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INCREMENTAL DRAG DUE TO GROOVES AND THREADS FOR KE PROJECTILES

AMEER G. MIKHAIL

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19. ABSTRACT (Continue on reverse if necessary and identify by block number) A prediction method for estimating the increment of increase in the drag coefficient due to the grooves on the surface of the body of a Kinetic Energy (KE) projectile is established based on wind tunnel data. Known existing methods give very poor results when compared to the experimental data. Twenty cases were considered for the present correlations and validation. The results achieved are very satisfactory. The present analysis considers parameters not included in the existing methods: the groove pitch (or alternately the groove depth); and multi-groove types residing on the same projectile body. The groove drag decreases very rapidly with the increase in Reynolds number and also decreases considerably with the increase in Mach number. The range of validity of the present expressions coincides with the usual KE projectile requirement, which establishes the Mach number range to be $3.5 < M < 5.5$ and $\alpha = 0^\circ$. <i>Subsonic projectile</i>					
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I. INTRODUCTION

There are certain types of flight vehicles, namely the Kinetic Energy (KE) projectiles, which must have relatively large body grooves to hold a sabot which functions as a bore-rider inside the gun tube and which separates at a short distance from the muzzle of the gun. These grooves, which usually extend over half of the total length of the projectile, increase the axial force coefficient and therefore the drag. The grooves also affect the aerodynamics of the fins of these projectiles; however, this effect will not be considered in this study.

In recent years, long KE projectiles (those with $L/D > 20$) have been receiving closer attention due to their terminal ballistics effectiveness. For these long penetrators, the groove drag is an increasingly important component of the total drag of the vehicle. It is for this class of projectiles that a method was needed for predicting the increase in the drag coefficient due to these grooves. Reference 1 offers wind tunnel results for this class of projectiles and represents a basis for validation of the models as well as a basis for a correlation study.

Because the groove drag is small for shorter projectiles ($L/D < 20$) at Mach number of 5, not many investigations have been carried out for its study. Also, because the effect of grooves on aerodynamics is only of interest to projectile designers, no investigations have been made in the corresponding missile aerodynamics studies. Therefore, most short KE projectiles are currently designed without any consideration to groove drag. Among the available literature on the subject, Reference 2 offers an expression for calculating ΔC_{Dg} , without any basis and without any reference to any experimental

data for validation. Reference 3 offers a more logically acceptable expression which is validated using four data points; but which is shown to provide much higher values (300% higher) in comparison with the present wind tunnel data. In addition, Reference 3 does not disclose the Reynolds number of the wind tunnel tests

With only poor prediction capabilities being available, the present work was pursued to fill a void and to provide a prediction capability based on the correlation of the experimental data of Reference 1 as well as other data provided by Reference 4. The results of this correlation will be compared to the predictions of both References 2 and 3.

The present correlations include the direct effect of the groove pitch (or the groove depth, since they are directly related for all standard threads). The present analysis also allows for multiple groove types on the same body. An example for this particular projectile configuration is shown in Figure 1 together with the general nomenclature.

The present work limits itself to zero angle of attack. Therefore, distinction between drag and axial force vanishes and the drag coefficient will be used throughout this report.

II. WIND TUNNEL DATA AND TEST

Reference 1 documents tests performed at the supersonic wind tunnel of the Naval Surface Weapons Center Laboratory at White Oak, Maryland. The tests were performed at Mach numbers 3.5, 4.0 and 5.0, with the majority of tests run at $M = 5.0$. The Reynolds number for the tests varied between 4.0 through 5.8×10^6 per foot ($12.2-19.0 \times 10^6$ per meter).

All the models had a diameter of 0.94 inch (23.9 mm). The models were sting-supported. Four lengths were tested, with L/D of 20.59, 25.59, 30.58 and 35.59. Both body-alone (without fins) smooth-surface and body-alone grooved-surface configurations were tested.

The groovings are made on two body sections, a forward and a rear section. The forward section is always grooved while the rear was either smooth, threaded or grooved. The notation (G/T) is used to denote Grooved forward section and Threaded rear section. Figure 2 displays the three basic combinations together with the totally smooth body. Tests were also performed with and without fins; however, the present work excludes the cases with fins in order to obtain a better evaluation of the effects of grooves on the axial force coefficient. The grooves were 8 per inch (i.e., pitch = 1/8 inch, or 3.17 mm) while the threads were 32 per inch (pitch = 1/32 inch, or 0.79 mm).

Four different noses were tested and they closely approximate the 10° cone (semi-vertex angle). They are two Sears-Haack noses and a bi-conical shape in addition to the 10° cone. In the present study, the nose configuration does not affect the results since only the drag increase between smooth and grooved configuration, ΔC_{D_g} , is considered. Most of the tests, however, were run with the Sears-Haack nose. This nose configuration and details of the forward and rear grooves are shown in Figure 2b.

Examination of the wind tunnel results of Reference 1 indicated that the measurements are not always perfect and the accuracy of the result may be affected. For example, the axial force coefficient for some cases did not exhibit symmetry for positive and negative α . Some obvious errors for other cases for the normal force existed, usually for negative α . With such long projectiles being supported by a sting from its base, possible rod bending may partially explain the source of asymmetry with respect to α .

Seventeen test cases were identified from the data of Reference 1 and they are listed in Table 1. The total pressure of the wind tunnel varied between 2.4 and 8.2 times the atmospheric value.

A second set of data was obtained from Reference 4. This data consisted of three tests at Mach numbers 3.49, 3.98 and 4.75 with a corresponding tunnel Reynolds number of 12, 14 and 18×10^6 per foot ($39.4, 45.9$ and 59.1×10^6 per meter), respectively. The tests were run at the Vought Corporation wind tunnel in Dallas, Texas. The wind tunnel model is shown in Figure 3, with grooving of four grooves per inch (pitch = 0.25 inch or 6.35 mm). Both smooth and grooved models were tested. The results and test conditions for these cases are listed in Table 2.

III. ANALYSIS AND CORRELATIONS

1. REVIEW OF EXISTING EXPRESSIONS

A quick review of existing expressions will be presented so that differences and similarities with the present work will be more clear.

Reference 2 provides an empirical relation in the form of:

$$\Delta C_{Dg} = [0.00025 M^{3.9} l_g] \cdot C_{D_{TSB}} \quad (1)$$

where ΔC_{Dg} is the incremental drag coefficient due to grooves, based on the projectile reference area $A_{ref} (= \pi D^2/4)$, where D is the reference diameter for the projectile

M is the free stream Mach number of the projectile

l_g is the length of the grooved portion, in calibers

$C_{D_{TSB}}$ is the body-alone total drag coefficient (including base drag) for the smooth body (i.e., without grooves) configuration.

Reference 3 provides the following expression for computing the incremental drag.

$$\Delta C_{Dg} = 1.6 \frac{l_g}{l_c} \cdot C_{D_{SFSB}} \quad (2a)$$

where l_c is the length of the cylindrical portion of the body, in calibers,

$C_{D_{SFSB}}$ is the drag coefficient due to skin friction of the smooth body of the cylindrical portion, l_c , of the body,

$$C_{D_{SFSB}} = (4d_c l_c) \cdot c_f \quad (2b)$$

The skin friction coefficient, c_f , is defined as:

$$c_f = \frac{\tau}{0.5 \rho_\infty V^2 A_{ref}},$$

where τ is the skin friction shear stress at the projectile surface

and d_c is the diameter of the cylindrical portion, in calibers.

Reference 3 did not specify the expression used for c_f , in its application to the projectile configuration shown in Figure 4. In this study, an expression for c_f for turbulent flow was used from Reference 5 as:

$$c_f = \frac{0.455(1 + 0.2M_\infty^2)^{-0.32}}{(\log_{10} R_{e_L})^{2.58}} \quad (2c)$$

where R_{e_L} is the Reynolds number based on the total length of the projectile.

2. PRESENT ANALYSIS

The present approach is similar in general to the data correlation approach used earlier in Reference 6. The influencing physical parameters are first identified then some of the wind tunnel data are used to determine the constants in the correlation. Finally, the rest of the data is used for validation of the obtained expressions.

a. First expression This expression is targeted only for "typical" KE projectiles as shown in Figure 1. More specifically, the "typical" KE projectile is defined as one with nose cone semi-vertex angle less than 9° , having a very small nose radius $< 0.03 D$, with no boattail, and having one main cylindrical diameter. This expression is not valid for KE projectiles shown in Figure 3. This expression is formulated in the form of Eq. (1) (of Reference 2), where the body-alone, smooth, total body drag coefficient is known, and only the incremental groove drag is required. This approach suits projectile designers where the groove drag is considered after all other design and shape requirements have been fulfilled.

First, considering the physical parameters that do affect the groove drag, one might write:

$$\Delta C_{D_g} = F(M, R_e, l_g, p, h, x_0)$$

where R_e is the Reynolds number per unit length,
 p is the groove pitch ($p = 1/n$ where n is the number of grooves per unit length),
 h is the depth of the groove,
 x_0 is the distance from the nose tip to the first groove.

The parameter x_0 was neglected based on the results of Reference 3 which indicated no measureable effect on ΔC_{D_g} of slightly varying x_0 . Also, the

parameter h was eliminated since both p and h are directly related for all standard groove geometries. Therefore, the expression will be a function of the pitch, p , which also can usually be more easily measured than h . The above expression is then simplified and written as:

$$\Delta C_{D_g} = [F(M, R_e, l_g, p)] \cdot C_{D_{TSB}}.$$

Based on the data for the 17 cases of Reference 1, the first correlation in the present work for typical projectiles is introduced as:

$$\Delta C_{D_g} = \left[2.05 \, TF_1 \cdot \frac{l_g}{L} \cdot MF_1 \cdot RF_1 \right] \cdot C_{D_{TSB}} \quad (3a)$$

where TF_1 is the Thread Factor defined as:

$$TF_1 = 0.84 + 0.117 \left(\frac{p}{0.031} \right) - 0.007 \left(\frac{p}{0.031} \right)^2 \quad (3b)$$

where p is the groove pitch in inches,

MF_1 is the Mach number Factor introduced as:

$$MF_1 = \frac{1}{M_\infty (1.453 - 0.067 M_\infty)} \quad (3c)$$

RF_1 is the Reynolds number Factor introduced as:

$$RF_1 = \left(\frac{4.2 \times 10^6}{R_e} \right)^{0.8} \quad (3d)$$

and R_e is the Reynolds number per foot.

For multiple grooves with different pitch on the same projectile, Eq. (3a) takes the form:

$$\Delta C_{D_g} = \left[2.05 \left(\frac{TF_{11} l_{g1} + TF_{12} l_{g2}}{L} \right) \cdot MF_1 \cdot RF_1 \right] \cdot C_{D_{TSB}} \quad (4)$$

where TF_{11} and TF_{12} are the thread factors for the threads of pitch p_1 and p_2 , respectively.

One can notice the large differences between the expression of Eq. (3a) and the expression of Eq. (1) of Reference 2. The latter does not include Reynolds number dependence or groove pitch dependence. Further, ΔC_{D_g} increases with the increase in Mach number, while in reality it decreases with Mach number. Also, Eq. (1) depends directly on l_g while Eq. (3a) is found to be more appropriate if l_g/L is used.

Finally, Eq. (1) is found, as will be shown in the results, to give results as high as 10 times the experimental value. Equation (3a) follows the data within $\pm 13\%$.

b. Second expression For non-typical KE projectiles, as shown in Figure 3, it is not adequate to use $\Delta C_{D_{TSB}}$ as an input. This is because large variations in nose or boattail angles can cause large changes in $\Delta C_{D_{TSB}}$ even though the groove drag on the bodies will not change.

Therefore, an approach like that of Reference 3, which relates only to the cylindrical part of the projectile (i.e., excluding the nose and the boattail), should be more appropriate. Therefore, this expression is targeted to be more general than that of Eq. (3).

Based on the data of References 1 and 4, the second correlation for the general projectile configuration was found to have the form:

$$\Delta C_{D_g} = \left[2.59 TF_2 \cdot \frac{l_g}{l_c} \cdot MF_2 \cdot RF_2 \right] \cdot C_{D_{SFSB}} \quad (5a)$$

where $TF_2 = \left(1 + 0.1 \left(\frac{p}{0.031} \right)^a \right) , \quad (5b)$

$$a = 0.528 \left(\frac{p}{0.031} \right)^{0.462} ,$$

$$MF_2 = M^{-[3.164 + 3.225 M^{0.4} - 0.156 M^{1.5}]} , \quad (5c)$$

$$RF_2 = \left(\frac{9 \times 10^6}{R_{e_L}} \right)^{0.6} , \quad (5d)$$

and $C_{D_{SFSB}}$ is as given by Eq. (2b).

For configurations with multiple thread types, Equation (5a) becomes:

$$\Delta C_{Dg} = 2.59 \cdot \left(\frac{TF_{21} l_{g1} + TF_{22} l_{g2}}{l_c} \right) \cdot MF_2 \cdot RF_2 \cdot C_{D_{SFSB}} \quad (6)$$

Comparison of the present expression of Eq. (5) and that of Eq. (2) (i.e., expression of Reference 3) reveals the following facts. First, the present analysis includes explicitly the effects of Mach number, groove pitch and Reynolds number. None of these effects is considered by Reference 3. Second, as will be shown, the results of applying the expression of Reference 3, in general, overpredict the wind tunnel data by 300%.

