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Anomalous Diffraction Approximation Limits

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ARL-TN-128

November 1998

19981123 155

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Adelphi, MD 20783-1197

ARL-TN-128

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Abstract

It has been reported in a recent article that the anomalous diffraction approximation (ADA) accuracy does not depend on particle refractive index, but instead is dependent on the particle size parameter. Since this is at odds with previous research, we thought these results warranted further discussion.

The anomalous diffraction approximation (ADA) of van de Hulst (1981) provides a method by which gross scattering properties (scattering efficiencies and albedo) can be rapidly obtained. The primary assumption used to derive the approximation is that the scattering particle is soft; i.e., $|m - 1| \ll 1$. In a recent article, Liu, Jonas, and Saunders (1996) reported that "the ADA accuracy depends mainly on the particle size parameter and is not sensitive to the condition of $|m - 1| \ll 1$." Since this is at odds with several recently published studies (Mitchell and Arnott, 1994; Ackerman and Stephens, 1987; Evans and Fournier, 1996; Chýlek and Klett, 1991; and Chýlek and Videen, 1994), we felt that this statement warranted further clarification. There are two points to consider: First, the ADA produces more accurate results in the geometrical-optics (short-wavelength) limit. And second, the accuracy is independent of the particle refractive index (as proposed by Liu, Jonas, and Saunders, 1996).

The first point (the accuracy increases in the geometrical-optics limit) is not surprising. It is well-known that for large particles, the extinction efficiency approaches 2. This is illustrated in figure 1(a), which shows the extinction efficiencies plotted as a function of size parameter $x = 2\pi r/\lambda$ for a sphere of radius r . The extinction approaches zero as the radius approaches zero. As the radius increases, bringing the sphere into the resonance region, structure appears in the Mie extinction curves, which is beyond the capacity of the ADA to replicate. As the radius further increases, the oscillations gradually die, approaching the final, geometrical-optics limit. Since the ADA also approaches the proper, geometrical-optics limit, it is no surprise that the extinction efficiency accuracy increases. The accuracy of the absorption efficiency can similarly be explained. As the particle size increases, any light incident upon the particle will be absorbed (assuming nonzero absorption and a soft particle). The absorption efficiency must therefore approach unity for an absorbing soft particle. The Mie and ADA absorption efficiencies are shown as a function of sphere size parameter in figure 1(b). For the soft particle ($m = 1.1 + 0.01i$) having some absorption, the absorption approaches unity in the geometrical-optics limit.

The peculiar aspect of the investigation by Liu, Jonas, and Saunders (1996) is their claim of the lack of any accuracy dependence on the refractive index. This aspect can be understood when we consider the refractive indices of the particles in their study. They concentrated on ice particles through much of the ultraviolet (UV), visible, and infrared (IR) spectra. The value of the ice refractive index changes drastically throughout this range (as demonstrated in fig. 1 of Liu et al's report). However, for only a

couple (relatively narrow) spectral bands, the condition $|m - 1| \ll 1$ holds ($\lambda \sim 2.8 \mu\text{m}, 10 \mu\text{m}$). In these bands, the deviations of the ADA results from those given by Mie theory are less than 10 percent; whereas, in the other regions, the errors are typically much greater than 10 percent and sometimes even over 100 percent of the actual value. Unfortunately, these small regions of applicability only represent a small percentage of the entire spectrum and apparently went unnoticed in their analysis. As illustrated in figure 2, a strong accuracy dependence exists on particle refractive index as long as the particle remains sufficiently soft. Figure 2 shows the percent errors in the (a) extinction and (b) absorption efficiencies as a function of size parameter x for three different refractive index values. For real refractive index $m_r = 1.1$, the errors decrease in the small- and large-wavelength regions and can be quite substantial (10 to 30 percent) in the resonance region. Figure 2 illustrates that as m_r is further increased, the percent of errors become increasingly large as to be impracticable. Indeed, this is because of the particles themselves being well outside the bounds on refractive index for which the approximation was derived. The ice refractive index has an even wider range of values than is illustrated in figure 2. The associated errors resulting from the ADA calculations are so large that they become meaningless.

