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1. REPORT DATE (DD-MM-YYYY)		2. REPORT TYPE		3. DATES COVERED (From - To)	
31-07-2008		REPRINT			
4. TITLE AND SUBTITLE Ion-Collision Emission Excitation Cross Sections for Xenon Electric Thruster Plasmas				5a. CONTRACT NUMBER	
				5b. GRANT NUMBER	
				5c. PROGRAM ELEMENT NUMBER 61102F	
				5d. PROJECT NUMBER 2303	
6. AUTHORS Jason D. Sommerville* Lyon B. King* Yu-Hui Chiu Rainer A. Dressler**				5e. TASK NUMBER RS	
				5f. WORK UNIT NUMBER A1	
7. PERFORMING ORGANIZATION NAME(S) AND ADDRESS(ES) Air Force Research Laboratory /RVBXT 29 Randolph Road Hanscom AFB, MA 01731-3010				8. PERFORMING ORGANIZATION REPORT NUMBER AFRL-RV-HA-TR-2008-1057	
9. SPONSORING/MONITORING AGENCY NAME(S) AND ADDRESS(ES)				10. SPONSOR/MONITOR'S ACRONYM(S) AFRL/RVBXT	
				11. SPONSOR/MONITOR'S REPORT NUMBER(S)	
12. DISTRIBUTION/AVAILABILITY STATEMENT Approved for Public Release; distribution unlimited.					
13. SUPPLEMENTARY NOTES Reprinted from <i>Journal of Propulsion and Power</i> , Vol. 24, No. 4 pp. 880-888, doi:10.2514/1.33657 *Michigan Technological University, Houghton, MI 49931 **Spectral Sciences, Inc., Burlington, MA 01803-3304					
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15. SUBJECT TERMS Xenon ion collision Emission excitation cross section Ion beam luminescence spectra Electric thruster Hall effect thruster					
16. SECURITY CLASSIFICATION OF:			17. LIMITATION OF ABSTRACT	18. NUMBER OF PAGES	19a. NAME OF RESPONSIBLE PERSON
a. REPORT	b. ABSTRACT	c. THIS PAGE			Yu-Hui Chiu
UNCL	UNCL	UNCL	UNL	10	19b. TELEPHONE NUMBER (Include area code)

20080819 339

Ion-Collision Emission Excitation Cross Sections for Xenon Electric Thruster Plasmas

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DOI: 10.2514/1.33657

Energetic xenon ion collisions with neutral xenon atoms play an important role in electric thruster plasma radiation at low electron temperatures. The respective emission excitation cross sections are necessary for the derivation of plasma parameters from the observed radiance. We present apparent emission excitation cross sections for near-infrared $5p^56p$ to $5p^56s$ transitions of neutral xenon impacted by singly and doubly charged xenon ions. The cross sections were measured over a laboratory energy-per-charge-number range of 100 to 900 eV, a range that covers typical Hall effect thruster discharge voltages. The cross sections are derived from ion beam luminescence spectra produced at single-collision conditions and at pressures for which radiation trapping effects were shown to be negligible. The Xe^+ cross sections are significantly higher than those of Xe^{2+} and increase with energy throughout the investigated range. The Xe^{2+} cross sections plateau at approximately 600 eV. The cross sections are incorporated in a collisional-radiative model. The calculations of near-infrared spectra demonstrate that the sensitivity of the model diagnostic with respect to electron temperature increases with ion energy.

Nomenclature

$B(\lambda)$	= uncorrected background spectrum
E	= kinetic energy
e	= electronic charge
$f(E)$	= energy distribution function
G	= geometric correction factor
I	= ion beam current
J_λ	= radiation intensity in photons per second for transition at λ
k_λ	= emission excitation rate coefficient at wavelength λ
M	= metastable enhancement factor
m	= atom or ion mass
n_0	= neutral density
n_e	= electron density
P	= target gas pressure
q	= species charge number
$R(\lambda)$	= uncorrected spectrum
S	= sensitivity of emission line intensity ratio to electron temperature.
$s(\lambda)$	= detector spectral sensitivity
T_e	= electron temperature
t	= exposure time
α_1	= singly charged ion fraction
$\Delta\lambda$	= peak integration width
λ	= wavelength

ξ_λ	= unscaled emission excitation cross section at wavelength λ
σ_λ	= emission excitation cross section at wavelength λ

I. Introduction

XENON is used as a propellant in electric spacecraft thrusters, including Hall effect thrusters (HETs) and gridded ion thrusters. These thrusters use electric and magnetic fields to create a plasma and expel the ions at high velocity to generate thrust [1]. The development and testing of these thrusters requires diagnostic techniques that can provide information on plasma properties such as species densities and electron temperature. Optical emission spectroscopy [2] is an attractive diagnostic because it is entirely passive and noninvasive; therefore, it guarantees that the plasma is not disturbed in the process of taking the measurement. Furthermore, it can be applied to remote observations or to examine regions of thruster plasma that are too harsh for electrostatic probes. Observed spectra of absolute line intensities can contain sufficient information to derive electron energy distributions and electron number densities [2–6]. Researchers studying HETs have applied optical emission spectroscopy to investigate general emissive properties [7–9], plasma fluctuations [10–12], transient operation [13], and thruster body erosion [14,15].