Since the second correlation expression (Eq. (5)) is more general, i.e., can be applied to all KE configurations (including the configuration of Figure 3), it is recommended over the first correlation expression. However, in cases where both expressions apply, the results of those two expressions should be very close to each other.

IV. RESULTS AND COMPARISONS

Equation (3) of the first expression was first applied to the configurations of Reference 1. Seventeen cases were computed. Also the expressions of References 2 and 3 were applied and compared with the experimental data.

Figures 5a-d show the variation of the incremental drag with the increase of Mach number for the four different L/D's of 20, 25, 30 and 35, respectively. Figure 5a shows the predictions of Reference 3 to be twice as high as the data, while those of Reference 2 are much higher. Also, the trend given by the expression of Eq. (1) of Reference 2 is in the wrong direction compared to the data. Figures 5b-d all show the same features as stated above and also indicate the accuracy of the present correlation.

Figures 6a-c show the effect of the length of the projectile (and hence the length of the grooved section) on the incremental drag, for each of the Mach numbers of 3.5, 4.0 and 5.0. The expressions of References 2 and 3 overpredict the results by more than 100% as seen in Figure 6a. The same results are repeated in Figures 6b and 6c. In all cases, the present correlation predicts the results very accurately. It is of interest to notice that in Figure 6c, Eq. (1) of Reference 2 overpredicts ΔC_{Dg} by more than ten times for L/D = 35.

Figures 7a-d focus on the variation of l_g/L , at Mach 5.0 where most of the tests of Reference 1 were made for the same four L/D's of 20, 25, 30 and 35. The same good agreement between present expression predictions and the experimental data persist. Also, one notices the considerable and consistent overprediction of the expressions of References 2 and 3 being evident.

Figures 8a-d re-cap the results as variation of the incremental drag with the groove type. The same agreement with the data persists for the present expressions while the expressions of Reference 2 and 3 consistently overpredict the results as shown in Figure 8a. Figures 8b-d gives the same trend described above.

A summary of the present predictions of Eq. (3), and the wind tunnel data of References 1 and 4 is given in Table 3.

The second expression, Eq. (5), was applied to the test case of Reference 4 where the first expression, Eq. (3), is not applicable. The results as shown in Figure 9 are in good agreement considering the complexity of the shape. It is to be noted that the present results underpredict the data. This underprediction is expected since the grooves are known to increase the boattail drag.⁷ This trend has also been shown in Reference 7 where serrated body ends were tested in wind tunnels at subsonic speeds, causing an increase in the base drag. The effect of grooves on boattail drag is not included in this study; however, this effect explains why the present predictions for this projectile configuration must be smaller than those of the wind tunnel.

The second expression, Eq. (5), was then applied to some of the cases of Reference 1 to see how closely the two correlation expressions compare. Nine of the seventeen cases were recomputed and the results were very similar. Those results and a comparison with the earlier results of the first expression, Eq. (3), are tabulated in Table 4.

A summary plot for the comparison of the present two prediction correlations, Eqs. (3) and (5), with the wind tunnel data is given by Figure 10. The predictions are shown to be in good agreement with the data.

It is felt that the most influential factor is the Reynolds number, followed by the Mach number. The incremental groove drag decrease significantly with increase in Reynolds number. It also increases rapidly with the decrease of Mach number towards the sonic value. The present correlations are only to be used for $M > 3.5$ since Mach number function was targeted for that region only. It is found that for $M < 3.5$ the Mach number function may greatly overpredict the expected answer.

V. SUMMARY AND CONCLUSIONS

A fast prediction capability for computing the incremental drag of Kinetic Energy projectiles due to surface grooves at supersonic speeds has been established.

Physical variables affecting the groove drag have been identified and studied with respect to their importance.

Two correlations with wind tunnel data are established for typical and general KE projectile configurations. The second correlation is recommended since it can be applied to many KE projectile configurations. The correlation expressions are valid only for the Mach range $3.5 < M < 5.5$ and for $\alpha = 0^\circ$. Any application for Mach numbers less than 3.5 will yield overpredicted results.

Comparison with other existing prediction expressions heavily favored the present correlations. The accuracy of these present correlations is $\pm 20\%$ while the existing prediction methods of References 2 and 3 can be as high as $+1000\%$ and $+300\%$, respectively.

Finally, the Reynolds number is expected to play a significant role. The incremental drag decreases significantly with the increase of Reynolds number. It is hoped that more tests will be available at the Reynolds number of 35×10^6 per foot rather than the present values of $4-6 \times 10^6$ per foot. These tests could be run in cryogenic wind tunnels and the results further used to validate the present correlations.

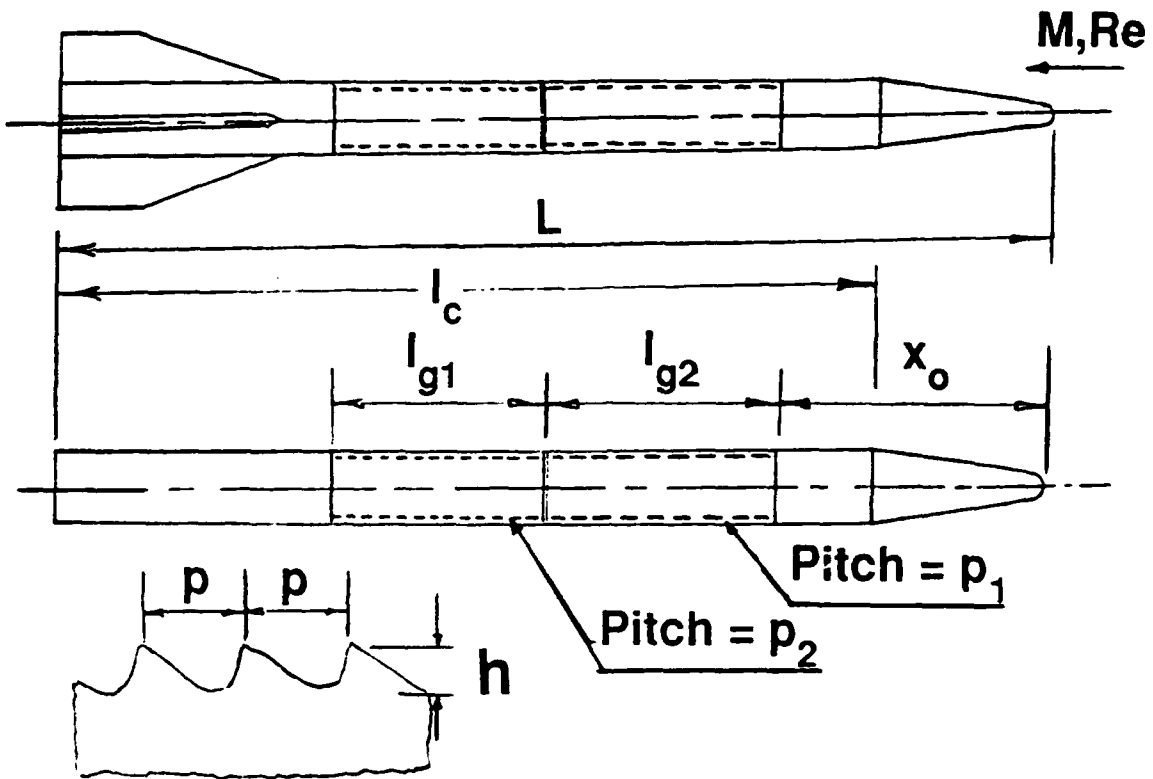


Figure 1. General configuration and nomenclature
for a Kinetic Energy projectile.

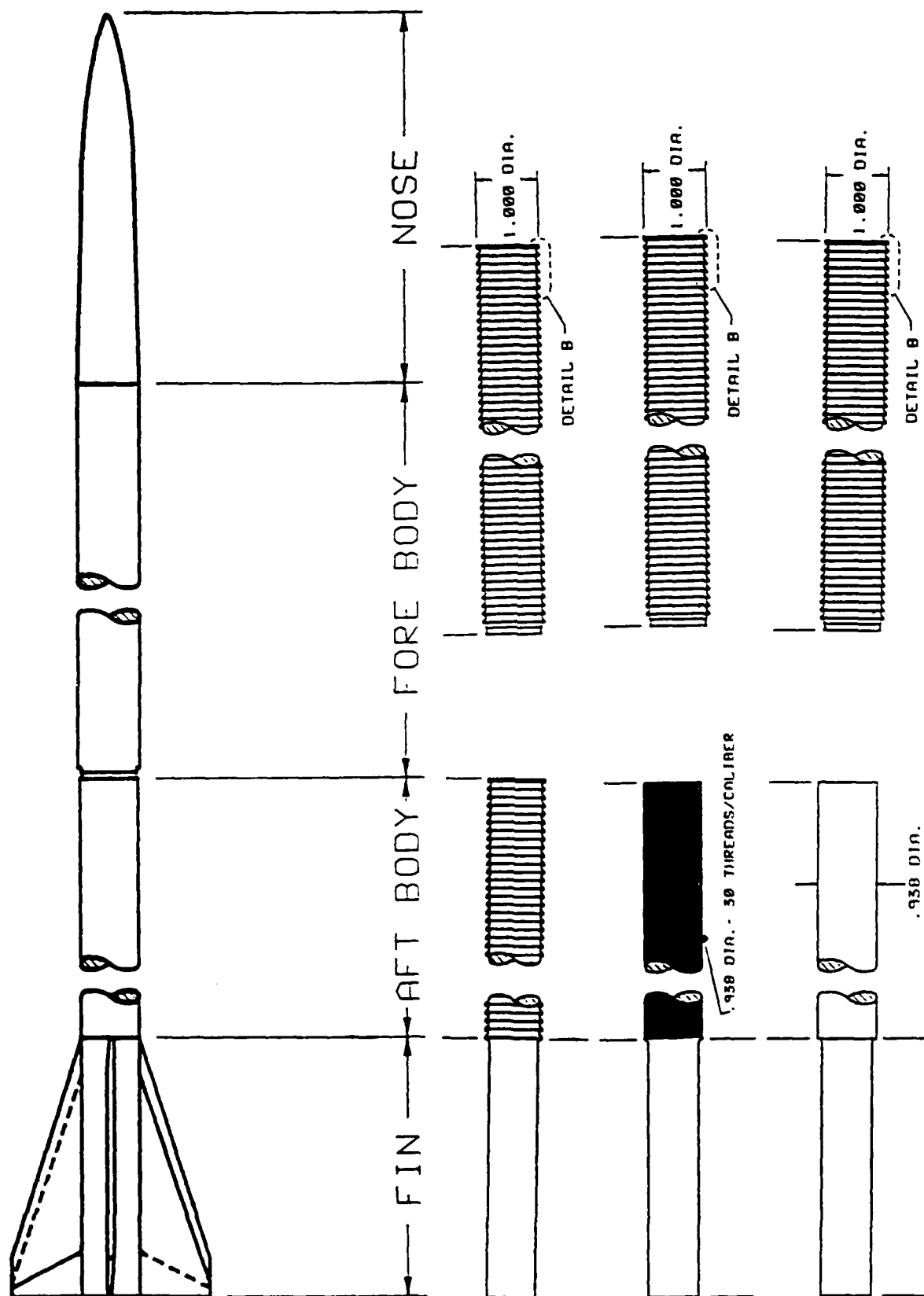
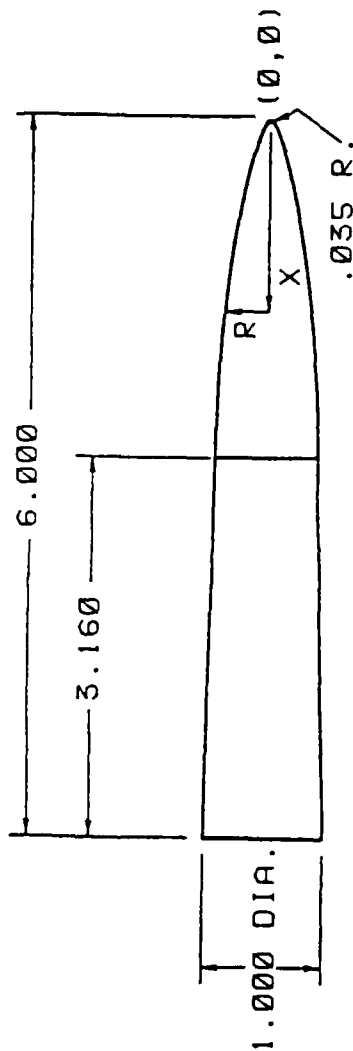
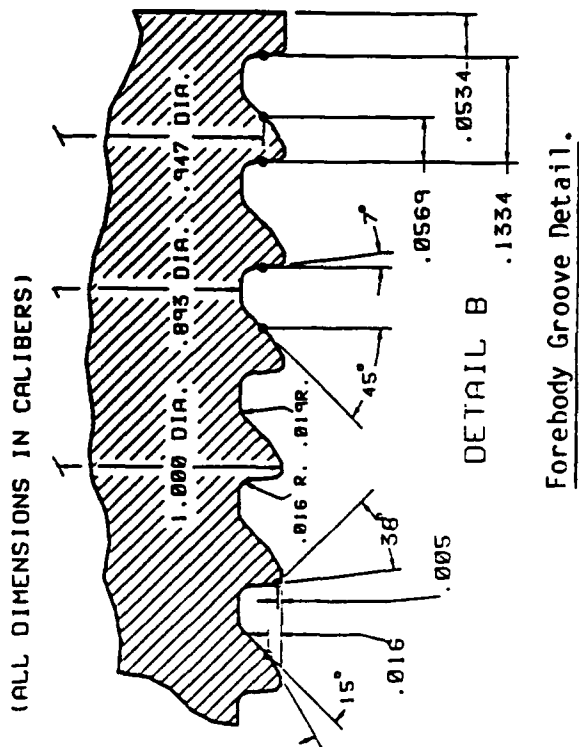
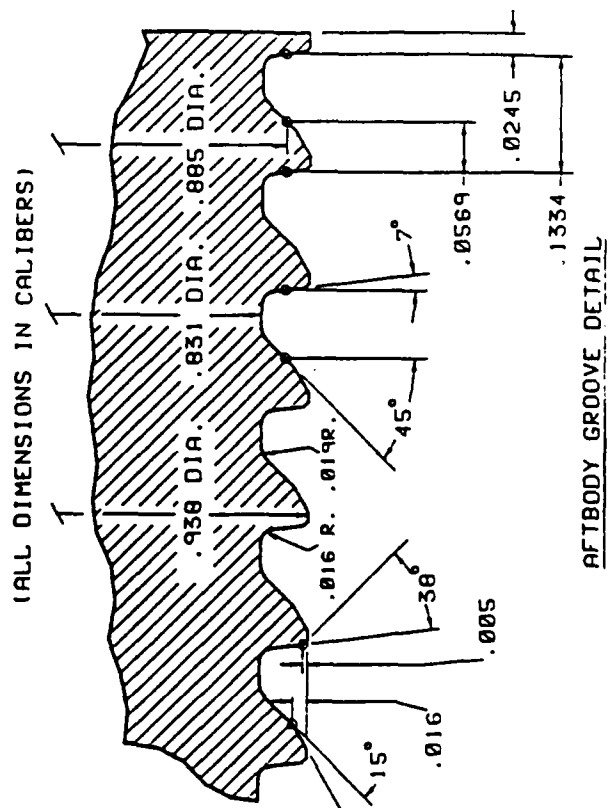


