

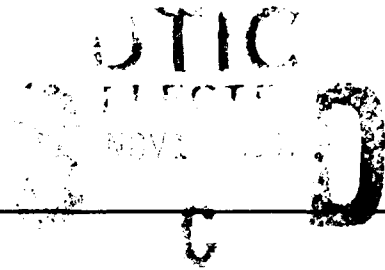
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TECHNICAL REPORT BRL-TR-3275

**BRL**



EXPERIMENTAL MEASUREMENT OF  
TAILBOOM STRAIN DURING  
CHARGE IGNITION

JOSEPH W. COLBURN  
ARTHUR A. KOSZORU

OCTOBER 1991

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# TABLE OF CONTENTS

	<u>Page</u>
LIST OF FIGURES.....	v
ACKNOWLEDGMENTS.....	vii
1. INTRODUCTION.....	1
2. SIMULATOR BACKGROUND.....	1
3. EXPERIMENTAL.....	2
3.1 Simulator Instrumentation.....	2
3.2 Granular Charge.....	4
3.3 Stick Charge.....	4
4. DATA ANALYSIS.....	4
4.1 Pressure Time Histories.....	5
4.2 Radar Data.....	6
4.3 Strain Records.....	7
5. CONCLUSIONS.....	10
6. REFERENCES.....	11
APPENDIX.....	13
DISTRIBUTION LIST.....	17

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## LIST OF FIGURES

Figure		Page
1.	Chamber Simulator Diagram.....	2
2.	Experimental Test Setup.....	3
3.	Granular Charge Breech Pressure.....	4
4.	Stick Charge Breech Pressure.....	5
5.	Granular Charge Displacement.....	6
6.	Granular Charge Velocity.....	6
7.	Stick Charge Displacement.....	6
8.	Stick Charge Velocity.....	6
9.	Granular Charge Strain.....	7
10.	Granular Charge, Strain Axis Direction.....	8
11.	Granular Charge, Strain Magnitude.....	8
12.	Stick Charge Strain.....	9
13.	Stick Charge, Strain Axis Direction.....	9
14.	Stick Charge, Strain Magnitude.....	10
15.	Modified Test Projectile and Simulator Chamber.....	11

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## 1. INTRODUCTION

Armor Piercing Fin Stabilized Discarding Sabot (APFSDS) projectiles which intrude significantly into the gun chamber have become commonplace in modern tank gun ammunition. At the Ballistic Research Laboratory (BRL), interior ballisticians have developed propelling charges which have continued to push the level of performance despite the increasing intrusion of projectile tailbooms into the gun chamber. Most propelling charge development studies have treated the tailboom as a solid obstacle to be worked around, instead of as a dynamic mechanism which requires "gentle" treatment to preserve the performance of the gun system.

This experiment used projectiles instrumented with strain gauges in transparent acrylic gun chamber simulators to compare the magnitudes of the transverse forces applied to a projectile tailboom during the ignition phase of two different propelling charges. The first charge was based on stick propellant. The second charge used granular propellant.

Stick charges, which are usually packed in a very uniform, radially symmetric fashion generally show little resistance to gas flow in the axial direction, but offer large resistance to radial gas flow (Minor 1988). Large portions of a stick charge are typically bound to the tailboom of an APFSDS projectile. This practice creates radially symmetric regions of ullage which enhance the flame spreading process and generally reinforce the symmetry of the ignition process. Without taking into account any asymmetries caused by the ignition system, the symmetry of the stick charge design yields the intuitive impression that the radial forces applied to the tailboom by the charge during the early ignition phase should be uniform.

Granular charges, which are usually packed in a random fashion, show a uniformly high resistance to gas flow in all directions. Because the propellant is not bound to the projectile it tends to settle during shipping and handling, forming unpredicted ullage regions in the charge. Ullage regions have a very low resistance to gas flow, and can produce unpredictable, asymmetric ignition paths in a charge. By intuition, asymmetric ignition increases the likelihood of unbalanced loading of the projectile tailboom.

## 2. SIMULATOR BACKGROUND

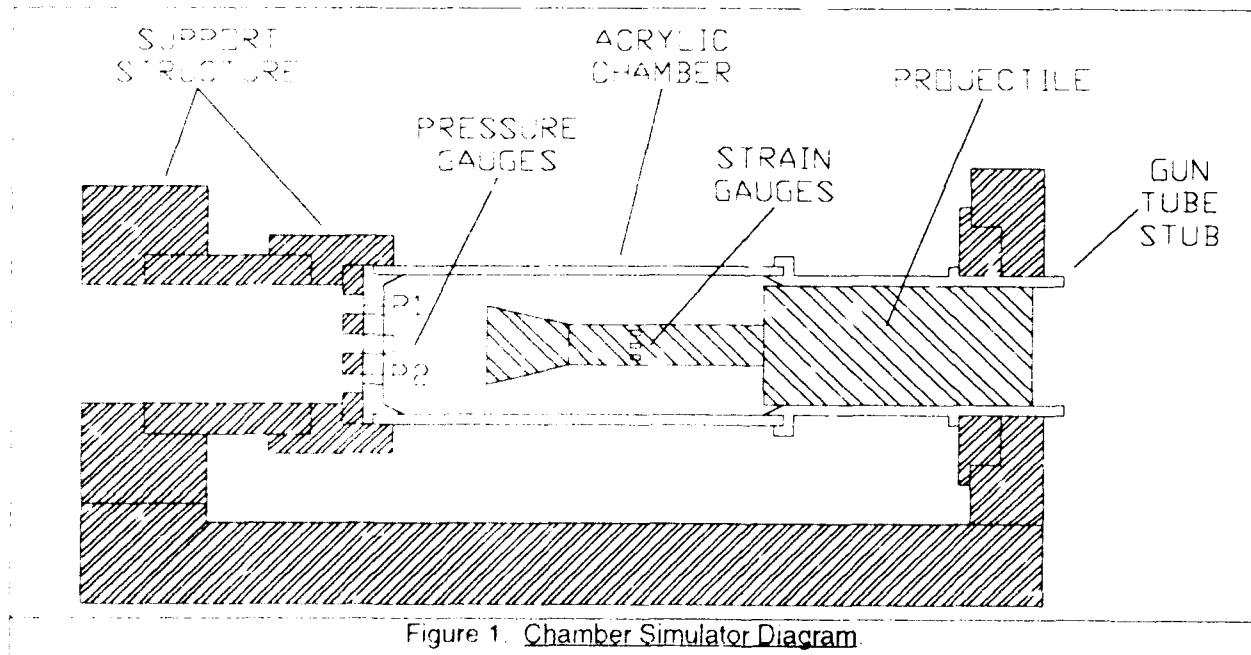
Gun chamber simulators are an important tool supporting the development of advanced large caliber gun ammunition at BRL. Simulators generally use a transparent tube, with inside dimensions duplicating the gun chamber under study. This allows the experimenter visual access to events in

the early stages of the interior ballistic cycle. These events include the functioning of the ignition system, flame spreading, and propellant bed motion and compaction. Recently, simulators have been used in the development of low vulnerability charges (Chang 1988), and in the analysis of flame spreading in howitzer charges (Minor and Horst 1986).