The extraction of quantitative plasma parameters from observed radiance relies on an emission model that describes the intensity of radiation from two or more transition lines as a function of plasma parameters such as electron temperature and species densities. Central to most of these models are emission cross section data. The cross sections are necessary to determine the emission excitation rate coefficient:

$$k_\lambda = \int_0^\infty \sqrt{\frac{2E}{m}} \sigma_\lambda(E) f(E) dE \quad (1)$$

where $f(E)$ is normalized. As Eq. (1) implies, cross sections are collision-energy-dependent. HET plasma electrons typically have a broad energy distribution; therefore, a model requires cross sections over a wide range of electron energies. In contrast, ion-energy distributions in HETs are narrow; typical full widths at half-maximum are about 15% of the acceleration voltage and are centered at energies about 10% below the corresponding acceleration voltage

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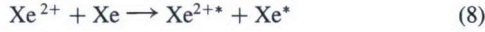
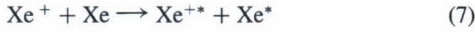
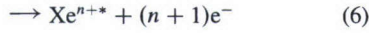
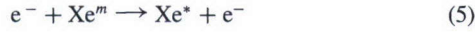
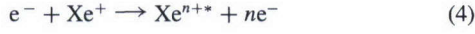
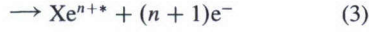
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[16,17]. Therefore, a single set of ion cross sections is sufficient to model an HET operating at a fixed discharge voltage, provided these cross sections change slowly over the range of ion energies present in the HET beam.

The model must incorporate cross sections for all processes that give rise to significant emission on lines within a spectral region of interest. A comprehensive list of xenon emission excitation processes that affect plasma radiation include



where an asterisk denotes excited emitting atoms or ions, and m denotes metastable atoms.

Until recently, a lack of accurate emission excitation cross sections for many of these processes limited the utility of optical emission spectroscopy in determining electron temperature and other plasma parameters in xenon-propelled electric thrusters. This problem was circumvented by applying a simple collisional-radiative model such as the coronal equilibrium model [2,3,18]. The corona model assumes that all excitation occurs from electron collisions with ground-state neutral atoms; that is, only process 2 is considered. This greatly reduces the need for cross section data. Bouchoule et al. [13], however, demonstrated experimentally that metastables play a fundamental role in the HET plasma radiance spectrum, and therefore processes 5 and 6 cannot be ignored. Meanwhile, Chiu et al. [19] reported cross sections for ion emission excitation (processes 8 and 9) for 300-eV ions, demonstrating that they are large for near-infrared lines. Karabadzhak et al. [20] developed a more rigorous collisional-radiative model (CRM) for the near-infrared spectrum of xenon thruster plasmas. In addition to electron collisions with ground-state neutrals (process 2), the model also considers excitation from metastable states (process 5) and excitation due to heavy-particle collisions using the cross sections of Chiu et al. [19]. The results of Karabadzhak et al. [20] demonstrated that metastables can enhance emission intensities of particular lines by factors as high as three, depending on the electron temperature. Karabadzhak et al. also showed that ion-neutral collisions are a significant excitation mechanism at low electron temperatures, such as may be encountered in the plume of an HET. To reduce the number of coupled state-to-state transition rate equations necessarily introduced by the added processes, the CRM makes use of apparent emission cross sections. These cross sections include emission contributions from cascade transitions from higher levels into the upper level as well as direct excitation to the upper level from the ground state.

The CRM of Karabadzhak et al. [20] applies to the emission lines associated with the lowest excited states of neutral xenon: specifically, the transitions from the $5p^56p$ (2p in Paschen notation) upper states to the $5p^56s$ (1s) lower states (Table 1). These lines have high intensities and occur in a relatively sparse region of the spectrum, thereby avoiding experimental problems due to line interference, and cross section data for many of the reactions are

Table 1 Xenon emission lines studied

Upper state	Branching ratio ^a	Lower state	λ , nm
2p ₁	0.720 ± 0.059	1s ₂	788.74
2p ₃	0.863 ± 0.050	1s ₂	834.68
2p ₅	0.998 ± 0.000	1s ₄	828.01
2p ₆	0.743 ± 0.041	1s ₅	823.16
2p ₆	0.243 ± 0.039	1s ₄	895.23
2p ₇	0.901 ± 0.020	1s ₄	916.27
2p ₇	0.099 ± 0.019	1s ₅	840.92
2p ₈	1.000 ± 0.000	1s ₅	881.94
2p ₉	0.620 ± 0.029	1s ₄	992.32
2p ₉	0.379 ± 0.046	1s ₅	904.54
2p ₁₀	0.975 ± 0.032	1s ₅	979.97

^aCalculated from electron-impact excitation cross sections given in [21].

available. Because the subject lines involve only emissions from neutral xenon, processes 3, 4, 6, and 9 may be ignored. Cross sections for process 2 have recently been measured over a broad range of energies [19,21]. The model includes an approximate method based on cascade population of metastable states to account for process 5.

Before this work, the CRM of Karabadzhak et al. [20] was limited to thrusters with discharge voltages of approximately 300 V. Although 300 V may be regarded as the most common thruster discharge voltage, HETs operate over a range of voltages from 100 V [22] to 1000 V [23–27]. Most gridded ion thrusters operate above 1000 V, but some work has been done on low-power and low-specific-impulse ion thrusters that used voltages in the 600–900-V range [28,29]. Given the demonstrated importance of ion-collision-induced emission excitation for 300-V discharge voltages and the possibility that the respective emission excitation cross sections depend strongly on energy, measurements of cross sections for processes 7 and 8 over a broad energy range are necessary to apply the diagnostic model to engines with other discharge voltages. To extend the applicability of the Karabadzhak et al. [20] CRM to a wider range of thruster operating points, we measured the apparent emission cross sections for singly and doubly charged ion impact of neutral xenon for the E/q range from 100 to 900 eV for the transitions listed in Table 1. The new cross sections are incorporated into the CRM of Karabadzhak et al., and the effect of HET discharge voltage on the predicted near-infrared xenon plasma radiation spectrum is discussed.