Figure 2. Wind tunnel test model for tests of Reference 1.

a. General configuration.



$$R = .432 * (1. - (1. - 2. * X / 5.995)^2)^{.75}$$

$$0.0 \leq X \leq 2.839$$

(ALL DIMENSIONS IN CALIBERS)

Sears-Haack nose (Sm SH).

Figure 2. Wind tunnel test model for tests of Reference 1.

b. Nose configuration and grooves geometry.

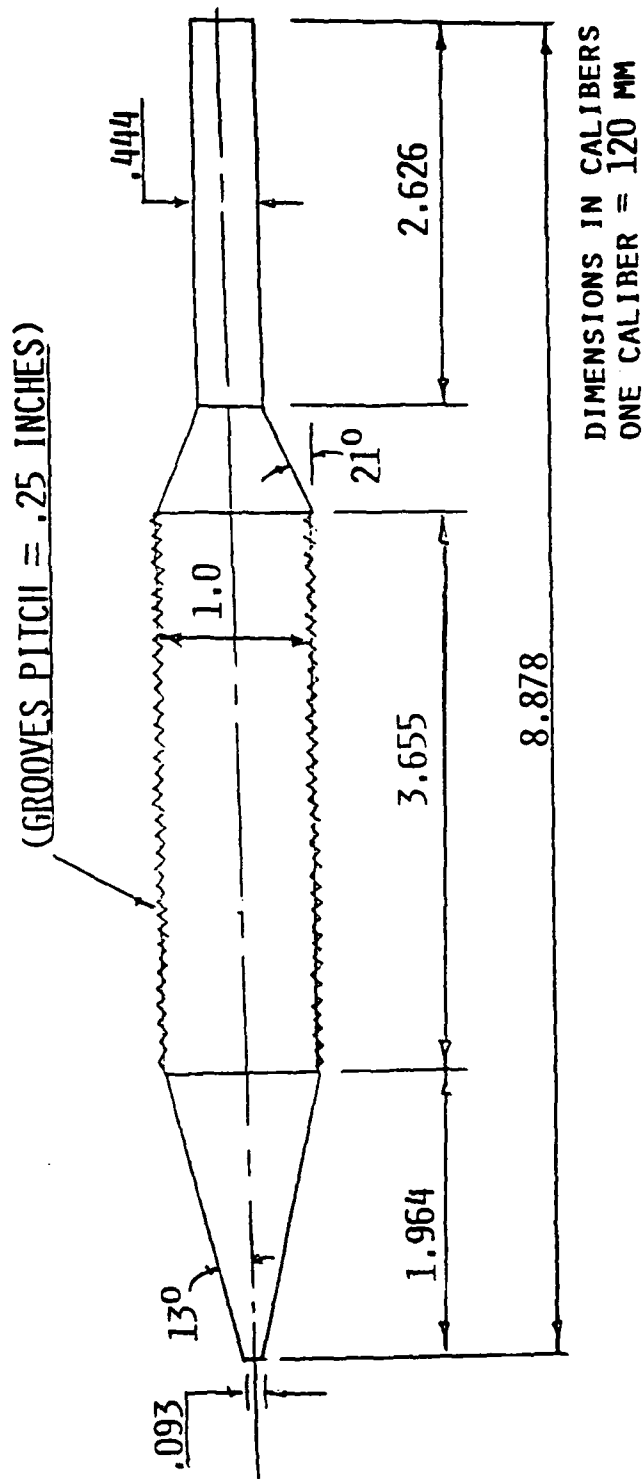


Figure 3. Wind tunnel test model for the tests of Reference 4.

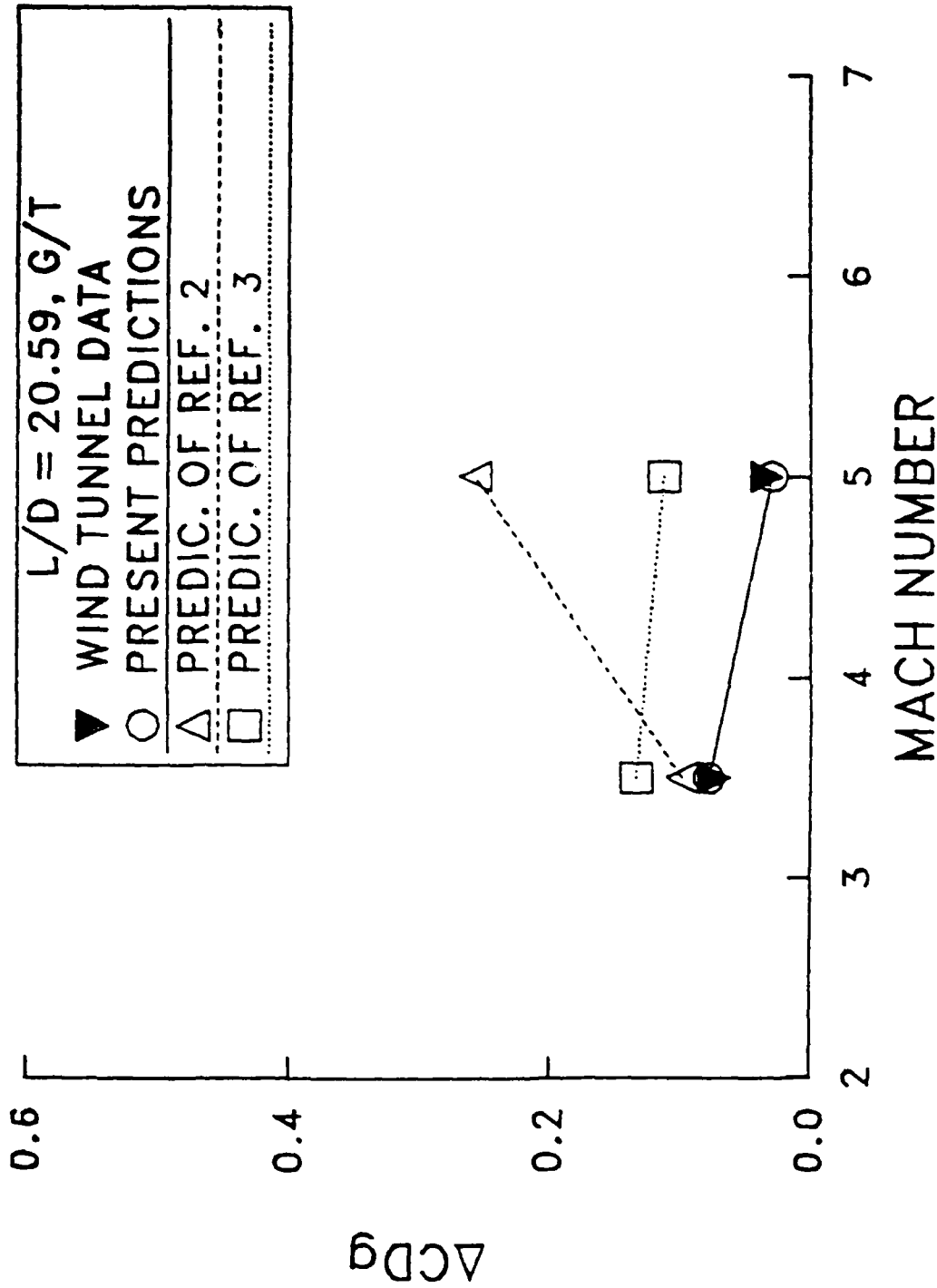


Figure 5. Results of present correlation: validation of ΔC_D with Mach number.

a. For $L/D = 20.59$.

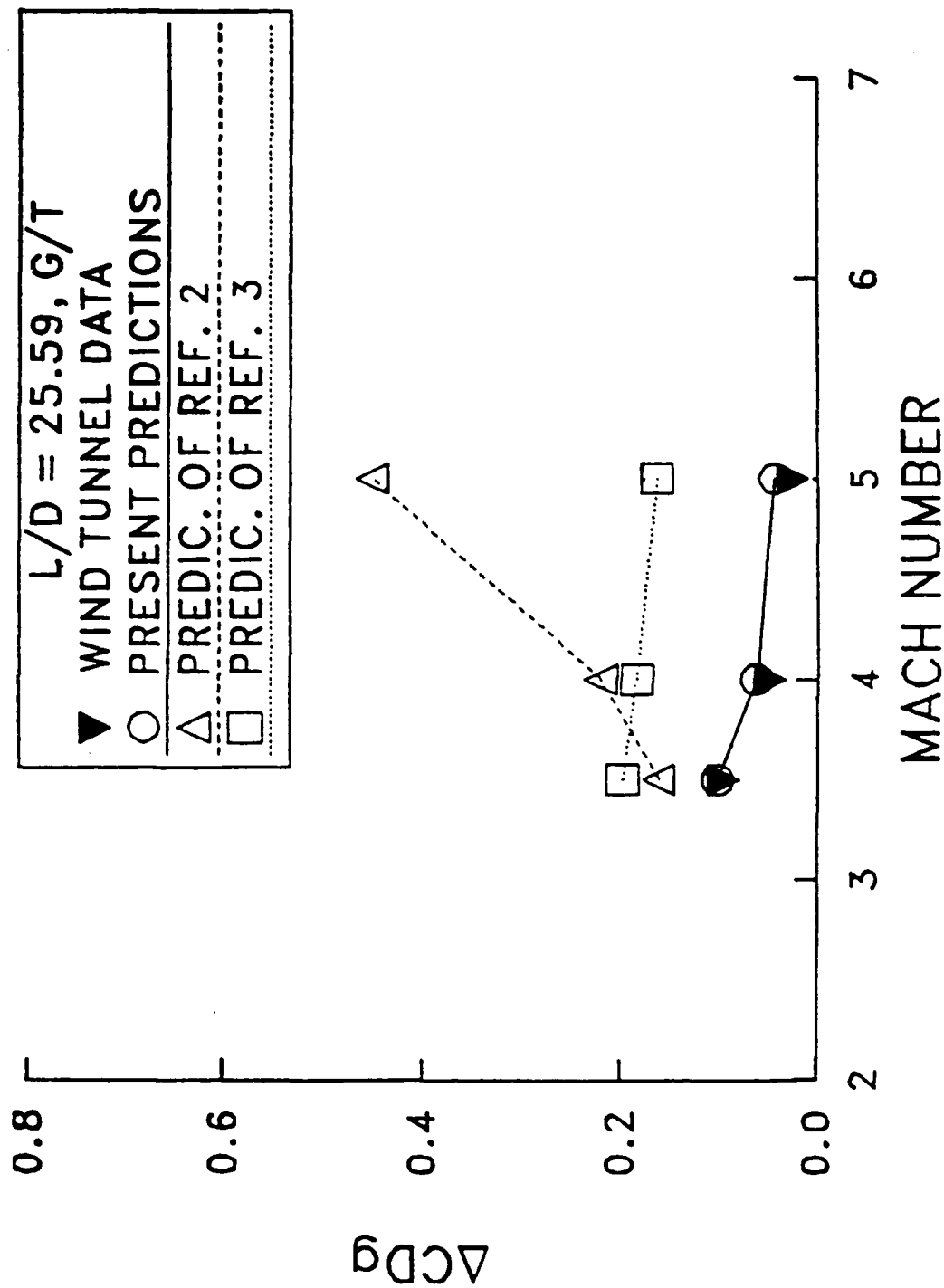


Figure 5. Results of present correlation: validation of ΔC_D with Mach number.

b. For $L/D = 25.59$.

