

REPORT DOCUMENTATION PAGE				Form Approved OMB NO. 0704-0188	
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1. REPORT DATE (DD-MM-YYYY) 15-08-2013		2. REPORT TYPE Book		3. DATES COVERED (From - To) -	
4. TITLE AND SUBTITLE InAs1-xSbx alloys with native lattice parameters grown on compositionally graded buffers: structural and optical properties			5a. CONTRACT NUMBER W911NF-12-2-0057		
			5b. GRANT NUMBER		
			5c. PROGRAM ELEMENT NUMBER 611102		
6. AUTHORS Ding Wang, Dmitry Donetsky, Youxi Lin, Gela Kipshidze, Leon Shterengas, Gregory Belenky, Wendy L. Sarney, Stefan P. Svensson			5d. PROJECT NUMBER		
			5e. TASK NUMBER		
			5f. WORK UNIT NUMBER		
7. PERFORMING ORGANIZATION NAMES AND ADDRESSES Research Foundation of SUNY at Stony Brook U Office of Sponsored Programs W-5510 Melville Library Stony Brook, NY 11794 -3362				8. PERFORMING ORGANIZATION REPORT NUMBER	
9. SPONSORING/MONITORING AGENCY NAME(S) AND ADDRESS(ES) U.S. Army Research Office P.O. Box 12211 Research Triangle Park, NC 27709-2211				10. SPONSOR/MONITOR'S ACRONYM(S) ARO	
				11. SPONSOR/MONITOR'S REPORT NUMBER(S) 62447-EL.3	
12. DISTRIBUTION AVAILABILITY STATEMENT Approved for public release; distribution is unlimited.					
13. SUPPLEMENTARY NOTES The views, opinions and/or findings contained in this report are those of the author(s) and should not be construed as an official Department of the Army position, policy or decision, unless so designated by other documentation.					
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15. SUBJECT TERMS InAsSb, compositionally graded buffer, MBE, infrared, minority carrier lifetime, reciprocal space mapping					
16. SECURITY CLASSIFICATION OF:			17. LIMITATION OF ABSTRACT UU	15. NUMBER OF PAGES	19a. NAME OF RESPONSIBLE PERSON Gregory Belenky
a. REPORT UU	b. ABSTRACT UU	c. THIS PAGE UU			19b. TELEPHONE NUMBER 631-632-8397

Report Title

InAs_{1-x}Sb_x alloys with native lattice parameters grown on compositionally graded buffers: structural and optical properties

ABSTRACT

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FRONTIERS IN ELECTRONICS

Selected Papers from the Workshop on
Frontiers in Electronics 2011 (WOFE-11)

San Juan, Puerto-Rico

18 – 21 December 2011

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Published by

World Scientific Publishing Co. Pte. Ltd.

5 Toh Tuck Link, Singapore 596224

USA office: 27 Warren Street, Suite 401-402, Hackensack, NJ 07601

UK office: 57 Shelton Street, Covent Garden, London WC2H 9HE

British Library Cataloguing-in-Publication Data

A catalogue record for this book is available from the British Library.

Selected Topics in Electronics and Systems — Vol. 53

FRONTIERS IN ELECTRONICS

Selected Papers from the Workshop on Frontiers in Electronics 2011 (WOFE-11)

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ISBN 978-981-4536-84-4

Printed in Singapore

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InAs_{1-x}Sb_x ALLOYS WITH NATIVE LATTICE PARAMETERS GROWN ON COMPOSITIONALLY GRADED BUFFERS: STRUCTURAL AND OPTICAL PROPERTIES

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GaInSb and AlGaInSb compositionally graded buffer layers grown on GaSb by MBE were used to develop unrelaxed InAs_{1-x}Sb_x epitaxial alloys with strain-free native lattice constants up to 2.1% larger than that of GaSb. The in-plane lattice constant of the strained top buffer layer was grown to be equal to the native, unstrained lattice constant of InAs_{1-x}Sb_x with given x . The InAs_{0.56}Sb_{0.44} layers demonstrated a photoluminescence (PL) peak at 9.4 μm at $T = 150$ K. The minority carrier lifetime measured at 77 K for InAs_{0.8}Sb_{0.2} was 250 ns.