There are two basic types of gun chamber simulators used in large caliber interior ballistic studies at BRL. One type is fabricated from transparent acrylic. This allows the use of high speed cameras to record ignition and flame spreading phenomena, as well as propellant bed motion and compaction. The other type of simulator chamber is made of wound fiberglass. This type withstands significantly more pressure before bursting, but it is only transparent to flash X-rays, and translucent to intense light sources (e.g., flame fronts). The limited transparency of the fiberglass wound chamber makes the clear acrylic chamber the most useful choice unless higher maximum pressures are required to observe the phenomena under study.

### 3. EXPERIMENTAL

3.1 Simulator Instrumentation. Generally, the simulator is instrumented with one or more pressure transducers in the simulator breech, and an array of pressure, acceleration, and force transducers at the rear of the projectile. Video and high speed film cameras can be used to capture the ignition event on film. Flash X-rays are normally taken before (static), and during (dynamic) the ignition to measure movement of the propellant bed or ignition system parts. A displacement measurement device is usually used to detect axial motion of the projectile before chamber failure.



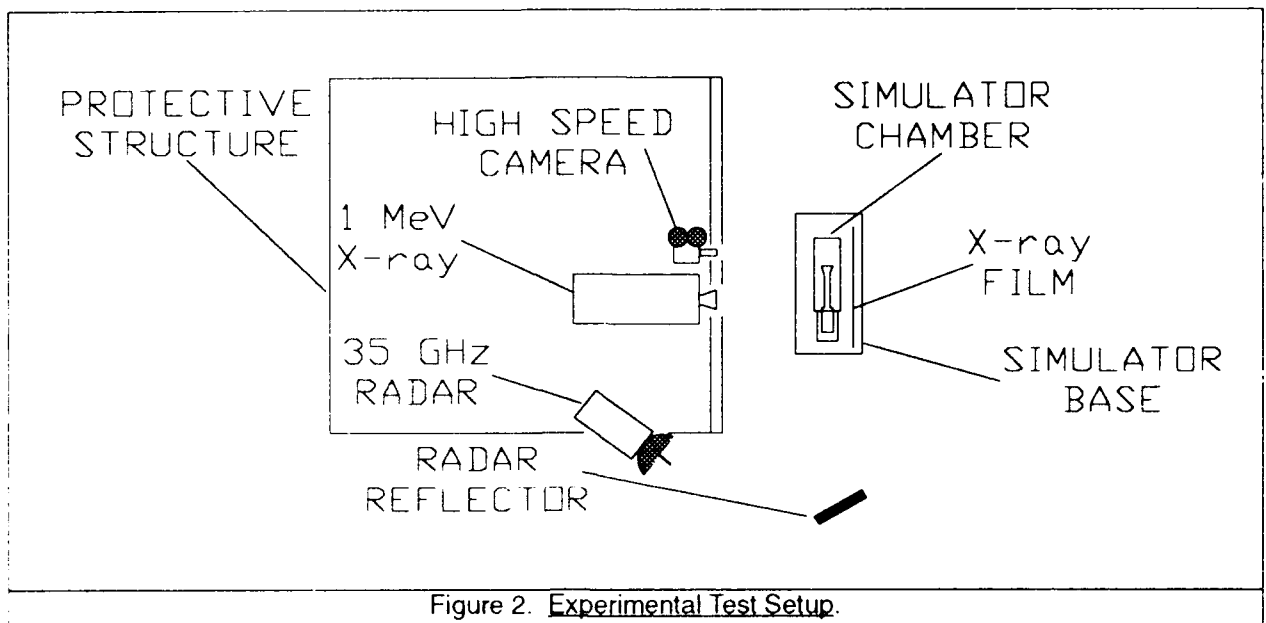


Figure 2. Experimental Test Setup.

Figure 1 presents a cross-sectional view of the transparent acrylic gun chamber simulator used in this experiment. The transparent chamber offered excellent visualization for the recording of the events which occurred inside. The chamber was capable of withstanding dynamic pressures in excess of 15 MPa (2170 psi) before rupturing. The rear end of the chamber was adapted to the base of a real cartridge and its forward end to a short gun barrel in which an instrumented projectile was loaded. The unit was mounted on a steel fixture.

Figure 2 illustrates the experimental arrangement for the simulator diagnostics. The instrumentation used included two pressure gages, six 350 ohm strain gauges, one Photec high speed 16-mm camera, one 1MeV X-ray head, and one 35 GHz TERMA microwave interferometer for measurement of projectile axial displacement.