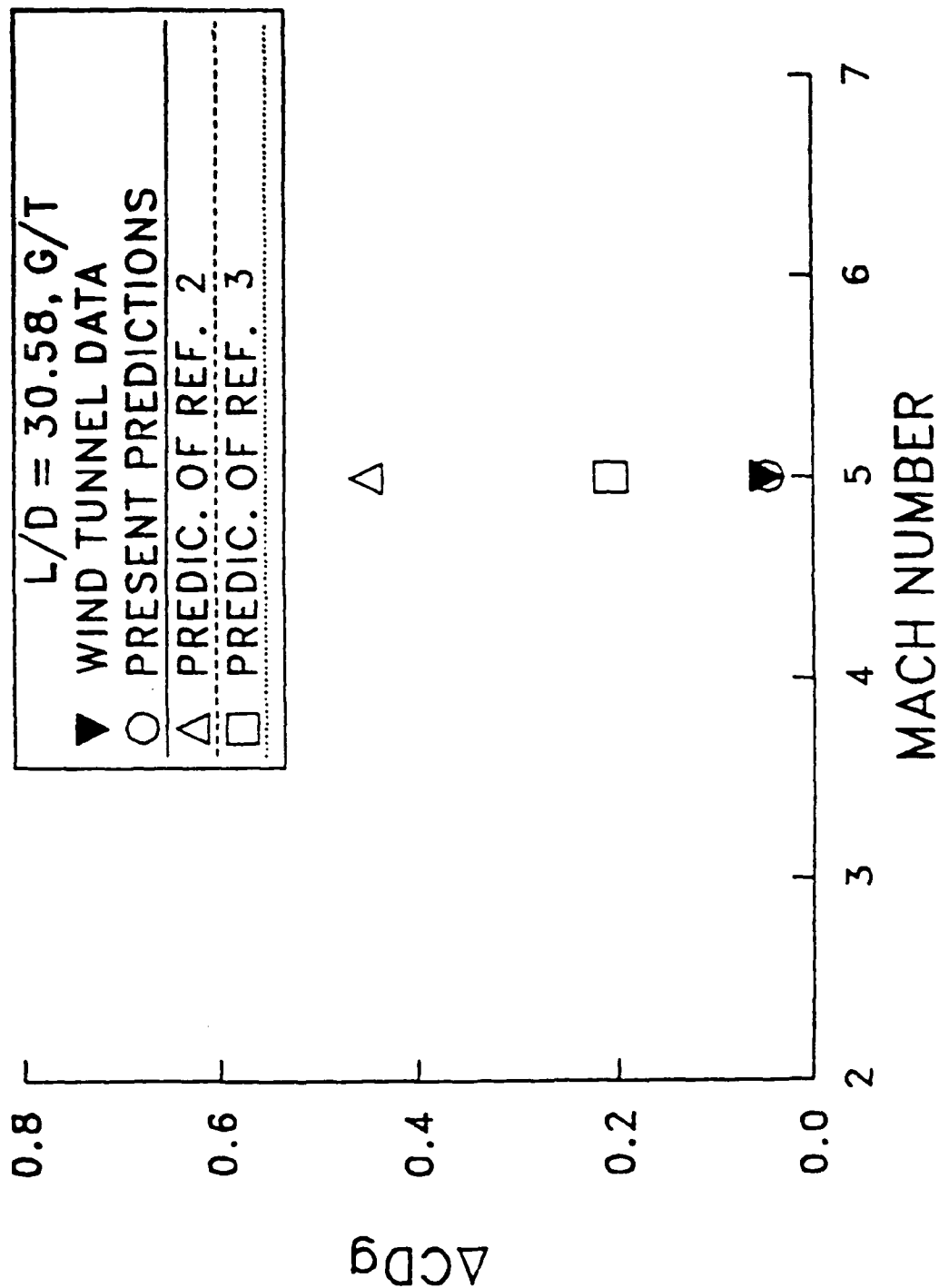


Figure 5. Results of present correlation: validation of ΔC_{Dg} with Mach number.

c. For $L/D = 30.58$.

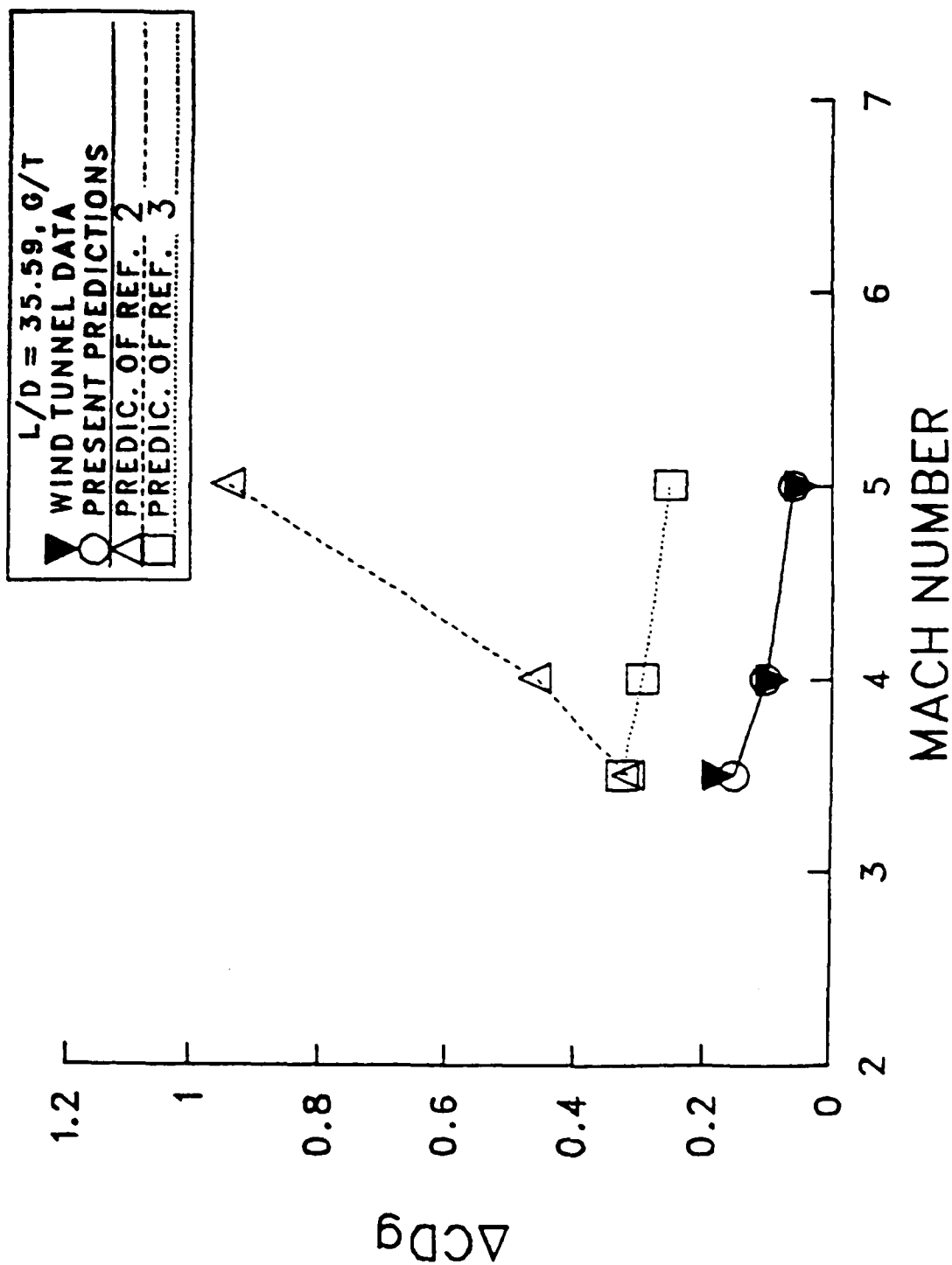


Figure 5. Results of present correlation: validation of ΔC_p with Mach number.

d. For $L/D = 35.59$.

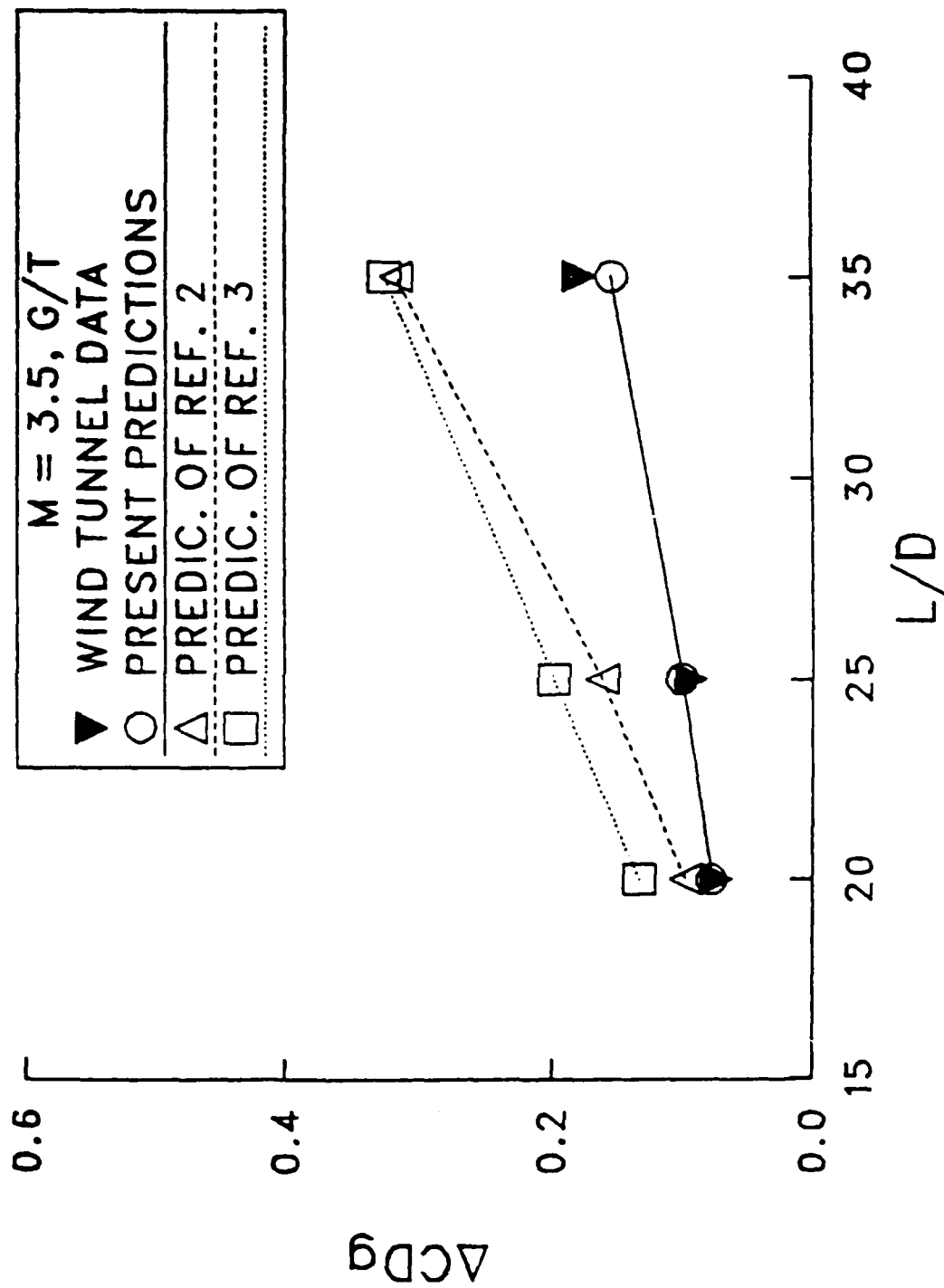


Figure 6. Result of present correlation: variation of ΔC_D with groove length, l_g/D

a. $M = 3.5$.