Keywords: InAsSb; compositionally graded buffer; MBE; infrared; minority carrier lifetime; reciprocal space mapping.

Introduction

GaSb based III-V materials are widely used in the development of mid- and long-wave infrared optoelectronic devices because of the narrow bandgap and the flexibility in forming heterojunctions with various types of band offsets, i.e. type I, type II staggered or type II broken gap. For device applications, heterostructures with a considerable thickness are preferably grown lattice-matched or nearly lattice-matched to the substrate. Therefore, the device design is restricted by the lattice parameters of commercially available III-V substrates. For example, in the case of GaSb based type I lasers, the content of As in InGaAsSb quantum wells must be high enough to satisfy the conditions of pseudomorphic growth, but high As content in quantum wells severely affects the device performance [1]. In principle, the problem can be addressed by the epitaxial growth of lattice-mismatched materials of the desired lattice parameters.

The key issue in mismatched epitaxy is to minimize the dislocations that penetrate through the epi-structures. In this work, we expand the lattice parameter of the GaSb substrate by growing linearly compositionally graded Ga(Al)InSb buffers, following the approach in [2-3]. The graded strain in the buffer layers facilitates the glide of threading dislocations and reduces the densities of dislocations that propagate through the buffer layer into the device [2]. High quality InAs_{1-x}Sb_x layers having non-tetragonally distorted, strain-free lattice parameters were grown on top of the buffer layers with thickness up to 1.5 μm .

$\text{InAs}_{1-x}\text{Sb}_x$ alloys are of special interest, because the bowing effect in the band gap E_g is dependent on the Sb composition, which allows the growth of layers having bandgaps narrower than that in InSb [4-15]. In the second part of the paper, we present the optical properties of non-distorted $\text{InAs}_{1-x}\text{Sb}_x$ alloys grown on linearly compositionally graded Ga(Al)InSb buffers. Strong PL was observed for $\text{InAs}_{1-x}\text{Sb}_x$ alloys in a wide temperature range. A relatively long carrier lifetime was obtained in $\text{InAs}_{0.8}\text{Sb}_{0.2}$ alloys from the PL response to modulated optical excitation.

Growth and Structural Characterization

The heterostructures were grown on GaSb substrates by solid-source molecular beam epitaxy using a Veeco GEN-930 system equipped with As and Sb valved cracker sources. Molecular beam fluxes were measured by an ion gauge positioned in the beam path. The substrate temperature was controlled by a pyrometer, which was calibrated using references such as the III to V enriched surface reconstruction transition, oxide desorption and the melting point of InSb. The compositionally graded 2–3.5 μm thick Ga(Al)InSb buffer layers were grown at temperatures ranging from 460 to 520°C. The growth temperature was maintained near 415°C for the InAsSb layers. The Sb incorporation was controlled by the adjustment of the relative pressure of the As and Sb group V elements as measured by the beam-flux-monitor. The growth rate was about 1 μm per hour. $\text{InAs}_{1-x}\text{Sb}_x$ layers with $x = 20, 30$ and 44% were grown on GaInSb and AlGaInSb buffers.

The defect distribution in linearly compositionally graded GaInSb and AlGaInSb buffers were characterized by cross-sectional TEM images. Figure 1 shows the XTEM images of structures with either laser or absorber layers grown on top of three different linearly graded buffer layers; including (a) GaInSb with top In content of 16%; (b) GaInSb with top In content of 30%; (c) AlGaInSb with top Al, Ga and In contents of 75, 0 and 25 %, respectively. The images were taken with a (220) bright field two-beam condition to emphasize the dislocations. In all three structures, the misfit dislocation network was confined in the bottom part ($\sim 1.5 \mu\text{m}$) of the graded buffers; the topmost portion of the buffers as well as the epi-structures grown onto the buffers is free from misfit dislocations. TEM results did not show any noticeable difference in the dislocation morphology of these two buffer layers or in the laser or absorber layer structures grown on top of them, both appear to be equally efficient in accommodating the misfit strain. From the images, we can estimate that the threading dislocation density is below 10^7 cm^{-2} in the $\text{InAs}_{1-x}\text{Sb}_x$ layers.

