiting cases pose questions regarding the actual physics governing the torsional structure: (1) do the clusters possess

free internal rotation between the cluster solute and solvent or do they possess torsional oscillations for which the lowest eigenstates behave more or less harmonically; and (2) in either case, what is the dependence of the intermolecular potential upon the relative orientation of the cluster constituents?

II. Experimental Procedures.

The apparatus and techniques employed in this study are largely the same as those previously reported.^{1,3} The $\text{ben}(\text{CD}_4)_1$ and $\text{ben}(\text{CF}_4)_1$ spectra are recorded using a pulsed supersonic molecular jet expansion in combination with two-color time-of-flight mass spectroscopy (2-color TOFMS). Two Nd^{+3} /YAG pumped dye lasers are used to probe the cluster spectra. An LDS 698 dye laser whose output is frequency doubled and mixed with the Nd^{+3} /YAG 1.064 micron fundamental is used to probe the clusters' $S_1 \leftarrow S_0$ transition in the isolated benzene 6_0^1 region ($^1\text{B}_{2u} \leftarrow ^1\text{A}_{1g}; \pi^* \leftarrow \pi$). Subsequent ionization of the clusters is accomplished using a R590 dye laser whose output is frequency doubled and mixed with the Nd^{+3} /YAG 1.064 micron fundamental. The ionization laser output is set at 45112 cm^{-1} . A 5% mixture of deuterated methane or carbon tetrafluoride in helium is placed in line with liquid benzene in a trap at room temperature. The three component mixture is expanded using a pulsed molecular jet nozzle having a 500 micron orifice while maintaining 100 psig backing pressure. Apparatus chamber pressure is maintained at or below 5×10^{-6} torr during the experiments.

III. Theoretical Considerations.

The NCAs are conducted using the same methods as described in our previous publication on vdW cluster vibronic structure.³ For $\text{ben}(\text{CD}_4)_1$, the calculated cluster ground state geometry, force field, and binding energy are taken as those of $\text{ben}(\text{CH}_4)_1$. Only the masses are changed for the deuterium/hydrogen substitutions in the NCA. For $\text{ben}(\text{CF}_4)_1$, the cluster ground state geometry and binding energy are calculated via intermolecular energy minimization employing the methods previously described.^{1,3} The NCA is conducted using

the modeled geometry and its corresponding force field. The potential parameters for the C-F ($A_{CF} = 6.08 \times 10^7 \text{ cm}^{-1}/\text{\AA}^2$; $C_{CF} = 1.13 \times 10^5 \text{ cm}^{-1}/\text{\AA}^6$) and H-F ($A_{HF} = 7.77 \times 10^6 \text{ cm}^{-1}/\text{\AA}$; $C_{HF} = 3.12 \times 10^4 \text{ cm}^{-1}/\text{\AA}$) atom-atom interactions are determined from parameters and methods described elsewhere.⁶

The 3D-HRRA centers upon analyzing the effects of the intermolecular potential on the cluster solvent rotational degrees of freedom. In the absence of a perturbing field from the cluster solute, the cluster solvent rotates freely in three dimensions against a fixed cluster solute framework. Application of a specified perturbation field results in the eventual restriction of this free rotation to torsional oscillation as the perturbation field magnitude is increased; that is, rotation ceases and vibrational oscillatory motion begins as the residence times of the solvent in the torsional potential well becomes longer.

In the 3D-HRRA, both cluster constituents are assumed internally rigid, their respective internal geometries remaining constant and at "equilibrium". The cluster solvents are taken as the rotating portions of the clusters since their zero field rotational constants are orders of magnitude greater than those of the cluster solute. The hindered solvent rotations are presupposed to contribute to the spectral features observed along with the intermolecular bending and stretching modes.

The 3D-HRRA involves setting up a molecule fixed coordinate system (x, y, z) and a space fixed coordinate system (ξ, η, ζ) on the cluster solvent as shown in Figure 1. Both systems have their origins at the nuclear center-of-mass of the cluster solvent. The molecule fixed coordinate system is chosen such that its principle axes lie along the C_2 rotational axes of the solvent tetrahedron (T_d point group symmetry). The space fixed coordinate system is chosen to be at the Euler angles $\theta = 45$ degrees, $\phi = 0$ degrees, and $\chi = 180$ degrees with respect to the molecule fixed coordinate system. The relative

orientation of these two systems is chosen so that the orientational dependence of the perturbation field can be expressed as a simple function of the curvilinear coordinates θ , ϕ , and χ .

An analytical expression for the potential governing the torsional motion can be derived by analyzing the solvent rotations in each of the rotational coordinates. Since a solvent rotation is periodic over 360 degrees, the potential can be expressed by an even function Fourier series written as⁷

$$V(q) = \frac{V_0}{2} + \sum_{p=1}^{\infty} A_p \cos(pq) \quad (1)$$

with q a function of the rotational coordinates and V_0 taken as the barrier height to internal rotation. The function q is determined by the potential dimensionality and the number of minima in the potential along each of the rotational coordinates. The first cosine term in the series describes the major form of the potential, and the rest of the terms in the series "fine tune" the potential shape. Since the coefficients A_p in these terms are assumed to be small, the series can be truncated at its first cosine term without sacrificing the general potential shape. Doing this yields the mathematically convenient potential form:

$$V(q) = \frac{V_0}{2} (1 - \cos(q)). \quad (2)$$

The symmetry of the cluster solvents can be employed to determine the dimensionality of and the number of minima in the potential and, hence, the function q . Rotating the cluster solvents against the benzene frame by 90 degrees in θ or χ , 270 degrees in χ , or 180 degrees in the ϕ coordinate results in equivalent maxima. Equivalent minima occur at 0 degrees rotation

in θ , ϕ , and χ and at 180 degrees rotation in θ and χ . The maxima have the same energy V_0 relative to the equivalent minima in all three rotational coordinates. Generalizing these relative displacements to all coordinate values, one finds that the potential magnitude depends upon the relative displacements of all three rotational coordinates as,

$$q = 2\theta + 2\chi + \phi. \quad (3)$$

Inserting eq. (3) into eq. (2) yields the three dimensional potential function used in the calculations:

$$V(\theta, \phi, \chi) = \frac{V_0}{2} [1 - \cos(2\theta + 2\chi + \phi)] \quad (4)$$

in which V_0 is the barrier height to internal rotation. The rotational wavefunctions chosen as the basis set for the calculation are the rigid rotor symmetric top wavefunctions which depend on the curvilinear coordinates and on the quantum numbers J , k and m . Under zero field conditions, these wavefunctions are solutions to the spherical top Schrödinger equation,⁸

$$-\frac{B}{\hbar^2} \hat{J}^2 |jkm\rangle = E_{\text{rot}}^J |Jkm\rangle \quad (5)$$

with eigenvalues corresponding to those of the spherical top cluster solvent. Upon application of the perturbation field, the eigenvalues are obtained via diagonalization of the energy matrix possessing elements of the form

$$\langle J' k' m' | -\frac{B}{\hbar^2} \hat{J}^2 + V(\theta, \phi, \chi) | J'' k'' m'' \rangle, \quad (6)$$

using eq. (1) for $V(\theta, \phi, \chi)$ with V_0 as an adjustable parameter. In the calculations performed, the only possible nonzero off-diagonal matrix elements occur with the following restrictions:

$$|J' - J''| \leq 2 \quad (7)$$

$$|k' - k''| = 2 \quad (8)$$

$$|m' - m''| = 1. \quad (9)$$

Diagonal matrix elements contain only the zero field spherical top energies and a $V_0/2$ potential term.

