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**Temperature Dependence of Species Concentrations
Near the Substrate During CVD Diamond Growth**

by

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Temperature dependence of species concentrations near the substrate during CVD diamond growth

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Measurements have been made of the temperature dependence of CH₃, CH₄, and C₂H₂ very near the substrate during filament-assisted diamond growth. CH₃ was detected using (2+1) resonance-enhanced multiphoton ionization (REMPI), and CH₄ and C₂H₂ concentrations were measured using sampling mass spectrometry. A strong dependence of the CH₃ REMPI signal on substrate temperature was observed, which at low temperatures may be characterized as having an activation energy of approximately 4 ± 1 kcal/mole. Methane and acetylene, on the other hand, are relatively independent of substrate temperature. These results are most likely due to recombination of methyl to methane in the cool gas layer near the substrate or on the surface at low substrate temperatures.

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I. INTRODUCTION

The technology of diamond film deposition has developed rapidly since the discovery that diamond could be grown efficiently under metastable conditions by chemical vapor deposition. Polycrystalline films can now be produced with optical, mechanical, and thermal properties approaching those of natural diamond. Nevertheless, many proposed applications for diamond require that several remaining problems be overcome, for example, the high growth temperatures needed, the incorporation of hydrogen and sp^2 carbon impurities, and the relatively low efficiency (and therefore high cost) of current growth processes.

One barrier to overcoming these problems is that the fundamental mechanisms of diamond growth and graphite suppression under CVD conditions are not well understood. Recent work in several laboratories has shown that the methyl radical CH_3 is the dominant growth species under typical conditions [1-5], but under some conditions diamond can grow from acetylene as well [4]. It has also been proposed that atomic carbon may contribute to diamond growth in some plasma environments [6, 7]. Mechanisms have been proposed for incorporation of carbon into the diamond lattice from either methyl [8-12] or acetylene [12-15].

One important test of a growth mechanism is whether it can reproduce the observed dependence of the growth rate on substrate temperature [16,17]. The measured temperature dependence, however, may have two sources: the intrinsic temperature dependence of the surface growth mechanism, and the variation with substrate temperature of the gas-phase reactants (e.g., CH_3 , C_2H_2 , H) at the surface. To compare theory with experiment, it is first necessary to measure the dependence of the gas-phase reactant concentrations at the surface on substrate temperature. This information is also useful for validating computational models of the diamond growth environment [7, 18-22].

In this paper, we examine the effect of substrate temperature on the gas-phase concentrations of CH_3 , CH_4 , and C_2H_2 very close to the substrate during filament-assisted diamond growth. The methyl radical is detected *in situ* using resonance-enhanced multiphoton ionization (REMPI), which has been shown previously to be a sensitive diagnostic for methyl under filament-assisted diamond growth conditions [23]. Methane and acetylene are measured by sampling mass spectrometry.

The results of this study show that the CH_3 concentration near the substrate is significantly depleted at substrate temperatures below 1000 K. The CH_4 and C_2H_2 concentrations near the substrate are only weakly dependent on substrate temperature. As discussed below, the most likely explanation for these results is either homogeneous recombination of methyl via $CH_3 + H \rightarrow CH_4$ occurring in the cool gas region near the substrate, or heterogeneous recombination on the diamond surface. However, surface reactions which produce CH_3 (particularly etching of sp^2 carbon by atomic hydrogen) may contribute as well to the observed temperature dependence.

II. EXPERIMENTAL

The experimental apparatus is shown schematically in Fig. 1. The hot filament reactor consists of a 2.0 inch OD five-way cross with water-cooled walls. The gas inlet and the electrical feedthroughs for the filament current and the ionization probe bias are in the top arm of the cross, and the electrical feedthrough for the substrate temperature control and the gas exhaust to the vacuum pump are in the bottom arm. A tee in the exhaust line diverts a small portion of the exhaust to a quadrupole mass spectrometer for on-line analysis during the REMPI measurements. Two opposing side arms are fitted with calcium fluoride windows to allow the laser beam inside the reactor.

Figure 1(b) shows the filament, substrate and ionization probe geometry in more detail. The substrate consists of a 5 mm wide, 20 mm long, and 125 μm thick molybdenum foil. The substrate is heated electrically using a DC current, which allows us to control its temperature independently of the filament temperature. The substrate temperature is measured with an icepoint-corrected chromel-alumel thermocouple (type K) spot-welded to the center of the substrate lower surface. The substrate is supported by two 1.6 mm diameter tungsten posts, which are in turn mounted on a water-cooled copper feedthrough.

The filament consists of a single, straight 200 μm diameter tungsten wire, approximately 25 mm long, and was typically kept at a distance of 7 to 15 mm from the substrate. The lowest attainable substrate temperature (with no power applied) was 900 K at 7 mm and 700 K at 15 mm, when the filament temperature was kept at 2200 K.

The filament brightness temperature is measured through a glass viewport on a side arm using a one-color disappearing-filament pyrometer. While we report here only the directly-measured brightness temperature, an emissivity correction can be applied to estimate the true filament temperature [24].

The ionization probe is a 250 μm platinum wire supported by a ceramic tube. A bead was formed on the probe end by a standard arc thermocouple welder. The probe bead is located 2 mm above the center of the substrate. The laser is focused just above the substrate, at a distance from the substrate of less than 0.5 mm. Since modeling studies [18] have shown that the chemical mean free path of CH_3 is approximately 1 mm under filament-assisted conditions, we expect the methyl concentration at the measurement position to be essentially the same as that at the film surface. Either a 20 or 50 cm focal length focusing lens is used, resulting in a beam waist of less than 40 μm .

To make the CH_4 and C_2H_2 measurements, the experimental setup was modified to accommodate a quartz microprobe for mass spectroscopic sampling. In the present apparatus it was not possible to accommodate both the microprobe and the ionization probe simultaneously, and so CH_4 and C_2H_2 concentrations were measured with the microprobe in separate experiments which replicated the conditions of the REMPI CH_3 measurements. To insure that the conditions were identical, the gas exhaust

composition was measured both during the REMPI and microprobe experiments, and used as a check that the conditions were the same. We found that changes in the gas composition at the substrate, as measured with the microprobe, were always reflected in the exhaust gas also, even though the measured absolute mole fractions were different.

The microprobe is a 6 mm diameter quartz tube, drawn down to a probe tip diameter of 150 μm . The pressure inside the probe is held at 0.3 torr, to quench chemistry other than radical recombination. This sampling method has been used successfully by Harris, Weiner and Perry under similar conditions [25]. The sampled gas flows past an orifice, which allows a small portion of the gas into the quadrupole mass spectrometer ionizer. Methane and acetylene are monitored from mass peaks 15 and 26, respectively. A calibration gas mixture of 0.3% CH_4 and 0.3% C_2H_2 in H_2 was used to calibrate the signal to yield absolute mole fractions.

The UV beam used for the REMPI measurements was produced by frequency-doubling the output of a YAG-pumped pulsed dye laser (Continuum YG 661 pumping a TDL 51) in an autotracking converter (INRAD). The dye used was DCM, dissolved in DMSO to extend its emission to 670 nm. The dye laser linewidth is approximately 0.08 cm^{-1} . After frequency doubling, a pulse energy of 1 mJ with a typical duration of 7–8 ns was obtained at the wavelength of the CH_3 REMPI line (333.5 nm).

The REMPI signal from the ionization probe was amplified and detected by a boxcar averager (Stanford Research Systems SR 250), using an integration gate width of 100 ns. The signal was then digitized and sent to a microcomputer for further processing.

Before any measurements were taken, the substrate was scratched with diamond paste and a diamond film was grown for at least 12 hours, using a gas flow of 100 sccm of 0.5% methane in hydrogen, a filament brightness temperature of 2200 K, and a substrate temperature of 1000 K. Under these conditions, a continuous diamond film with typical crystal size of approximately 1 μm was always observed to grow.

III. VALIDATION OF REMPI DETECTION SCHEME

A typical survey ionization spectrum is shown in Fig. 2. This spectrum was taken immediately above the substrate, using a gas mixture of 0.5% CH_4 in H_2 . The main features in this spectrum are the CH_3 (2+1) REMPI signal at 333.5 nm and the atomic carbon (2+1) REMPI signal at 320.2 nm, which results from photolytically-produced C. We were not able to identify the broad peak near 326 nm. Another feature is the high background ionization level, which is a function of laser energy. The two broad peaks near 319 and 329 nm are close to CH_3 assignments. Celii *et al.* [23] observed, however, that the CH_3 $3p^2A_2'' 2_1^1$ transition is hidden under an unassigned peak at 329.1 nm. We did not observe the other unattributed peak at 332.1 nm observed by Celii *et al.*, even at the highest filament temperatures in our system.

The peak at 333.5 nm results from the vibrationless (0-0) two-photon absorption from the X^2A_2'' electronic ground state of CH_3 to the $3p^2A_2''$ Rydberg state, followed

by an incoherent single-photon ionization. The observed peak is the strong Q branch of the band, whose rotational structure is completely unresolved, due to effect of predissociation on the rotational linewidth.

All results reported below are obtained from the peak of this transition, correcting for background ionization. To relate this signal to the CH₃ concentration, several factors must be considered: (1) the electron collection must be saturated, (2) the collection efficiency must be independent of changes in parameters to be varied (i.e., temperature), and (3) the Boltzmann distribution of population among the vibrational and rotational levels of the lower electronic state must be accounted for.

