



AFRL-ML-WP-TP-2007-523

**SELF-PUMPED PHOTOREFRACTIVE GRATINGS IN
Fe:KNbO₃ (PREPRINT)**

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**Hardened Materials Branch
Survivability and Sensor Materials Division**

JANUARY 2006

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REPORT DOCUMENTATION PAGE				<i>Form Approved</i> OMB No. 0704-0188	
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1. REPORT DATE (DD-MM-YY) January 2006		2. REPORT TYPE Conference Paper Preprint		3. DATES COVERED (From - To)	
4. TITLE AND SUBTITLE SELF-PUMPED PHOTOREFRACTIVE GRATINGS IN Fe:KNbO ₃ (PREPRINT)				5a. CONTRACT NUMBER In-house	
				5b. GRANT NUMBER	
				5c. PROGRAM ELEMENT NUMBER 62102F	
6. AUTHOR(S) Dean R. Evans (AFRL/MLPJ) G. Cook (Universal Technology Corporation) M.A. Saleh (UES, Inc.) J.L. Carns (General Dynamics Information Technology, Inc.)				5d. PROJECT NUMBER 4348	
				5e. TASK NUMBER RG	
				5f. WORK UNIT NUMBER M08R1000	
7. PERFORMING ORGANIZATION NAME(S) AND ADDRESS(ES) Hardened Materials Branch (AFRL/MLPJ) Survivability and Sensor Materials Division Materials and Manufacturing Directorate Wright-Patterson Air Force Base, OH 45433-7750 Air Force Materiel Command, United States Air Force ----- Universal Technology Corporation 1270 North Fairfield Road Dayton, OH 45432				8. PERFORMING ORGANIZATION REPORT NUMBER AFRL-ML-WP-TP-2007-523	
9. SPONSORING/MONITORING AGENCY NAME(S) AND ADDRESS(ES) Air Force Research Laboratory Materials and Manufacturing Directorate Wright-Patterson Air Force Base, OH 45433-7750 Air Force Materiel Command United States Air Force				10. SPONSORING/MONITORING AGENCY ACRONYM(S) AFRL/MLPJ	
				11. SPONSORING/MONITORING AGENCY REPORT NUMBER(S) AFRL-ML-WP-TP-2007-523	
12. DISTRIBUTION/AVAILABILITY STATEMENT Approved for public release; distribution unlimited.					
13. SUPPLEMENTARY NOTES Conference slideshow submitted to the Proceedings of the Third International Photorefractive Workshop. The U.S. Government is joint author of this work and has the right to use, modify, reproduce, release, perform, display, or disclose the work. PAO Case Number: AFRL/WS 06-0322, 07 Feb 2006.					
14. ABSTRACT <ul style="list-style-type: none"> • High gain confirmed in off-axis geometries for Fe:KNbO₃ • Mismatch between theory and experiment for mid-range crystal angles, especially for the a-c plane • Large apparent variation in the effective trap density with crystal angle • Modified theory gives a good fit to experimental data • Mechanism for trap density anisotropy is unclear 					
15. SUBJECT TERMS Photorefractive, Potassium Niobate, Two-Beam Coupling, Focal Plane					
16. SECURITY CLASSIFICATION OF:			17. LIMITATION OF ABSTRACT: SAR	18. NUMBER OF PAGES 20	19a. NAME OF RESPONSIBLE PERSON (Monitor) Dean R. Evans 19b. TELEPHONE NUMBER (Include Area Code) N/A
a. REPORT Unclassified	b. ABSTRACT Unclassified	c. THIS PAGE Unclassified			

Self-Pumped Photorefractive Reflection Gratings in Fe:KNbO_3



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Outline



- Potassium niobate as a photorefractive
- Theory
- Off-axis geometries
- Experiments
- Results
- A controversial suggestion.....
- Discussion
- Conclusion (confusion?)