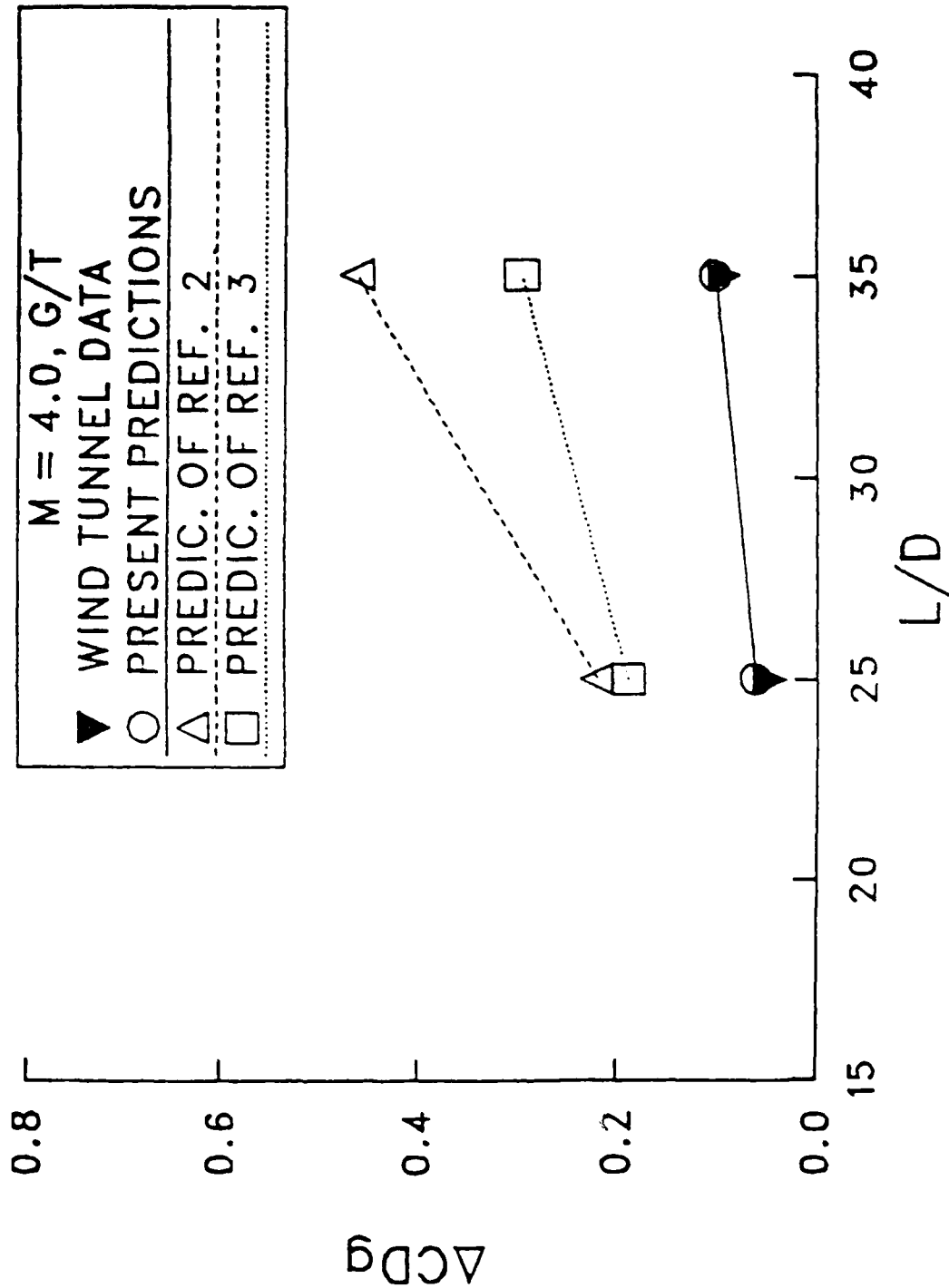


Figure 6. Result of present correlation: variation of ΔC_{Dg} with groove length, l_g/D

b. $M = 4.0$.

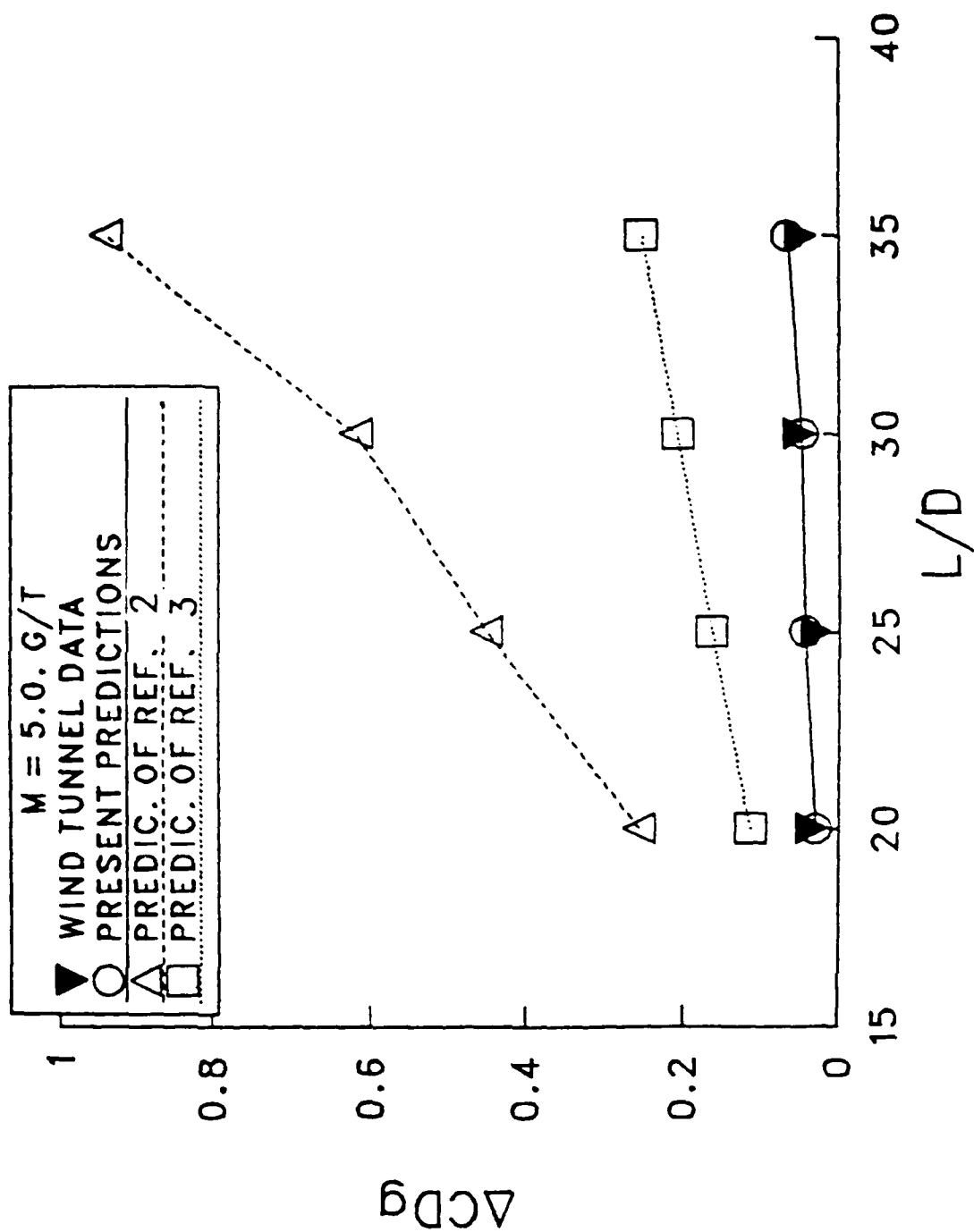


Figure 6. Result of present correlation: variation of ΔC_D with groove length, l_g/D

c. $M = 5.0$.

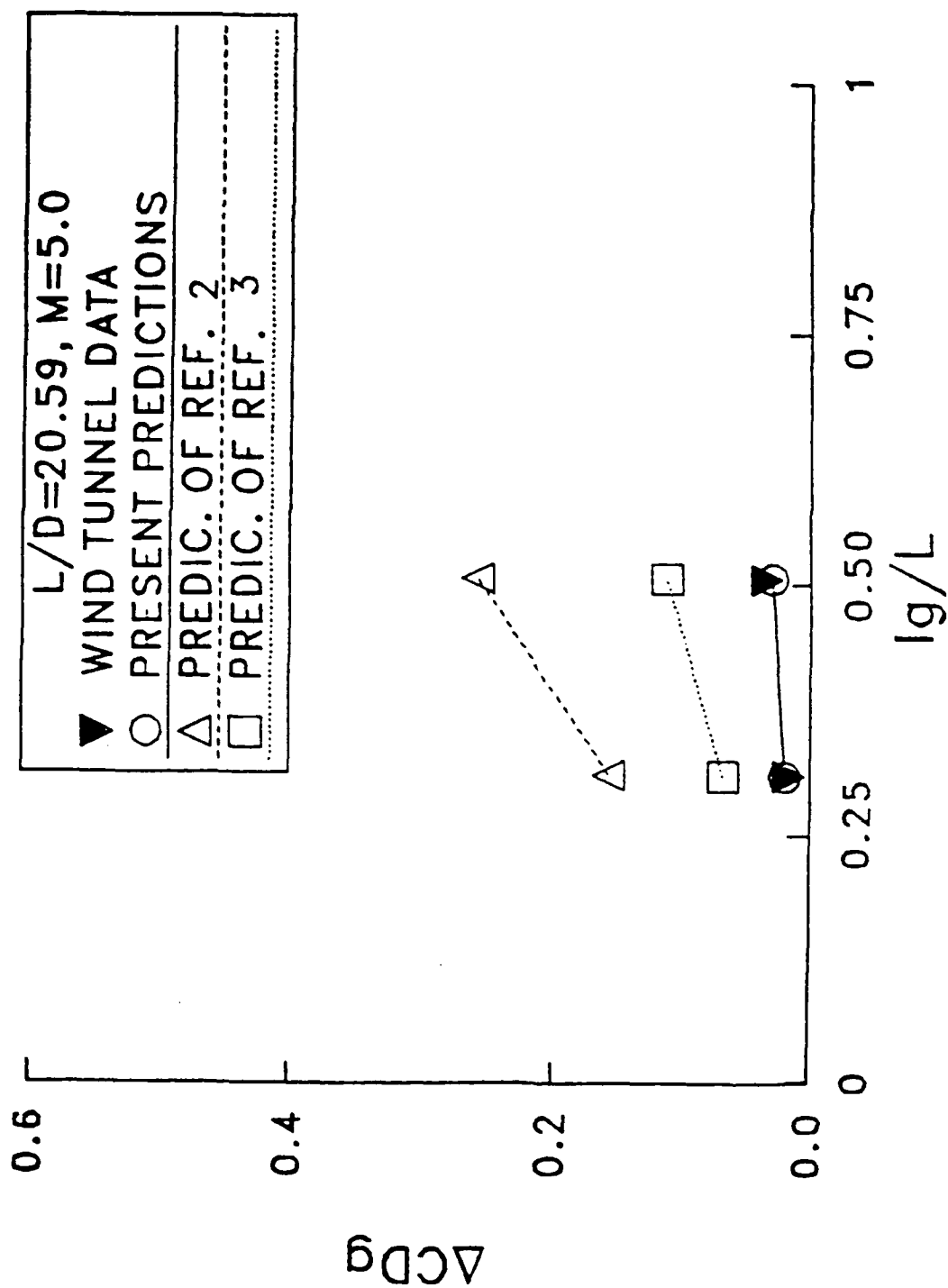


Figure 7. Result of present correlation: variation of ΔC_{Dg} with lg/L .

a. $L/D = 20.59, M = 5.0$.