The resulting matrix is diagonalized for a selected value of V_0 yielding eigenvalues corresponding to the solvent torsional eigenstates at the specified perturbation. The perturbation is varied until a reasonable fit with the experimental spectrum is obtained. The calculations are performed on a Cyber 205 computer using a basis set consisting of 680 wavefunctions to ensure convergence of the lowest eigenstates at their proper eigenvalues. Only the torsional structures of $\text{ben}(\text{CH}_4)_1$ and $\text{ben}(\text{CD}_4)_1$ are calculated since their respective rotational constants are large enough that the coupling of the rotational levels does not require an extremely large basis set to insure convergence. Matrix elements are determined via numerical integration using a non-adaptive integration routine. The matrix is prediagonalized into a tridiagonal form using orthogonal similarity transformations⁹ and diagonalized using an implicit QL method.¹⁰ The rotational constants used for methane and deuterated methane are 5.2 cm^{-1} , and 2.6 cm^{-1} , respectively.

IV. Results.

A. $\text{Ben}(\text{CD}_4)_1$.

Fig. 2 and Table 1 present the $\text{ben}(\text{CD}_4)_1$ vdW cluster spectrum in the region between 38552.4 cm^{-1} and 38678.6 cm^{-1} using 2-color TOFMS. The

cluster 6_0^1 transition is "red shifted" by 41.2 cm^{-1} with respect to the benzene 6_0^1 . The bathochromic shift is comparable to that of $\text{ben}(\text{CH}_4)_1$ (41 cm^{-1}) indicating that the binding energy difference between S_1 and S_0 for $\text{ben}(\text{CD}_4)_1$ is nearly identical to that of $\text{ben}(\text{CH}_4)_1$. Ten pronounced intermolecular vibronic features are observed to the "blue" of the cluster 6_0^1 . As in the case of the $\text{ben}(\text{CH}_4)_1$ cluster, no features are observed in the symmetry forbidden benzene 0_0^0 region. Thus, the cluster geometry must possess at least a three-fold axis of symmetry.

The calculated $\text{ben}(\text{CD}_4)_1$ geometry (fig. 3) is assumed to be the same as that calculated for $\text{ben}(\text{CH}_4)_1$. The geometry possesses C_{3v} point group symmetry. In this geometry, the CD_4 center-of-mass lies $3.47 \text{ \AA}$ above the benzene molecular plane along the three-fold rotational axis. The ground state binding energy in this configuration is calculated at 540 cm^{-1} . Using the cluster "red shift" of 41.2 cm^{-1} , the excited state binding energy is calculated to be 581 cm^{-1} .

The NCA reveals six vdW vibrations (fig. 3 and Table 1), two being two-fold degenerate. The ground state vibrational energies are 75 cm^{-1} for the vdW stretch $s_z(a_1)$, 15 cm^{-1} for the bends $b_{xy}(e)$, and 20 cm^{-1} ($t_z(a_2)$) and 64 cm^{-1} ($t_{xy}(e)$) for the vdW torsions. The eigenvector normal modes (fig. 3) transform in an identical fashion to those calculated for $\text{ben}(\text{CH}_4)_1$. The vdW stretching mode transforms as the translation of the cluster constituents away from one another along the z (3-fold) axis. The vdW bends transform as some combination of cluster constituent translations in opposite directions perpendicular to the 3-fold axis in the xy plane. One vdW torsion transforms as a rotation of the cluster constituents about the 3-fold axis in opposite directions. The remaining two vdW torsions transform as rotations about orthogonal axes perpendicular to the 3-fold axis.

The 3D-HRRA results are also included in Table 1 for the lowest eigenstates in the torsional manifold. The torsional mode structure is calculated

for $V_0 = 300 \text{ cm}^{-1}$ and $B = 2.6 \text{ cm}^{-1}$. This perturbation results in a reasonable fit for the torsional features observed experimentally. The torsional "zero point energy" is 75 cm^{-1} . Two distinct torsional manifolds result from the calculations: one manifold has eigenvalues grouped quartically (nearly four fold degenerate located at approximately 21 cm^{-1} , 46 cm^{-1} , and 75 cm^{-1} above the zero point energy) and the other manifold has eigenvalues grouped octally (nearly eight fold degenerate located at approximately 68 cm^{-1} above the torsional zero point energy).

For comparison, the observed vibronic features for $\text{ben}(\text{CH}_4)_1$ are reproduced in Table 1 along with the results of the NCA and the 3D-HRRA. The NCA results for $\text{ben}(\text{CH}_4)_1$ are those reported previously.³ The 3D-HRRA torsional structure is calculated for $V_0 = 300 \text{ cm}^{-1}$ and $B = 5.2 \text{ cm}^{-1}$. This perturbation is chosen since it is assumed that the barrier to internal rotation is nearly identical in the two systems as they only differ by isotopic substitution and have the same electronic structure. Additionally, using the same potential barrier of electronic origin for both CH_4 and CD_4 clusters provides a check on the validity of the model in predicting the torsional mode structure of the clusters. For $\text{ben}(\text{CH}_4)_1$ the torsional "zero point energy" is 96 cm^{-1} . As in the case of the $\text{ben}(\text{CD}_4)_1$ calculations, two distinct torsional manifolds result: one manifold with eigenvalues grouped quartically at about 29 cm^{-1} , 54 cm^{-1} and 73 cm^{-1} above the torsional zero point energy and the other manifold with eigenvalues grouped octally at about 89 cm^{-1} above the zero point.

B. $\text{Ben}(\text{CF}_4)_1$.

Fig. 4 and Table 2 present the $\text{ben}(\text{CF}_4)_1$ cluster spectrum recorded using 2-color TOFMS in the region between 38578.6 cm^{-1} and 38702.6 cm^{-1} . Unlike the $\text{ben}(\text{CD}_4)_1$ and $\text{ben}(\text{CH}_4)_1$ cluster 6_0^1 transitions, the $\text{ben}(\text{CF}_4)_1$ cluster 6_0^1 is blue shifted by 6.1 cm^{-1} with respect to the benzene 6_0^1 . The small hypsochromic shift indicates that the binding energies in the S_0 and S_1 states are

nearly identical with the ground state binding energy being slightly greater. The relative displacement between the two potential surfaces is also small since only 5 vdW transitions are observed and their intensities decrease abruptly at about 50 cm^{-1} above the cluster 6_0^1 origin. No cluster spectrum is observed in the forbidden benzene 0_0^0 region indicating that the cluster possesses at least a three-fold rotation axis. This result is not unexpected since the same observation is made for the $\text{ben}(\text{CH}_4)_1$ and $\text{ben}(\text{CD}_4)_1$ systems.

The ground state configuration and vdW eigenvalues/eigenvectors are shown in fig. 5 and Table 2. Only one minimum energy geometry is calculated for the cluster. The geometry has C_{3v} point group symmetry with the CF_4 center-of-mass at 3.43 Å above the benzene molecular plane along the 3-fold axis. The cluster ground state binding energy is 1064 cm^{-1} .

The NCA reveals six vdW vibrations. Their ground state vibrational energies are 69 cm^{-1} for the vdW stretch $s_z(a_1)$, 11 cm^{-1} for the vdW bends $b_{xy}(e)$, and 13 cm^{-1} ($t_z(a_2)$) and 36 cm^{-1} ($t_{xy}(e)$) for the vdW torsions. The eigenvector normal modes transform as the translational and rotational representations of the C_{3v} point group as indicated and in the same manner as those for $\text{ben}(\text{CH}_4)_1$ and $\text{ben}(\text{CD}_4)_1$.

V. Discussion.

In comparing the calculated ground state vdW vibrational structure and experimental vdW vibronic structure, we assume that the intermolecular potential surfaces of the clusters studied are identical in both the S_1 and S_0 electronic states. Additionally, we utilize the group theoretical arguments developed in our previous publication³ on vdW cluster vibronic structure to assign and understand the observed cluster spectra.