Potassium niobate as a photorefractive



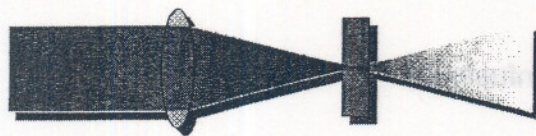
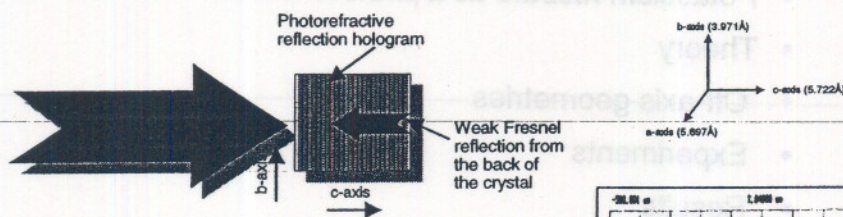
- High trap density
 - ↳ Allows efficient counter-propagating gratings to be written
- High sensitivity
 - ↳ Fast response times
- Broad spectral response
 - ↳ 400nm - ~700nm (with Fe doping)
 - ↳ 400nm - >700nm? (with Ni doping)
- Difficult to grow reproducibly
 - ↳ Program under way to fix this (looks very promising!)



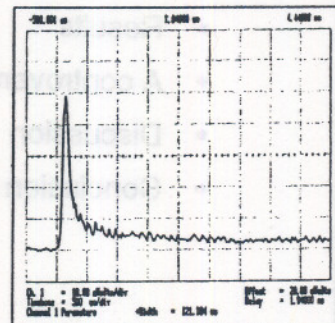
Experiment



- Self-pumped two beam coupling



Focal plane geometry





Theory



• Optical fields

$$\left. \begin{aligned} E_s(z,t) &= \frac{1}{2} A_s(z,t) \exp i(-kz - \omega t) + c.c. \\ E_p(z,t) &= \frac{1}{2} A_p(z,t) \exp i(+kz - \omega t) + c.c. \end{aligned} \right\}$$

• Intensity fringes

$$I(z,t) = (I_s + I_p) \left(1 + \frac{A_s A_p^*}{I_s + I_p} \exp(2ikz) + c.c. \right)$$

$$\left. \begin{aligned} \frac{dN_d^+}{dt} &= s(I + I_{Dark})(N_d - N_d^+) - \gamma_r n N_d^+ \\ \frac{dn}{dt} &= \frac{dN_d^+}{dt} + \frac{1}{e} \frac{dJ}{dz} \\ J &= e\mu n E_{sc} + \mu k_B T \frac{dn}{dz} + sI(N_d - N_d^+) e\delta \\ \epsilon_s \frac{dE_{sc}}{dz} &= e(N_d^+ - N_a^- - n) \end{aligned} \right\}$$

(Kiev group/Kukhtarev material equations)

A_s, A_p are the slowly varying amplitudes of the electric fields



Theory



• Solving the Kiev group/Kukhtarev's equations for the space charge field gives:

$$\frac{a}{\tau_{di}} \frac{\partial E_{sc}}{\partial t} + b E_{sc} + c m = 0$$

• Where:

$$m = \frac{\sqrt{I_s I_p}}{I_s + I_p + I_{Erasure}} \exp(-i\varphi),$$

$$a = 1 + \frac{E_d}{E_m} - i \frac{E_0}{E_m},$$

$$b = 1 + \frac{E_d}{E_q} - i \left(\frac{E_0 + (N_a / N_d) E_{pv}}{E_q} \right),$$

$$c = E_0 + E_{pv} + i E_d$$

$$\tau_{di} = \epsilon_s \gamma_r N_a / e \mu s (I_p + I_s + I_{Erasure})(N_d - N_a)$$

$$E_d = \frac{2\pi k_B T}{e\Lambda}$$

$$E_q = (1 - N_a / N_d) e N_a / 2k\epsilon_s$$

$$E_{pv} = \gamma_r N_a^- \delta / \mu$$

$$E_m = \gamma_r N_a / (\mu K)$$



Theory



- The space charge field modifies the refractive index through the linear Pockels effect
- Substituting the modulated index into the optical wave equation gives the coupled equations for the intensities and phase:

$$\frac{\partial I_p}{\partial z} = -\alpha I_p - \frac{2\pi n^3 r_{\text{eff}}}{\lambda} \sqrt{I_p I_s} \text{Im}(E_{sc} \exp(i\varphi))$$

$$\frac{\partial I_s}{\partial z} = +\alpha I_s - \frac{2\pi n^3 r_{\text{eff}}}{\lambda} \sqrt{I_p I_s} \text{Im}(E_{sc} \exp(i\varphi))$$

$$\frac{\partial \varphi}{\partial z} = \frac{\pi n^3 r_{\text{eff}} (I_p - I_s)}{\lambda \sqrt{I_p I_s}} \text{Re}(E_{sc} \exp(i\varphi))$$