II. Experiment

The measurements were performed on the ion beam luminescence apparatus of the U.S. Air Force Research Laboratory at Hanscom Air Force Base. The high-vacuum apparatus, schematically shown in Fig. 1, was previously described by Chiu et al. [19], who measured cross sections for 300-eV-per-charge ions. In summary, a glow-discharge ion source generates a xenon ion beam. The beam passes through a Wien velocity filter that selects the desired ionization level. The beam is then focused into a collision cell filled with a xenon target gas. After passing through the cell, the beam current is collected by a Faraday cup. To assess the divergence of the ion beam, the beam current to the Faraday cup and the back wall of the cell are monitored with electrometers, and the sum of the currents is used to quantify the beam current. As the beam enters the Faraday cup, it passes through a retarding potential analyzer (RPA), which is used to verify the beam energy before the luminescence experiment is run. The target gas pressure is monitored with a capacitance manometer. Both target pressure and beam current are continuously monitored throughout each experiment and the time-averaged values are used. The face of a fiber optical bundle collects luminescence photons produced in the cell and the fibers guide the photons out of the vacuum chamber to a spectrograph equipped with a charge-coupled device detector.

The spectral sensitivity of the spectrograph was calibrated by comparing the recorded spectrum of a halogen-tungsten lamp with the Planck distribution for a body at the temperature of the lamp, as measured by a pyrometer. The temperature of the lamp was found to

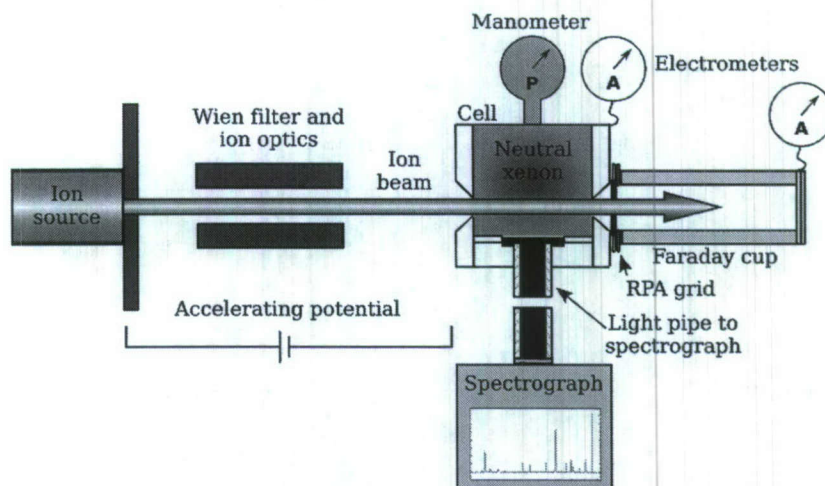


Fig. 1 Schematic representation of the ion excitation cross section experiment.

be 3065 K, close to the manufacturer's specified nominal temperature of 3200 K. The intensity in photons per unit time of a line at wavelength λ was calculated by first correcting the raw spectra for background light, the spectral sensitivity of the detection system and the exposure time, and then integrating over a fixed spectral width:

$$J_{\lambda} = \int_{\lambda-\frac{1}{2}\Delta\lambda}^{\lambda+\frac{1}{2}\Delta\lambda} s(\lambda') \frac{R(\lambda') - B(\lambda')}{t} d\lambda' \quad (10)$$

The integration width was chosen to approximate the full width at half-maximum of the spectral peaks. The widths were 0.6 and 1 nm for experiments run with 50- and 100- μm entrance slits, respectively.

The cross sections are derived from the spectral measurements of a given transition line according to

$$\sigma_{\lambda} = \frac{qeJ_{\lambda}}{GPI} \quad (11)$$

All experiments were conducted with a cell pressure of 1.5 mtorr. Chiu et al. [19] showed that radiation trapping effects are insignificant for ion excitation of the studied emission lines between the pressures of 0.1 and 2.0 mtorr. The geometric correction factor G accounts for the effective solid angle of observation and detector efficiency. It was experimentally determined by comparing the $\text{Xe}^+ + \text{Xe}$ reaction at 300 eV with the results previously obtained at this energy on the same apparatus [19]. We define the unscaled cross section as Eq. (11) with $G = 1$:

$$\xi_{\lambda} = \frac{qeJ_{\lambda}}{PI} \quad (12)$$

If σ_i are the values reported by Chiu et al. [19] for different transitions i , and ξ_i are the respective measured unscaled cross sections, then G is given by the average:

$$G = \left\langle \frac{\xi_i}{\sigma_i} \right\rangle \quad (13)$$

The G factor is used to scale all measurements, yielding the absolute cross sections.

There are two major contributions to error in the experiment. The largest contribution arises from the uncertainty in the G factor, and the estimate for this contribution is given by the standard deviation of the individual ratios averaged in Eq. (13), which yields a 30% error. The second major contribution is the uncertainty in the spectral-sensitivity calibration. We estimate this error by the largest deviation between the presently determined line ratios for lines with common upper states and the average of values derived from previously published data, as shown in Table 2. The data listed under experiment 1 are derived from measurements of electron-impact excitation cross sections by Fons and Lin [21]. Experiment 2 refers to the ratios of Einstein A coefficients listed by Radzig and Smirnov [30]. The theoretical values from Aymar and Coulombe [31] are the ratios of A coefficients that were calculated with a central field model. The largest deviation of the present ratios from the average of these values is 22.8%. We therefore conservatively estimate a 25% spectral-sensitivity calibration error for all lines. Inspection of the data suggests that the spectral-sensitivity calibration error most strongly affects the 916- and 905-nm transitions. Combining the G -factor and spectral-sensitivity contributions, the error on all measurements at E/q of 200 eV and greater is estimated at 40%. For lines on which multiple measurements were conducted, the standard deviations were checked and all were within this 40% estimate. For the 100-eV measurements, the errors are estimated at 60%, except for the 841-nm line, for which the error is estimated at 85%. Note that this line is not used in the CRM of Karabadzha et al. [20], because it has the same upper state ($2p_7$) as the 916-nm line, which is significantly more intense. The higher errors at 100 eV are due to smaller signal levels combined with less stable beam operation and high divergence at this low energy.