A. $\text{Ben}(\text{CD}_4)_1$.

The calculated ground state vdW vibrations (NCA) of $\text{ben}(\text{CD}_4)_1$ and the experimental vibronic spectrum is shown in fig. 2 and Table 1. As previously, the assignments are made by direct comparison between the calculations and the experimental vibronic spectra. The intense feature at 22.5 cm^{-1} to the blue of

the cluster 6_0^1 corresponds to the t_z torsion calculated at 20 cm^{-1} . Thus this feature and its observed overtones at 48.4 cm^{-1} , 70.5 cm^{-1} , and 100 cm^{-1} are assigned to 6_0^1 vdW combination bands $6_0^1 t_z(a_2)_0^1$, $6_0^1 t_z(a_2)_0^2$, $6_0^1 t_z(a_2)_0^3$, and $6_0^1 t_z(a_2)_0^4$. The features at 32 cm^{-1} and 62.7 cm^{-1} to the blue of the cluster 6_0^1 correspond to the first and third overtones of the 2-fold degenerate vdW bending modes whose fundamental is calculated at 15 cm^{-1} . They are assigned to the 6_0^1 combination bands $6_0^1 b_{xy}(e)_0^2$, and $6_0^1 b_{xy}(e)_0^4$, respectively. The feature at 58.1 cm^{-1} to the blue of the cluster 6_0^1 is assigned to a 2-fold degenerate t_{xy} vdW torsions/cluster 6_0^1 combination band, $6_0^1 t_{xy}(e)_0^1$. The t_{xy} torsion fundamental is calculated at 64 cm^{-1} . With this assignment and that of the t_z torsion, the feature at 84 cm^{-1} is assigned to the $6_0^1 t_{xy}(e)_0^1 t_z(a_2)_0^1$ combination band. The vdW stretch is calculated at 75 cm^{-1} . This mode is identified in the cluster vibronic spectrum at 72.9 cm^{-1} to the blue of the cluster 6_0^1 origin and is thus assigned as $6_0^1 s_z(a_1)_0^1$. Under this scheme, we assign the features at 93.3 cm^{-1} and 97.5 cm^{-1} as $6_0^1 s_z(a_1)_0^1 b_{xy}(e)_0^1$ and $6_0^1 s_z(a_1)_0^1 t_z(a_2)_0^1$ combination bands.

The torsional structure resulting from the 3D-HRRA for $\text{ben}(\text{CD}_4)_1$, (Table 1), confirms the torsional assignments made using the NCA. The quartically grouped torsional levels at 21 cm^{-1} , 46 cm^{-1} and 75 cm^{-1} correspond to those associated with the t_z torsion and its overtones in the cluster spectrum at 22.5 cm^{-1} , 48.4 cm^{-1} and 70.5 cm^{-1} . The octally grouped torsional levels at 68 cm^{-1} correspond to the t_{xy} torsion assigned at 58.1 cm^{-1} .

The correspondence between these levels and those calculated using the NCA can be understood as follows. In the 3D-HRRA, four symmetrically equivalent minima exist in the torsional potential surface. These minima correspond to the four ways of placing the solvent tetrahedron upon the solute with a tetrahedral face toward the solute molecular plane. If the barrier between these minima is infinitely high, penetration of the local wavefunctions through the barrier

separating the potential minima does not take place. In this case, each "well" contains eigenstates corresponding to intermolecular torsions which occur with small amplitude about each potential well minimum. Since the four "potential wells" are identical in shape and depth, a four-fold "structural" degeneracy exists in which all four potential wells contain identical torsional structure. Thus, for example, a nondegenerate torsional eigenstate actually has a four-fold structural degeneracy, etc.

Considering the actual situation in which the potential barrier is finite, tunneling occurs and the "structural" degeneracy is lifted via interaction of the local wavefunctions through the potential barrier. The eigenstate splitting due to this tunneling may, or may not, be observed depending on the experimental resolution and the relative difference between the eigenstate energy and the barrier height. Splitting of the "structural" degeneracy in the lower portion of the potential well is minimal unless the barrier is low. Based on the calculations, energy level splittings of the first few sets of eigenstates should not be observed unless the barrier is below 150 cm^{-1} .

At moderate barriers (about 300 cm^{-1}), the eigenstates in the lower portion of the well are nearly degenerate and behave more or less harmonically. For all practical purposes, we can assume these levels to be degenerate. The torsional level structure can then be determined from the eigenstates in one of the minima. Thus, the 3D-HRRA torsional structure calculation simplifies into the NCA. Physically, the cluster can be considered at least "semi-rigid" in the respect that it has a definable equilibrium configuration. The potential energy barrier separating one minimum from the others in the potential surface is large and may be of the order of the cluster binding energy.

The vibronic structure in both the $\text{ben}(\text{CH}_4)_1$ and $\text{ben}(\text{CD}_4)_1$ spectra suggests that the systems are more or less rigid. None of the observed vibronic state energies follow a free rotor formalism for which the energy level structure

is described by eq. (2). If the systems behaved nonrigidly, the free rotor eigenstates would lie at approximately $2.6J(J+1) \text{ cm}^{-1}$ and $5.2J(J+1) \text{ cm}^{-1}$ for $\text{ben}(\text{CD}_4)_1$ and $\text{ben}(\text{CH}_4)_1$, respectively, and a $\Delta J = \pm 1$ selection rule would govern the transitions.

In the $\text{ben}(\text{CH}_4)_1$ and $\text{ben}(\text{CD}_4)_1$ systems, the t_z torsional mode shifts by 17.6 % upon deuterium substitution of the cluster solvent. Both the NCA and 3D-HRRA models predict a 28.6% frequency shift for t_z , while the free rotor model would predict a 50% frequency shift for t_z upon deuteration. Considering that only one mode is taken into account and the mode couplings may be different in the protonated and deuterated clusters, we conclude that torsional tunneling does not take place in the lower levels of the $\text{ben}(\text{CH}_4)_1$ and $\text{ben}(\text{CD}_4)_1$ cluster potential wells and that the clusters can be considered to be rigid.

The negligible isotopic red shift in the b_{xy} bends also suggests that the cluster is rigid and that the entire observed spectrum is not solely due to internal rotation. The same rationale holds for the vdW stretch: its observation also dispels the notion that only rotor modes occur in the spectrum. No experimental isotopic shift can be determined for the stretching mode, however, since it is not observed in the $\text{ben}(\text{CH}_4)_1$ system. Theoretically, the mode should red shift by 8.5% upon deuteration of the cluster solvent. The $\text{ben}(\text{CH}_4)_1$ stretch should then be at about 80 cm^{-1} in the vibronic spectrum based upon the observed stretch in the $\text{ben}(\text{CD}_4)_1$ system. (The $\text{ben}(\text{CH}_4)_1$ stretch is calculated by the NCA to be at 82 cm^{-1}).

The $\text{ben}(\text{CD}_4)_1$ cluster spectrum is also substantially richer than the $\text{ben}(\text{CH}_4)_1$ spectrum in the respect that both the vdW stretch and t_{xy} torsions are observed along with well developed t_z torsion and bend progressions. This is probably due to more favorable Franck-Condon factors in the $\text{ben}(\text{CD}_4)_1$ case resulting from the isotopic substitution.

As is the case for $\text{ben}(\text{CH}_4)_1$, H-T coupling influences the vibronic intensities (selection rules) in the $\text{ben}(\text{CD}_4)_1$ spectrum. The nontotally symmetric modes t_z , t_{xy} and b_{xy} all appear in the spectrum with $\Delta v = \pm 1, \pm 2, \pm 3, \dots$ selection rules as can be seen in fig. 2 and Table I.