Theory



- In steady state the space charge field and coupled equations reduce to:

$$E_{sc}(z) = \frac{-(E_0 + iE_d + E_{pv})m(z)}{1 + \frac{E_d}{E_q} - i \left(\frac{E_0}{E_q} + \frac{N_a E_{pv}}{N_d E_q} \right)}$$

$$\frac{dI_p}{dz} = -\alpha I_p - \Gamma \frac{I_p I_s}{I_p + I_s + I_{\text{Erasure}}}$$

$$\frac{dI_s}{dz} = +\alpha I_s - \Gamma \frac{I_p I_s}{I_p + I_s + I_{\text{Erasure}}}$$

Where

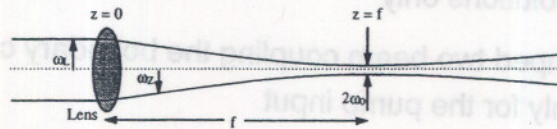
$$\Gamma = \frac{2\pi}{\lambda} n^3 r_{\text{eff}} \text{Im}(E_s)$$



Theory



• Focusing



$$\frac{dI_p}{dx} = -\alpha I_p - \Gamma \frac{I_p I_s}{I_p + I_s + I_{Erasure}} - \frac{2(z-f)I_p}{z_R^2 + (z-f)^2}$$

$$\frac{dI_s}{dx} = +\alpha I_s - \Gamma \frac{I_p I_s}{I_p + I_s + I_{Erasure}} - \frac{2(z-f)I_s}{z_R^2 + (z-f)^2}$$



Theory



• Piezoelectric/photoelastic contributions^{1,2}

$$r_{ij}^{eff} = r_{ijk}^S \hat{n}_k + p_{ijkl}^E \hat{n}_l A_{km}^{-1} B_m$$

$$A_{ik} = C_{ijkl}^E \hat{n}_j \hat{n}_l$$

$$B_i = e_{kij} \hat{n}_k \hat{n}_j$$

r_{ijk}^S is the clamped EO tensor

C_{ijkl}^E is the elastic stiffness tensor

p_{ijkl}^E is the effective elasto-optic tensor

e_{kij} is the piezoelectric tensor

• Scalar effective EO coefficient:

$$r_{eff} = \hat{n}_p \cdot r_{ij}^{eff} \cdot \hat{n}_s$$

- 1) M. Zgonik, K. Nakagawa, P. Günter, "Electro-optic and dielectric properties of photorefractive BaTiO₃ and KNbO₃", J. Optical Society of America B, vol. 12, no. 8, pp 1416-1421, 1995.
- 2) M. Zgonik, R. Schlessler, I. Biaggio, E. Voit, J. Tscherry, P. Günter, "Materials constants of KNbO₃ relevant for electro- and acousto-optics", J. Applied Physics, vol. 74, no. 2, pp 1287-1297, 1993.



Theory



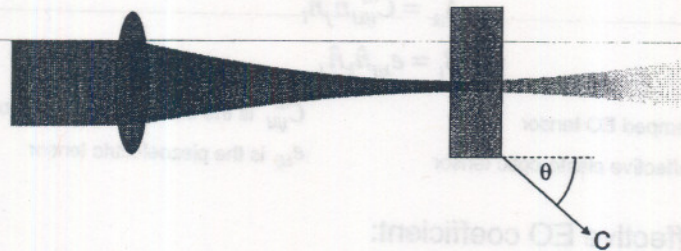
- No closed form solution to the coupled equations
- Numerical solutions only
- For self pumped two beam coupling the boundary conditions are known only for the pump input
- Iterative shoot and match methods are required