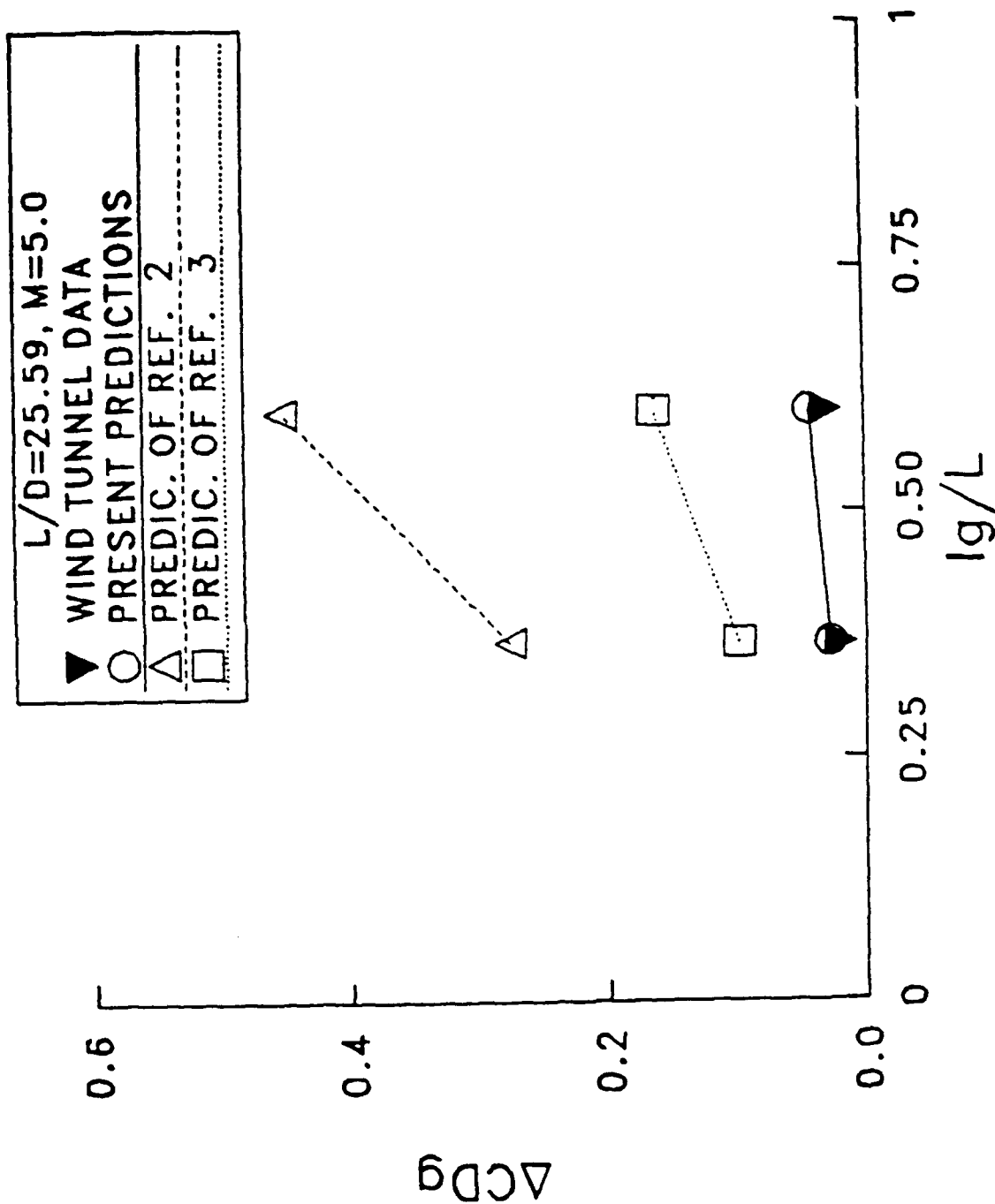


Figure 7. Result of present correlation: variation of ΔC_{Dg} with lg/L .

b. $L/D = 25.59, M = 5.0$.

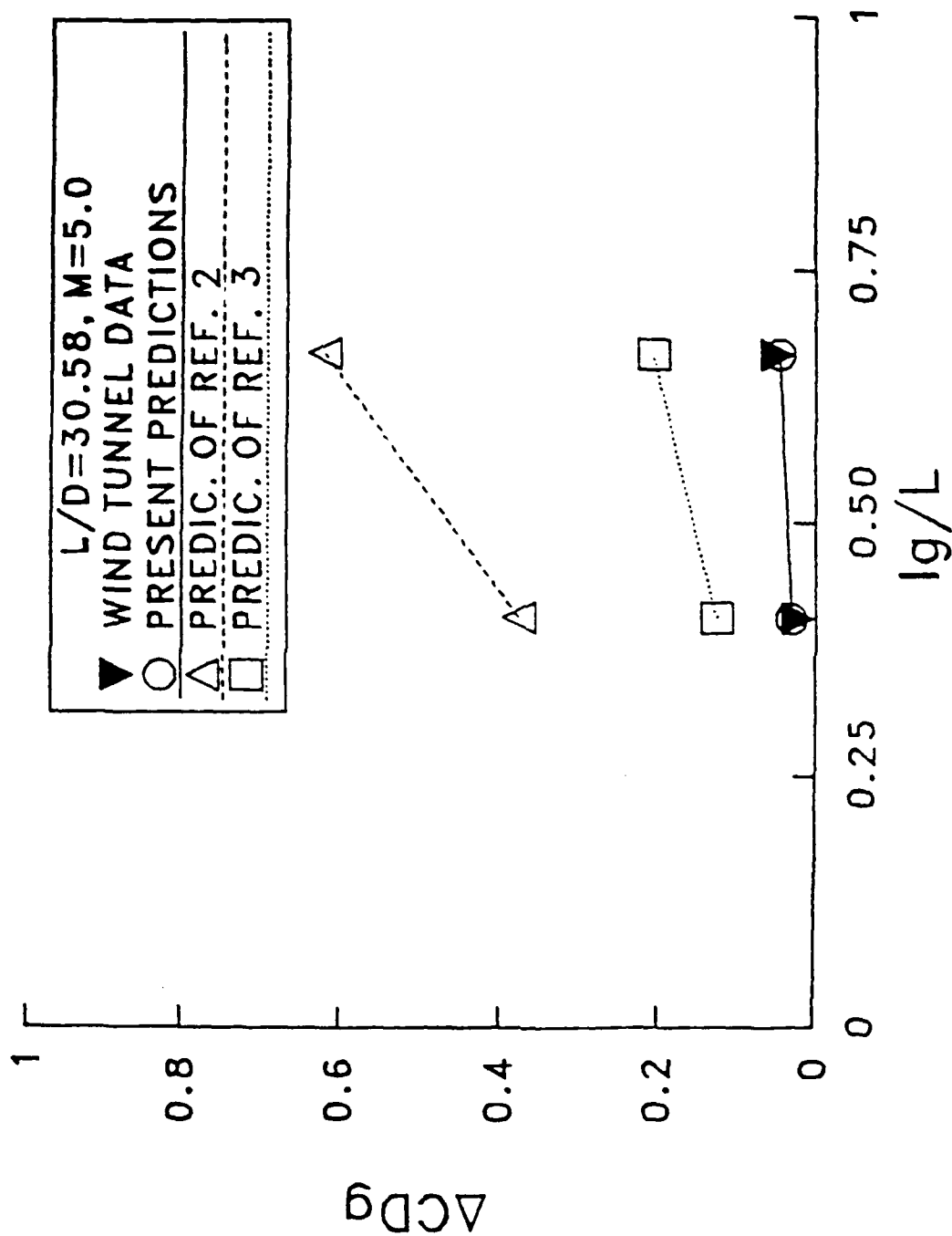


Figure 7. Result of present correlation: variation of $\Delta C D g$ with lg/L .

c. $L/D = 30.58, M = 5.0$.

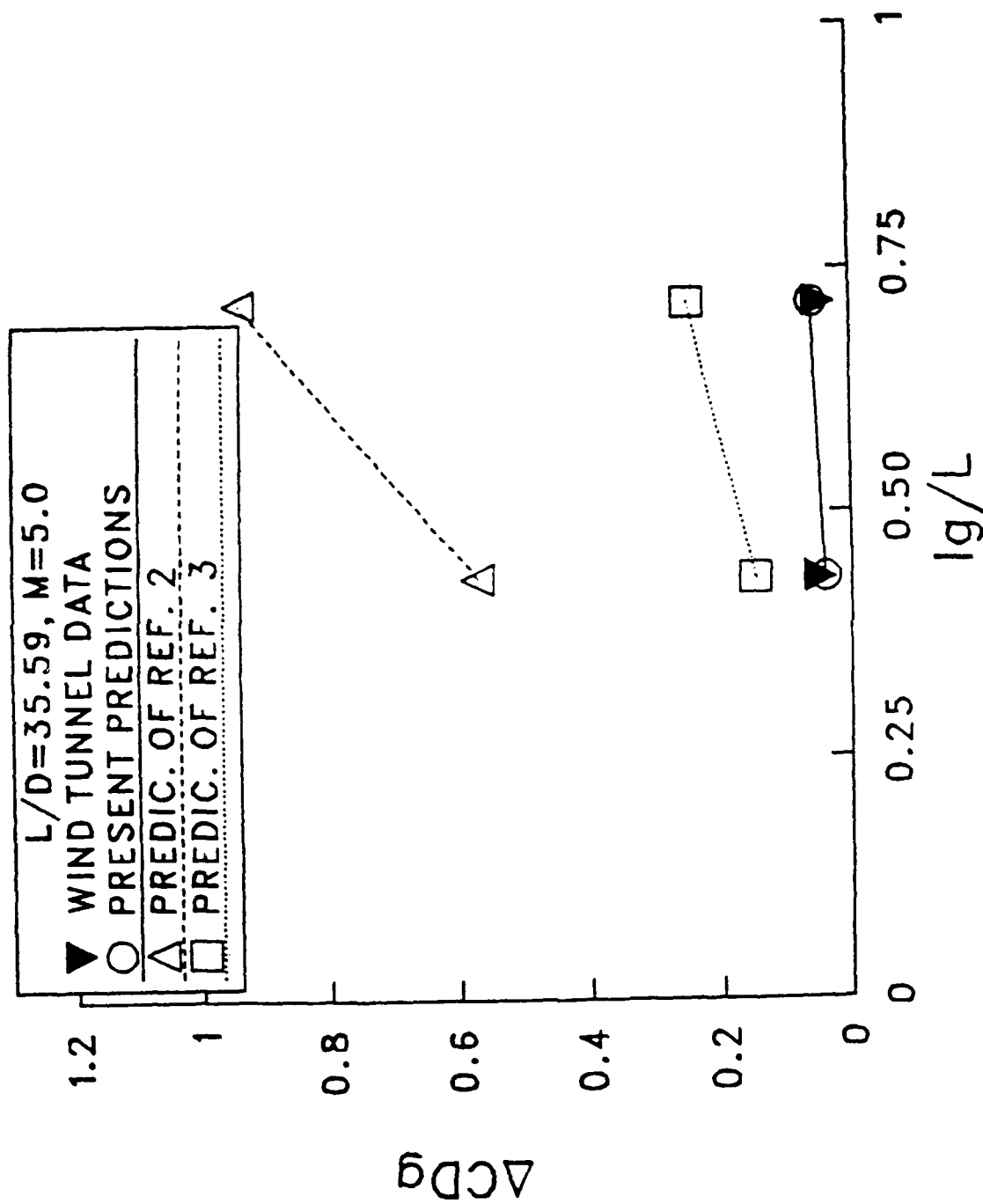


Figure 7. Result of present correlation: variation of ΔC_D with lg/L .

d. $L/D = 35.59, M = 5.0$.

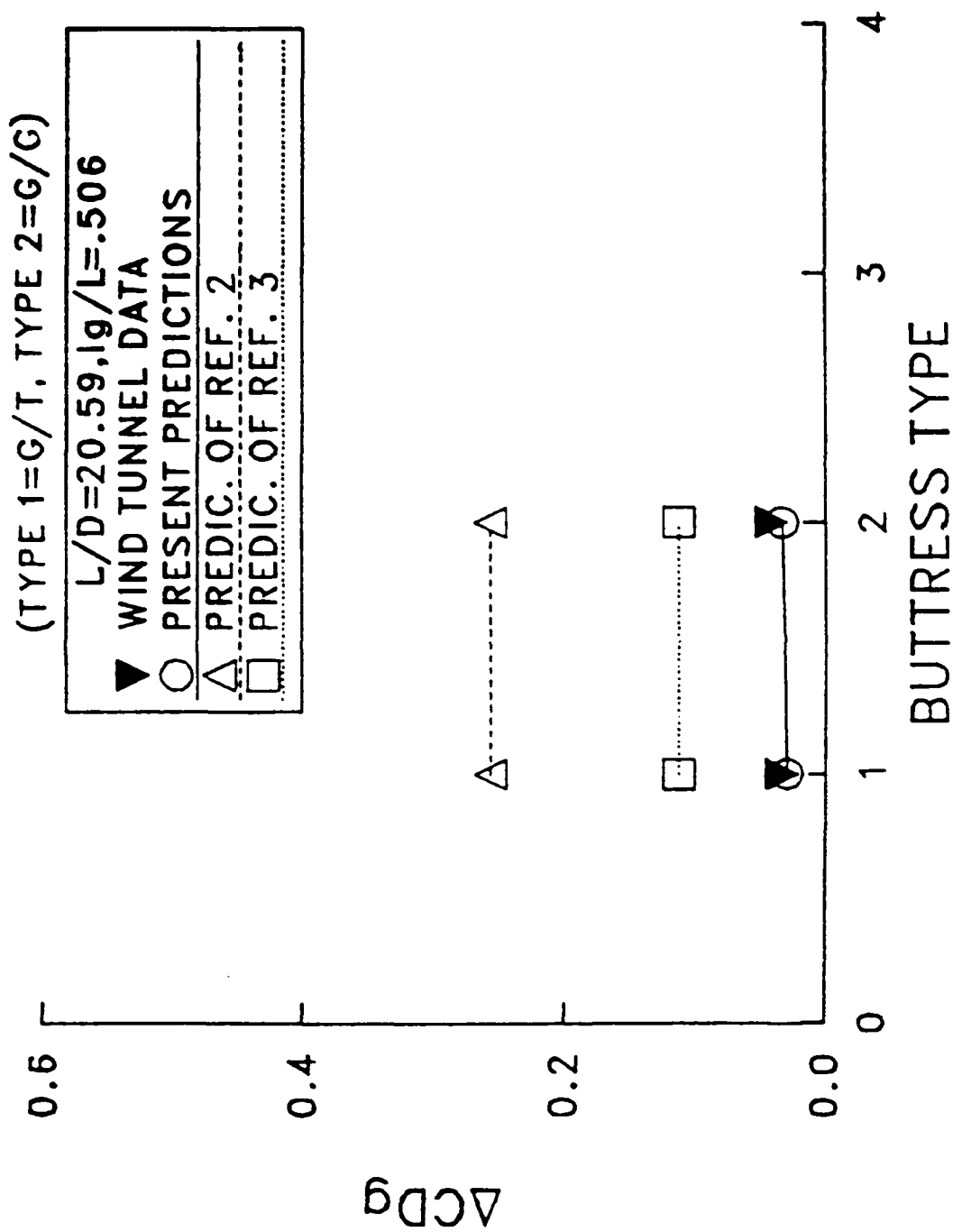


Figure 8. Result of present correlation: variation of ΔC_D with Buttruss type.