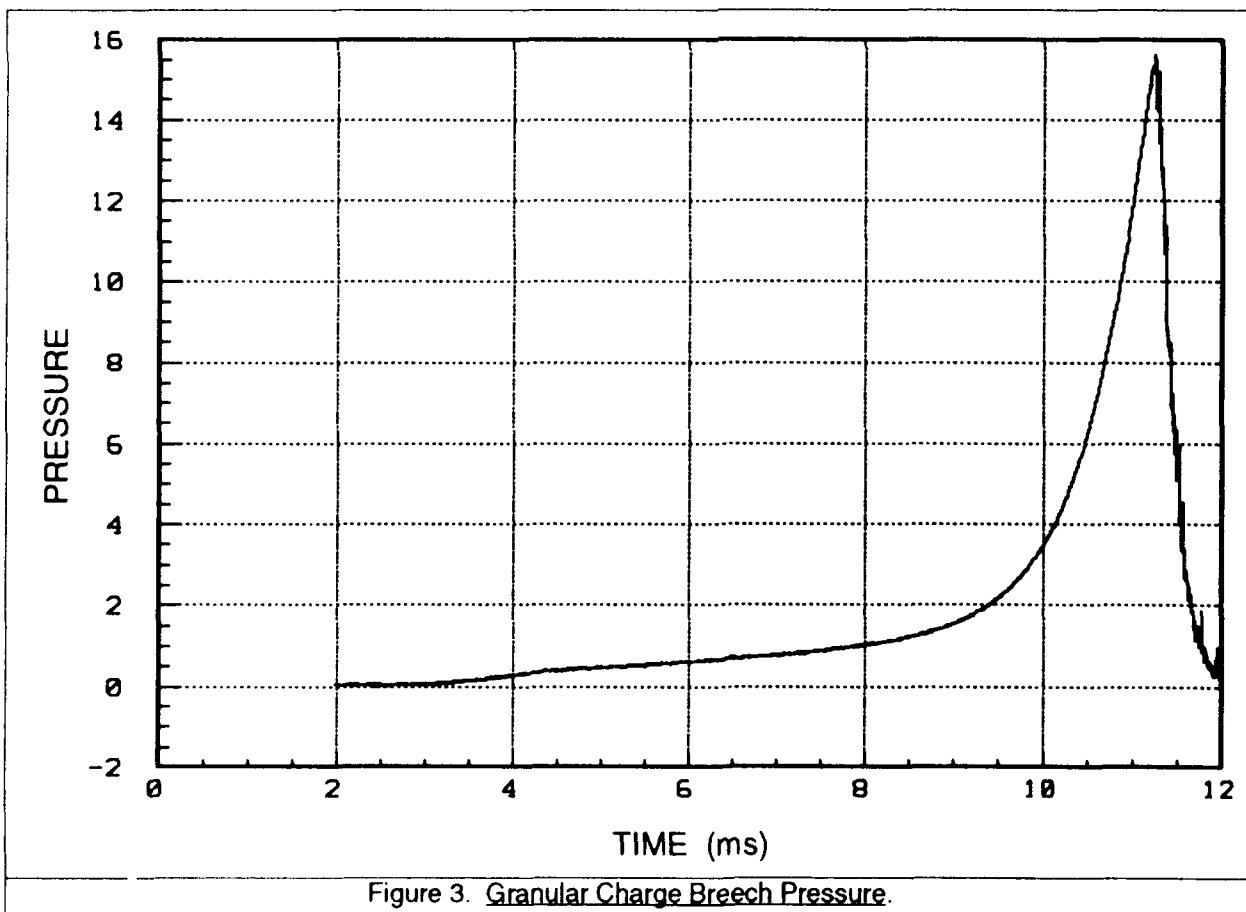
The pressure gages (quartz PCB Model 113A23) monitored the pressure at the breech end of the chamber. The strain gauges monitored the stress on the projectile tailboom from six radially symmetric locations at the same axial position. The high speed camera monitored flame spreading in the entire chamber. The framing rate of the cameras was set at 5000 frames per second. The X-ray head was positioned on one side of the chamber and a cassette containing Kodak XR-5 film was on the opposite side to record propellant bed motion and compaction. The radar signal was reflected onto the nose of the projectile to record projectile motion as a function of time. The resolution of the 35 GHz radar was 0.4287 cm per interference fringe. Data acquisition and data reduction were performed by the Telemetry Acquisition Reduction and Plotting System (TARPS) at BRL's Sandy Point large caliber firing facility. Data were filtered by a 90 kHz low pass filter to eliminate aliasing before being digitized at 1 MHz.

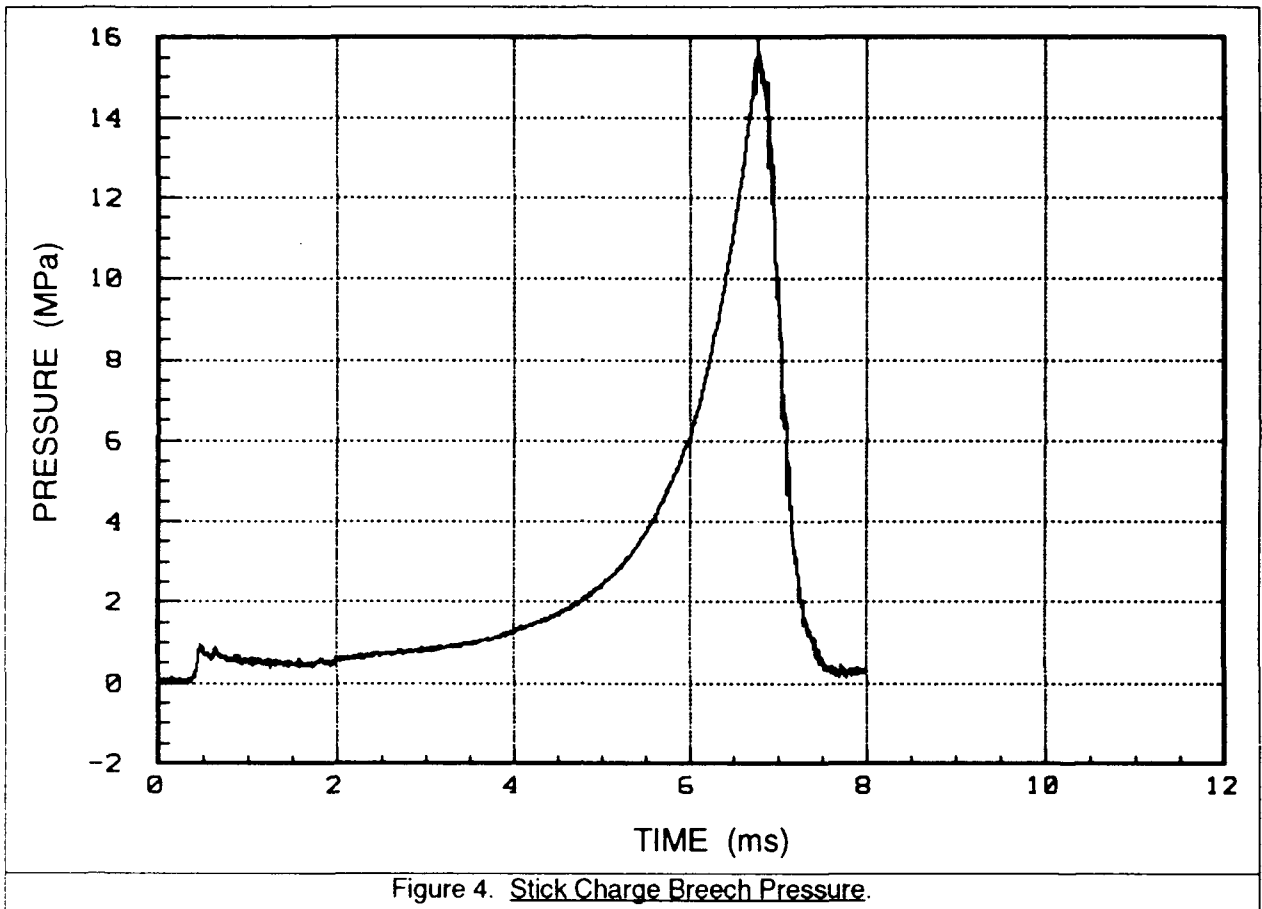
3.2 Granular Charge. The charge for the granular firing consisted of a modified XM125 bayonet primer and 6.58 kg (14.5 lbs) of JA2, multi-perforated, hexagonal propellant. Before firing, the charge was rotated about its long axis to form a uniform 2 cm ullage at the top of the chamber. This was an attempt to create an asymmetry in the path of flame spreading.