III. Results

Figure 2 shows $2p$ -line emission excitation cross sections for $\text{Xe}^+ + \text{Xe}$ collisions as a function of ion energy (laboratory frame). In the case of the $2p_6$, $2p_7$, and $2p_9$ upper levels, two lines are

Table 2 Comparison of line ratios

Upper state	λ_1/λ_2	Exp. 1 ^a	Exp. 2 ^b	Theory ^c	Average	Present	% deviation
$2p_6$	823 nm/895 nm	3.06 ± 0.653	2.87 ± 0.406	2.36	2.76	2.65 ± 0.211	-4.10
$2p_7$	916 nm/841 nm	9.10 ± 1.95	10.8 ± 1.53	14.0	11.3	13.0 ± 4.71	15.0
$2p_9$	992 nm/905 nm	1.64 ± 0.350	1.60 ± 0.226	1.81	1.68	1.30 ± 0.178	-22.8

^aCalculated from electron-impact excitation cross sections given in [21].

^bCalculated from Einstein A coefficients given in [30].

^cCalculated from Einstein A coefficients given in [31].

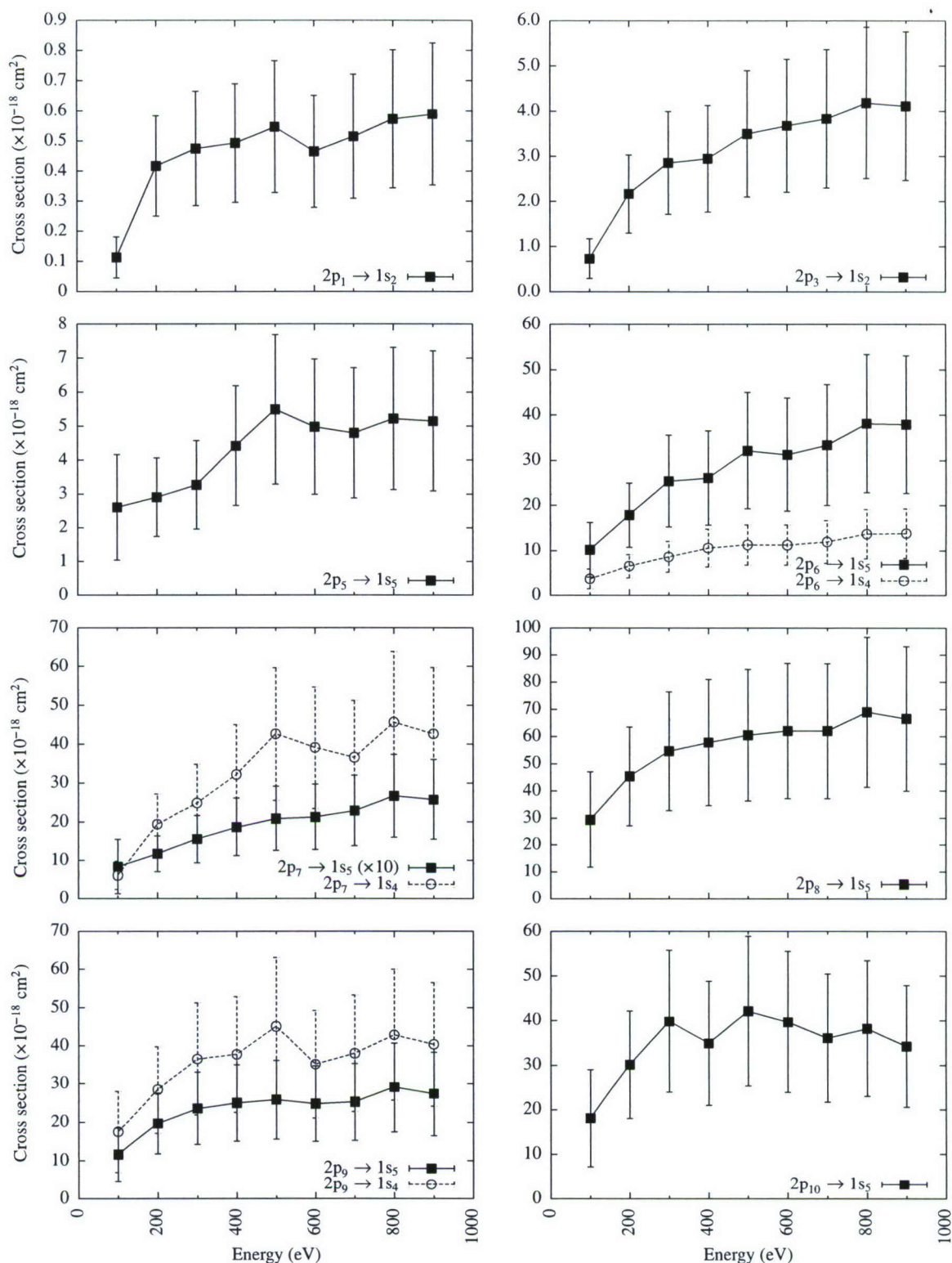


Fig. 2 $\text{Xe}^+ + \text{Xe}$ emission excitation cross sections as a function of ion energy.

produced in the near-infrared spectrum, and the emission excitation cross sections for both lines are shown. Except for the 980-nm line originating in the $2p_{10}$ level, all emission excitation cross sections are shown to increase with ion energy throughout the measured ion-energy range. In the case of the 980-nm line, the cross section becomes constant with energy above 300 eV. The cross section values are listed in Table 3. The largest cross section of $69 \times 10^{-18} \text{ cm}^2$ is observed for the $2p_8$ line at an energy of 800 eV,

whereas the maximum electron-induced emission excitation cross section is approximately $38 \times 10^{-18} \text{ cm}^2$ at 15 eV for the same line [21].