a. $L/D = 20.59, M = 5.0.$

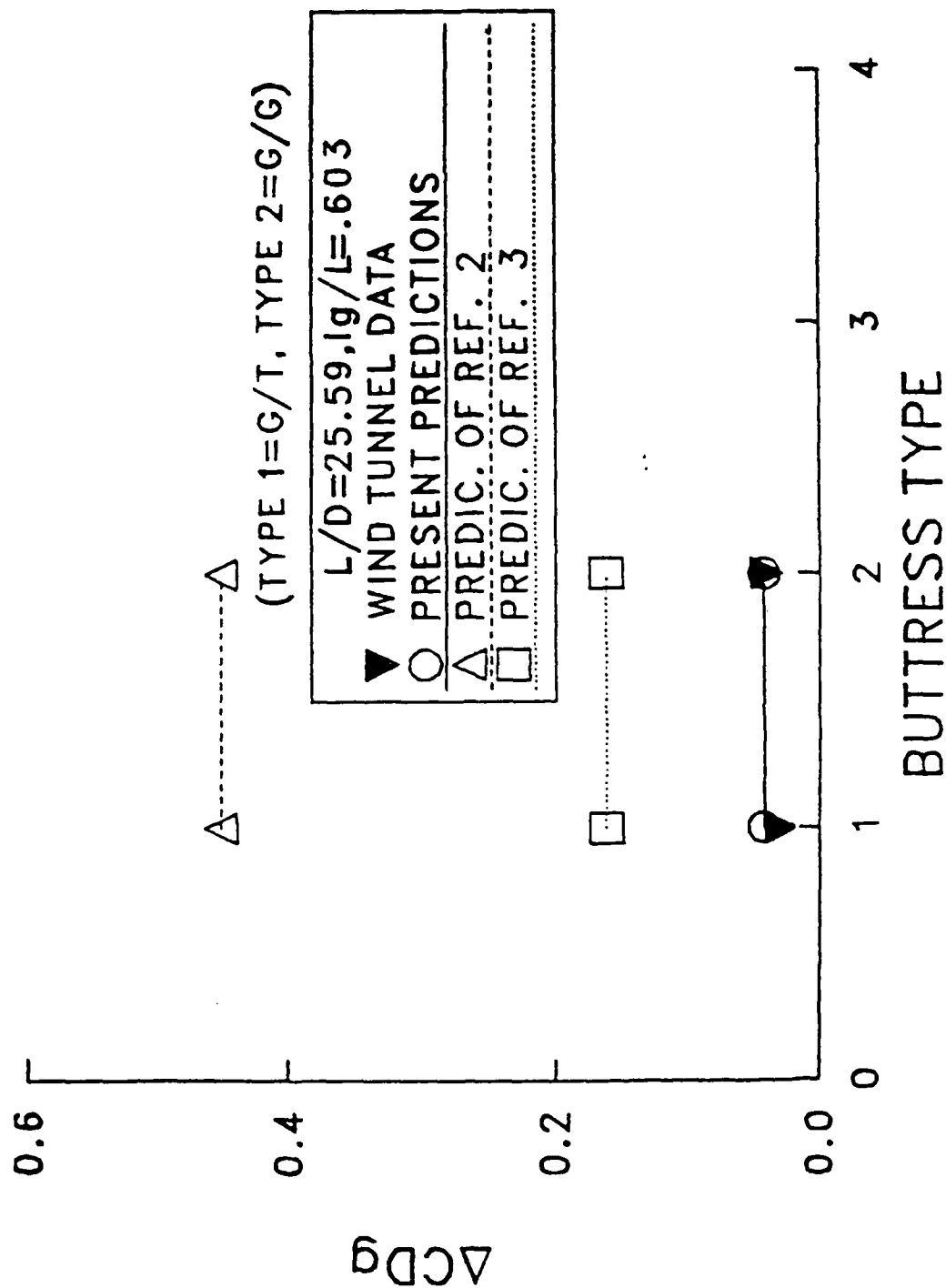


Figure 8. Result of present correlation: variation of ΔC_D with Buttruss type.

b. $L/D = 25.59, M = 5.0.$

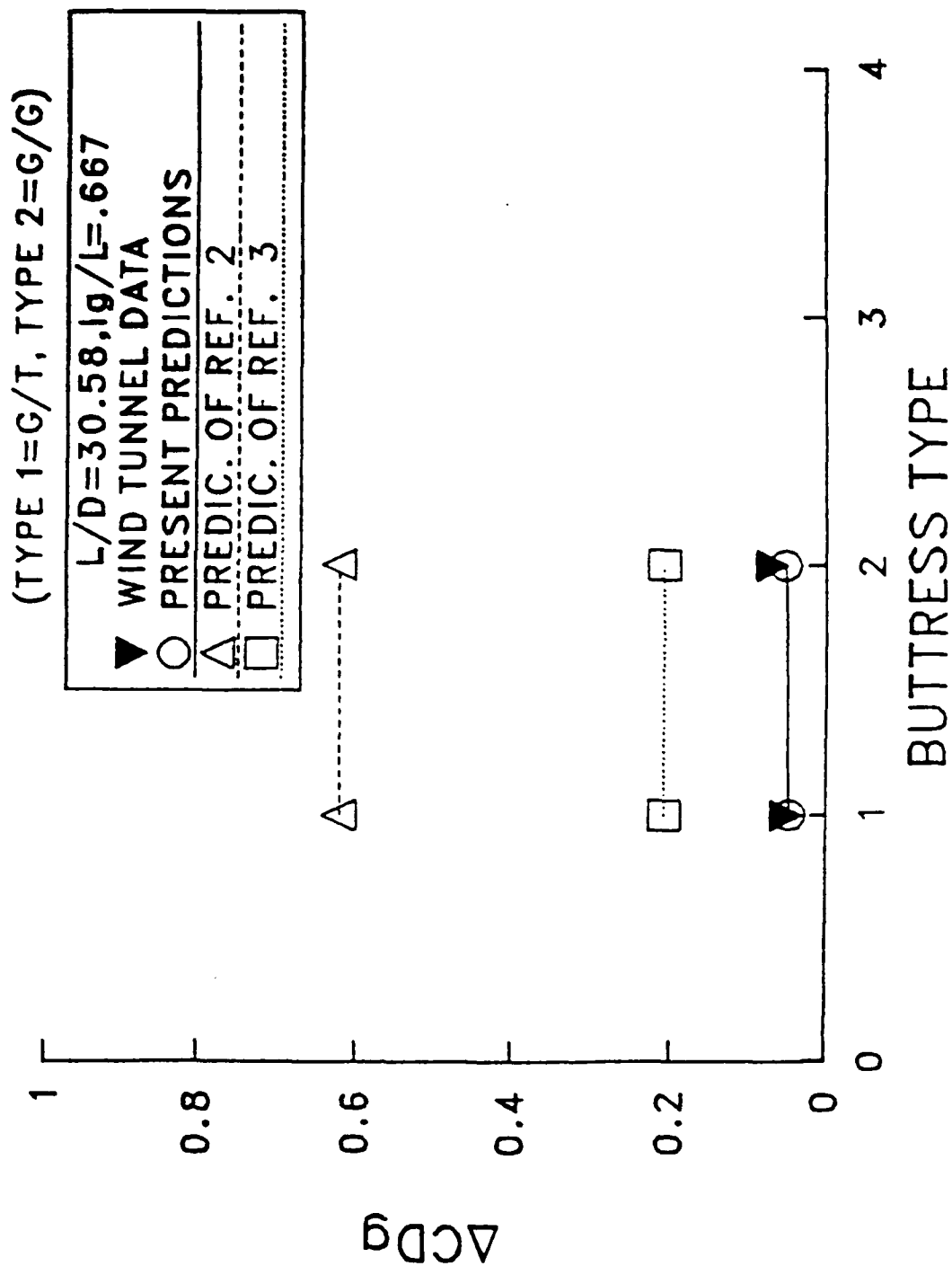


Figure 8. Result of present correlation: variation of ΔC_D with Buttress type.

c. $L/D = 30.58, M = 5.0.$

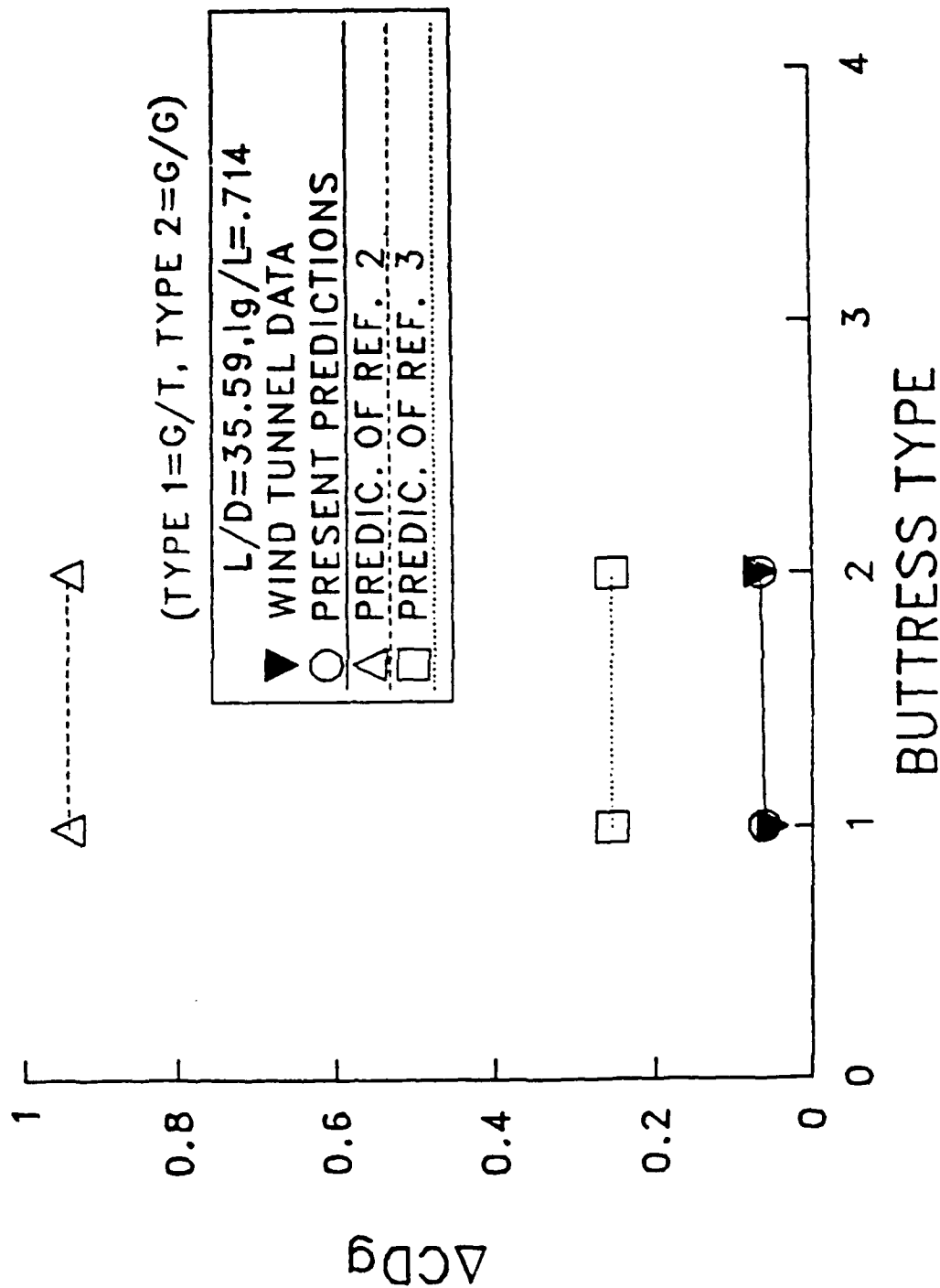


Figure 8. Result of present correlation: variation of ΔC_D with Buttruss type.