Obtaining saturated electron collection depends primarily on shaping the electric field near the probe appropriately. The biasing scheme used in these measurements is shown in Fig. 1(b). The substrate is grounded and the probe and filament are positively biased. The filament bias was maintained at 20 V less than the probe bias. The electrical noise due to thermionic electron emission was filtered when the filament biasing power supply was used. Its output shunt capacitance was found to substantially reduce the thermionic emission noise, and no further shielding of the probe from the filament was necessary. This biasing scheme resulted in saturated electron collection for various experimental conditions, as evidenced by the plot of collected current vs. bias voltage shown in Fig. 3. For the measurements discussed below, the probe bias voltage was set at 100 V.

To verify the relationship of the REMPI signal to the concentration of the probed species near the surface, a small amount of argon was bled into the reactor and the (3+1) REMPI signal at 314.4 nm from the Ar $3p^6\ ^1S \rightarrow 3p^54s\ 4s'[1/2]^oJ = 1$ transition was measured as a function of filament and substrate temperatures, under the same conditions used for the CH₃ REMPI measurements. Since argon is non-reactive, we expect the Ar concentration to vary inversely with gas temperature, due to the change in gas density with temperature. A similar procedure has been used previously by Smyth and Tjossem [26] to calibrate the REMPI collection efficiency in flame experiments.

Since the $4s'$ transition is dipole-allowed, retroreflection of the beam to form a standing wave geometry is necessary to avoid cancellation of the REMPI signal by competing third harmonic generation [27]. This was accomplished using two 20 mm focal length lenses and one mirror. Due to the sensitivity of the Ar (3+1) REMPI signal to the precise overlap of the beams at the focus, as well as to the stronger dependence of the signal on laser pulse energy (which was stable only to 5%), the fluctuation in the Ar (3+1) REMPI signal is larger than that we observed for the methyl (2+1) signal. Nevertheless, the Ar REMPI signal shows the expected $1/T$ dependence on substrate temperature to within experimental error, as seen in Fig. 4.

Since the CH₃ REMPI signal results from the $v = 0$ vibrational ground state, it measures directly only the population in this state. To relate this to the total CH₃ population, we assume

$$[\text{CH}_3] \propto SQ_{\text{vib}}(T) \quad (1)$$

where S is the measured REMPI signal, Q_{vib} is the vibrational partition function, and T is the substrate temperature. We estimate that the gas temperature at the measurement location is within 50 K of the substrate temperature. The partition function is calculated from the known frequencies and degeneracies of the vibrational modes of CH_3 ($\nu_1 = 2992.6 \text{ cm}^{-1}$, $\nu_2 = 606 \text{ cm}^{-1}$, $\nu_3 = 3158.9 \text{ cm}^{-1}$, $\nu_4 = 1394.6 \text{ cm}^{-1}$, $g_1 = g_2 = 1$, $g_3 = g_4 = 2$), using measured and calculated vibrational energy levels for the first 5 levels of the ν_2 mode [28, 29] and the harmonic oscillator approximation otherwise. Over the temperature range of interest here (700–1300 K), Q_{vib} varies from 1.5 to 3.4, resulting in more than a factor of 2 correction in the temperature dependence data. The large value of Q_{vib} results primarily from the low-frequency out-of-plane vibrational mode ν_2 , which becomes significantly populated at these temperatures.

In principle, a similar partition function correction should be applied to the rotational levels as well, to account for the fact that the laser pumps only some rotational lines. However, this is not feasible in practice since the Q branch is rotationally unresolved. Instead, we assume, like previous investigators [30], that the unresolved rotational manifold of the Q branch is relatively temperature-independent. This is supported by the spectra in Fig. 5, taken at substrate temperatures of 800 K and 1200 K, which show no significant differences in the shape of the manifold. In this case, the signal at the peak is proportional to the area under the curve to within experimental error. However, we cannot at present rule out the possibility that the appearance of the manifold is affected by J-dependent predissociation. In this case, there may be some temperature dependence to the REMPI signal which we do not account for here.

IV. RESULTS

A. Filament temperature dependence

While the focus of this work is the effect of substrate temperature on species concentrations, some measurements of the filament temperature dependence were also made. The variations with filament temperature of the CH_3 REMPI signal and the CH_4 and C_2H_2 mole fractions near the substrate are shown in Fig. 6, for an initial gas composition of 0.5% CH_4 in H_2 . All measurements were performed with the substrate temperature fixed at 1000 K. The conversion of methane to acetylene increases sharply above a filament brightness temperature of 2100 K, accompanied by a sharp increase in the methyl REMPI signal. The CH_4 and C_2H_2 results are similar to those reported previously for a gas composition of 3% CH_4 in H_2 [24]. Also, Celii and Butler [23] report a similar variation of CH_3 with filament temperature. The present results show conclusively that the sharp changes seen in CH_4 , C_2H_2 , and CH_3 all occur over the same filament temperature range, and occur even for the more typical composition used for diamond growth of 0.5% CH_4 in H_2 .

Sommer and Smith [24] have argued that the sharp change in gas composition as

the filament temperature is increased results from filament poisoning at low temperatures, due to a graphite layer coating the filament. For a composition of 0.5% CH₄ in H₂, the carbon-hydrogen equilibrium phase diagram shows that graphite is stable below about 2170 K [24], and our measurements are consistent with this interpretation. (The actual filament temperature is somewhat higher than the brightness temperature, but when the filament is graphite-coated the emissivity is approximately 0.9, so the difference is not large.)

As shown in Fig. 6, once the graphite layer desorbs from the filament, the CH₃, CH₄, and C₂H₂ signals saturate, and are apparently unaffected by further increases in filament temperature. This is in contrast to atomic hydrogen, which several studies have shown to continue to increase with increasing filament temperature [31–33].

The effect of gas flow rate on the methane and acetylene mole fractions is shown in Fig. 7. While the CH₄ mole fraction in the exhaust gas was measured to be 0.5% at low filament temperatures, near the substrate it is always less than this value, and shows a dependence on gas flow rate. This behavior is consistent with thermal diffusion (the Soret effect), which will cause heavy species to diffuse preferentially to cold regions of the reactor near the walls, depleting the hot filament/substrate region of hydrocarbons. The conversion of methane to acetylene is seen to be relatively unaffected by gas flow rate, since for both cases the change in the CH₄ signal and the magnitude of the C₂H₂ signal are nearly the same.

B. Substrate temperature dependence

The substrate temperature dependence of the CH₄ and C₂H₂ mole fractions is shown in Fig. 8. At $T_{fil} = 2100$ K (where the filament is still partially graphite-coated and CH₄ concentrations are still high) some conversion of methane to acetylene occurs as the substrate temperature is raised; this may however be due to an increase in the average gas temperature, rather than direct conversion on the substrate. (The possibility that increasing the substrate temperature could affect the filament temperature was examined by monitoring the filament power, voltage drop, and brightness temperature. Since these were found to be independent of substrate temperature, we conclude that the conversion of methane to acetylene is not due to a slight increase in the filament temperature.) At $T_{fil} = 2200$ K, where the graphite layer has largely been removed from the filament, no effect of substrate temperature on acetylene is seen, and only a small decrease in the methane mole fraction is observed.

The substrate temperature has a larger effect on the near-substrate methyl radical concentration, however, as indicated in Fig. 9. The relative CH₃ concentration near the substrate is shown as a function of substrate temperature for two different flow rates (100 sccm and 400 sccm) and filament temperatures (2100 K and 2200 K). In each case, the methyl concentration rises as the substrate temperature is increased from 700 to 1000 K, then increases more slowly or begins to decrease at higher substrate temperatures. The change in the CH₃ REMPI signal observed for increasing

substrate temperature was completely reversible when the substrate temperature was lowered, and was very repeatable from run to run. Also, the results do not appear to be reactor-specific, since we found the same behavior in REMPI measurements carried out in a much-larger bell jar reactor as in the small 5-way cross.

The substrate temperature dependence of the uncorrected CH_3 signal measured near the substrate is compared in Fig. 10 with the signal observed 4 mm above the substrate. Since the gas temperature away from the substrate is not measured, no partition function correction could be applied to this data. These results indicate that the fall-off in methyl seen near the substrate at low temperatures is not observed further away. We therefore conclude that it is due to chemistry occurring in the gas phase in the immediate vicinity of the substrate, or else directly on the surface.

An Arrhenius plot of the data of Fig. 9 is shown in Fig. 11. At temperatures below about 1000 K, the data fall on reasonably straight lines, and can all be described by an activation energy of approximately 4 kcal/mole. We may compare this value to reported activation energies for the diamond growth rate. Kondoh *et al.* [16] have measured growth rates which peak at about 1200 K, and decline for higher temperatures. They report an activation energy in the temperature range of 1000–1200 K of 22–24 kcal/mole. Chu *et al.* [17] have measured the homoepitaxial growth rates on (100), (111), and (110) faces over the temperature range of 950–1250 K. They report activation energies of 8, 12, and 18 kcal/mole for homoepitaxial growth on the (100), (111), and (110) faces, respectively, although the temperature dependence appears to be larger in the range 950–1000 K, and they assign an activation energy of 50 kcal/mole in this region. This comparison indicates that the temperature dependence of the methyl radical concentration at the surface contributes to the measured activation energies for growth, but is less important than other factors.