B. $\text{Ben}(\text{CF}_4)_1$.

Comparison of the calculated ground state vdW vibrations (NCA) of $\text{ben}(\text{CF}_4)_1$ and the experimental cluster spectrum is presented in fig. 4 and Table II. The intense feature at 17.4 cm^{-1} to the blue of the $\text{ben}(\text{CF}_4)_1$ 6_0^1 corresponds to the t_z torsion calculated at 13 cm^{-1} . Thus, this feature and its first overtone at 35.4 cm^{-1} are assigned to the 6_0^1 vdW combination bands $6_0^1 t_z(a_2)_0^1$ and $6_0^1 t_z(a_2)_0^2$. The shoulder feature at 15.7 cm^{-1} to the blue of the cluster 6_0^1 corresponds to the vdW 2-fold degenerate b_{xy} bend fundamental calculated at 11 cm^{-1} . Taking this feature as the combination band $6_0^1 b_{xy}(e)_0^1$ leads to the assignment of the b_{xy} bend first overtone $6_0^1 b_{xy}(e)_0^2$ at 32.3 cm^{-1} . Finally, the spectral feature at 39.5 cm^{-1} in the $\text{ben}(\text{CF}_4)_1$ spectrum corresponds to the t_{xy} torsion calculated at 36 cm^{-1} . Thus, the feature is assigned to the combination band $6_0^1 t_{xy}(e)_0^1$.

The $\text{ben}(\text{CF}_4)_1$ spectrum, like the $\text{ben}(\text{CD}_4)_1$ and $\text{ben}(\text{CH}_4)_1$ spectra, suggests that the system is rigid. Since the rotational constant for CF_4 is small (ca. 0.18 cm^{-1}), the free rotor energy level structure should appear at about $.37 \text{ cm}^{-1}$ intervals. This structure is not observed. Instead, the spectrum possesses oscillatory torsional structure commensurate with the NCA theoretical predictions. These theoretical and experimental results demonstrate the rigidity of the systems.

The $S_1 \leftarrow S_0$ excitation of the $\text{ben}(\text{CF}_4)_1$ cluster involves very little change in cluster geometry. In this spectrum, the progression intensities decrease dramatically at energies greater than the 6_0^1 cluster origin. In the lowest two vdW mode progressions, the intensities decrease approximately 13%

between the mode fundamentals and the first overtones. Furthermore, the high energy vdW stretch calculated at 69 cm^{-1} is not observed. At these high energies, the Franck-Condon factors must be very small.

The observation of the nontotally symmetric vdW b_{xy} bends and the t_z torsion progressions with a $\Delta v = \pm 1, \pm 2 \dots$ selection rule suggests that interelectronic state mixing (H-T coupling) is an important contributor to the mode intensity mechanism. As in the $\text{ben}(\text{CH}_4)_1$ and $\text{ben}(\text{CD}_4)_1$ cases, H-T coupling becomes apparent in the low lying vdW modes of the $\text{ben}(\text{CF}_4)_1$ system. In fact, the interelectronic state mixing is substantial in the $\text{ben}(\text{CF}_4)_1$ system and can be demonstrated by the observation of the well defined bend fundamental at 15.7 cm^{-1} and by the observation of the t_{xy} torsion fundamental (39.5 cm^{-1}). Both of these latter features should not be observed if H-T coupling is not present.

VI. Summary and Conclusions.

Two-color TOFMS and supersonic molecular jet techniques have been employed to study the $S_1 \leftarrow S_0$ vibronic spectra of $\text{ben}(\text{CD}_4)_1$ and $\text{ben}(\text{CF}_4)_1$ vdW clusters. These studies reveal detailed information regarding the geometries, the intermolecular energetics, and the physical nature of the vdW interactions present in the systems. The experimental observations demonstrate that the clusters are at least semi-rigid systems possessing unique equilibrium geometries and that the intermolecular motion present in the systems is oscillatory. Through comparison of these spectra with those of $\text{ben}(\text{CH}_4)_1$ previously studied in this laboratory, the "isotopic" shifts observed in the intermolecular vibronic structures demonstrate that all the low-lying intermolecular eigenstates are nearly harmonic. They are not admixtures of oscillatory intermolecular bends and stretches and free intermolecular rotation as would occur if the clusters behaved nonrigidly.

ment of the intermolecular motion. Both the NCA and 3D-HRRA models indicate that the intermolecular torsional motion is oscillatory and that the motion is constrained by an orientationally dependent intermolecular potential. Detailed analysis of the intermolecular interaction reveals that the torsional motion is governed by a three dimensional potential possessing torsional barriers on the order of the cluster binding energy. At this barrier magnitude, torsional tunneling in the lower portions of the potential is minimal and the residence time of the cluster constituents in a particular well minimum is substantially longer than the timescale of experimental observation.

The assignment of geometry and identification of the intermolecular modes using the theoretical models has also proven useful in understanding the physics governing the spectroscopic properties of the clusters. The observation of non-totally symmetric intermolecular mode fundamentals and combination bands in the vibronic spectra suggest that interelectronic state mixing (H-T coupling) is an important factor in the overall intensity mechanism governing the transitions. Additionally, the interelectronic state mixing is most important for the low-lying intermolecular bending and torsional modes.

Both the NCA and the 3D-HRRA adequately model the intermolecular modes in the systems studied; however, the NCA is the more useful here due to its mathematical simplicity and passive treatment of all six intermolecular modes in a single calculation. Unlike the 3D-HRRA, the NCA does not require detailed a priori knowledge of the potential surface along each coordinate of motion. Furthermore, the large amount of energy level coupling occurring in the 3D-HRRA at moderate barriers suggests that the model is best used in systems possessing free or nearly free internal rotation.

The rigidity of these clusters can be qualitatively understood by considering the nature of the solvent rotational dynamics. Any rotational state for which $J \neq 0$ has a displacement in the θ coordinate: ϕ and χ displacements may or may not be present depending on the values of the k and m quantum numbers. Given the geometry of the system then, low barrier one-dimensional rotation (solvent rotating around the cluster three-fold C-X bond axis) is not possible. Basically the two dimensional hindering in $\text{ben}(\text{CX}_4)_1$ clusters is insufficient to produce the one dimensional rotation occurring in "attached tops" such as toluene.¹¹

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TABLE I

vdW spectral features in $\text{ben}(\text{CD}_4)_1$ and $\text{ben}(\text{CH}_4)_1$ 6_0^1 regions
and calculated ground state vdW modes (refer to fig. 2).

Ben(CD ₄) ₁ Energy ^{a)}				Ben(CH ₄) ₁ Energy ^{a)}			
Observed	NCA	3D-HRRA ^{c)}	Assignment ^{b)}	Observed	NCA	3D-HRRA ^{c)}	Assignment ^{b)}
0(38567.4)	0	0(4)	6_0^1	0(38567.6)	0	0(4)	
16.1	15(τ_{xy})		$6_0^1 b_{xy}^1$	16.1	16(τ_{xy})		
22.5	20(τ_z)	21(4)	$6_0^1 \tau_{zo}^1$	27.3	28(τ_z)	29(4)	
32			$6_0^1 b_{xy}^2$	32.3			
				48.4			
48.4		47(4)	$6_0^1 \tau_{zo}^2$	51.4		54(2)	
58.1	64(τ_{xy})	68(8)	$6_0^1 \tau_{xyo}^1$				
62.7			$6_0^1 b_{xyo}^4$	64.6			
70.5		75(4)	$6_0^1 \tau_{zo}^3$	73.5		73(4)	
72.9	75(τ_z)		$6_0^1 \tau_{zo}^1$		82(τ_z)		
84		80(12)	$6_0^1 \tau_{xyo}^1 \tau_{zo}^1$				
					89(τ_{xy})	84(8)	
93.3			$6_0^1 b_{xyo}^1$				
97.5			$6_0^1 \tau_{zo}^1$				
100			$6_0^1 \tau_{zo}^4$				

a) Energies are reported in cm^{-1} relative to the 6_0^1 cluster origins.

b) vdW mode representations as per fig. 2.

c) Values in parentheses represent "structural" degeneracy calculated in 3D-HRRA.

d) From references 1 and 3.