The surface morphology was characterized by atomic force microscopy (AFM) in tapping mode (AFM Dimension V). Cross-hatched patterns with crossing lines along the [110] crystallographic directions were observed on all structures. However, structures grown on AlGaInSb buffers showed better surface morphology. Figure 2 (a) and (b) show the AFM amplitude images measured over a 50 by 50 μm area for samples



Fig. 1. Cross-sectional TEM images of heterostructures: (a) GaInSb with top In content of 16%; (b) GaInSb with top In content of 30%; (c) AlGaInSb with top Al, Ga and In contents of 75, 0 and 25 %, respectively.



Fig. 2. AFM amplitude images measured over a 50 by 50 μm area for samples: (a) GaInSb buffer and (b) AlGaInSb buffer.

with $\text{InAs}_{0.8}\text{Sb}_{0.2}$ layers. The surface roughness is both nearly twice as shown in the image of the surface roughness, i.e. increasing the Sb content as indicated by surface roughness.

Strain relaxation was characterized by reciprocal space mapping (RSM) of the (335) Bragg reflection. The RSM of the (335) Bragg reflection of the AlGaInSb buffer layer is larger than that of the GaInSb buffer layer. The RSM of the (335) Bragg reflection of the topmost section of the

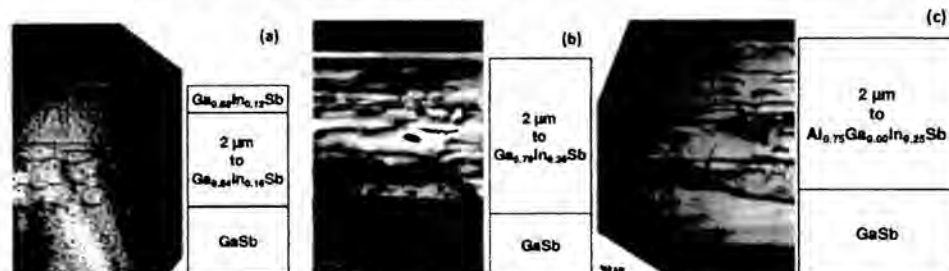


Fig. 1. Cross-sectional TEM images of samples with 2 μm thick linearly graded buffers grown on GaSb substrates: (a) GaInSb with top In content of 16% - mismatch accommodated 0.9%; (b) GaInSb with top In content of 30% - mismatch accommodated 1.4%; (c) AlGaInSb with top Al, Ga and In contents of 75, 0 and 25% - mismatch accommodated 1.4%.

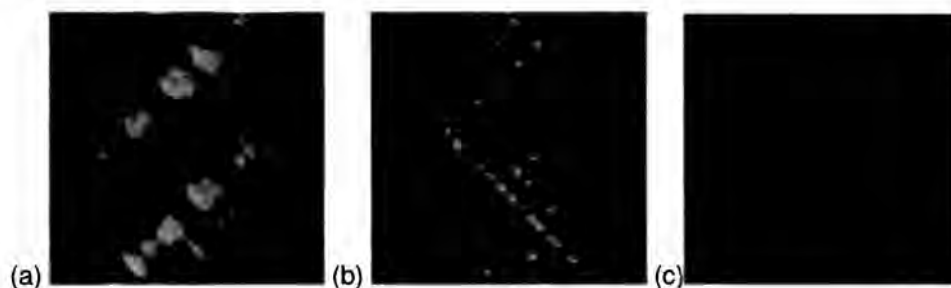


Fig. 2. AFM amplitude images measured over 50 by 50 μm area for samples with InAs_{0.8}Sb_{0.2} layer grown on (a) GaInSb buffer and (b) AlGaInSb buffer; (c) shows the enlarged image (3 by 3 μm) of sample (b).