d. $\frac{L/D = 35.59, M = 5.0.}{}$

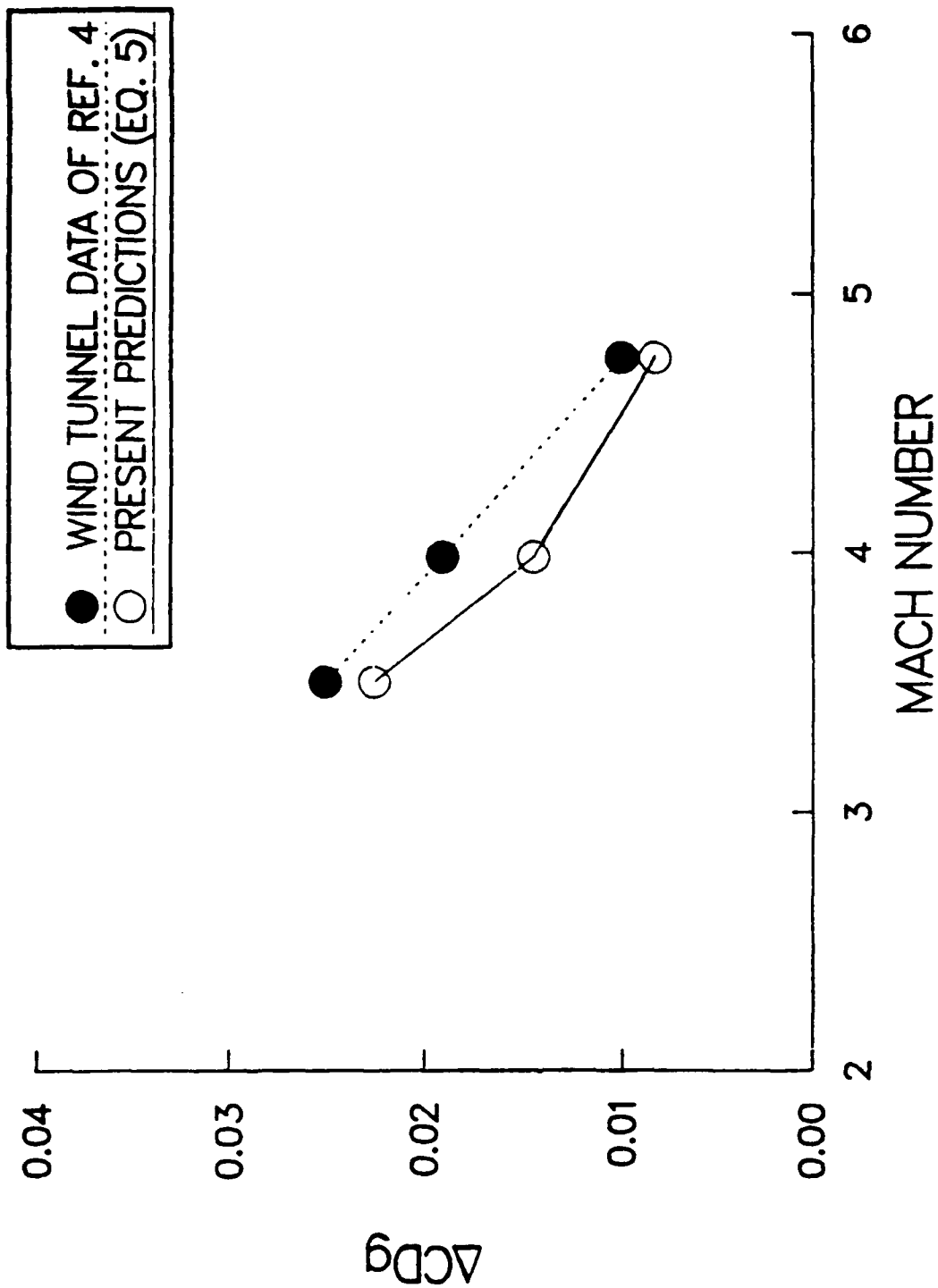


Figure 9. Validation of present predictions against data of Reference 4.

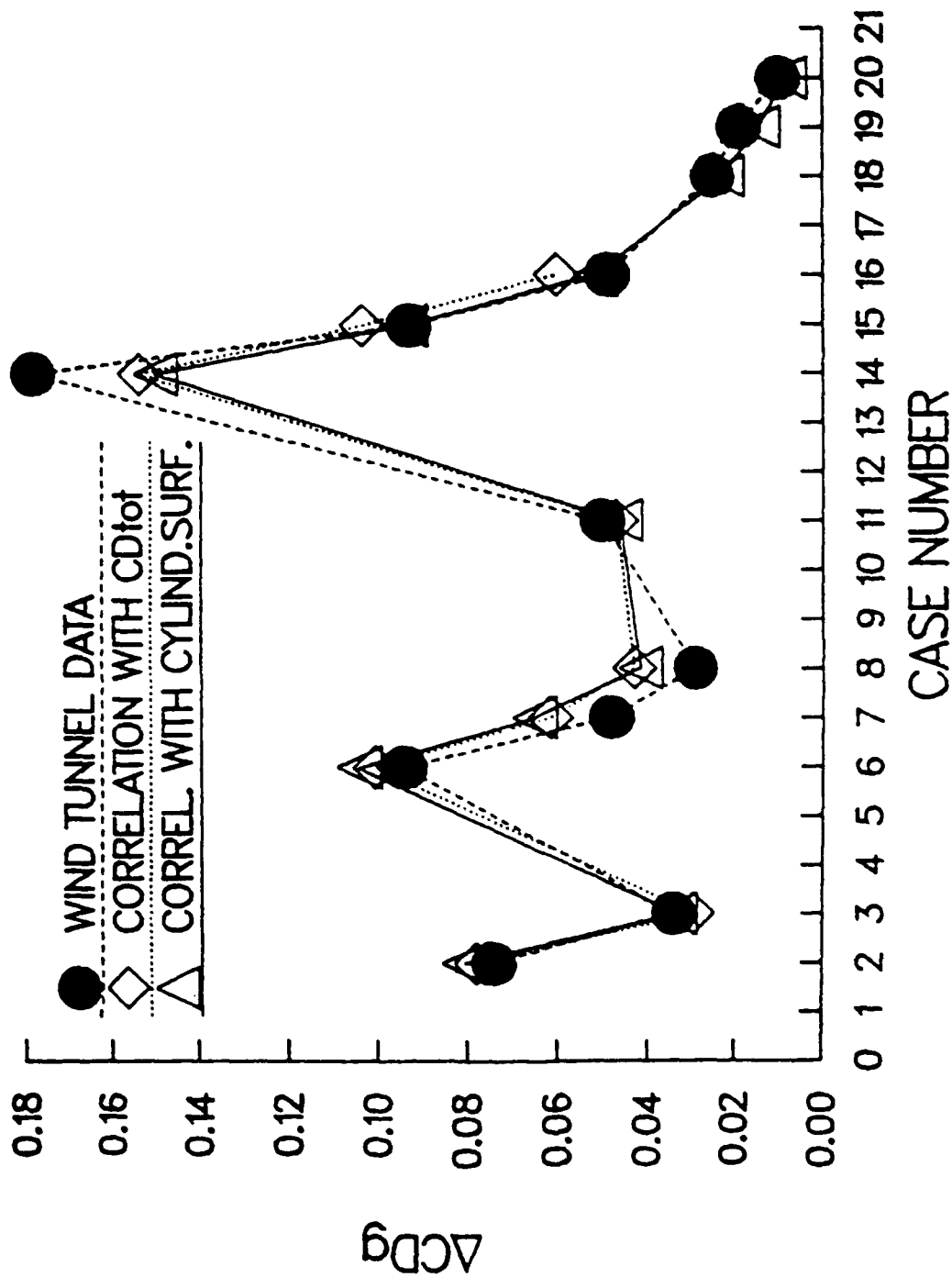


Figure 10. Comparison of the results of the present two correlations.

TABLE 1. Case Designation and Test Conditions of Reference 1.

Case No.	Total length, L (caliber)	Length and Type* of Threads		Cylindrical length, l_c (caliber)	Mach Number	R_e per ft $\times 10^6$	$R_{eL} \times 10^6$ (based on L)
		l_{g_1} (caliber)	l_{g_2} (caliber)				
1	20.59	6.366(G)	-	10.417	5.0	5.79	9.32
2		6.366(G)	4.051(T)		3.5	4.15	6.68
3		6.366(G)	4.051(T)		5.0	5.65	9.09
4		6.366(G)	4.051(G)		5.0	5.56	8.94
5	25.59	9.392(G)	-	15.420	5.0	5.68	11.36
6		9.392(G)	6.028(T)		3.5	4.19	8.38
7		9.392(G)	6.028(T)		4.0	5.23	10.45
8		9.392(G)	6.028(T)		5.0	5.41	10.81
9		9.392(G)	6.028(G)		5.0	5.88	11.76
10	30.58	12.408(G)	-	20.412	5.0	5.71	13.65
11		12.408(G)	8.004(T)		5.0	5.71	13.64
12		12.408(G)	8.004(G)		5.0	5.76	13.77
13	35.59	15.424(G)	-	25.420	5.0	5.65	15.72
14		15.424(G)	9.996(T)		3.5	3.79	10.55
15		15.424(G)	9.996(T)		4.0	4.49	12.48
16		15.424(G)	9.996(T)		5.0	5.70	15.83
17		15.424(G)	9.996(G)		5.0	5.81	16.15
* TYPE: (G) for grooves (8 grooves per inch) (T) for threads (32 threads per inch)							

TABLE 2. Case Designation and Test Conditions of Reference 4.

Case No.	Mach Number	R_e per ft $\times 10^{-6}$	C_D		ΔC_{D_g} Wind Tunnel
			Grooved Body	Smooth Body	
18	3.49	12	0.280	0.255	0.025
19	3.98	14	0.227	0.208	0.019
20	4.75	18	0.196	0.186	0.010

TABLE 3. Comparison of Present Predictions with Data of References 1 and 4.

Case No.	C_D Wind Tunnel		ΔC_{Dg}	
	Grooved Body	Smooth Body	Wind Tunnel	Present Predictions Eq. (3)
1	0.213	0.197	0.016	0.0193
2	0.384	0.310	0.074	0.0779
3	0.231	0.197	0.034	0.0294
4	0.238	0.197	0.041	0.0331
5	0.256	0.236	0.020	0.0275
6	0.434	0.340	0.094	0.1005
7	0.328	0.280	0.048	0.0615
8	0.265	0.236	0.029	0.0429
9	0.277	0.236	0.041	0.0434
10	0.282	0.258	0.024	0.0312
11	0.308	0.258	0.050	0.0469
12	0.327	0.258	0.069	0.0509
13	0.298	0.248	0.050	0.0407
14	0.578	0.400	0.178	0.1532
15	0.450	0.357	0.093	0.1035
16	0.347	0.297	0.049	0.0603
17	0.354	0.298	0.056	0.0656
18	0.280	0.255	0.025	Not Applicable
19	0.227	0.208	0.019	
20	0.196	0.186	0.010	

TABLE 4. Comparison of the Present Two Correlations.

Case No.	ΔC_{Dg}		
	Wind Tunnel	Present Predictions Eq. (3)	Present Predictions Eq. (5)
2	0.074	0.0779	0.0805
3	0.034	0.0294	0.0334
6	0.094	0.1005	0.1040
7	0.048	0.0615	0.0650
8	0.029	0.0429	0.0415
11	0.050	0.0469	0.0461
14	0.178	0.1532	0.1494
15	0.093	0.1035	0.0937
16	0.049	0.0603	0.0514
18	0.025	Not Applicable	0.0226
19	0.019		0.0144
20	0.010		0.0083

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LIST OF SYMBOLS

A_{ref}	= projectile reference area, $(\pi D^2/4)$
C_A	= axial force coefficient
C_D	= drag coefficient, drag force/ $(0.5\rho_\infty V^2 A_{ref})$
$C_{D_{TSB}}$	= body-alone <u>total</u> drag coefficient (including base drag) for the <u>smooth body</u> (i.e., without grooves) configuration
$C_{D_{SFSB}}$	= drag coefficient due to <u>skin friction</u> of the <u>smooth body</u> of the cylindrical portion, l_c , of the body
c_f	= skin friction coefficient
d_c	= diameter of the cylindrical portion, in calibers
D	= reference diameter of the projectile
h	= groove depth
KE	= kinetic energy
l_c	= length of the cylindrical portion of the body, in calibers
l_g, l_{g1}, l_{g2}	= axial length of the grooved portion of the projectile body
L	= reference length of the projectile, usually the total length except as otherwise noted
M	= projectile Mach number
p	= groove pitch
Re	= Reynolds number per unit length
Re_L	= Reynolds number based on the total length of the projectile
V	= projectile velocity
x_0	= distance from the nose tip to the first groove

Greek Symbols

α	= angle of attack
ρ_∞	= free stream air density
τ	= shear stress at the wall

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