Figure 3 shows 2p-line emission excitation cross sections for $\text{Xe}^{2+} + \text{Xe}$ collisions as a function of ion energy. Again, transitions that share a common upper state are plotted together. All cross sections appear to plateau at 600 eV. The cross section values are listed in Table 4. The 789-nm Xe^{2+} cross section at 200 eV was

Table 3 Xe⁺ + Xe emission excitation cross sections (10⁻¹⁸ cm²)^a

Energy, eV	789 nm	823 nm	828 nm	835 nm	841 nm	882 nm	895 nm	905 nm	916 nm	980 nm	992 nm
100	0.112 ^b	10.1 ^b	2.60 ^b	0.729 ^b	0.832 ^c	29.3 ^b	3.67 ^b	11.5 ^b	6.05 ^b	18.1 ^b	17.4 ^b
200	0.416	17.8	2.90	2.16	1.16	45.3	6.49	19.6	19.4	30.1	28.4
300	0.475	25.4	3.27	2.85	1.55	54.6	8.55	23.6	24.8	39.7	36.5
400	0.493	26.1	4.42	2.94	1.86	57.8	10.5	25.0	32.1	34.8	37.7
500	0.547	32.1	5.49	3.49	2.08	60.4	11.2	25.8	42.5	42.0	45.0
600	0.464	31.2	4.97	3.67	2.12	62.0	11.2	24.8	39.0	39.6	35.1
700	0.515	33.3	4.79	3.82	2.28	62.0	11.9	25.2	36.5	36.0	38.0
800	0.573	38.1	5.22	4.18	2.66	68.9	13.6	29.1	45.6	38.1	42.8
900	0.589	37.9	5.14	4.11	2.57	66.5	13.7	27.3	42.6	34.1	40.3

^aAll measurements have a 40% uncertainty unless otherwise marked.^b60% uncertainty^c85% uncertainty

unobtainable due to the low signal-to-noise ratio for that peak. It is denoted by a dash in the table. Except for the 789-nm line, all cross sections are smaller than their Xe⁺-induced counterparts by roughly an order of magnitude. The largest cross section is approximately 13×10^{-18} cm², observed for the 882-nm line at 600 eV.

IV. Discussion

The presently determined energy dependencies of emission excitation cross sections in Xe⁺ + Xe and Xe²⁺ + Xe collisions differ significantly from those reported for e⁻ + Xe collisions [19,21], in which the cross section maxima are observed much closer to the emission excitation threshold. From a purely classical perspective, the observation of higher heavy-particle emission excitation cross sections than those observed for electrons is surprising because the relative nuclear motion is substantially slower than the motion of the atomic electrons at all of the investigated energies, making momentum transfer between a heavy particle and an electron unlikely. The observed heavy-particle emission excitation cross sections must thus be the consequence of a nonclassical, nonadiabatic mechanism [32,33]. The observed increasing trend with energy is fully consistent with a nonadiabatic excitation mechanism.

In the case of the Xe²⁺ + Xe system, the emission excitation cross sections reach a plateau near ion energies of 600 eV (300 eV per charge unit). For both charge states, charge transfer (two-electron charge transfer in the Xe²⁺ + Xe system) is likely to play an important role in the excitation mechanism, and the excited atom may be associated with either the original ion or the target atom. This conclusion can be made because charge exchange occurs effectively at large interatomic distances in regions of weak interaction, whereas the electronic transitions leading to radiance occur at short range in regions of strong interaction.

Given the growing Xe⁺-induced emission cross sections with energy, it can be anticipated that the ionic contribution to the observed emission spectrum of a Xe-propelled HET grows with discharge voltage. Furthermore, the excitation rate scales with collision velocity, leading to an additional enhancement of ionic contributions with energy. As reported by Chiu et al. [19] and Karabadzahk et al. [20], a significant factor contributing to the sensitivity of the near-infrared spectrum to electron temperature in plasmas with low electron temperatures is the significant difference between the ion-induced and low-energy electron-induced near-infrared spectra. With the higher ion-induced excitation rates at higher discharge voltages, the sensitivity of the diagnostic proposed by Karabadzahk et al. will be further enhanced. As the electron temperature is increased, the number of electrons with energies exceeding the excitation threshold increases, and the high-energy ion contribution becomes less significant.

By studying emissions simulated by the CRM of Karabadzahk et al. [20], which incorporates the present ion-collision excitation cross sections, we can quantify the contributions of ion-induced emission over a range of electron temperatures and ion energies, the latter of which approximately corresponds to HET discharge voltage.