Off-axis concept



- Optical gain is potentially much higher away from the c-axis



- c-axis is best for Fe:LiNbO₃ owing to the huge PV effect
- The same is NOT true for Fe:KNbO₃

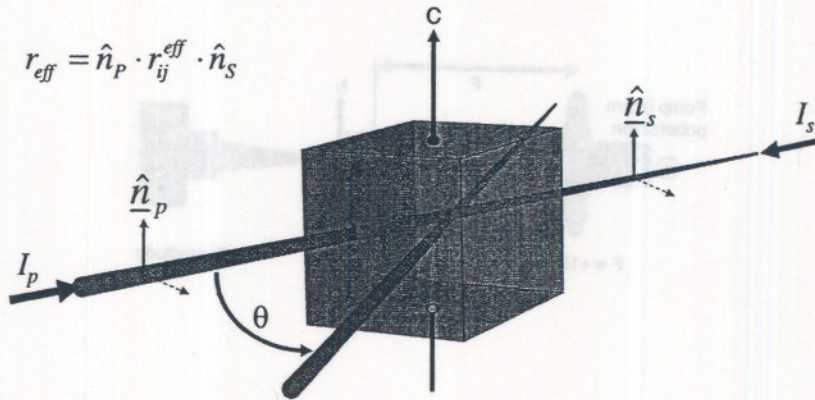


Off-axis effective electro-optic coefficient



- Beam axis rotation about the c-axis:

$$r_{eff} = \hat{n}_P \cdot r_{ij}^{eff} \cdot \hat{n}_S$$



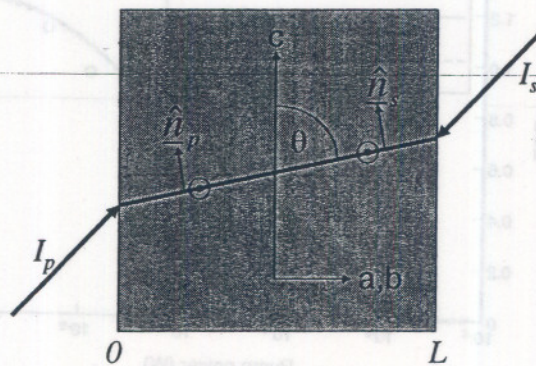
r_{eff} is zero for all polarizations



Off-axis effective electro-optic coefficient



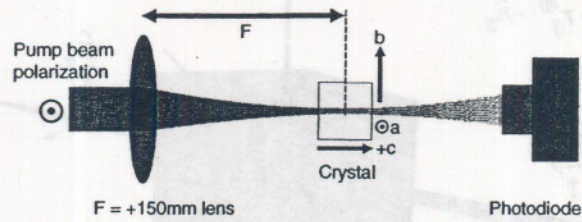
- Beam axis rotation in the a-c or b-c planes:



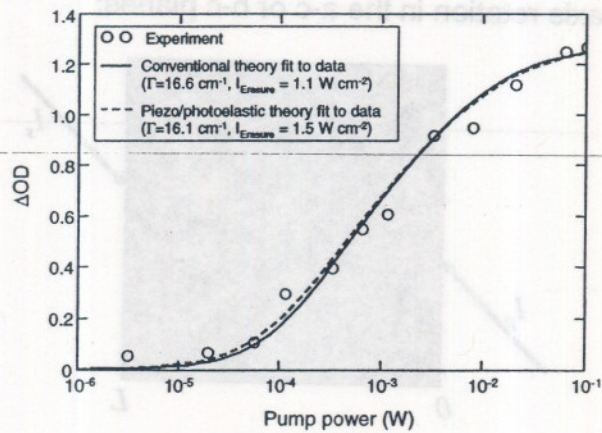
$$r_{eff} = \hat{n}_P \cdot r_{ij}^{eff} \cdot \hat{n}_S$$