with InAs_{0.8}Sb_{0.2} layer grown on a (a) GaInSb buffer and a (b) AlGaInSb buffer. The undulation amplitude and period in sample (a) were about 10 nm and 9 μm , respectively, both nearly twice as much the $\sim 5\text{ nm}$ and $\sim 5\text{ }\mu\text{m}$ measured for sample (b). Figure 2 (c) shows the image of sample (b) measured over 3 by 3 μm area; the root mean square surface roughness, i.e., in between of the dips in cross-hatch pattern, was below 1 nm. Increasing the Sb content led to larger peak-to-peak variations in the cross-hatch pattern, as indicated by surface roughness up to 10 nm for the InAs_{0.56}Sb_{0.44} samples.

Strain relaxation of the structures was examined using high-resolution X-ray diffraction reciprocal-space mapping (RSM) at the symmetric (004) and asymmetric (335) Bragg reflections. Figure 3 presents a set of RSM measurements for a structure consisting of a 1 μm InAs_{0.8}Sb_{0.2} layer grown on a 2 μm linearly compositionally graded AlGaInSb buffer layer. The native lattice constant of the InAs_{0.8}Sb_{0.2} layers is about 0.8% larger than that of GaSb. The native lattice constant of the buffer layer changed from that of GaSb to that of Al_{0.75}Ga_{0.13}In_{0.12}Sb with a strain ramp rate about 0.6% per μm . The topmost section of the graded buffer with Al_{0.75}Ga_{0.13}In_{0.12}Sb composition had a native

effect in the band gap E_g layers having bandgaps, we present the optical compositionally graded in a wide temperature $b_{0.2}$ alloys from the PL

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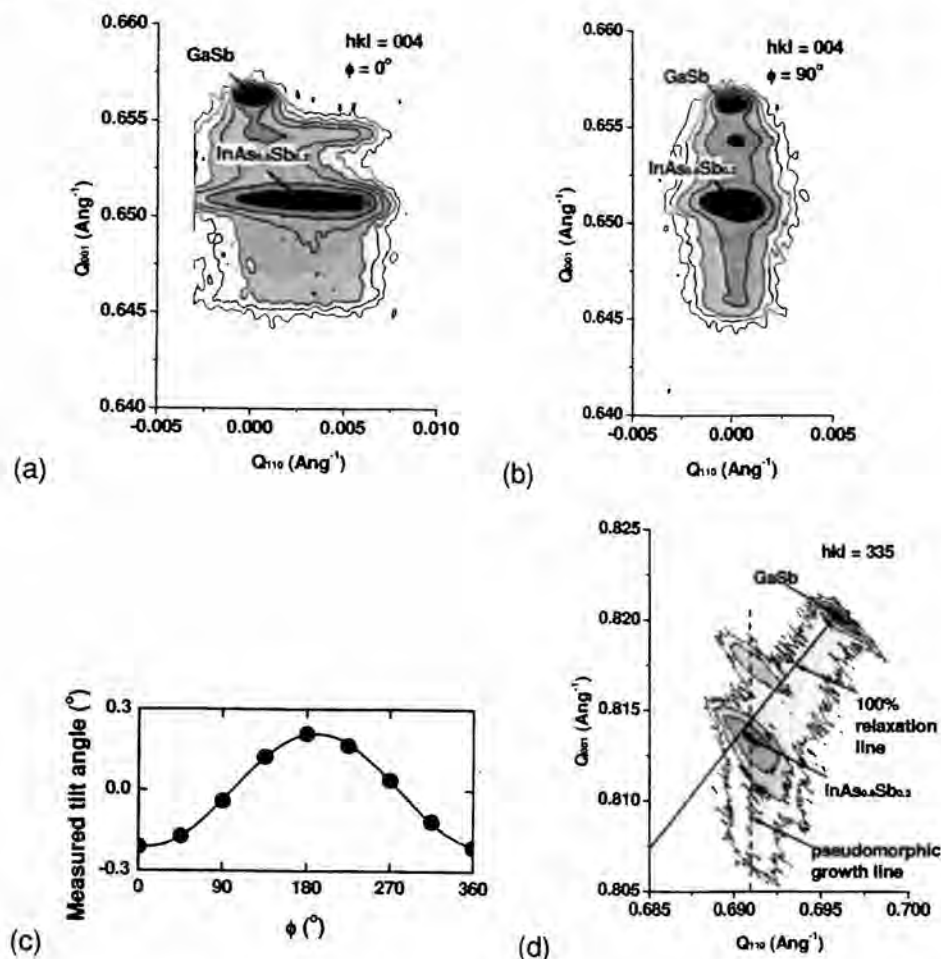