In this CRM, k_{λ_e} , k_{λ_1} , and k_{λ_2} are the rate coefficients for electron, singly charged ion, and doubly charged ion induced emissions, respectively, and the emission excitation rate for the line emitting at λ obeys the relation

$$J_{\lambda} \propto n_0 n_e \left[k_{\lambda_e}(T_e) + \alpha_1 k_{\lambda_1}(E_1) + \frac{1 - \alpha_1}{2} k_{\lambda_2}(E_2) \right] [1 + M(T_e, \alpha_1, \lambda)] \quad (14)$$

Note that k_{λ_e} is a function of the electron temperature, k_{λ_1} and k_{λ_2} are functions of the singly and doubly charged ion energies E_1 and E_2 , and $M(T_e, \alpha_1, \lambda)$ is an intensity-enhancement factor that accounts for population of the 1s₅ metastable state through radiative transitions out of the 2p₆₋₁₀ levels. Note that if we set M , k_{λ_1} , and k_{λ_2} to zero, then Eq. (14) reduces to the coronal equilibrium model.

Because of inaccuracies in the emission excitation cross sections, as well as in some of the assumptions pertaining to metastable excitation, the CRM is considered to be more accurate if a spectrum of the full set of the near-infrared lines is simulated and fit to the observation through the variation of plasma parameters (i.e., electron temperature and ion charge ratio). However, for the purpose of clarity, we focus on the emission rates of only two lines, which allows us to discuss the ratio of the two rates rather than a best fit for all of the lines. Because Karabadzahk et al. [20] demonstrated the significant sensitivity of the 823/828-nm intensity ratio to electron temperature, we shall use these lines. Furthermore, these two lines are attractive because of the spectral proximity that limits the effect of errors in the spectral-sensitivity calibration. Note that the M factor must be retained in the following analysis because the 823-nm line is coupled to the 1s₅ metastable state. Therefore, additional cross sections measured in this work that are associated with other transitions into the 1s₅ state must be incorporated in the two-line analysis via the M factor (see Karabadzahk et al. [20] for details on the M factor).

Figure 4 shows the fraction of emission due to ion-impact excitation of the 823- and 828-nm lines as a function of electron temperature for the nine ion energies studied. HET plumes typically exhibit electron temperatures between 1–2 eV in the far field (greater than 100 mm) and approximately 10 eV close to the exit plane of the thruster [34,35]. For xenon plasmas with 300-eV ions, the contribution to the 823-nm line intensity by ion collisions is approximately 10% at an electron temperature of 5 eV and rises rapidly with decreasing electron temperature. The contribution is particularly strong as the electron temperature approaches 2 eV. (Because of certain approximations, the CRM is inaccurate below 2 eV.) As the ion energy increases, so does the fraction of the emission produced by ion-impact processes. In fact, for the 900-eV ions, the fraction of emission due to ion-impact excitation does not drop below 10% until the electron temperature is greater than 8.5 eV. This behavior is typical of all of the modeled lines involving transitions into the 1s₅ metastable state. The 828-nm line shows the lowest fraction of emission due to ion-induced excitation of all of the investigated lines. However, even here, the fraction is significant for electron temperatures between 2 and 4 eV, particularly for the higher