Determination of Γ and $I_{Erasure}$



Determination of Γ and $I_{Erasure}$



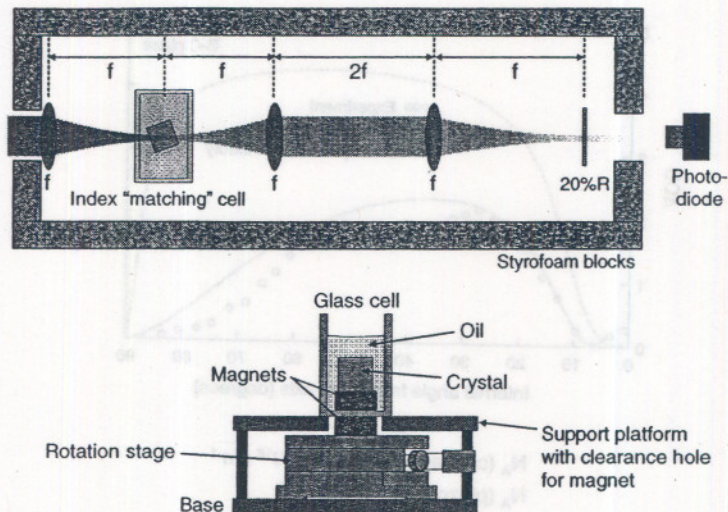
$$N_A (\text{conventional}) = 2.1 \times 10^{16} \text{ cm}^{-3}$$
$$N_A (\text{piezo}) = 5.0 \times 10^{16} \text{ cm}^{-3}$$



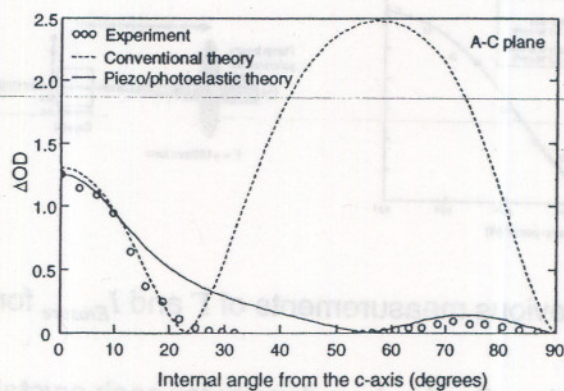
Measuring the off-axis gain in contra-directional Fe:KNbO₃



• Experiment:



Off-axis gain results

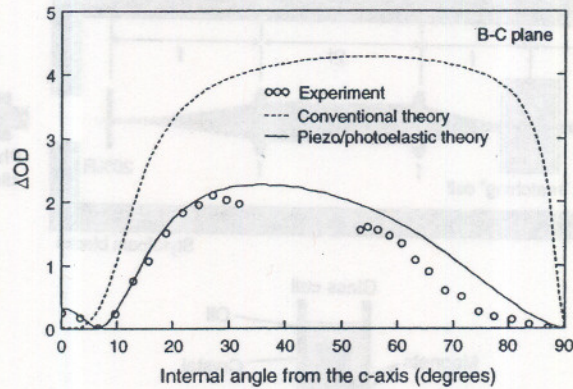


$$N_A (\text{conventional}) = 2.1 \times 10^{16} \text{ cm}^{-1}$$

$$N_A (\text{piezo}) = 5.0 \times 10^{16} \text{ cm}^{-1}$$



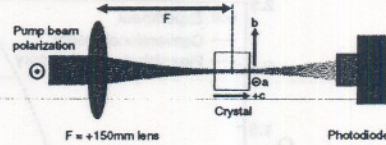
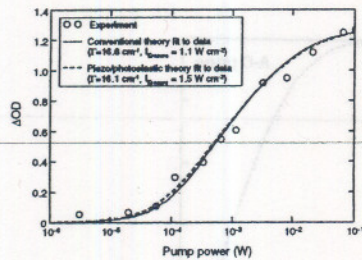
Off-axis gain results



N_A (conventional) = $2.1 \times 10^{16} \text{ cm}^{-1}$
 N_A (piezo) = $5.0 \times 10^{16} \text{ cm}^{-1}$



A controversial suggestion....