Fig. 3. (a) Symmetric (004) RSM taken at the azimuth angle emphasizing the tilt in the epi-layers; (b) (004) RSM taken at the azimuth angle minimizing the tilt in the epi-layers; (c) dependence of the measured tilt angle as a function of the azimuth angle; (d) asymmetric (335) RSM taken at azimuth angle equal to 90° . Solid line denotes the location of 335 reflexes corresponding to fully relaxed material with lattice parameter gradually increasing from that of GaSb. Dashed line denotes the location of 335 reflexes of the material with further increasing native lattice parameter but grown pseudomorphically to the top of fully relaxed section.

lattice constant about 1.3% larger than that of GaSb, but due to compressive strain, the in-plane lattice constant is equal to the native constant of the bulk InAs_{0.8}Sb_{0.2}. When the final structure was grown, the InAsSb layer was sandwiched between Al_{0.75}Ga_{0.13}In_{0.12}Sb carrier confinement layers to assist photoluminescence experiments.

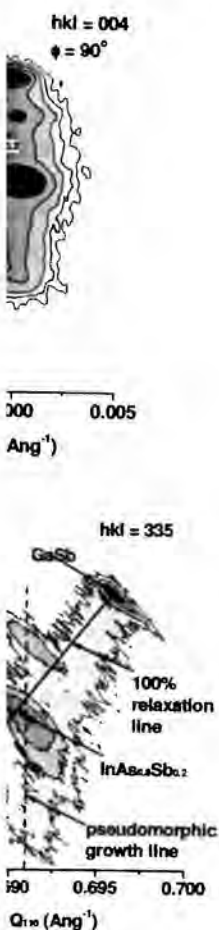
The symmetric reflection revealed the tilt present in the epi-structure. Figure 3 (a) and (b) shows the RSMs obtained near the symmetric (004) reflection at two azimuth angles, namely (a) $\phi = 0^\circ$ and (b) $\phi = 90^\circ$, corresponding to two perpendicular [110] crystallographic directions. The tilt angle projected to the measurement plane is

determined from the layers. As shown in the bottom part of the figure, the projected tilt angle is 0.2° in the direction of the [110] crystallographic axis.

Asymmetric (335) RSMs were taken in order to characterize the InAs_{0.8}Sb_{0.2} layer. Figure 3 (d) shows the asymmetric (335) RSM taken at $\phi = 90^\circ$, i.e., with the minimum tilt. The transition from the relaxed to the strained state is observed. For illustrative purposes, the relaxed section observed relaxation in this section and within our experimental range. The topmost section of the buffer layer is InAs_{0.8}Sb_{0.2} layer pseudomorphically relaxed to the lattice constant of the GaSb substrate, i.e., below 0.1%; therefore, the InAsSb reflection is not observed. The strained auxiliary layer is used for calibration purposes.