TABLE II

vdW spectral features in $\text{ben}(\text{CF}_4)_1 6_0^1$ region and calculated
ground state vdW modes (refer to fig. 4).

$\text{ben}(\text{CF}_4)_1$ a) Energy		Assignment b)	
Observed	NCA		
0(38614.7)	0	6_0^1	
15.7	11 (b_{xy})	$6_0^1 b_{xy}^1$	
17.4	13 (t_z)	$6_0^1 t_{zo}^1$	
32.3		$6_0^1 b_{xy}^2$	
35.4		$6_0^1 t_{zo}^2$	
39.5	36 (t_{xy})	$6_0^1 t_{xy}^1$	
	69 (s_z)		

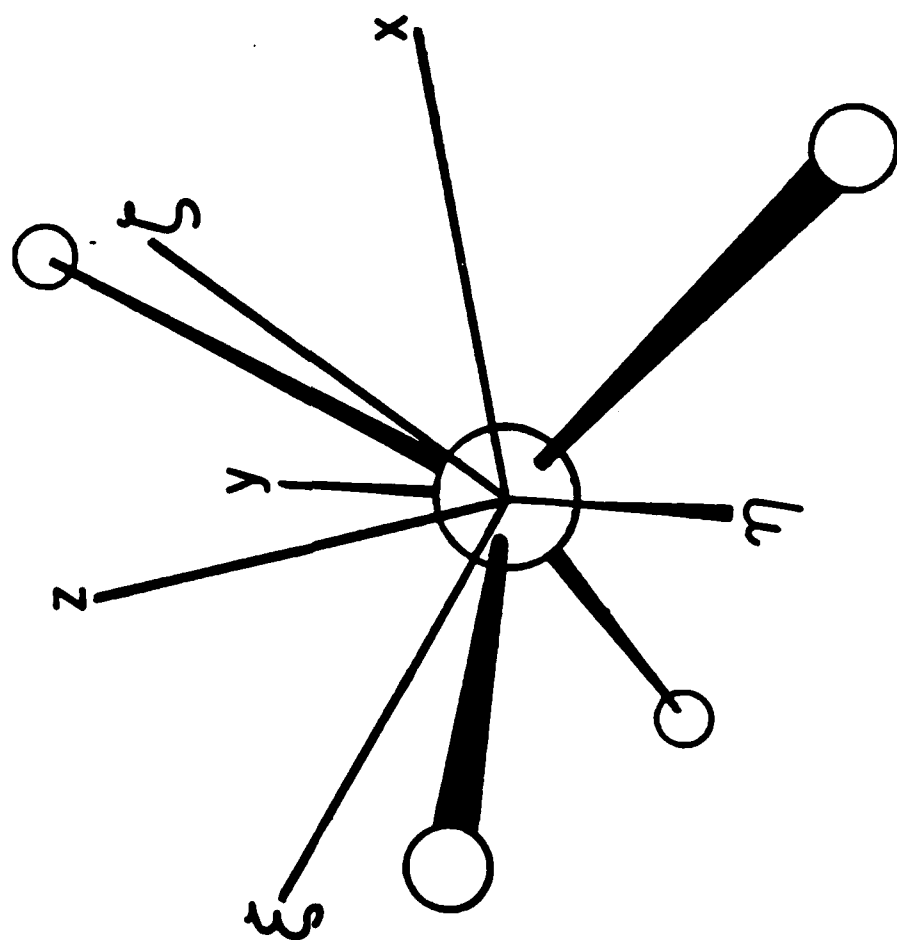
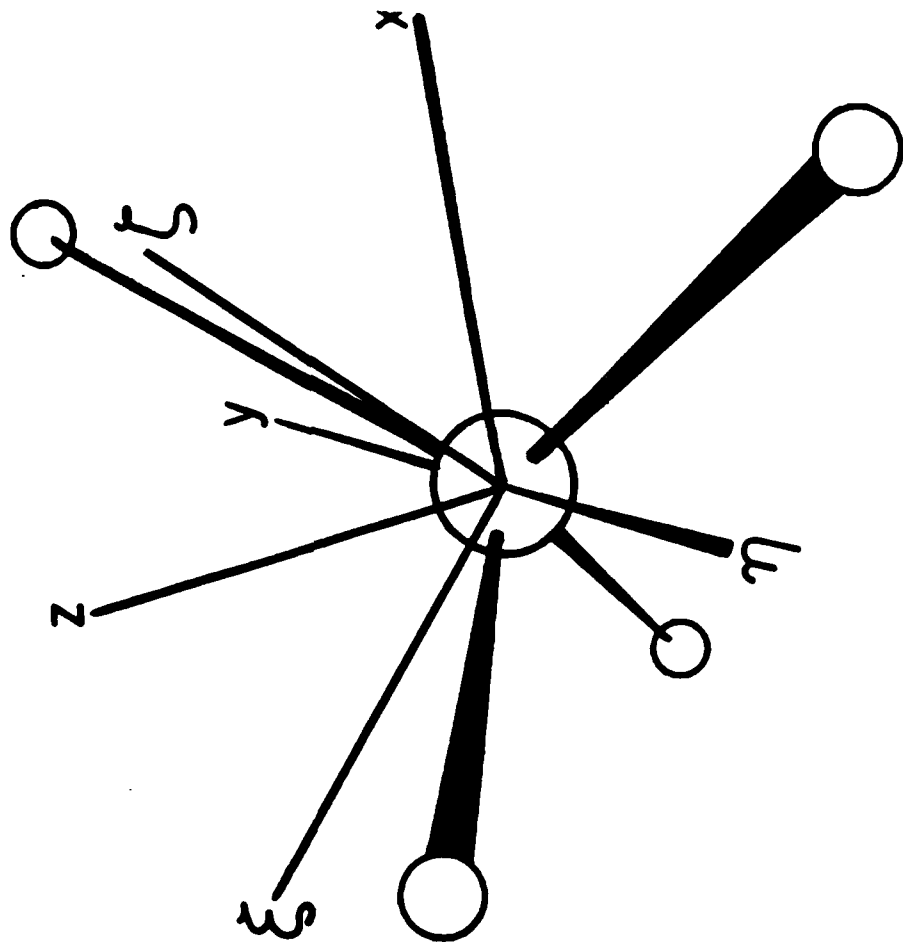
a) Energies are reported in cm^{-1} relative to the 6_0^1 cluster origin.