A similar dependence of the CH_3 concentration on substrate temperature is observed for different filament-substrate separations, as seen in Fig. 12. In all cases, the rise with temperature is approximately the same below 1000 K, and a saturation or decline is seen above 1000 K. The larger absolute signal at 7 mm compared to 11 mm indicates a larger CH_3 concentration closer to the filament, in agreement with expectation and previous measurements [23], although variations in REMPI collection efficiency with filament-substrate distance may affect the relative magnitude of the signal at the three distances.

The effect of different gas compositions on the CH_3 temperature dependence is shown in Fig. 13. At low temperatures, the rate of increase of CH_3 is very similar for the three compositions shown here. However, there is a pronounced shift in the location of the peak of the CH_3 concentration as the hydrocarbon content of the gas increases. For the 5% case, CH_3 declines at higher substrate temperatures. The methyl concentration is also seen to be approximately proportional to the initial methane fraction. As the CH_4 fraction is increased to 1% from 0.5%, the measured CH_3 concentration increases by about a factor of two below 1000 K, but somewhat less than this at higher temperatures.

V. DISCUSSION

The methyl concentration at the surface is determined both by chemistry, which may occur either in the gas-phase near the substrate or on the surface, and transport from further away in the gas. To properly interpret this data requires a numerical simulation which can account for all of these effects. A detailed comparison of this data with model predictions will be presented elsewhere. Here we only indicate the most likely factors influencing the methyl concentration at the substrate.

Modeling studies [18, 19] have shown that away from the substrate the methyl concentration under typical filament-assisted conditions is set by the partial equilibrium of the rapid reaction



The equilibrium constant for this reaction is nearly independent of temperature. Thus, if this reaction controls the CH_3 concentration, we have that $[\text{CH}_3] \propto [\text{CH}_4][\text{H}]$.

The observed results could be explained if reaction (2) is in partial equilibrium near the substrate and the atomic hydrogen concentration decreases at low temperatures. However, both simulations [19] and measurements [34] show that the partial equilibrium of reaction (2) breaks down near the surface. Also, the atomic hydrogen concentration at the surface is governed by a balance between diffusion to the surface and surface recombination; a lower H concentration at low substrate temperatures would require that the surface recombination coefficient of H on diamond be larger at low temperatures than at high temperatures. This behavior would be contrary to the known temperature dependence of H recombination on similar surfaces, such as quartz and glass [35, 36].

A second type of gas-phase reaction which might account for the present results is a radical-radical recombination reaction, such as



or



Under typical filament-assisted growth conditions, methyl consumption by reaction (4) is much more rapid than by reaction (3), due to the large H atom superequilibrium, and in fact reaction 4 is the dominant homogeneous loss mechanism for CH_3 . Since the rate for this reaction has a negative temperature exponent (faster at lower temperatures) [37], it could account for the observed methyl depletion at low substrate temperatures.

It is also possible that CH_3 could recombine heterogeneously on the diamond surface. Also, methyl is known to be a primary etch product of H-attack on amorphous carbon [38, 39]. Etching of non-diamond carbon by atomic hydrogen is widely thought

to be an essential aspect of CVD diamond growth, resulting in suppression of non-diamond phases. Since etching of carbon films by atomic hydrogen typically has a small positive activation energy (e.g., 5.3 ± 1.9 kcal/mole [40]), if this process occurs rapidly enough it could result in production of methyl at higher temperatures, leading to a temperature dependence similar to that observed here.

VI. SUMMARY

The filament- and substrate-temperature dependence of the CH_3 , CH_4 , and C_2H_2 concentrations near the substrate during filament-assisted diamond growth has been determined. The dependence on filament temperature is consistent with previous studies, and clearly shows that a reduction in filament activity at low temperatures occurs even for the dilute 0.5% CH_4 in H_2 gas composition. The substrate temperature studies show that CH_4 and C_2H_2 are relatively independent of substrate temperature at sufficiently high filament temperatures. On the other hand, CH_3 near the substrate exhibits a distinct temperature dependence which is qualitatively similar for all gas compositions, filament-substrate separations, flow rates, and filament temperatures. The most likely explanation for this behavior is recombination of CH_3 to CH_4 at low substrate temperatures, either through homogeneous chemistry occurring near the substrate, or heterogeneous chemistry on the substrate. These results suggest that depletion of growth species due to recombination may need to be considered in designing low-temperature processes for CVD diamond growth.

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FIGURES

FIG. 1. (a) Schematic of experimental setup. (b) close-up showing filament, substrate, ionization probe, laser beam, and biasing scheme.

FIG. 2. Typical survey ionization spectrum. Dashed line shows laser energy variation.

FIG. 3. Dependence of CH_3 REMPI signal on probe bias voltage for three substrate temperatures.

FIG. 4. Dependence of the Ar $4s'$ (3+1) REMPI signal on substrate temperature. The dashed curve shows a fit to a T^{-1} variation expected due to ideal gas expansion.

FIG. 5. Rotational manifold of the CH_3 REMPI Q branch. The solid curve shows the measured profile at a substrate temperature of 800 K, and the dashed curve shows the profile at 1200 K. Both curves are normalized to equal area, and some smoothing has been done to allow a better comparison between the curves.

FIG. 6. Measured CH_3 REMPI signal (arbitrary units) and CH_4 and C_2H_2 mole fractions vs. filament brightness temperature for 0.5% CH_4 in H_2 and $T_s = 1000\text{K}$.

FIG. 7. Methane and acetylene mole fractions vs. filament temperature for two different gas flow rates. Solid symbols: CH_4 mole fraction; open symbols: C_2H_2 mole fraction; squares: 100 sccm; circles: 400 sccm.

FIG. 8. Substrate temperature dependence of CH_4 and C_2H_2 . (a) $T_{fil} = 2100\text{K}$; (b) $T_{fil} = 2200\text{K}$. Circles: C_2H_2 ; triangles: CH_4 .

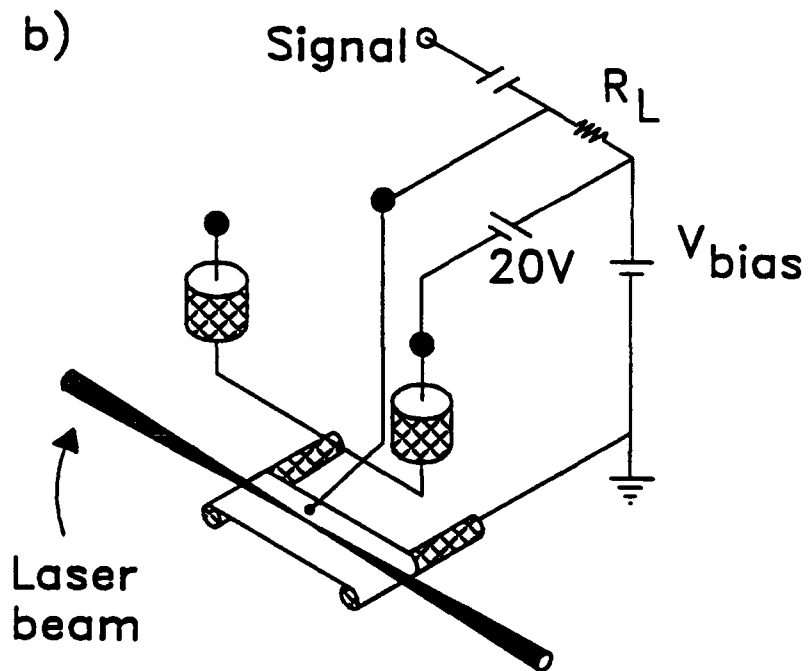
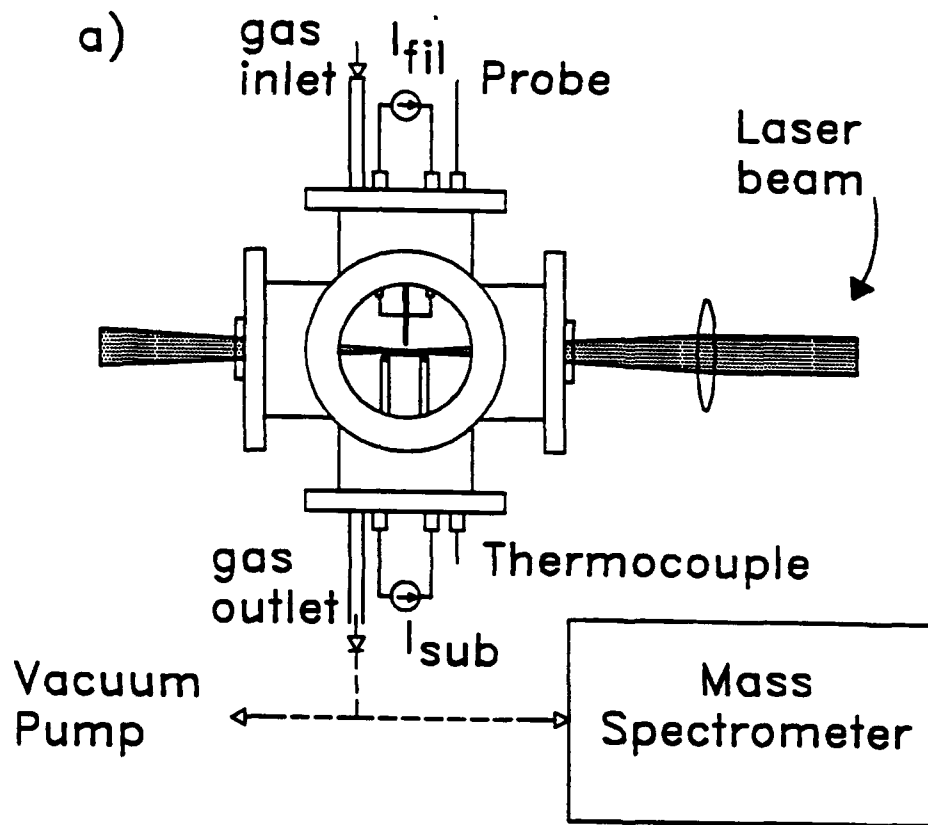
FIG. 9. Dependence of the CH_3 relative concentration on substrate temperature. Triangles: $T_{fil} = 2200$ K, 400 sccm; circles: $T_{fil} = 2200$ K, 100 sccm; squares: $T_{fil} = 2100$ K, 400 sccm.