3.3 Stick Charge. The charge for the stick firing consisted of an M123 primer, a doughnut shaped bag containing 100 g (0.22 lbs) of black powder around the primer, and 6.4 kg (14.1 lbs) of JA2 stick propellant. The majority of the stick propellant (4.13 kg (9.1 lbs)) was stacked symmetrically around the tailboom (long axis of the propellant parallel with the long axis of the projectile) and secured with packing tape. The remaining 2.27 kg (5 lbs) of propellant was stacked in a cylindrical shape, taped together, and placed in the chamber behind the projectile. X-ray data showed that the 2.27 kg charge increment was offset from the main portion of the charge by 1 cm before firing.

#### 4. DATA ANALYSIS

The following data are by no means generic to the types of charges used. The intent of this two round study was to verify the viability of the experimental techniques and to point out methods of improving future experiments.





4.1 Pressure Time Histories. Figure 3 shows the pressure time history for the granular round. The chamber ruptured at 15.5 MPa (2250 psi), 11.2 ms after application of the firing pulse. The pressure rise appeared to be smooth, without noticeable aberration. The bayonet primer vented into the center of the propellant bed so the breech mounted pressure gauges did not detect a pressure spike at the time of primer function. Since there were no pressure gauges at the forward end of the chamber no pressure wave analysis could be attempted. For purposes of clarity only the data from one pressure transducer are shown in the plots. In all cases the data from the other transducer are nearly identical to the data shown.

Figure 4 shows the pressure time history for the stick round. The chamber ruptured at 15.5 MPa (2250 psi), 6.8 ms after application of the firing pulse. A spike was observed on the pressure time history corresponding to the functioning of the M123 primer and combustion of the black powder bag. The ullage region at the rear of the propelling charge, where the ignition elements were located, was adjacent to the pressure transducers mounted in the stub base. This arrangement allowed the pressure transducers to respond immediately to the ignition pulse.

Except for the evidence of primer function on the stick charge pressure time history the only significant difference between the granular and stick charge pressure traces was the time delay before chamber rupture. It is possible that the low axial flow resistance of the stick charge allowed it to ignite more quickly, and burst the simulator 4.4 ms before the granular charge.

4.2 Radar Data. Figures 5 and 6 show the distance and velocity histories of the granular round. The projectile had moved 3.25 cm and was traveling at 15 m/s at the time of chamber rupture (11.2 ms). Figures 7 and 8 show the distance and velocity histories of the stick round. The projectile had moved 3.40 cm and was traveling at 30 m/s at the time of chamber rupture (6.8 ms).

The higher projectile velocity attained in a shorter time by the stick charge shows the effects of the lower resistance to axial gas flow as compared to the granular charge. It should also be noted that, in this case, greater acceleration places more stress on the projectile.

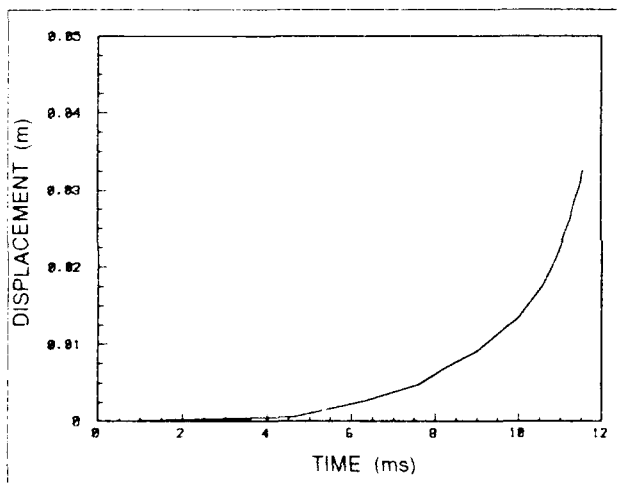


Figure 5. Granular Charge Displacement.

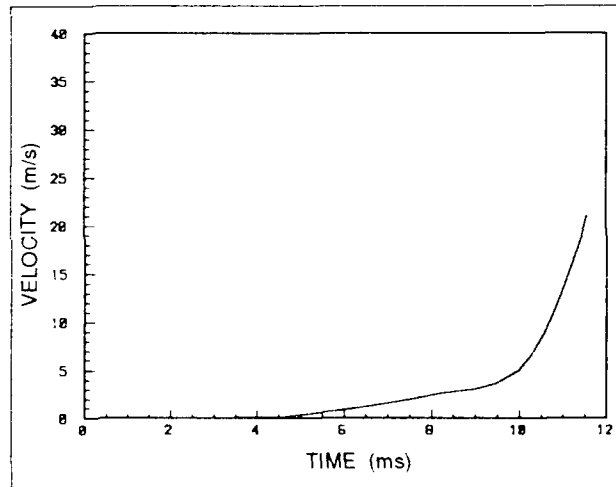


Figure 6. Granular Charge Velocity.

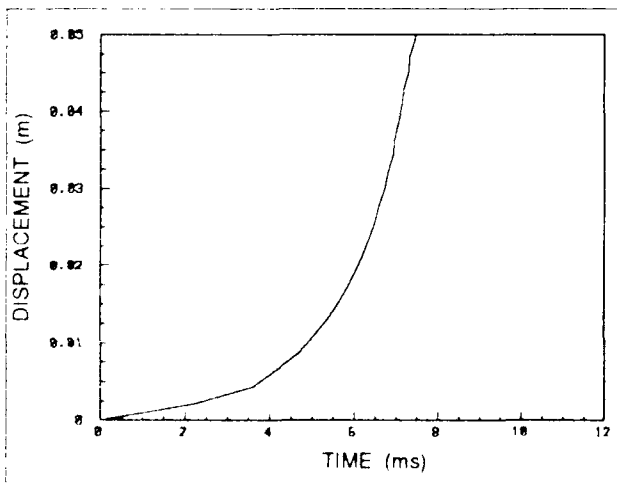


Figure 7. Stick Charge Displacement.