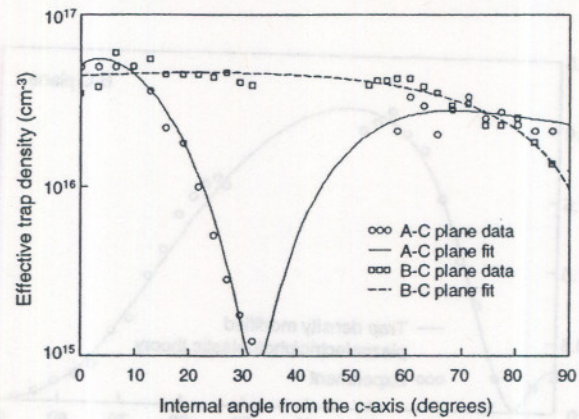


- Repeat previous measurements of Γ and I_{Erasure} for all crystal angles
- Calculate the effective trap density for each crystal angle

"Conventional" transmission/reflection grating method cannot be used reliably for trap density measurements owing to admittance angular restrictions at large crystal angles and competition from rear Fresnel reflection.



Effective trap density variations?

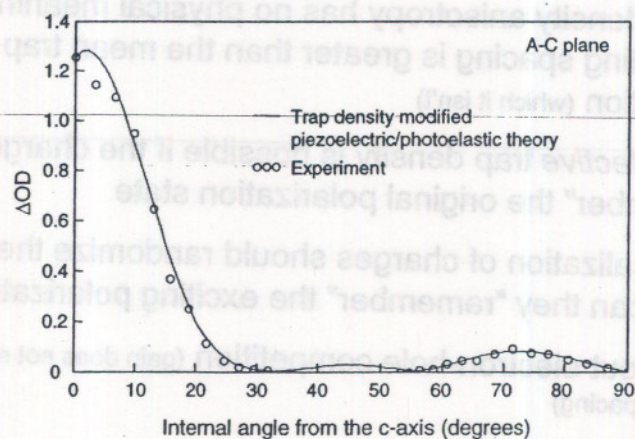


$$N_{A(AC)} = -9.09 \times 10^{12} \theta^6 + 3.03 \times 10^{15} \theta^5 - 3.85 \times 10^{17} \theta^4 + 2.26 \times 10^{19} \theta^3 - 5.67 \times 10^{20} \theta^2 + 3.21 \times 10^{21} \theta + 5 \times 10^{22}$$

$$N_{A(BC)} = -6 \times 10^{16} \theta^3 + 1 \times 10^{20} \theta + 4.42 \times 10^{22}$$

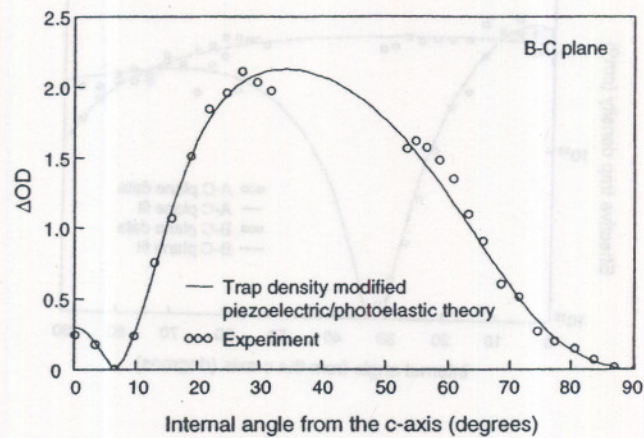


Theory modified for apparent trap density variations





Theory modified for apparent trap density variations



Discussion



- Trap density anisotropy has no physical meaning unless the grating spacing is greater than the mean trap separation (which it isn't)
- An *effective* trap density is possible if the charges “remember” the original polarization state
- Delocalization of charges should randomize the charge state (can they “remember” the exciting polarization?)
- Rule out electron-hole competition (gain does not reverse with grating spacing)



Summary



- High gain confirmed in off-axis geometries for Fe:KNbO₃
- Mismatch between theory and experiment for mid-range crystal angles, especially for the a-c plane
- Large apparent variation in the effective trap density with crystal angle
- Modified theory gives a good fit to experimental data
- Mechanism for trap density anisotropy is unclear