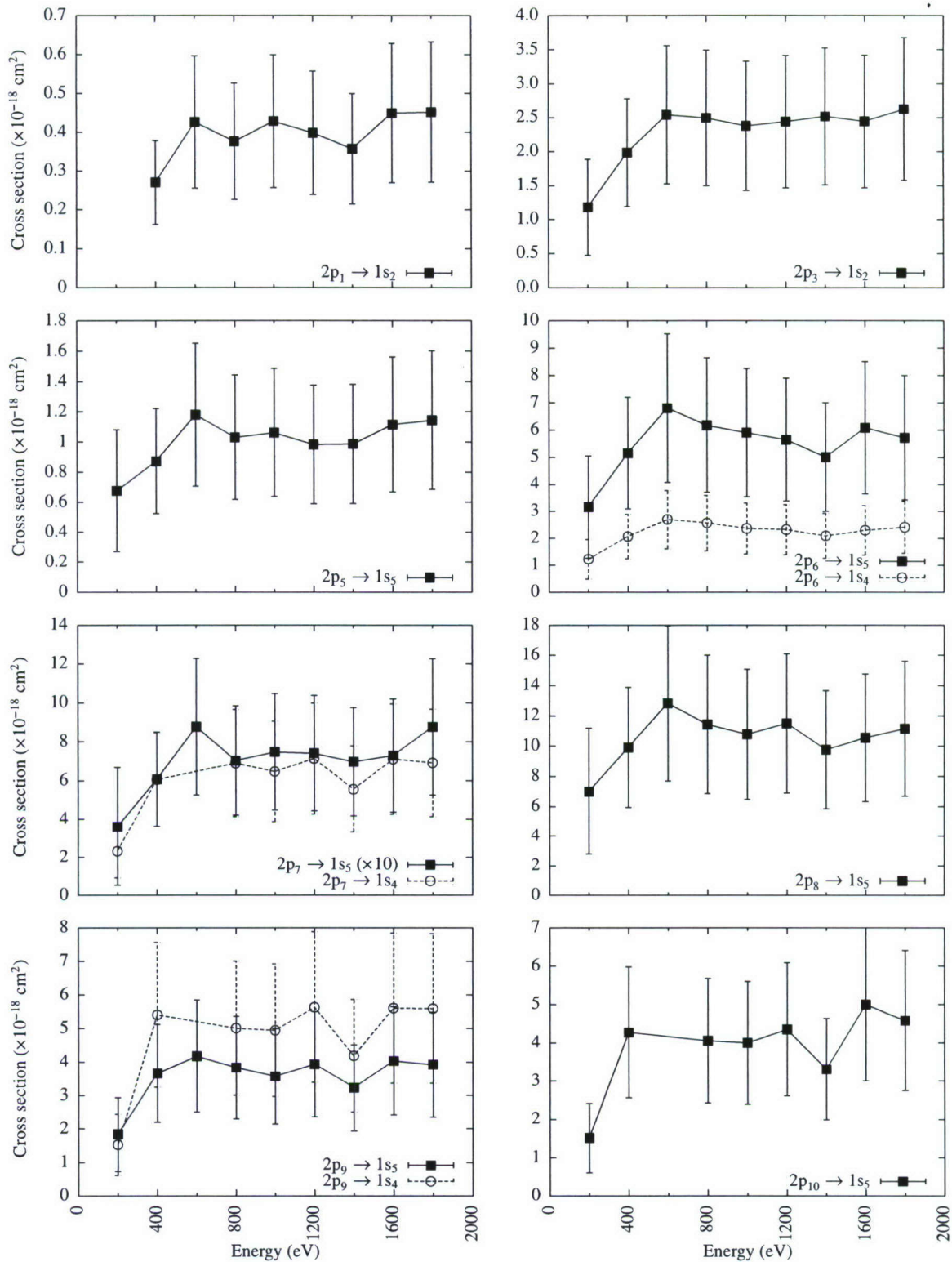


Fig. 3 $\text{Xe}^{2+} + \text{Xe}$ emission excitation cross sections as a function of ion energy.

ion energies. Clearly, ion-induced emissions account for a significant portion of the luminescence on the xenon near-infrared lines in HET plumes.

Figure 5 shows the modeled electron temperature versus the ratio of the emission intensity of the 823-nm line to the intensity of the 828-nm line for four different ion energies. Given an experimentally measured intensity ratio, this two-line model can predict the electron temperature if the ion energy is known. Because of the high ion-induced emission cross sections, particularly for Xe^+ , failure to

account for this excitation mechanism leads to increasingly erroneous results as the electron temperature is reduced or the ion energy is increased. This is further illustrated in Fig. 6, which compares the electron temperature that would be deduced from selected emission ratios using a coronalike CRM that neglects ion-induced excitation, with the results of the CRM incorporating the presently determined ion emission excitation cross sections. The failure to include ion-induced emission leads to increasingly erroneous results as the ion energy increases. These errors are

Table 4 $\text{Xe}^{2+} + \text{Xe}$ emission excitation cross sections (10^{-18} cm^2)^a

Energy, eV	789 nm	823 nm	828 nm	835 nm	841 nm	882 nm	895 nm	905 nm	916 nm	980 nm	992 nm
200	—	3.15 ^b	0.676 ^b	1.17 ^b	0.360 ^c	6.99 ^b	1.22 ^b	1.83 ^b	2.30 ^b	1.51 ^b	1.52 ^b
400	0.270	5.14	0.872	1.98	0.606	9.90	2.06	3.65	6.04	4.27	5.40
600	0.425	6.80	1.17	2.54	0.877	12.8	2.69	4.17	7.88	6.75	6.26
800	0.376	6.17	1.03	2.49	0.702	11.4	2.56	3.83	6.89	4.05	5.00
1000	0.427	5.90	1.06	2.37	0.748	10.7	2.36	3.56	6.46	4.00	4.94
1200	0.398	5.64	0.982	2.43	0.740	11.4	2.31	3.93	7.12	4.35	5.62
1400	0.356	5.00	0.986	2.51	0.696	9.76	2.09	3.22	5.55	3.31	4.18
1600	0.449	6.08	1.11	2.44	0.728	10.5	2.29	4.03	7.09	4.99	5.60
1800	0.452	5.71	1.14	2.62	0.876	11.1	2.40	3.92	6.90	4.57	5.58

^aAll measurements have a 40% uncertainty unless otherwise marked.^b60% uncertainty.^c85% uncertainty.

particularly problematic at low electron temperatures, for which the relative deviation is largest.

For example, suppose a ratio of intensities in a certain location of the plume of an HET operating with a discharge voltage of 800 V is found to be $J_{823}/J_{828} = 3.0$. This is a reasonable value, because Karabadzahk et al. [20] found ratios in HETs ranging between 1.1 and 4.5 from the discharge region to the plume of an anode-layer HET, and Matlock et al. [36] found ratios between ~ 1 and ~ 8 in the plume of a Busek BPT-200. If the coronalike model is used, an electron temperature of 4.4 eV is determined. If, however, the full Karabadzahk et al. [20] CRM is used, an electron temperature 5.8 eV is determined. The coronalike model underpredicts the electron temperature by about 25% in this case.

The sensitivity in the determination of the electron temperature from the ratio of two lines can be described by how rapidly that ratio changes with electron temperature:

$$S(T_e, E_i) \equiv \left| \frac{\partial(J_1/J_2)}{\partial T_e} \right|_{T_e} \quad (15)$$

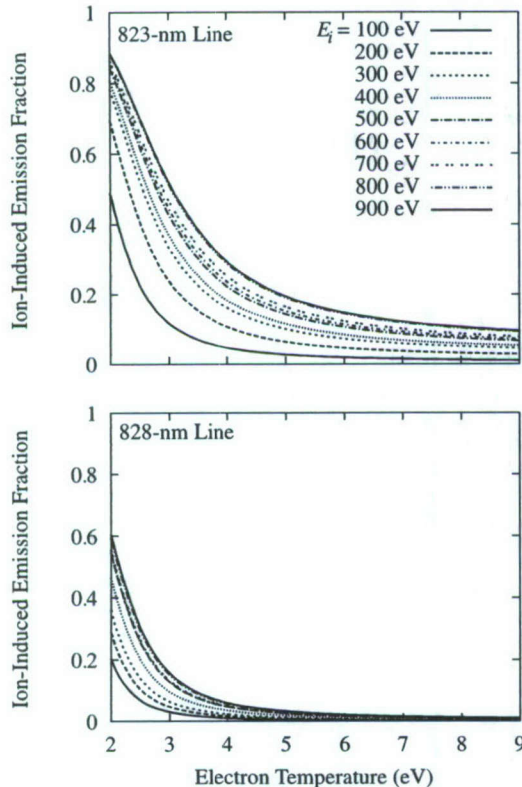


Fig. 4 Fraction of emission due to ion-induced excitation for the 823- and 828-nm lines as a function of electron temperature for each of the ion energies measured in this work.