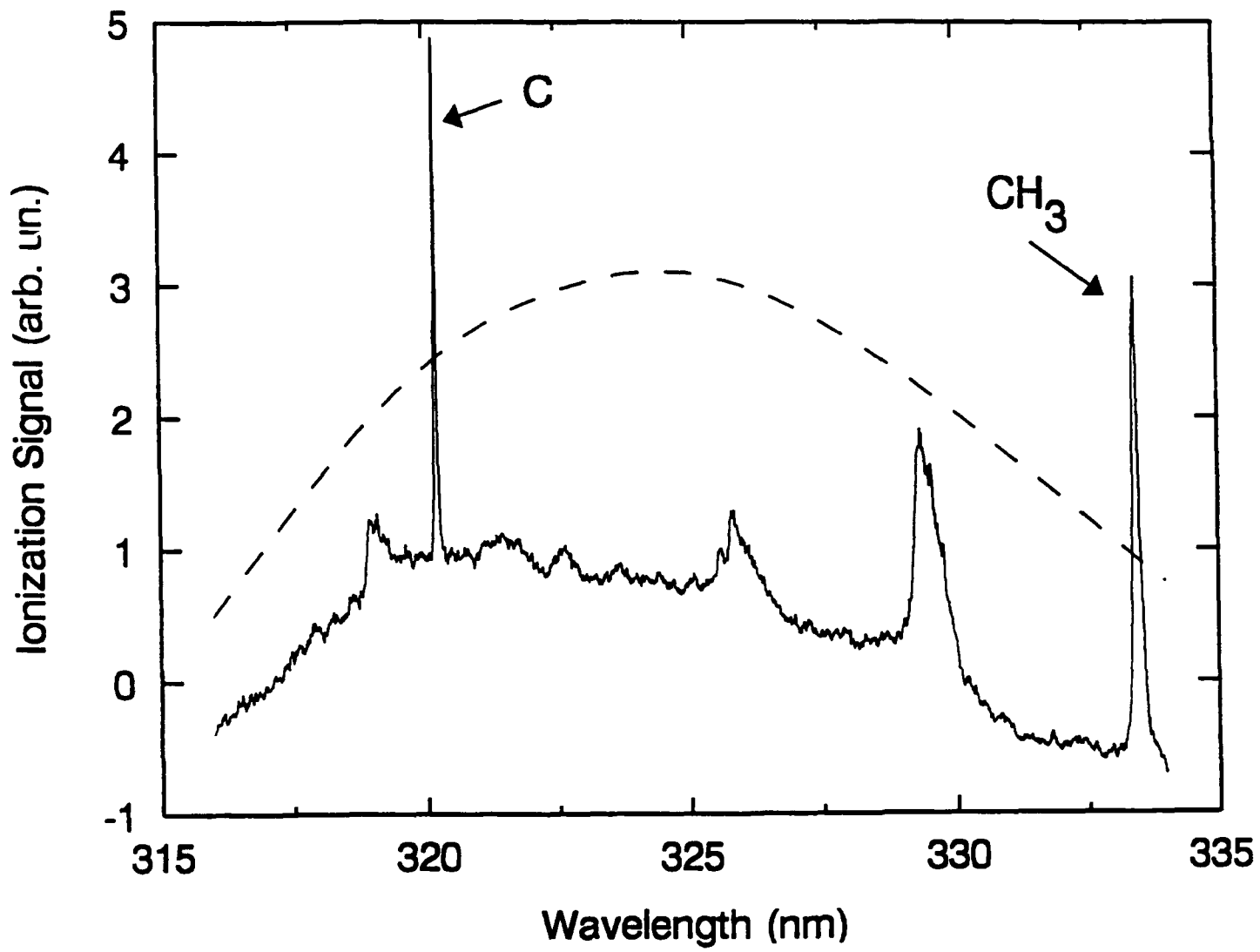
FIG. 10. Comparison of uncorrected REMPI signal vs. substrate temperature just above substrate to that 4 mm away.

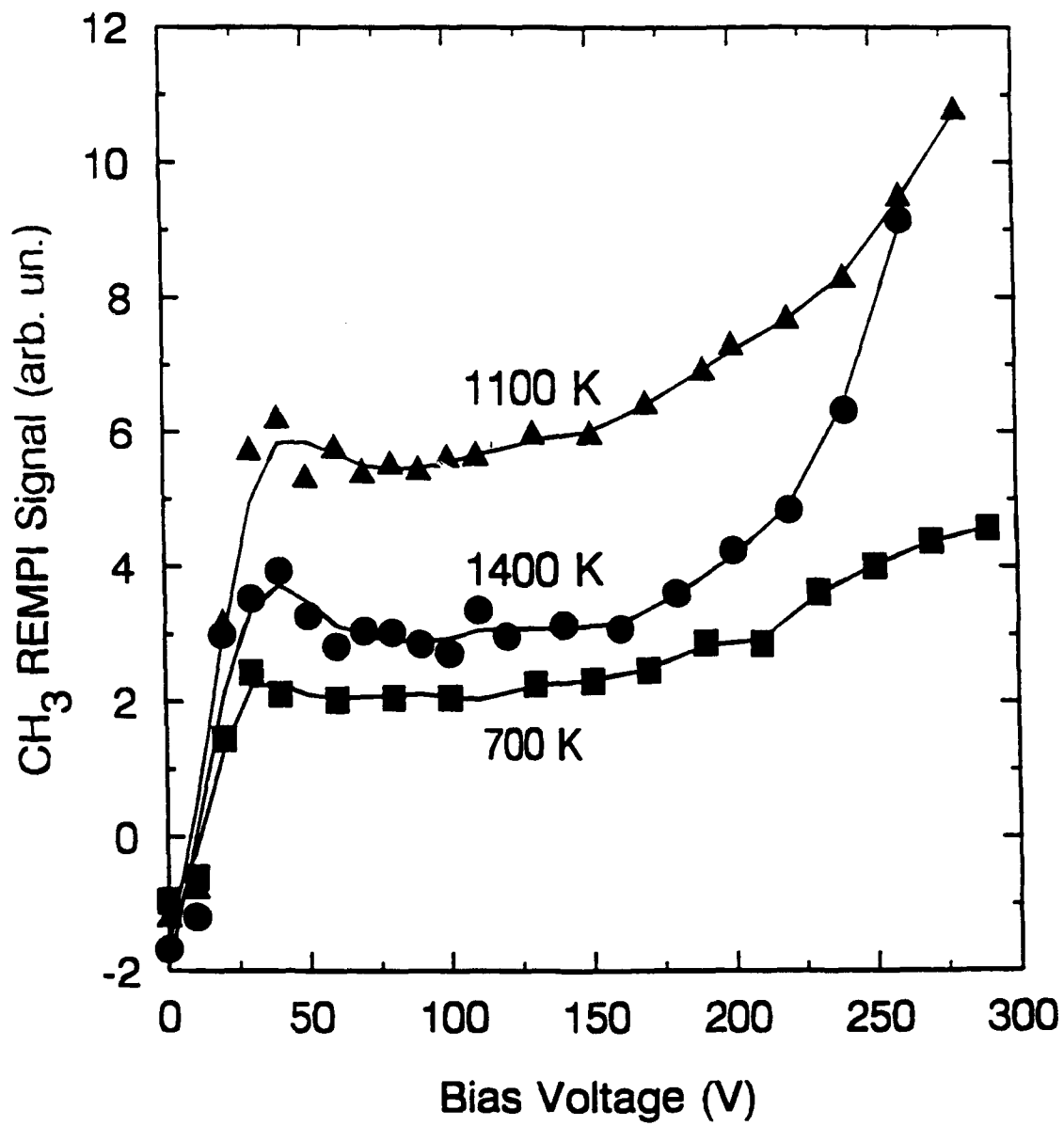
FIG. 11. Arrhenius plot of data in Figure 9. The low temperature fall-off of CH_3 may be characterized by an activation energy of approximately 4 kcal/mole.