Optical Characterization

The PL and absorption spectra were measured using a spectrometer equipped with a 1200 W laser and collected in a temperature range of 10–300 K. The PL spectra show the PL spectra at 13, 150 and 300 K. The InAs_{0.56}Sb_{0.44} layer is grown on InAs/GaSb superlattice. The PL spectra show 300 periods of InAs/GaSb superlattice. The AlAsSb carrier confinement layer is grown on the InAs/GaSb superlattice. The width at half maximum intensity is determined from the PL spectra.



the epi-layers; (b) (004) of the measured tilt angle equal to 90°. Solid line lattice parameter gradually the material with further axed section.

pressive strain, the As_{0.8}Sb_{0.2}. When the n Al_{0.75}Ga_{0.13}In_{0.12}Sb

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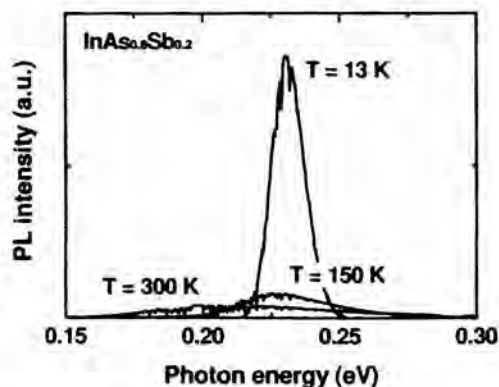
determined from the horizontal peak separation between the GaSb substrate and the epi-layers. As shown in Figure 3 (a), the tilt angle increases as the thickness increases in the bottom part of the graded buffer, and stops increasing in the consequent layers. The bottom part of the buffer layers is near completely relaxed, as will be shown later, suggesting that tilting is associated with the process of strain relaxation. Figure 3 (c) plots the projected tilt angle as a function of several azimuth angle ϕ . We estimate the tilt angle to be 0.2° in the direction about 10° away from the [110] direction ($\phi = 90^\circ$).

Asymmetric (335) RSM reflexes were measured at four different azimuth angles in order to characterize the degree of relaxation of the graded buffer layer and to confirm that the InAs_{0.8}Sb_{0.2} layer is lattice-matched to the topmost part of the graded buffer. Figure 3 (d) shows one of the (335) RSMs measured at an azimuth angle equal to 90°, i.e., with the minimum tilting effect. The shift visible in the (335) RSM corresponds to the transition from the strain relaxed to the pseudomorphic section of the graded buffer. For illustrative purposes, the solid line corresponds to a 100% relaxed square lattice. The observed relaxation is close to 100%. After the tilt angle is accounted for, the degree of relaxed in this section of the graded buffer can be estimated as 95%, i.e., nearly 100%, and within our experimental error. The pseudomorphic growth of the dislocation-free topmost section of the buffer layer is apparent from the (335) scan since the reflex from the buffer layer is nearly vertical (dashed line in Figure 3 (d)). The reflection from the InAs_{0.8}Sb_{0.2} layer is located at the turning point and on the same vertical line as the pseudomorphic section of the buffer, which confirms lattice matching to the in-plane lattice constant of the graded buffer layer. The amount of strain in the InAs_{0.8}Sb_{0.2} layer is below 0.1%; therefore, no strain relaxation is expected. The reflection located above the InAsSb reflection in both the (004) and (335) RSM corresponds to a pseudomorphically strained auxiliary AlGaSb layer (~150 nm) that was grown on top of the InAsSb layer for calibration purposes.

Optical Characterization

The PL and absorption spectra were measured with a Fourier-transform infrared (FTIR) spectrometer equipped with a liquid-nitrogen cooled HgCdTe detector with a cut-off wavelength of 12 μm . The PL was excited by either a 970 nm laser diode or a Nd:YAG laser and collected by reflective optics. PL was observed from all structures in a wide temperature range, up to room temperature from samples with 20% Sb. Figure 4 (a) shows the PL spectra from 1- μm thick InAs_{0.8}Sb_{0.2} layer grown on an AlGaInSb buffer at 13, 150 and 300 K. Figure 4 (b) presents the PL spectra measured from a 1- μm thick InAs_{0.56}Sb_{0.44} layer grown on an AlGaInSb buffer and a 1.8- μm thick long-wave InAs/GaSb superlattice grown on a GaSb substrate. The superlattice structure consists of 300 periods of InAs and GaSb layers with the cell period of 63 Å enclosed within 20-nm AlAsSb carrier confinement layers lattice-matched to GaSb. The spectral widths (full-width at half maximum) for the three samples were similar, about 11 meV. The PL intensities from both InAs_{0.56}Sb_{0.44} and superlattices were comparable at 13 K while