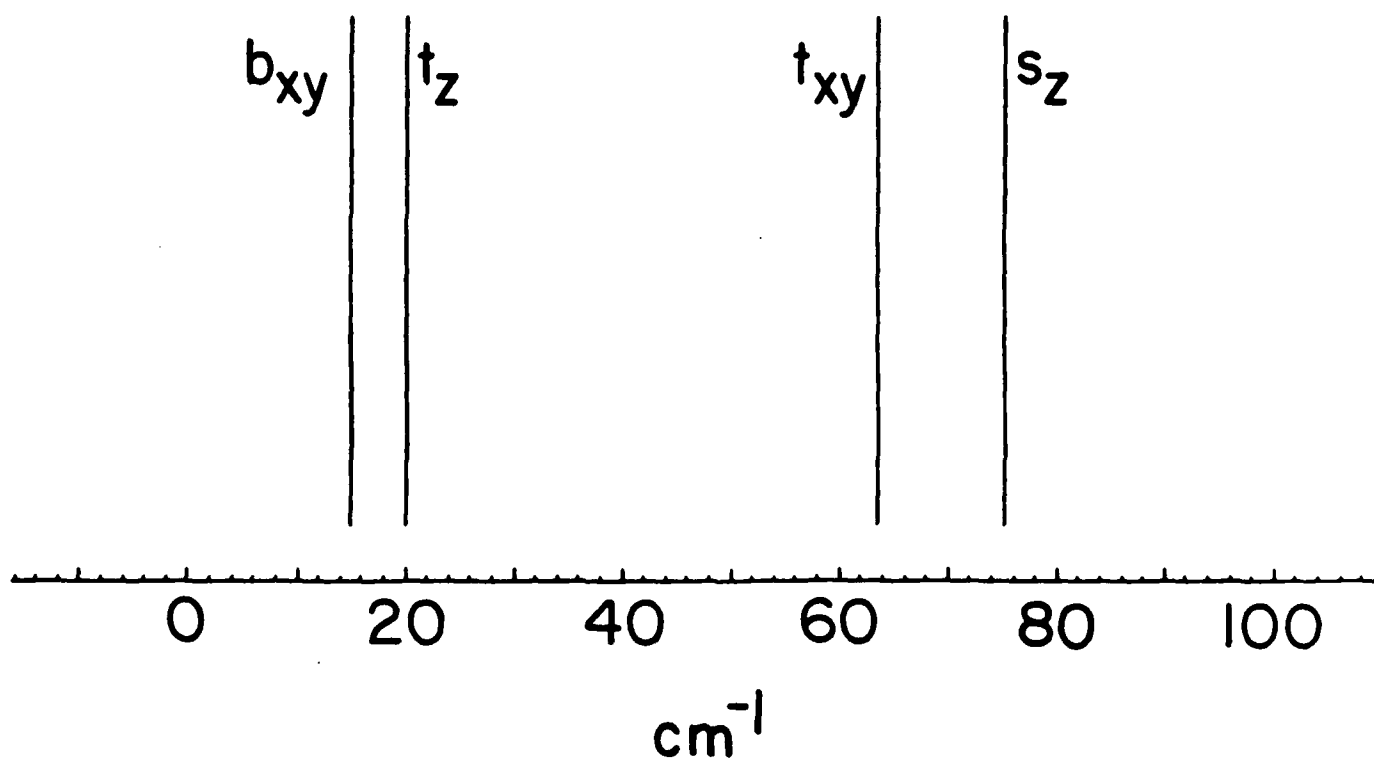
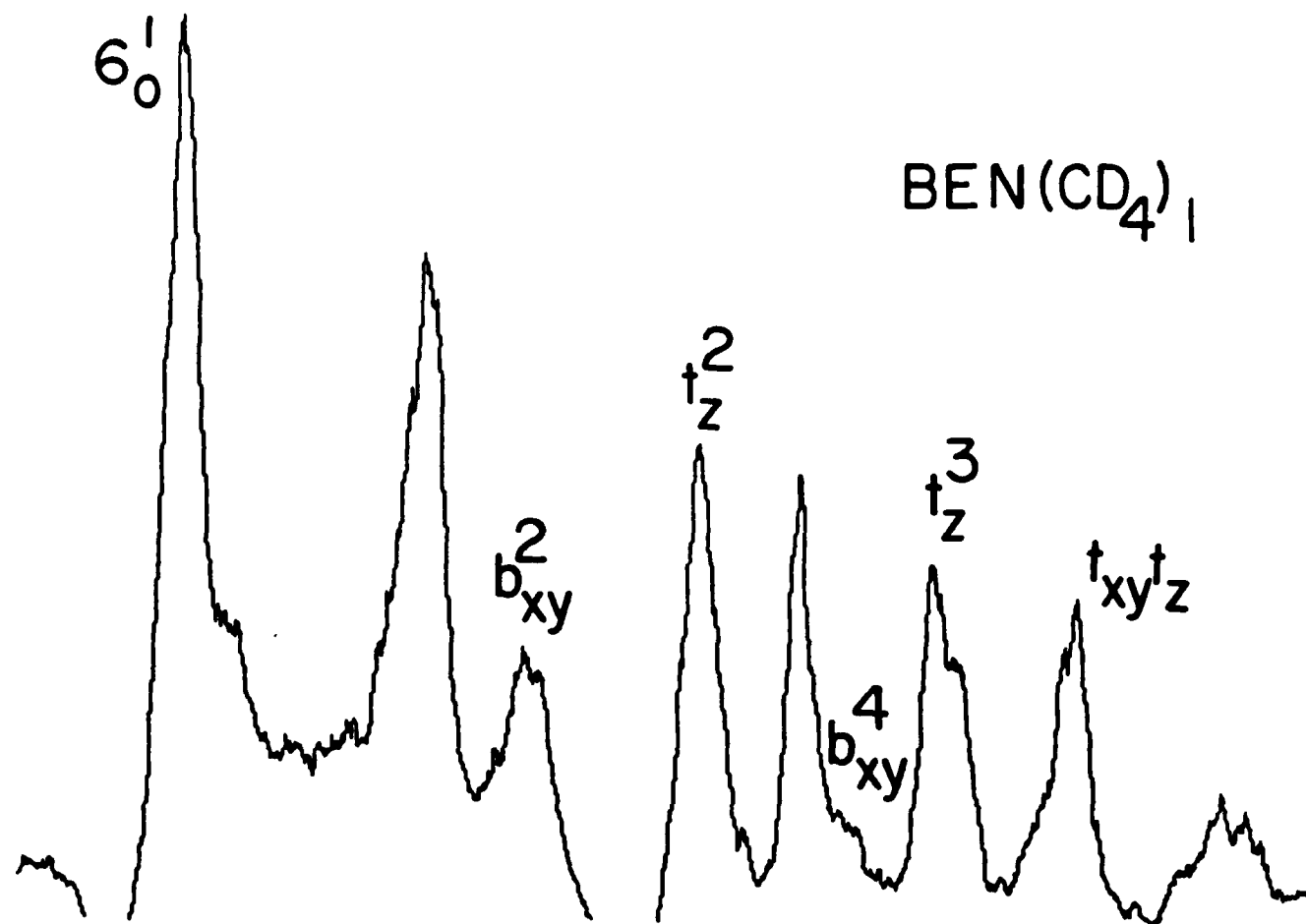
b) vdW mode representations as per fig. 4.

FIGURE CAPTIONS

- Figure 1 Stereoscopic projection of solvent tetrahedron showing relative orientation of the molecule-fixed (x,y,z) and space-fixed (ξ,η,ζ) coordinate systems used in the 3D-HRRA. The (x,y,z) coordinate system is located at $\theta = 45^\circ$, $\phi = 0^\circ$, and $\chi = 180^\circ$ from the (ξ,η,ζ) coordinate system.
- Figure 2 Two-color TOFMS $S_1 \leftarrow S_0$ spectrum and calculated ground state vdW modes (NCA) of $\text{ben}(\text{CD}_4)_1$. Energy scale is relative to $\text{ben}(\text{CD}_4)_1$ 6_0^1 transition (38567.4 cm^{-1}). Feature positions and assignments as per Table I and fig. 3.
- Figure 3 Calculated ground state minimum energy configuration (a) and NCA eigenvalue/eigenvector vdW modes (b) - (g) for $\text{ben}(\text{CD}_4)_1$. Cluster symmetry is C_{3v} with an equilibrium intermolecular distance of $3.47 \text{ \AA}$. Eigenvectors are normalized and displayed at 2x magnification ($2 \text{ \AA}$ total displacement).
- Figure 4 Two-color TOFMS $S_1 \leftarrow S_0$ spectrum and calculated ground state vdW modes (NCA) of $\text{ben}(\text{CF}_4)_1$. Energy scale is relative to $\text{ben}(\text{CF}_4)_1$ 6_0^1 transition (38614.7 cm^{-1}). Feature positions and assignments as per Table II and fig. 5.