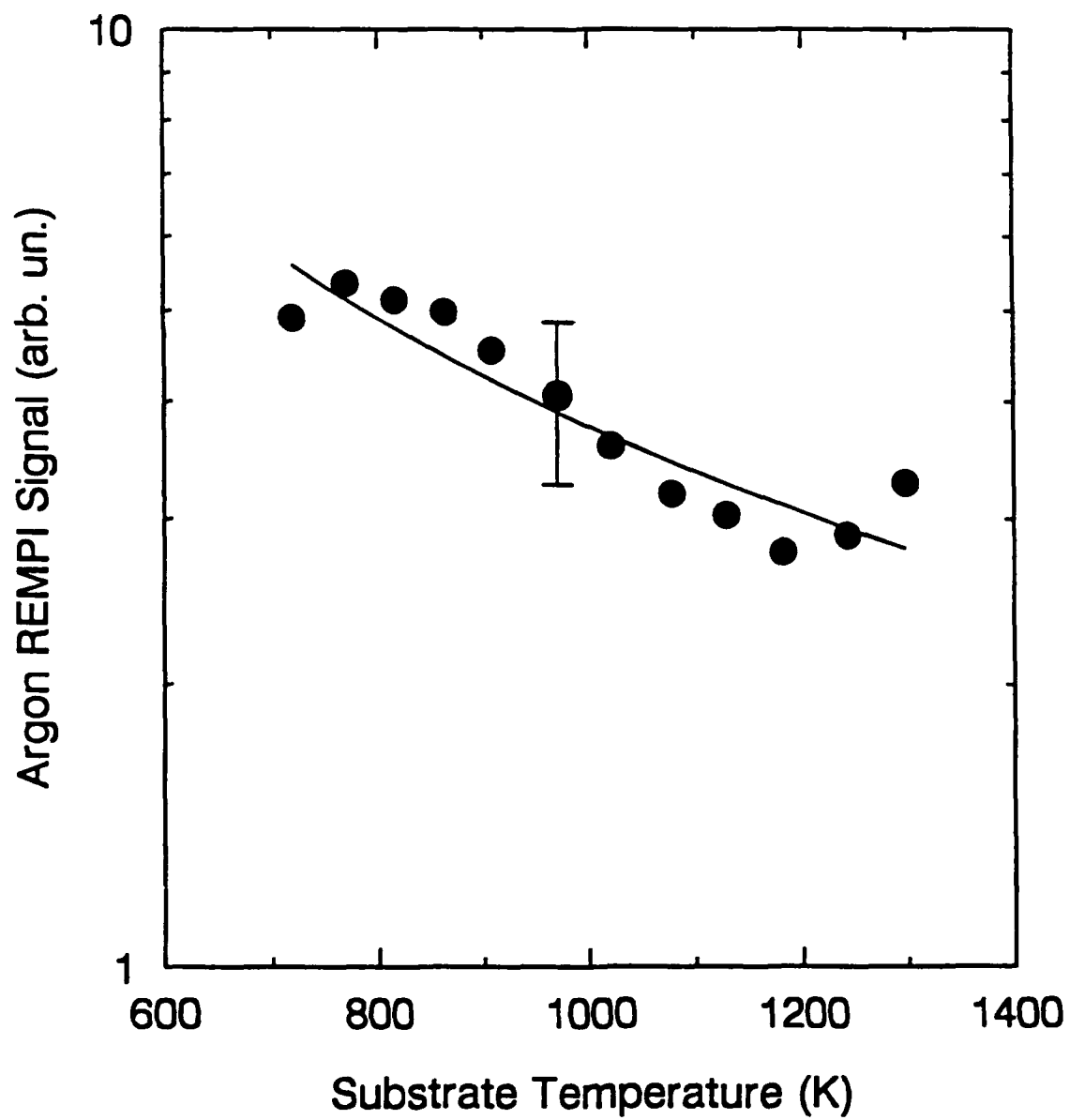
FIG. 12. Relative CH_3 concentration at the substrate vs. substrate temperature for three different substrate to filament distances.

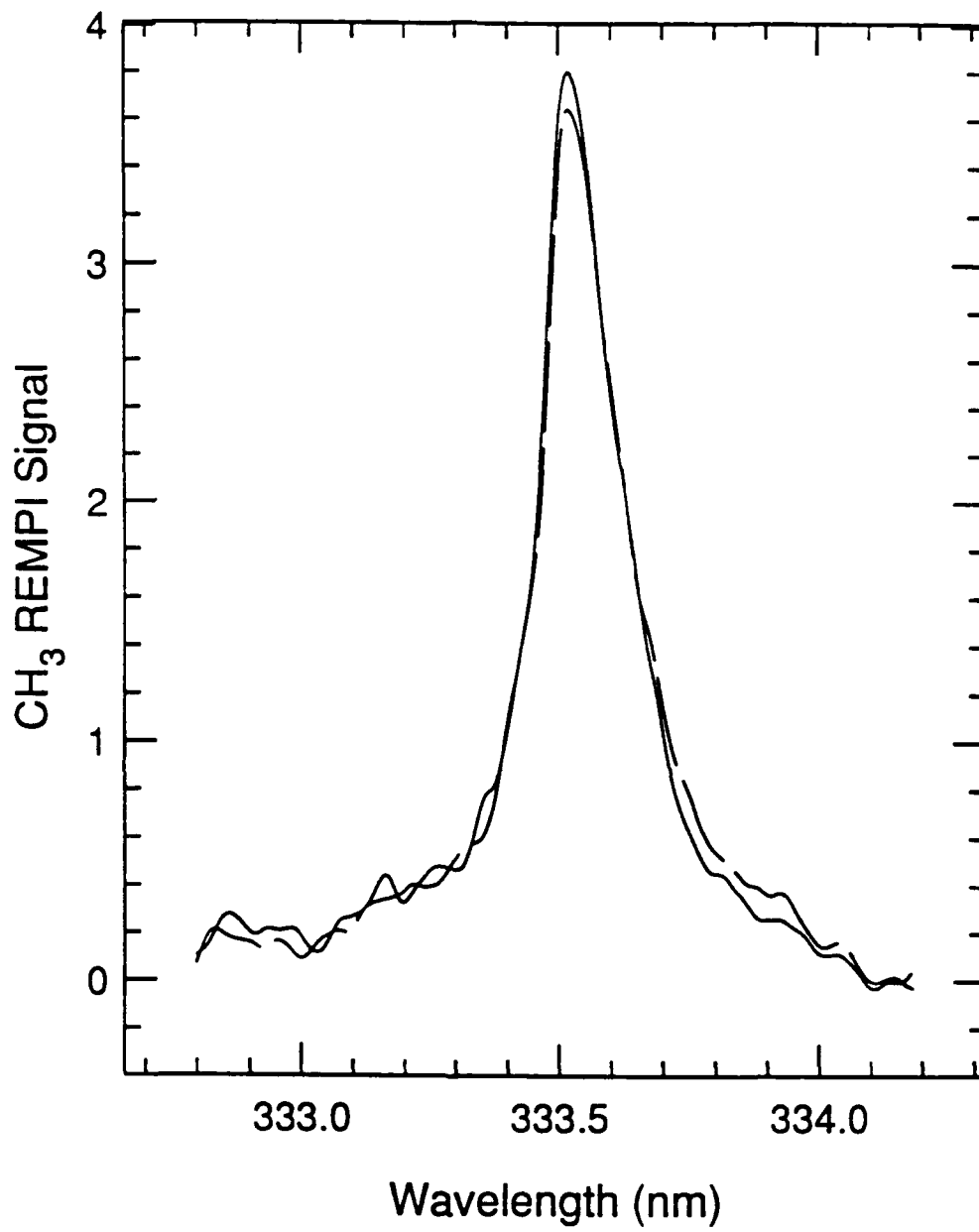
FIG. 13. Relative CH_3 concentration at the substrate vs. substrate temperature for three different methane percentages in the initial gas mixture. The 5% curve has been scaled to fit on the same plot.