Because the ion and low-energy electron-induced near-infrared spectra are quite different, ion contributions add to the sensitivity of the CRM in determining the electron temperature. Figure 7 shows the sensitivity 823/828 nm emission ratio as a function of ion energy. For all temperatures, the sensitivity increases with ion energy. Particularly at low electron temperatures, the sensitivity of the J_{823}/J_{828} ratio shows a dramatic increase as the ion energy is

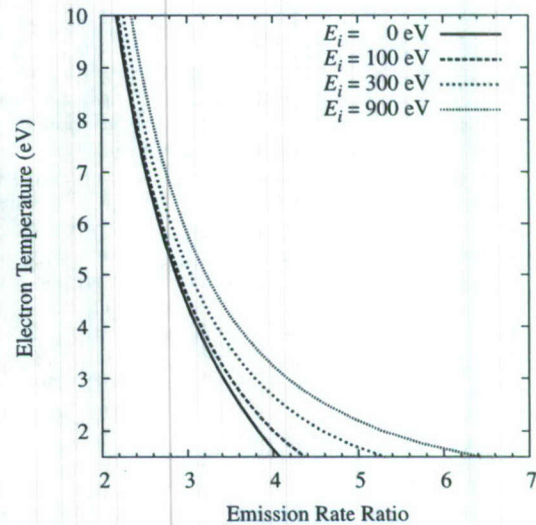


Fig. 5 Emission rate ratio J_{823}/J_{828} as a function of electron temperature for select ion energies.

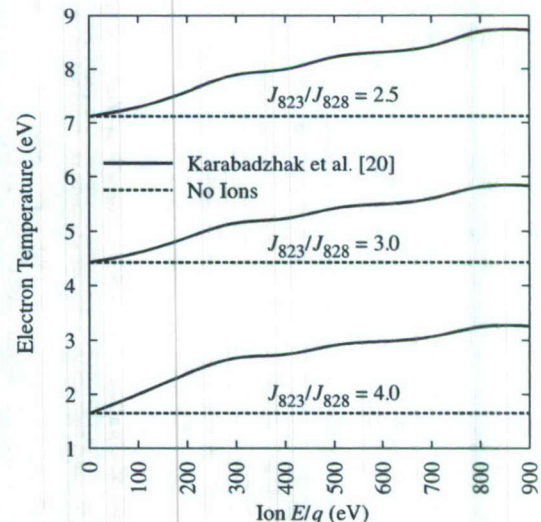


Fig. 6 Ion-energy dependence of derived electron temperature for select emission rate ratio J_{823}/J_{828} .

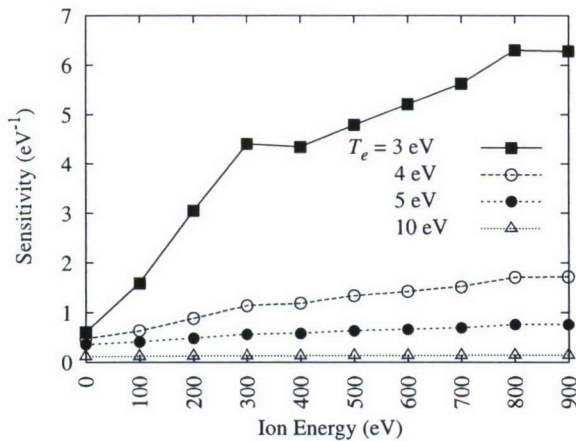


Fig. 7 Sensitivity of the 823- to 828-nm emission ratios to T_e as a function of ion energy at select electron temperatures.

increased. It is fortuitous for electric propulsion research that the proper incorporation of ion-induced emissions into a CRM not only improves the model's accuracy in determining electron temperature, but also improves its sensitivity.

V. Conclusions

We report experimental emission excitation cross sections for 11 emission lines of neutral xenon impacted by Xe^+ and Xe^{2+} across an E/q range from 100 to 900 eV. The cross sections are incorporated in a xenon plasma collisional-radiative model that includes emission excitation by energetic ions. Inclusion of an ion emission excitation mechanism is required to accurately model HET plasmas, in which ion-induced radiance may have comparable intensity with that produced by electron collisions, particularly at high ion energy and low electron temperature. Failure of a CRM to include the ion-induced emissions can result in errors in the determination of the electron temperature in the plume of an HET on the order of 25%. The present cross sections extend the applicability of the Karabadzhak et al. [20] CRM to a broader range of HET operating voltages. The upgraded model can now also be applied to plasmas of low-specific-impulse gridded ion thrusters.

Additional measurements of ion-induced emission excitation cross sections for xenon would further extend the availability and usefulness of the Karabadzhak et al. [20] and other CRMs. In particular, the extension of the range of the Xe^+ - and Xe^{2+} -induced cross sections to 2000 eV per charge number would open these diagnostics to the bulk of gridded ion thrusters. Furthermore, measurements of the emission excitation cross sections of ionic lines through the processes 4 and 6 would improve the accuracy of the model at high electron temperatures and could also provide estimates of absolute electron, neutral, and ion densities.

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