Figure 5 Calculated ground state minimum energy configuration (a) and
NCA eigenvalue/eigenvector normal modes (b) - (g) for $\text{ben}(\text{CF}_4)_1$.
Cluster symmetry is C_{3v} with an equilibrium intermolecular
distance of 3.43 Å. Eigenvectors are normalized and displayed
at 2x magnification (2 Å total displacement).





BEN(CD4)1

C_{3v}

(a)



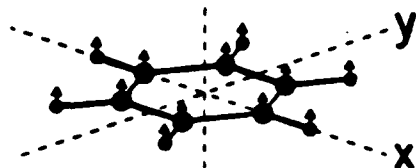
B.E. = 540 cm^{-1}

(b)

z

75 cm^{-1}

s_z stretch (a_1)

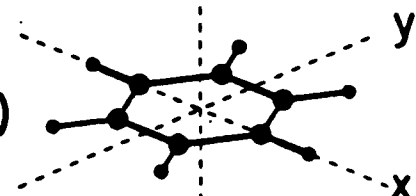


(c)

z

20 cm^{-1}

t_z torsion (a_2)

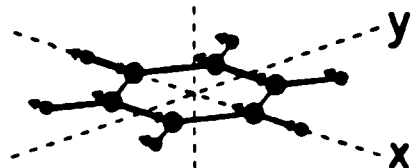


(d)

z

15 cm^{-1}

b_x bend (e)

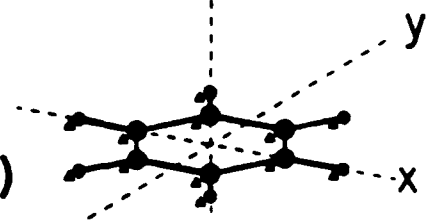


(e)

z

15 cm^{-1}

b_y bend (e)

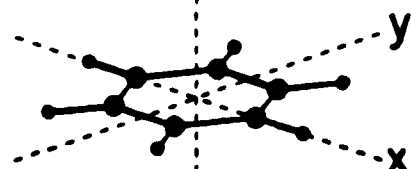


(f)

z

63 cm^{-1}

t_x torsion (e)

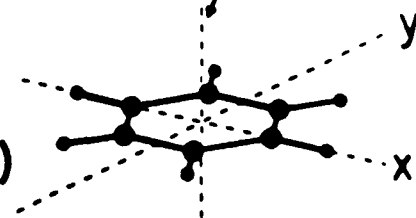


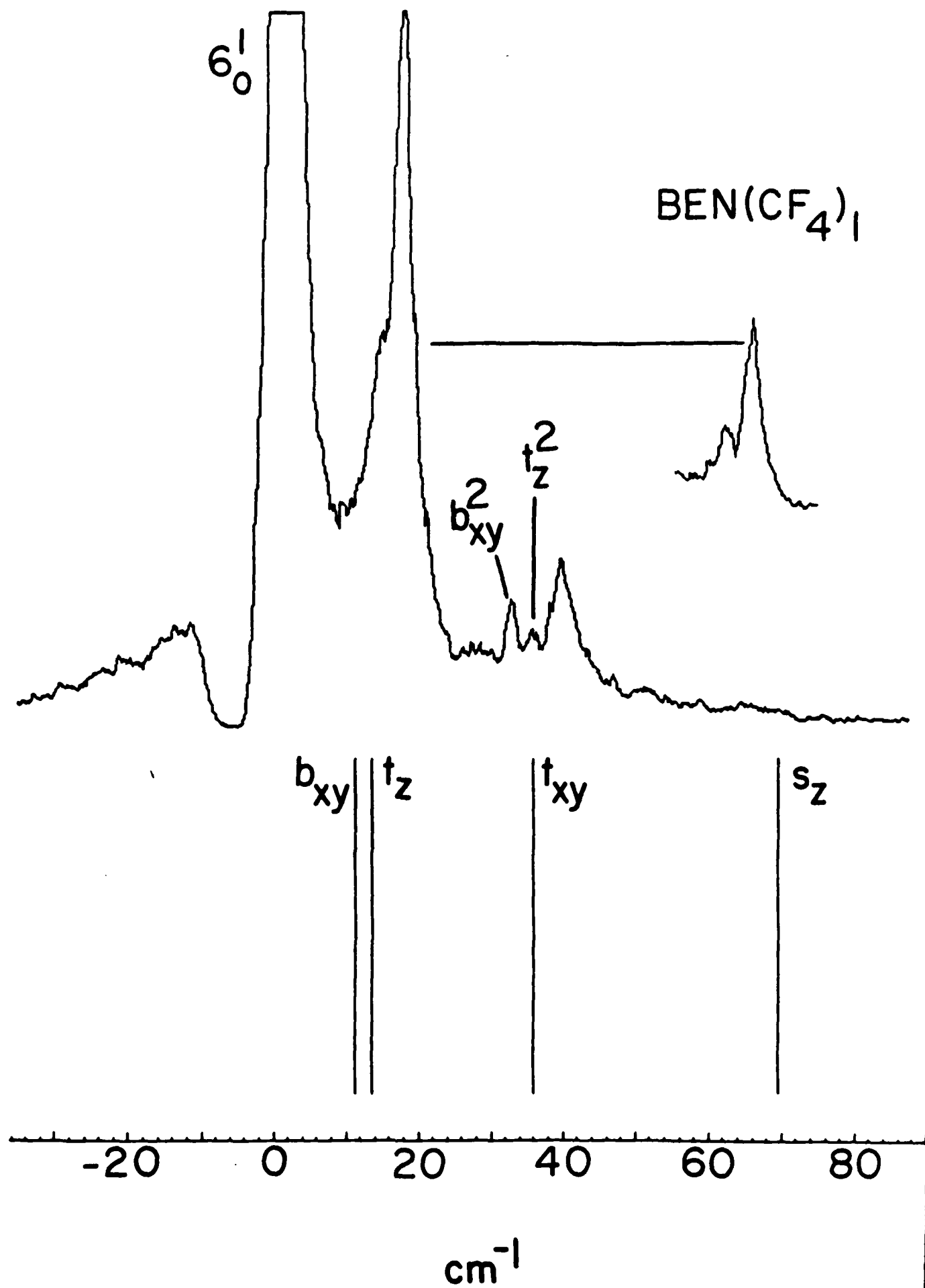
(g)