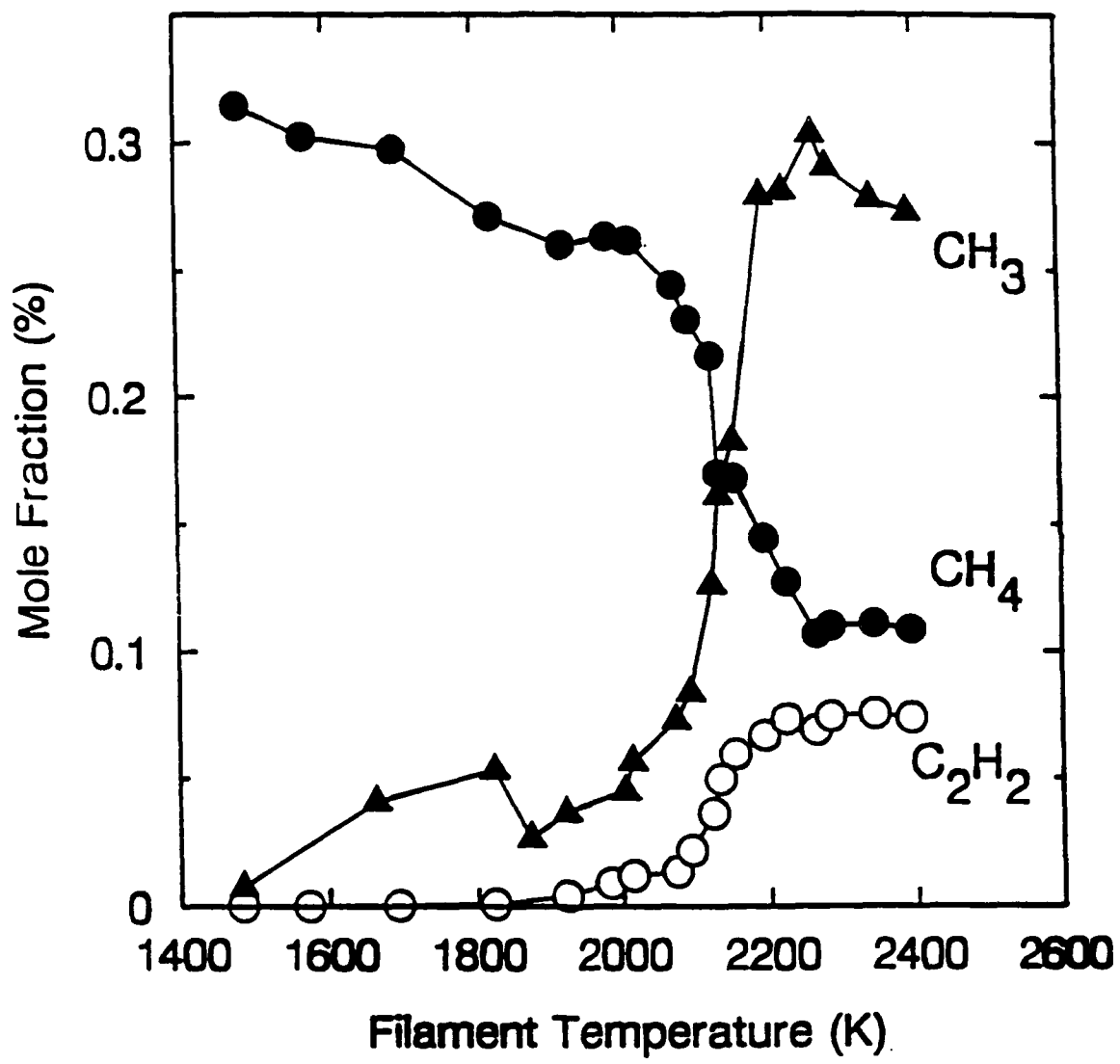


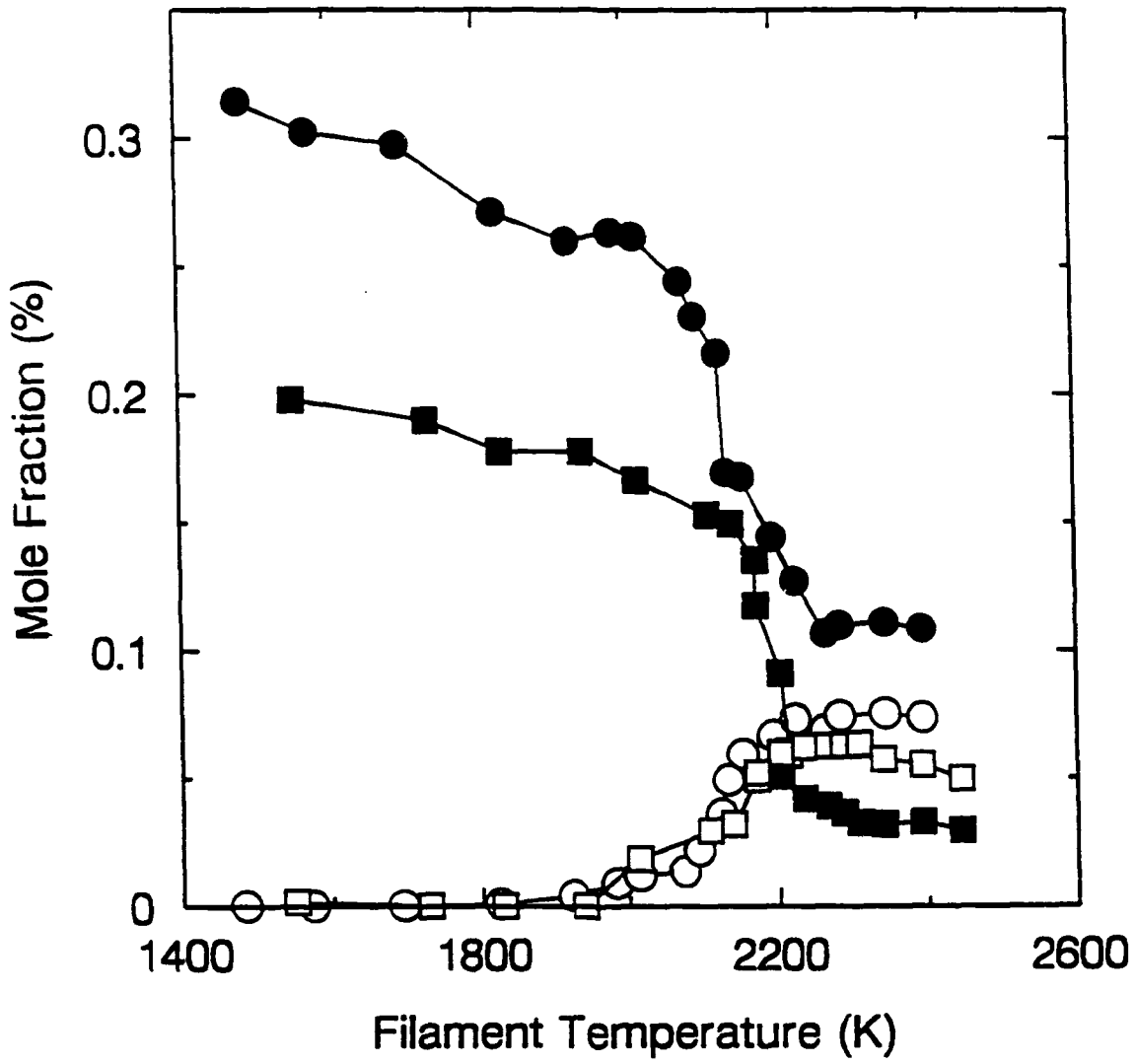


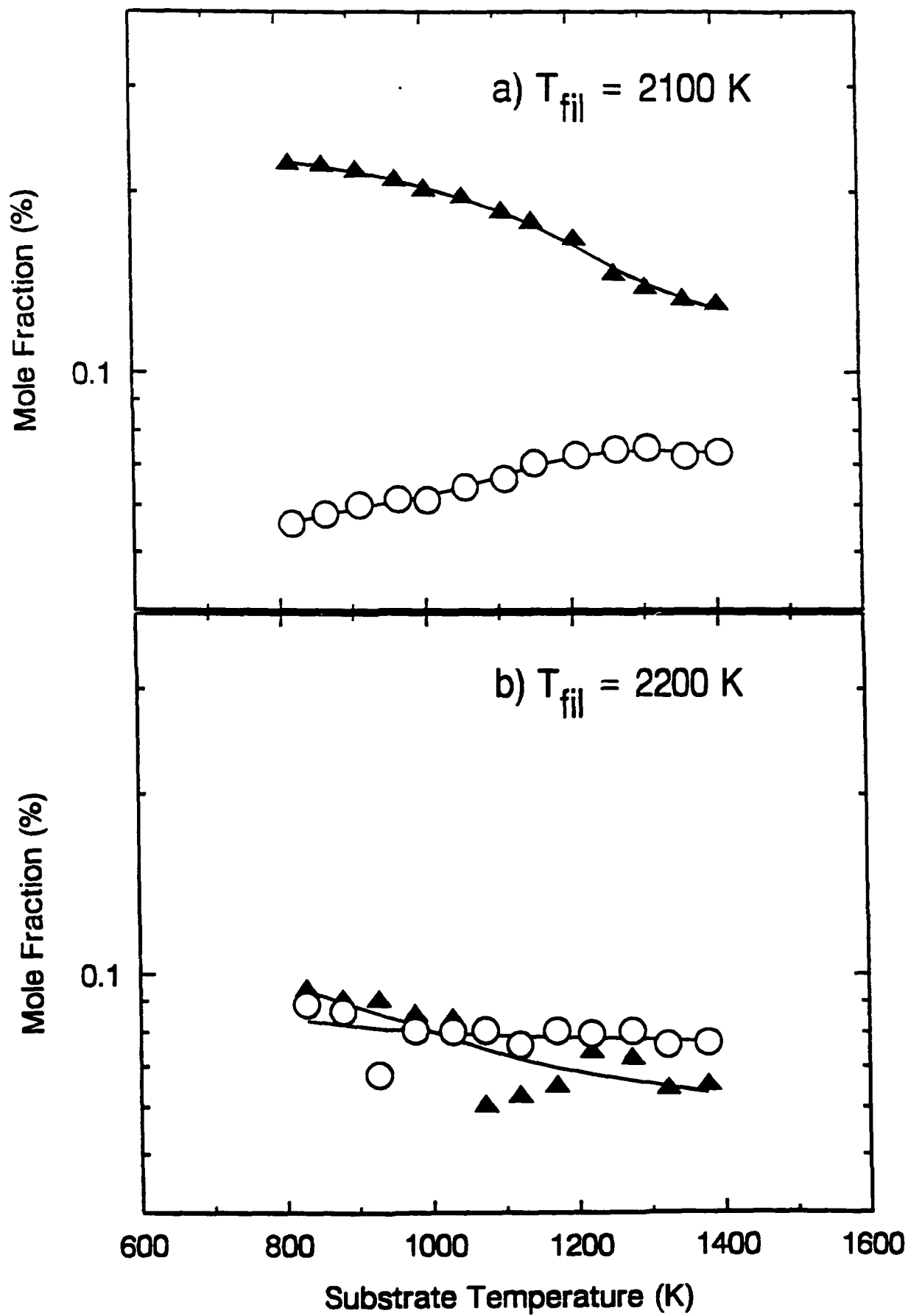