(a)



(b)

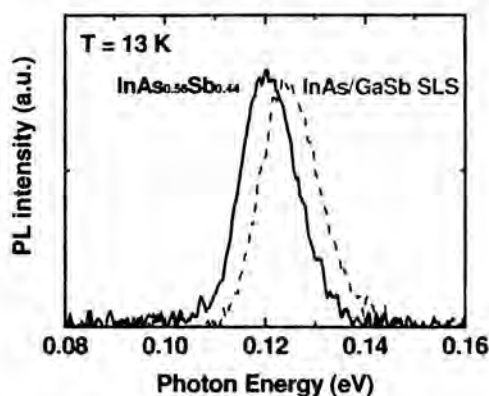


Fig. 4. (a) PL spectra from InAs_{0.8}Sb_{0.2} sample grown on AlGaInSb buffer at 13 K, 150 K and 300 K under an excitation power of 0.5 W. (b) PL spectra from InAs_{0.56}Sb_{0.44} layer grown on AlGaInSb buffer and long-wave InAs/GaSb superlattices grown on GaSb substrate at 13 K under an excitation of 0.1 W. The PL was excited by a Nd:YAG laser with a beam diameter of about 0.5 mm.

intensities drops much faster in InAsSb sample at elevated temperatures. Considering the challenge of creating adequate hole confinement in the As-rich alloys, the faster drop of the PL intensity with temperature can be explained by the increased diffusion of the excess carriers out of the InAsSb layer.

The absorption spectra were measured for the InAsSb layers with Sb compositions of 20% and 30% grown on GaInSb buffers. The absorbance was determined from transmission measurements taking into account the multiple reflections. The absorption spectrum was derived by subtracting the absorbance of the heterostructure with the epi-layers and the substrate. The transmission of the substrate was determined using the

same sample after the sample with 20% Sb was thinned. It was lapped down to the GaSb substrate at longer wavelengths. The GaSb substrate thickness of 55 μm was determined by transmission for the substrate. Both absorption and PL spectra are presented in Figure 5. The PL was associated

Fig. 5. Absorption and PL spectra of InAs_{0.8}Sb_{0.2} sample excited by a 970-nm laser. The PL near 0.29 eV was caused

same sample after the epi-layers were removed by polishing. The substrate of the sample with 20% Sb was thinned to 300 μm . The substrate of the heterostructure with 30% Sb was lapped down to near 50 μm thickness because of high free carrier absorption in the GaSb at longer wavelengths. The latter was determined to be 140 cm^{-1} near $\lambda = 8\text{ }\mu\text{m}$ for the GaSb substrates with Te doping level of $3 \times 10^{18}\text{ cm}^{-3}$. The sample with the measured thickness of 55 μm had near 50% transmission at this wavelength, as compared to a 2% transmission for the 300 μm -thick substrate. The free carrier absorption in the thin substrate was determined by a fit based on the absorption measurements for the thicker substrate. Both absorption and PL spectra measured for the two samples at 150 K are presented in Figure 5. The PL peak energy matched the absorption edge indicating that the PL was associated with the band-to-band recombination.