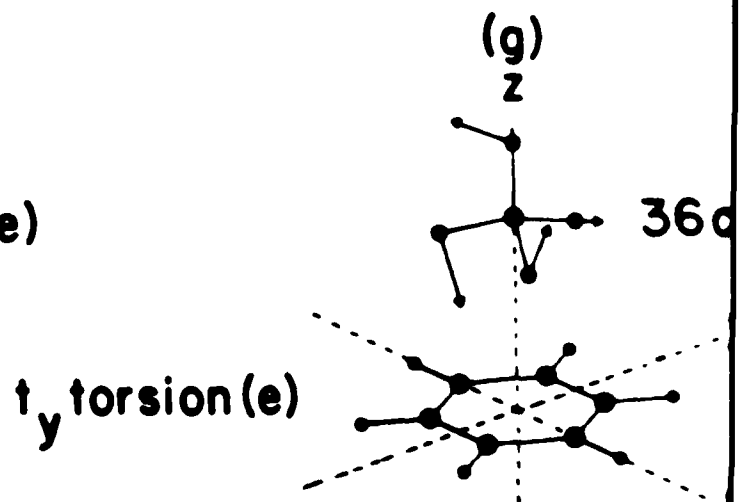
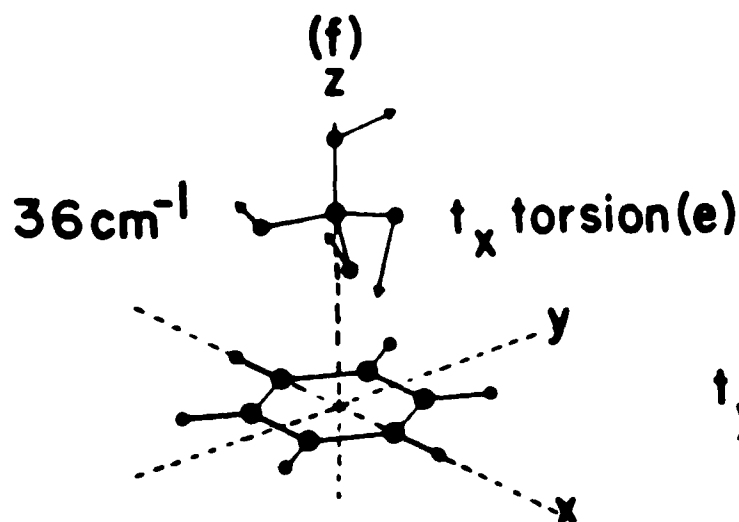
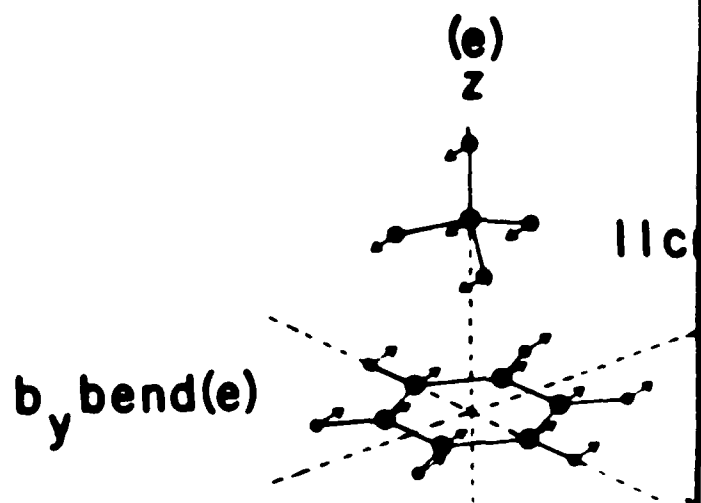
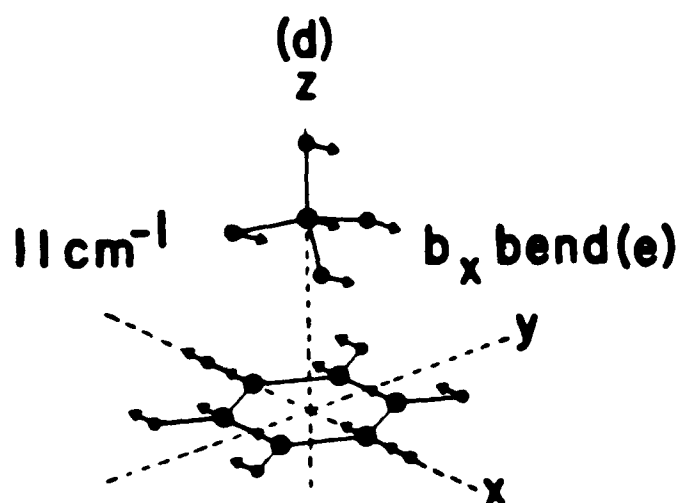
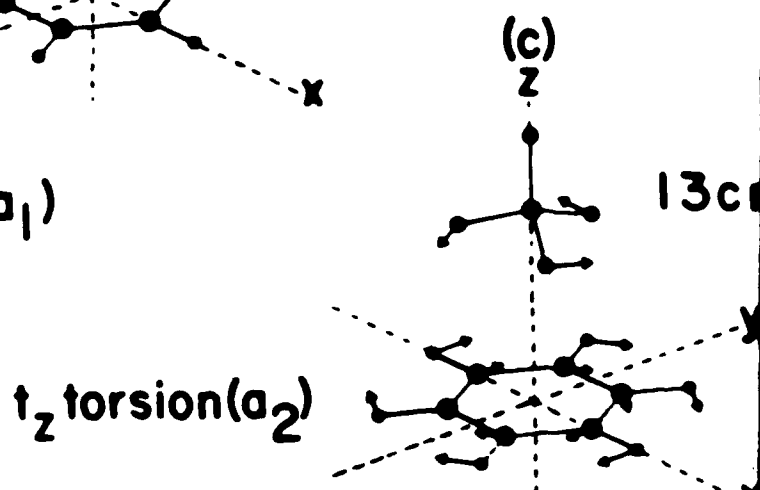
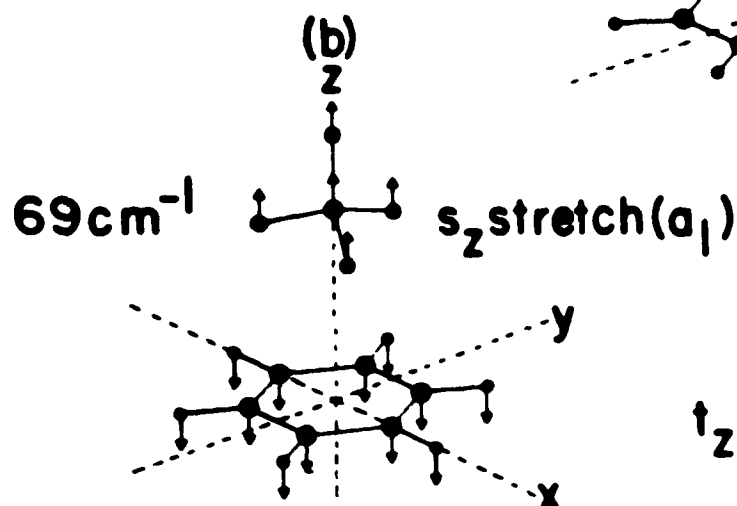
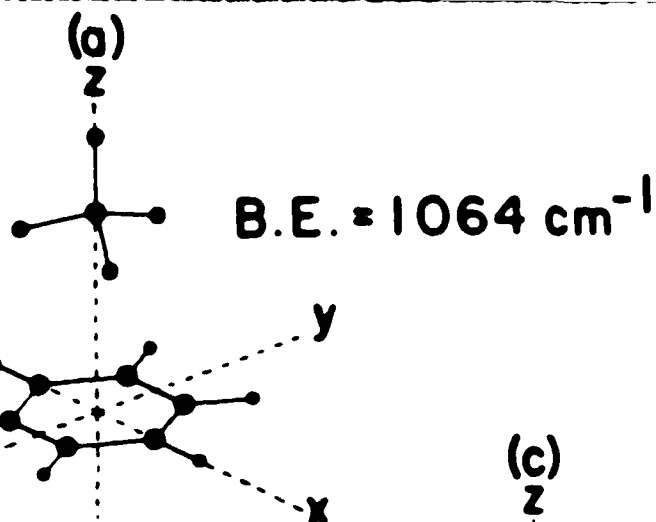
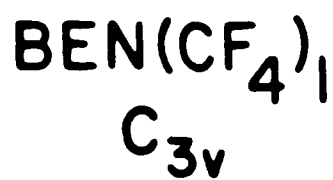
z

63 cm^{-1}

t_y torsion (e)







END

8-87

DTIC