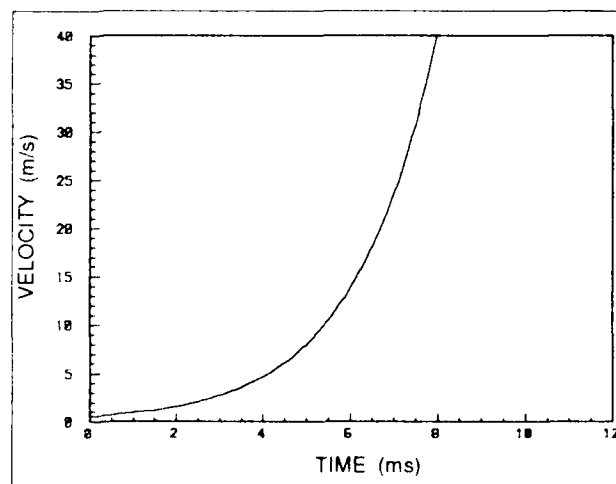
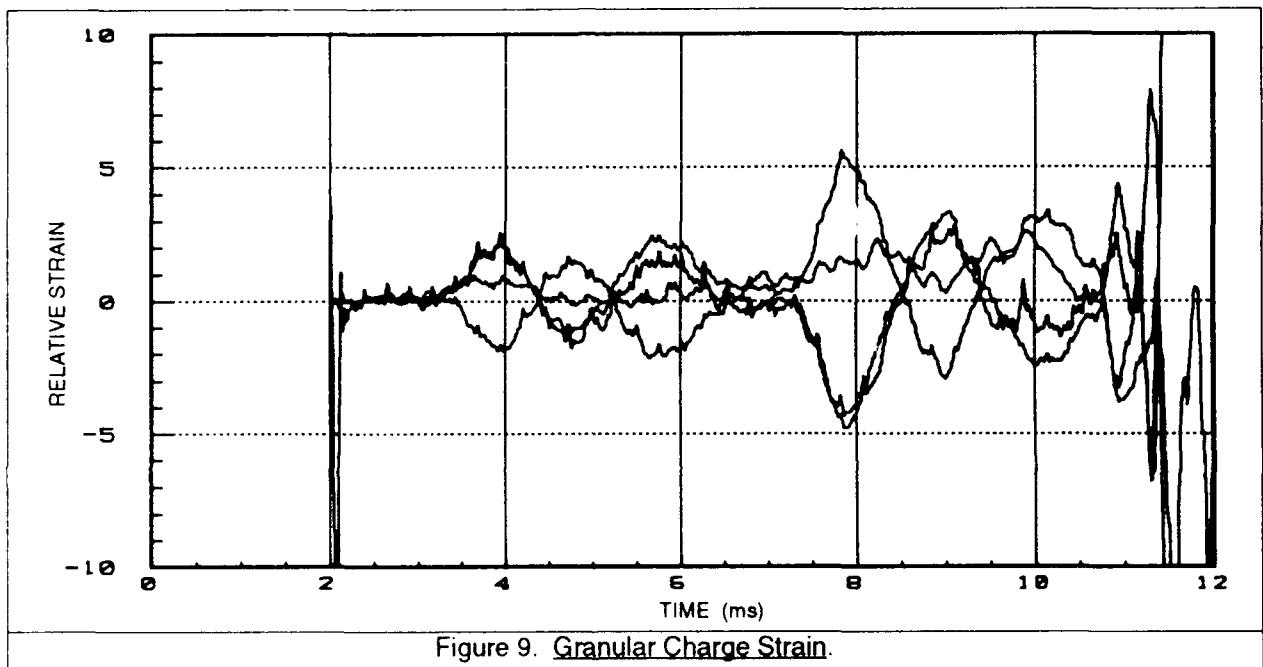


Figure 8. Stick Charge Velocity.

4.3 Strain Records. Figure 9 shows the strain histories for the granular charge. Data shown for diametrically opposite strain gauges account for approximately equal, but opposite peaks. The first major peak (at 4 ms) shows the beginning of oscillatory strains on the tailboom. These strains were probably due to propellant bed compaction resulting from primer function. Notice that the strains have settled at 7 ms, indicating that the propellant bed was fully compacted. The peak at 8 ms corresponds to a change in the slope of the pressure time history (Figure 3) which indicated ignition of the main charge. This was the beginning of a second set of strain oscillations in the tailboom which terminated at chamber rupture.



It was desired to evaluate the strain data to determine the angular location of the peak strain value, and the magnitude of the peak strain. It was acknowledged that this analysis would only be valid for the axial location common to the strain transducers under study.

If it is assumed that the magnitude of axial strain in the tailboom of a projectile falls off linearly with respect to angle, then the magnitude and angle of the principal strain axis of the tailboom can be determined at any point in time by evaluating the strain histories from two transducers which were not on the same radial axis. The analysis is executed by expressing the two measured strain data sets in terms of a projection of a strain of unknown magnitude which is at a maximum at an unknown angle from each strain measurement position. Since the angle between the strain measurement positions is known we are left with two equations and two unknowns. The simultaneous solution for the axis angle is substituted back into the original equations to determine the

magnitude of strain on that axis. This calculation is performed for each point in the strain-time histories. Appendix A details the mathematical manipulations.

Figures 10 and 11 display calculated values for the orientation of the axis of maximum strain (parallel to the bending axis of the tail boom) and the absolute magnitude of the strain that would have been measured on that axis if a transducer had been placed there. Each plot contains three separate data sets. These correspond to evaluation of the  $0^\circ$ ,  $60^\circ$  pair, the  $0^\circ$ ,  $120^\circ$  pair, and the  $60^\circ$ ,  $120^\circ$  pair of strain records. If the axis of maximum strain had not moved during the test the plot of axis direction versus time would appear as a square wave, with the angle shifting by  $180^\circ$  each time the strain passes through zero. As shown for the granular charge the axis of maximum strain seems to have moved from horizontal oscillations, to near vertical oscillations, then switched back to horizontal just before the chamber ruptured. It should be noted that if at any time the axis of maximum strain was perpendicular to the angular midpoint between the transducer pair under

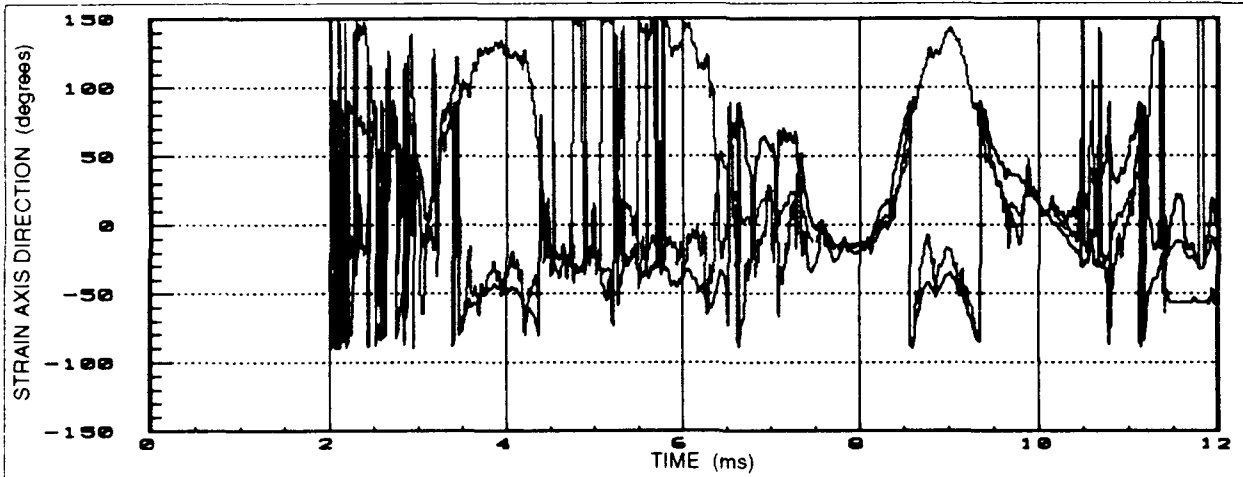


Figure 10. Granular Charge, Strain Axis Direction.