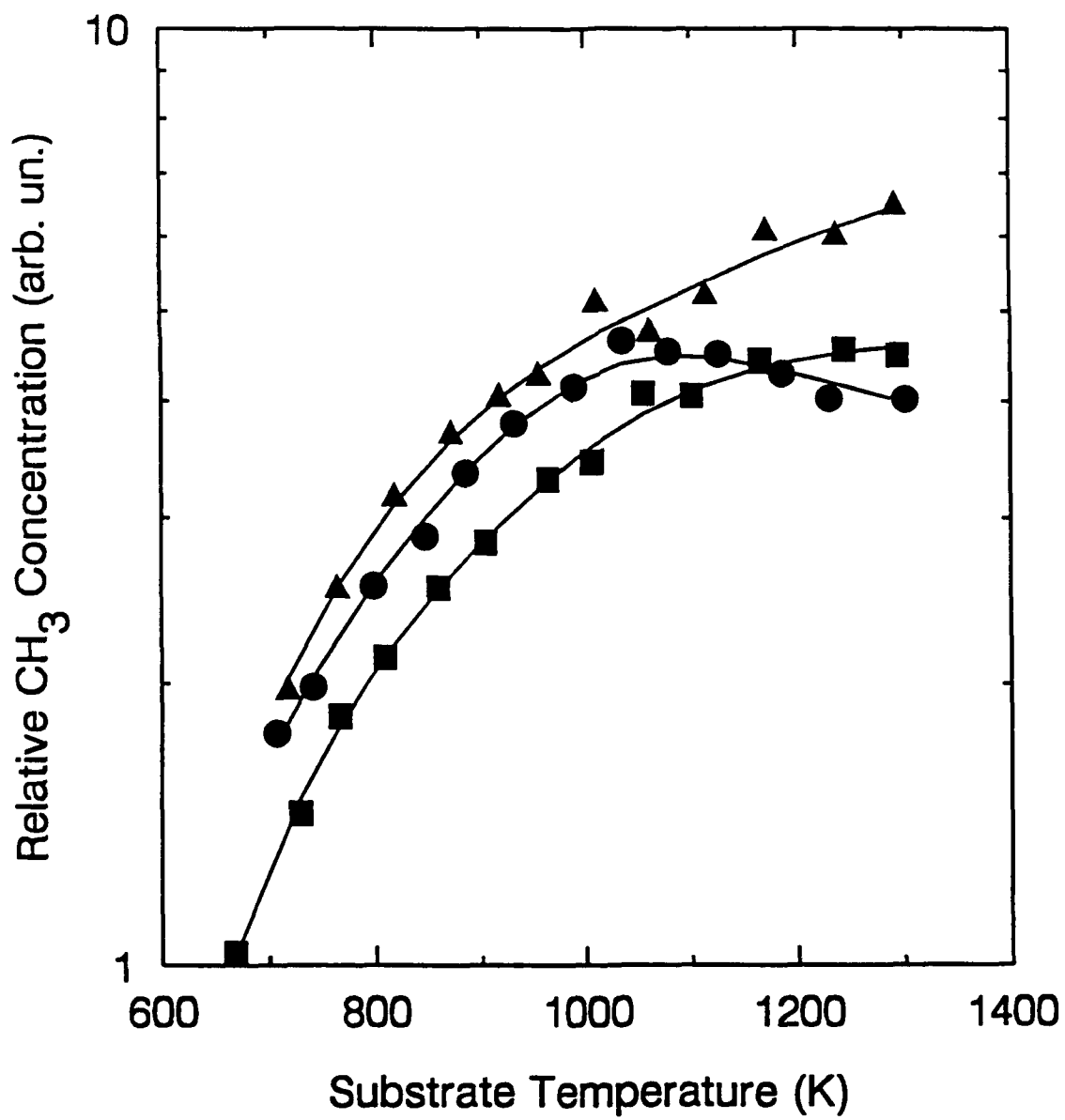


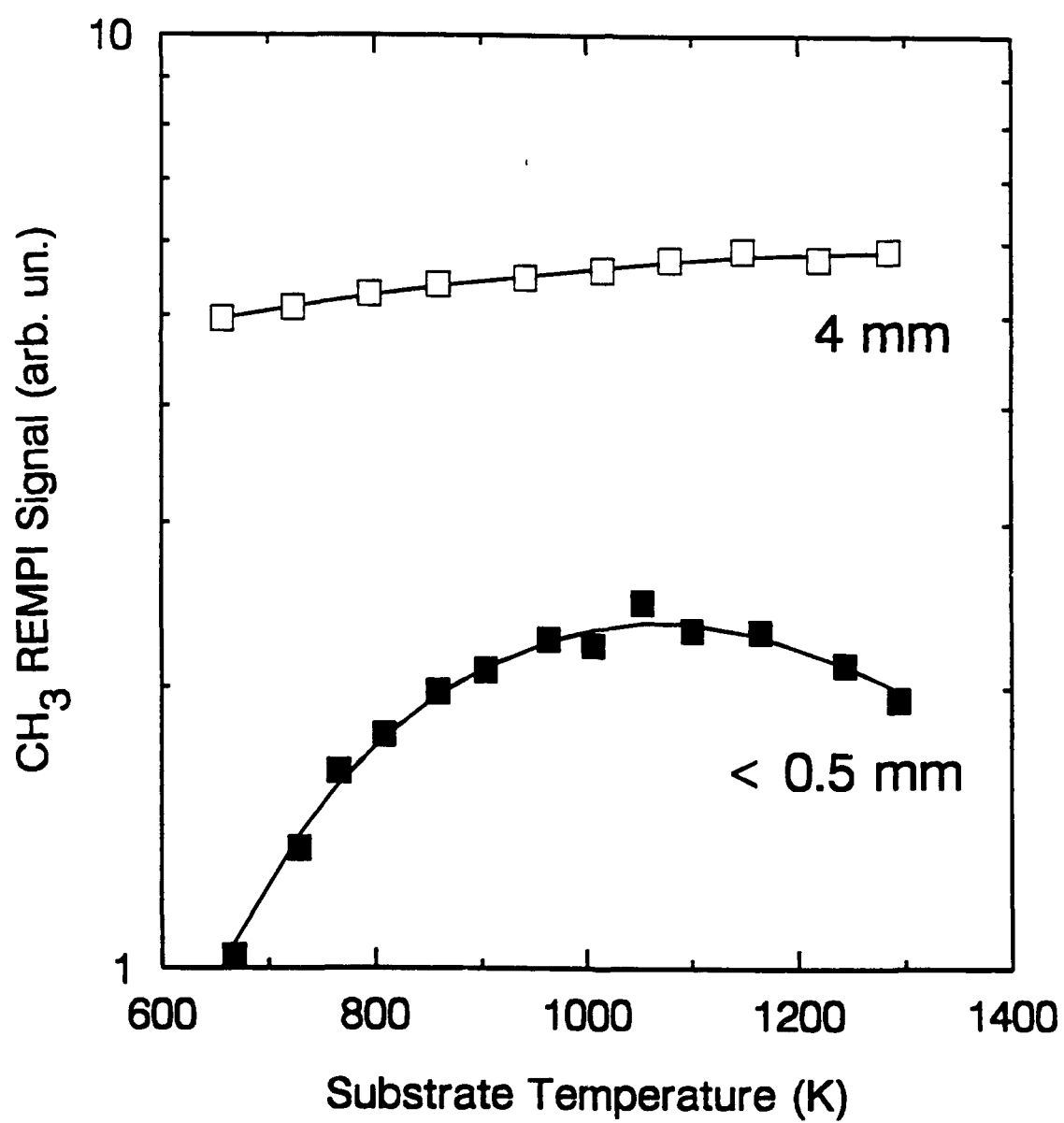


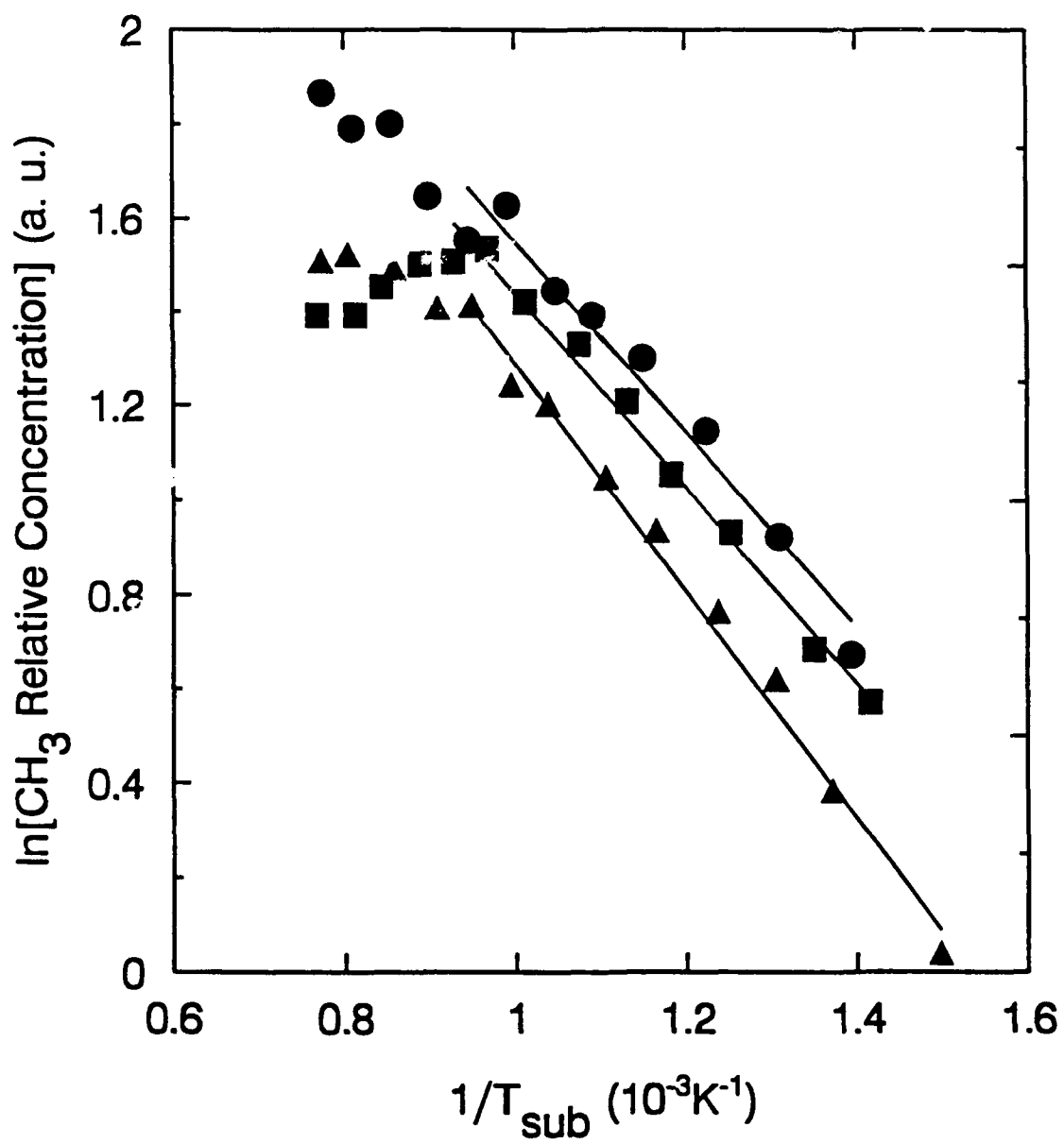