Figure 1. Mie and ADA
 (a) extinction and
 (b) absorption efficiencies
 as a function of sphere size
 parameter for three
 different refractive indices.
 Note that the three ADA
 absorption efficiencies
 overlap for spheres having
 the same imaginary part of
 the refractive index.

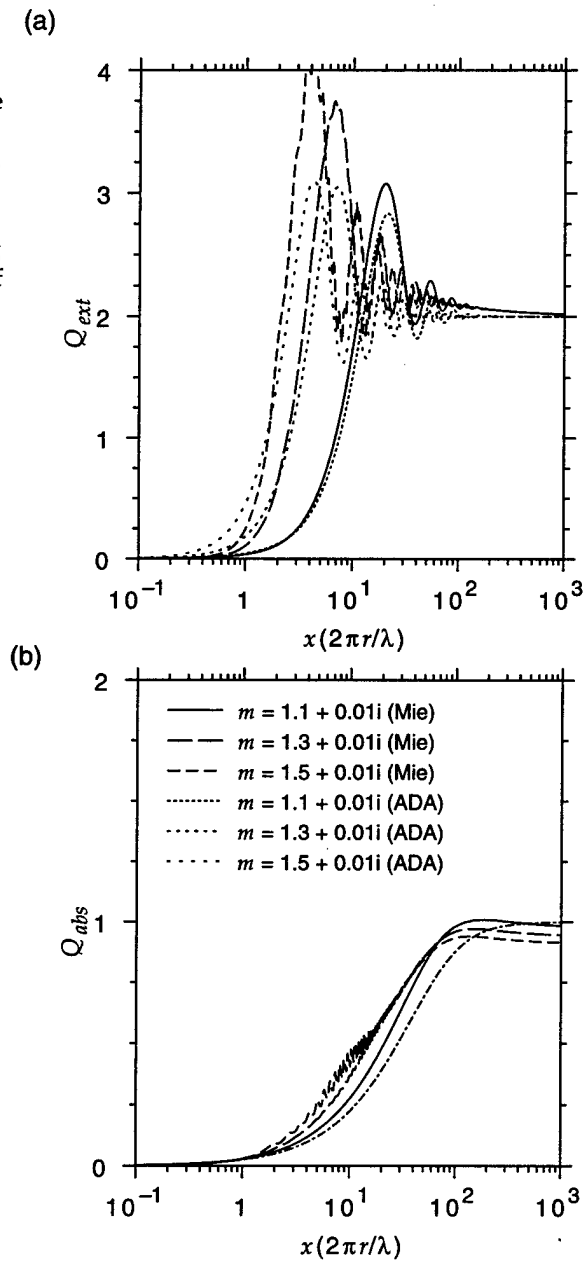
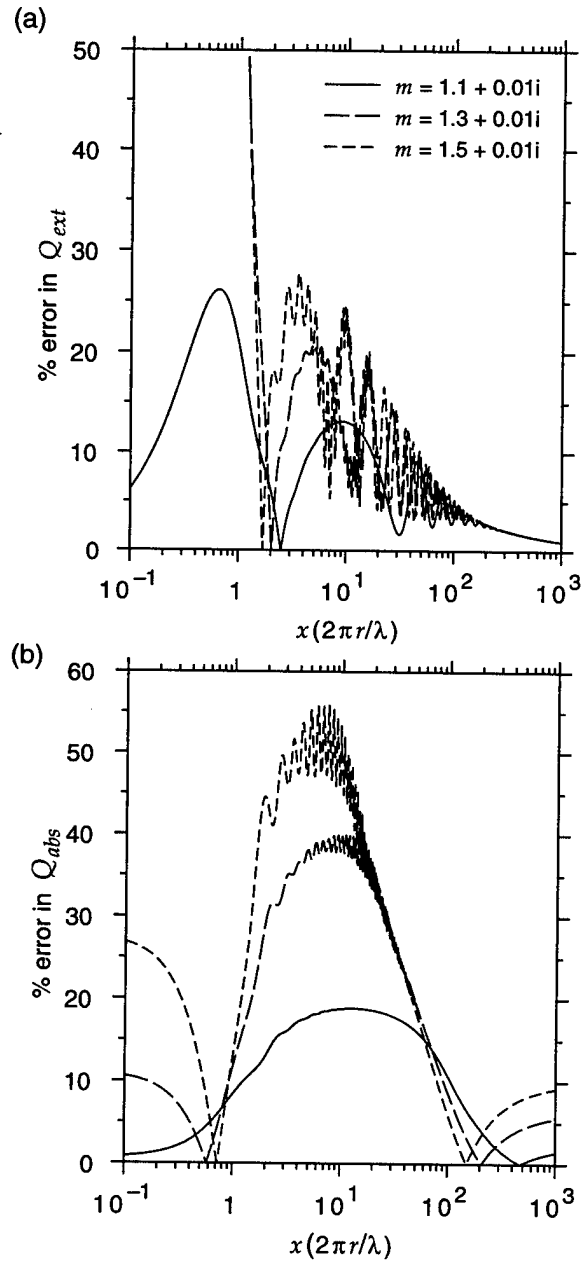