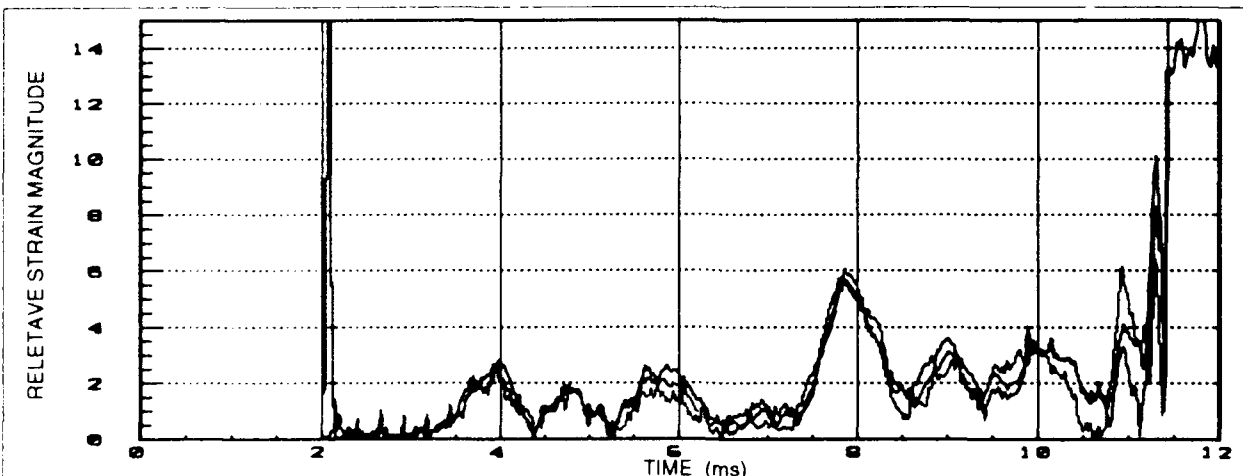


Figure 11. Granular Charge, Strain Magnitude.

evaluation, the axis direction calculation wouldn't show a zero crossing at that instant because it was on a null point. This occurred several times in the evaluation of the 60°, 120° pair. That trace stands out in Figure 10 because it remains above the 90° level while the other two traces show direction reversals.

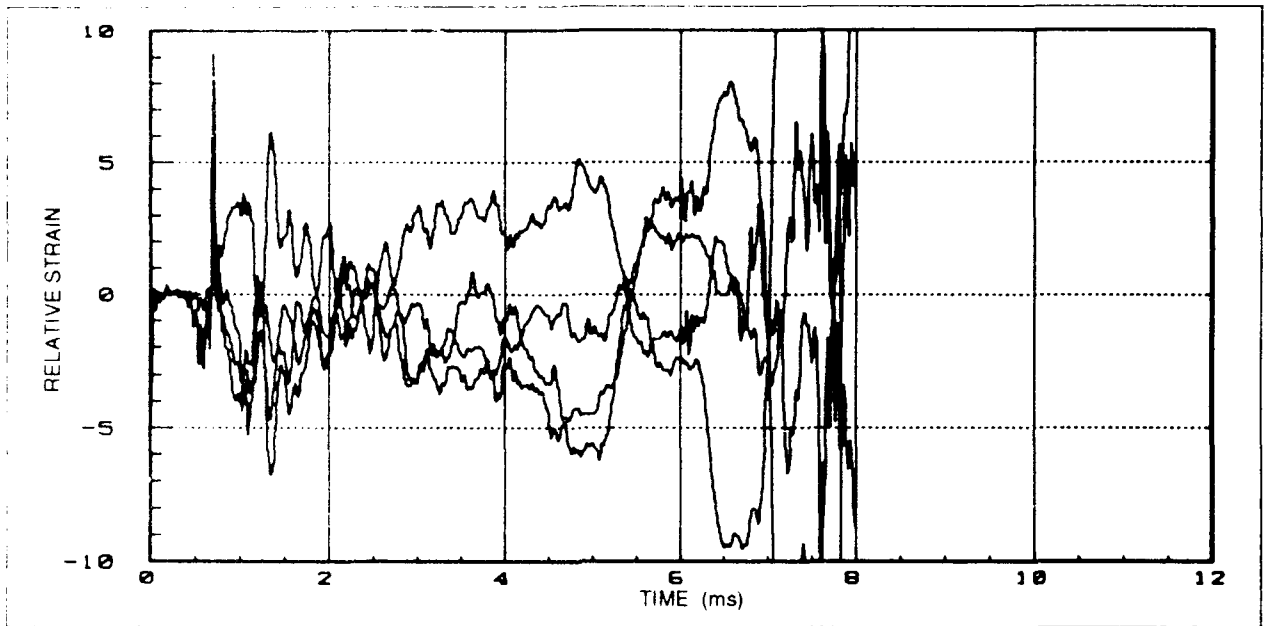


Figure 12. Stick Charge Strain.

Figure 12 shows the strain histories for the stick charge. The very large peak at 1 ms was electronic noise as indicated by the fact that all of the strain gauges responded in the same direction. Overall, the stick charge strain history showed oscillatory behavior similar to the granular charge, but at a much higher frequency. The first significant peak (1.5 ms) was probably due to the impact of the 2.27 kg charge increment on the main charge. The strain peak at 5 ms corresponded to the

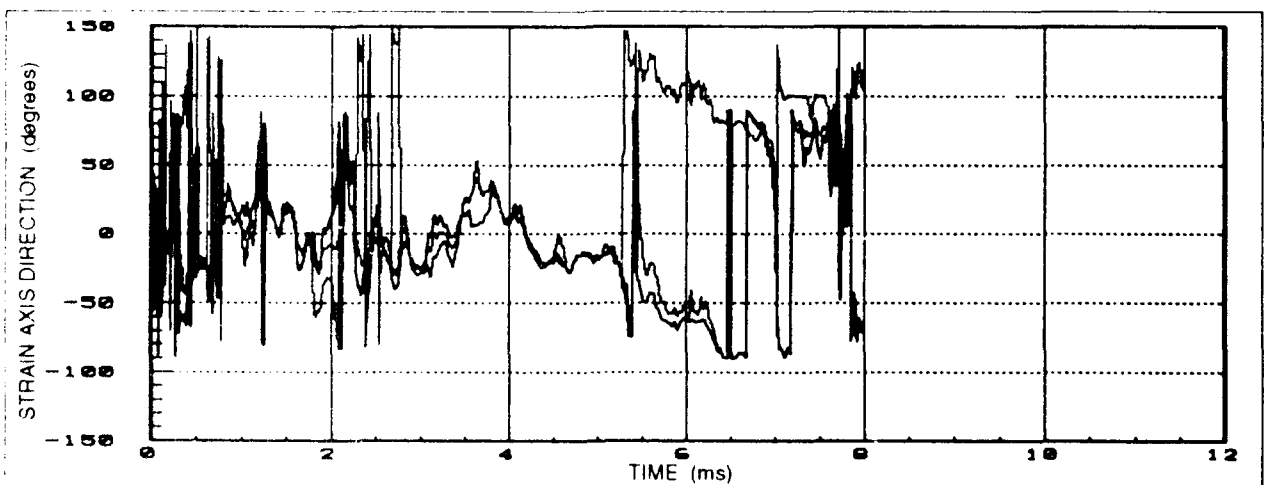


Figure 13. Stick Charge, Strain Axis Direction.