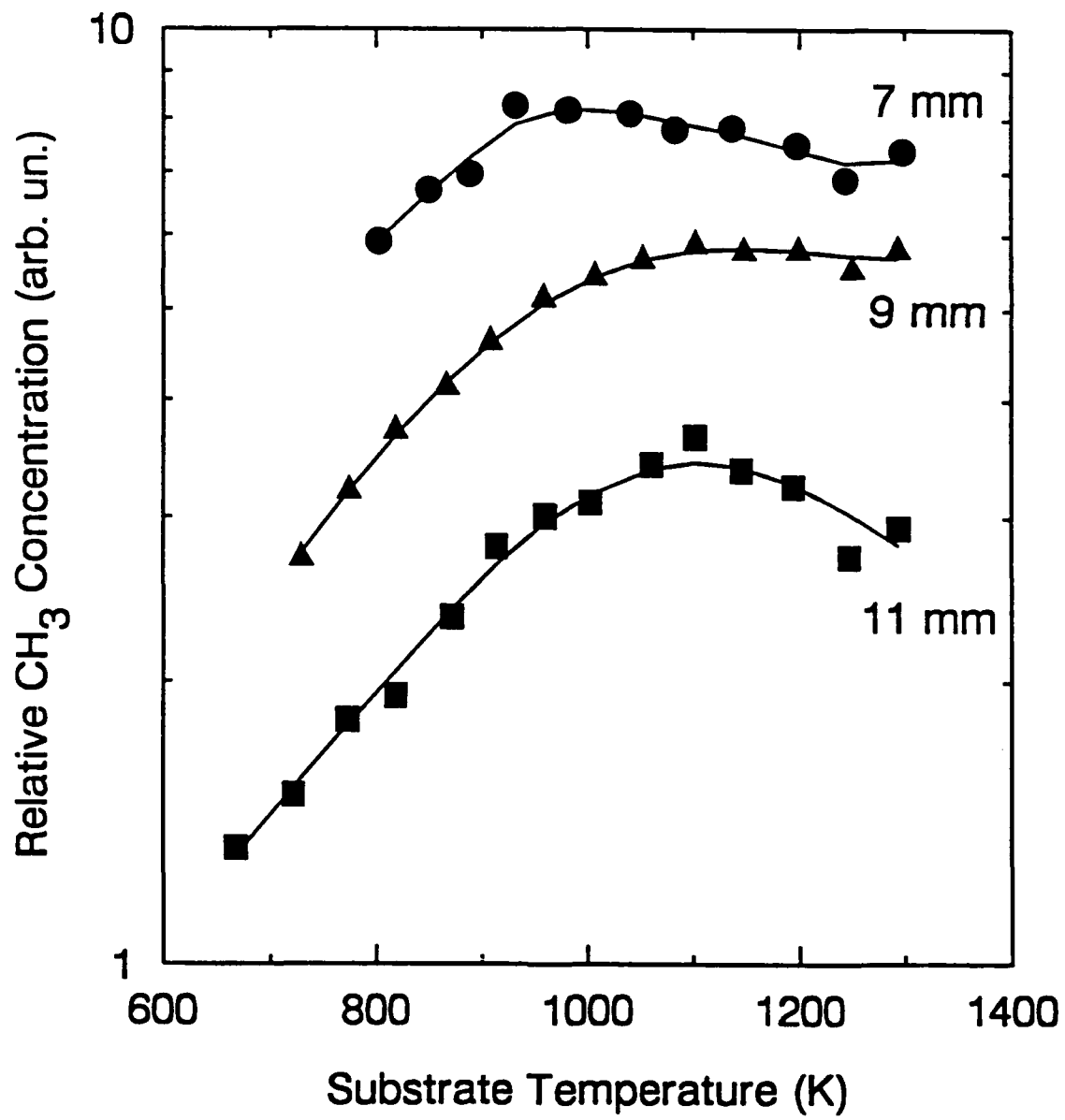


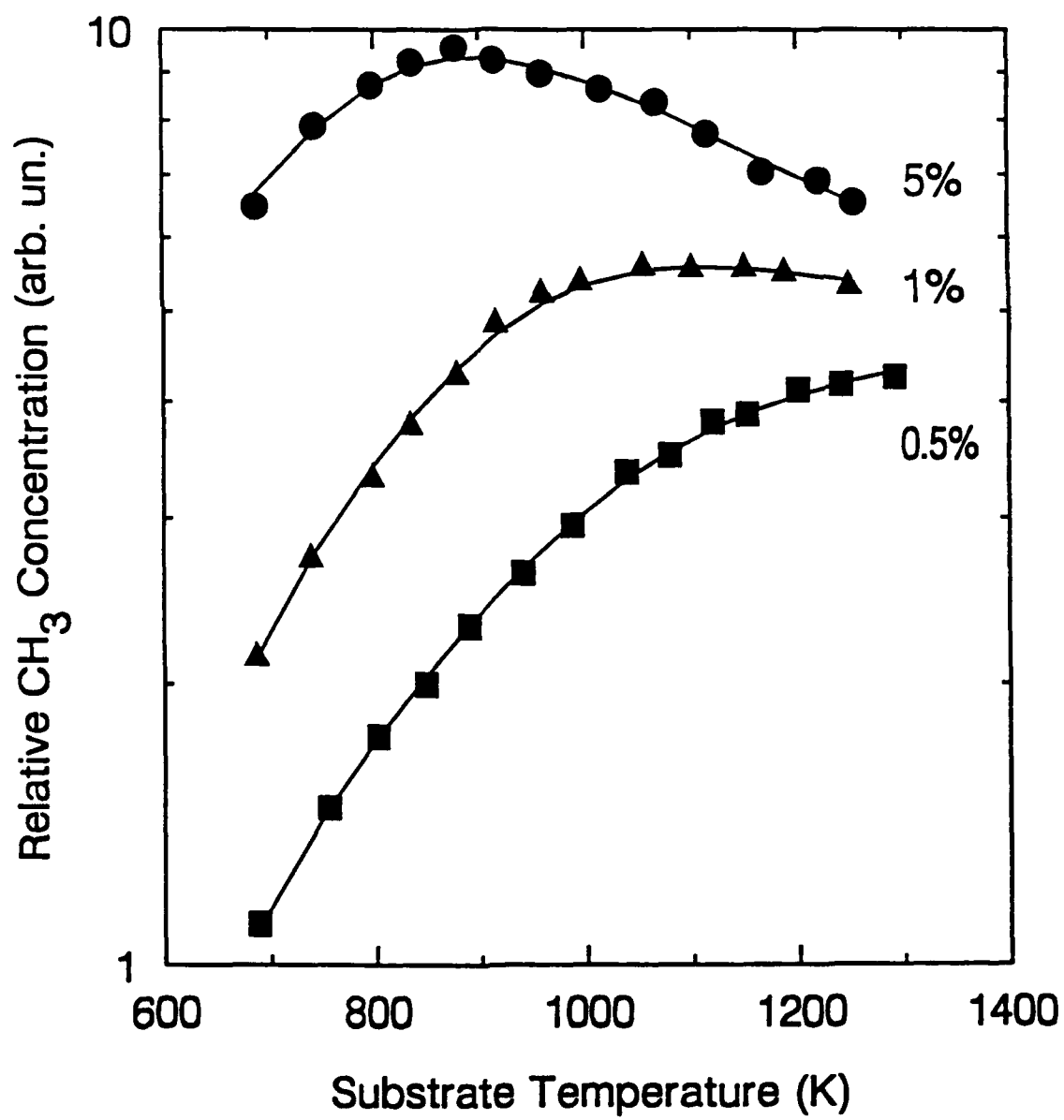












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