Figure 2. Percent error in (a) Q_{ext} and (b) Q_{abs} defined as the absolute value of the difference between the Mie and ADA extinction efficiencies divided by the Mie extinction efficiency.



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REPORT DOCUMENTATION PAGE			<i>Form Approved OMB No. 0704-0188</i>	
Public reporting burden for this collection of information is estimated to average 1 hour per response, including the time for reviewing instructions, searching existing data sources, gathering and maintaining the data needed, and completing and reviewing the collection of information. Send comments regarding this burden estimate or any other aspect of this collection of information, including suggestions for reducing this burden, to Washington Headquarters Services, Directorate for Information Operations and Reports, 1215 Jefferson Davis Highway, Suite 1204, Arlington, VA 22202-4302, and to the Office of Management and Budget, Paperwork Reduction Project (0704-0188), Washington, DC 20503.				
1. AGENCY USE ONLY (Leave blank)		2. REPORT DATE November 1998	3. REPORT TYPE AND DATES COVERED Final, 1 Oct 97 to 1 June 98	
4. TITLE AND SUBTITLE Anomalous Diffraction Approximation Limits			5. FUNDING NUMBERS DA PR: B53A PE: 61102A	
6. AUTHOR(S) Gordon Videen (ARL), Petr Chýlek (Department of Physics and Oceanography, Dalhousie University)				
7. PERFORMING ORGANIZATION NAME(S) AND ADDRESS(ES) U.S. Army Research Laboratory Attn: AMSRL-IS-EE email: videen@atm.dal.ca 2800 Powder Mill Road Adelphi, MD 20783-1197			8. PERFORMING ORGANIZATION REPORT NUMBER ARL-TN-128	
9. SPONSORING/MONITORING AGENCY NAME(S) AND ADDRESS(ES) U.S. Army Research Laboratory 2800 Powder Mill Road Adelphi, MD 20783-1197			10. SPONSORING/MONITORING AGENCY REPORT NUMBER	
11. SUPPLEMENTARY NOTES ARL PR: 7FEJ70 AMS code: 61110253A11				
12a. DISTRIBUTION/AVAILABILITY STATEMENT Approved for public release; distribution unlimited.			12b. DISTRIBUTION CODE	
13. ABSTRACT (Maximum 200 words) It has been reported in a recent article that the anomalous diffraction approximation (ADA) accuracy does not depend on particle refractive index, but instead is dependent on the particle size parameter. Since this is at odds with previous research, we thought these results warranted further discussion.				
14. SUBJECT TERMS scattering, ice crystals			15. NUMBER OF PAGES 12	
			16. PRICE CODE	
17. SECURITY CLASSIFICATION OF REPORT Unclassified	18. SECURITY CLASSIFICATION OF THIS PAGE Unclassified	19. SECURITY CLASSIFICATION OF ABSTRACT Unclassified	20. LIMITATION OF ABSTRACT UL	