slope change in the pressure time history shown in Figure 4. This slope change indicated the ignition of the main charge. The large strain peak after 6 ms in Figures 12 and 14 occurred after chamber rupture and was not considered significant. Figures 13 and 14 show the calculations for principal strain axis and strain magnitude as detailed previously. The data set pairs evaluated for this round were the  $0^\circ$ ,  $60^\circ$  pair, the  $0^\circ$ ,  $300^\circ$  pair, and the  $60^\circ$ ,  $300^\circ$  pair.

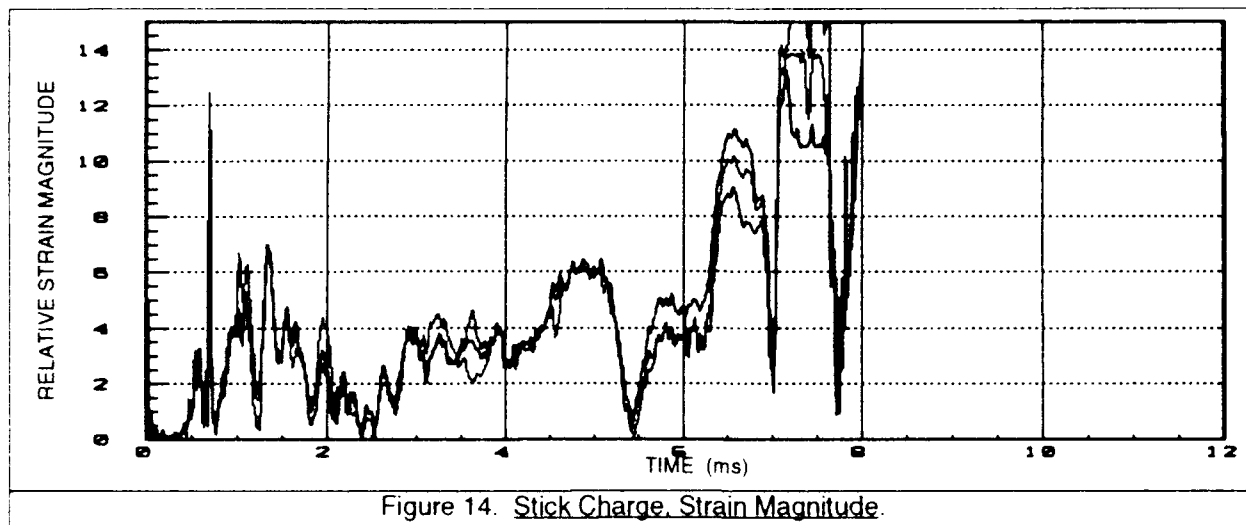


Figure 14. Stick Charge, Strain Magnitude.

## 5. CONCLUSIONS

The combination of plastic chamber simulators and strain gauge instrumented projectiles provides an opportunity to make unique measurements of charge/projectile interactions during the early phase of charge ignition. This type of system could be used to evaluate candidate propelling charges for detrimental effects on projectiles, as well as studying ignition anomalies caused by charge geometry. An instrumented tail boom can be a sensitive, yet robust transducer for data collection in the interior ballistic environment.

The success of these test firings indicates that this method of projectile instrumentation is indeed a viable method of recording transverse forces on projectile tailbooms in the early phase of charge ignition.

Although the limited scope of this test prevents the drawing of any general conclusions about the differences in propellant/tailboom interactions during the ignition phase of granular and stick charges, it does show that each charge has measurably different interactions with the tailboom. Future applications of this technique should be optimized to study specific projectile/charge systems or particular ignition asymmetries in each type of charge.

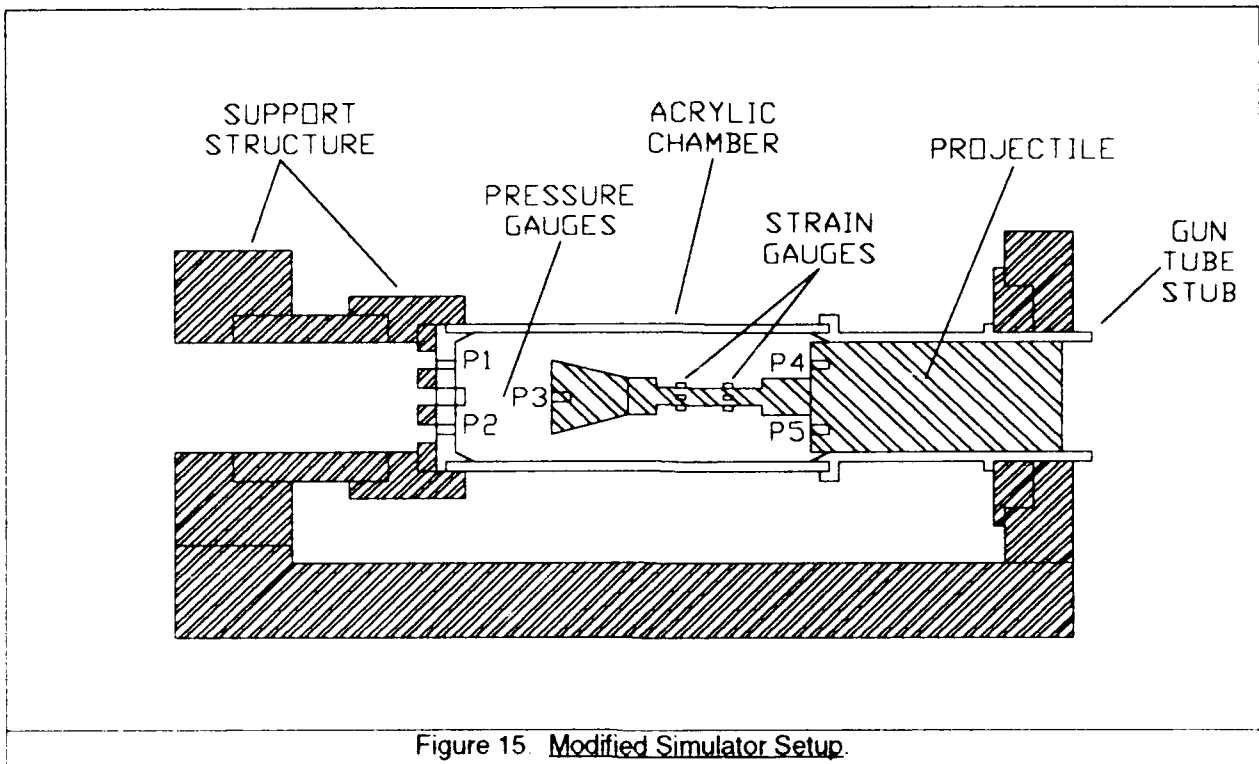


Figure 15. Modified Simulator Setup.

To accurately compare data between different charges and loading asymmetries a common projectile should be used. Figure 15 shows a sketch of a projectile which would be better suited for such a purpose. The pressure gauges mounted in the projectile would allow analysis of ignition induced pressure waves. Locating the strain gauges on a thinned down section of the tailboom would make the system more sensitive to asymmetric forces. An increased number of strain gauges combined with finite element analysis of the test projectile could allow the strain data from a generic projectile to be applied to specific systems under study.

6. REFERENCES

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Minor, T.C. "Flamespreading Processes in Obliquely Loaded Stick Propellant Beds." BRL-TR-2894, U.S. Army Ballistic Research Laboratory, Aberdeen Proving Ground, MD, March 1988.

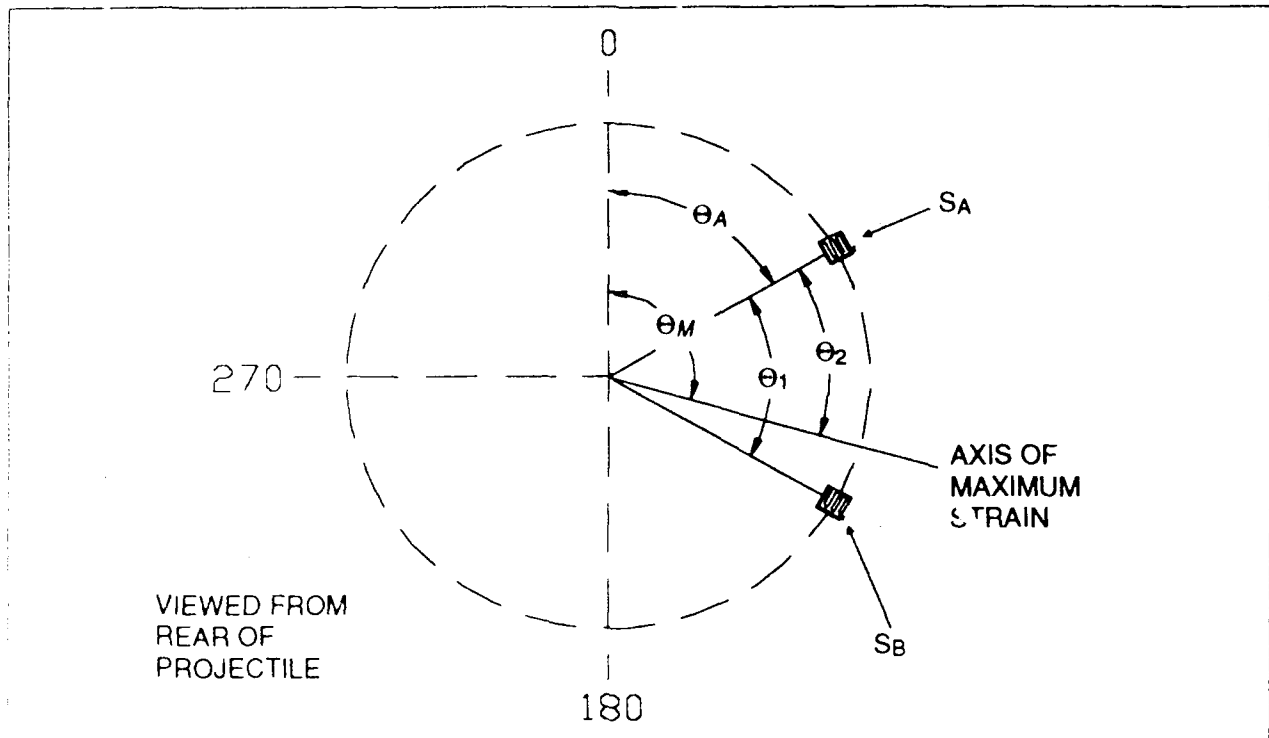
Minor, T.C., and A.W. Horst, "Theoretical and Experimental Investigation of Flamespreading Processes in Combustible-Cased, Stick Propellant Charges." BRL-TR-2710, U.S. Army Ballistic Research Laboratory, Aberdeen Proving Ground, MD, February 1986.

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**APPENDIX.**

**PROCEDURE FOR REDUCING STRAIN DATA**

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Equations A.1 through A.8 detail the procedure used to produce the data shown in Figures 10, 11, 13, and 14 in this report. The reader is referred to the diagram above to explain the angular references made in the equations.

Given:  $M_A$  and  $M_B$  are instantaneous strain magnitudes measured at non-radially-aligned transducer locations  $S_A$  and  $S_B$ . The angle between the two transducers is defined as  $\theta_1$ , the angle between  $S_A$  and the axis of maximum strain is defined as  $\theta_2$ , and the angle from top-dead-center to  $S_A$  is defined as  $\theta_A$ . The computed magnitude of maximum strain is represented as  $M_M$ , and the angle of the axis of maximum strain is  $\theta_M$ .

Our purpose is to solve for  $M_M$  and  $\theta_M$

Given:

$$\theta_M = \theta_A + \theta_2 \quad (\text{A.1})$$

$$M_A = M_M \cos \theta_2 \quad (\text{A.2})$$

$$M_B = M_M \cos(\theta_2 - \theta_1) \quad (\text{A.3})$$

Therefore:

$$\frac{M_B}{M_A} = \frac{\cos(\Theta_2 - \Theta_1)}{\cos\Theta_2} \quad (\text{A } 4)$$

$$\frac{M_B}{M_A} = \frac{\cos\Theta_2 \cos\Theta_1 + \sin\Theta_2 \sin\Theta_1}{\cos\Theta_2} \quad (\text{A } 5)$$

$$\frac{M_B}{M_A} = \cos\Theta_1 + \sin\Theta_1 \tan\Theta_2 \quad (\text{A } 6)$$

$$\arctan\left(\frac{\frac{M_B}{M_A} - \cos\Theta_1}{\sin\Theta_1}\right) = \Theta_2 \quad (\text{A } 7)$$

Therefore:

$$\Theta_M = \arctan\left(\frac{\frac{M_B}{M_A} - \cos\Theta_1}{\sin\Theta_1}\right) + \Theta_A \quad (\text{A } 8)$$

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