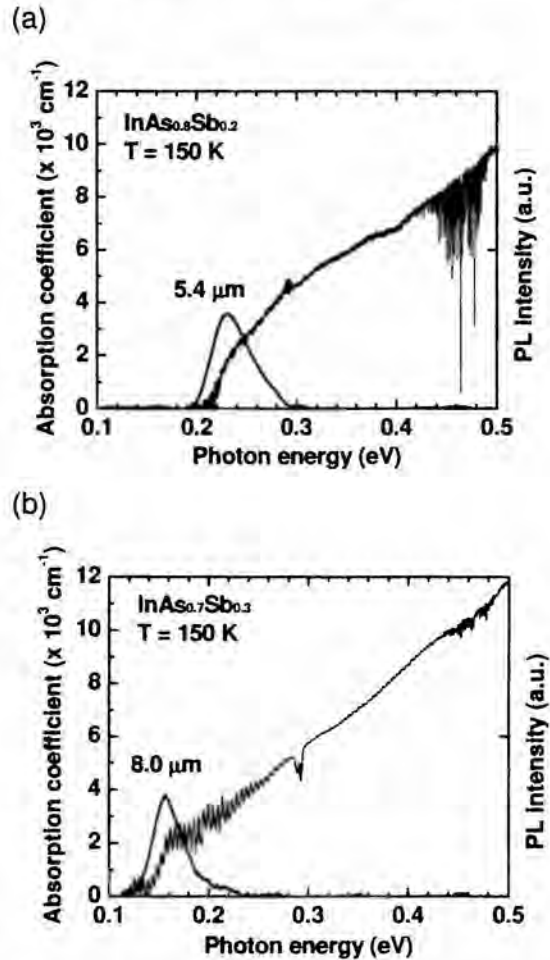


Fig. 5. Absorption and PL spectra measured at 150 K for (a) $\text{InAs}_{0.8}\text{Sb}_{0.2}$ and (b) $\text{InAs}_{0.7}\text{Sb}_{0.3}$. The PL was excited by a 970-nm laser diode at a power of 400 mW; the excitation area was $2.5 \times 10^{-3}\text{ cm}^2$. The distortion near 0.29 eV was caused by CO_2 absorption.

The energy positions of PL maxima at $T = 13$ K versus Sb composition x in the InAsSb layers are presented in Figure 6. The positions of PL maxima were used to determine the bandgaps and the bowing parameter, which was about 0.9 eV, considerably greater than the recommended value of 0.7 eV [15]. The lower value of bowing reported previously was based on measurements in materials grown without control of the strain relaxation. The observed difference in the bowing between the 0.9 eV determined in this work and the 0.7 eV reported in literature can be explained by the absence of residual strain in the InAsSb epitaxial layers.

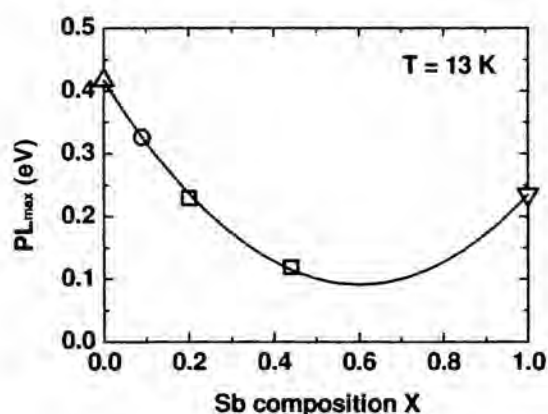


Fig. 6. Dependence of 13K PL maxima on composition X in InAs_{1-x}Sb_x epitaxial layers: InAs epilayer grown on InAs substrate (triangle), InAsSb_{0.08} epilayer grown lattice matched to GaSb substrate (circle), InAs_{1-x}Sb_x epilayers grown on AlGaInSb buffers (squares) on GaSb substrate, InSb epilayers grown on InSb substrate (inverted triangles).

Carrier lifetime measurements for the 1- μ m thick InAs_{0.8}Sb_{0.2} layer grown on AlGaInSb buffer layer were performed at $T = 77$ K using optical modulation response technique [16]. To minimize the effects of carrier separation on the carrier lifetime in undoped InAsSb layers, the Al_{0.25}Ga_{0.70}In_{0.05}Sb barriers were doped with Be to the level of $1 \times 10^{17} \text{ cm}^{-3}$. The dependence of the carrier lifetime on the excitation power is shown in Figure 7. Simulation of the band diagram showed that the minority holes remain confined under low excitation (inset of Figure 7 (b)). The carrier lifetime was determined from the PL response to a small signal modulation of excitation in the frequency domain. The PL response spectra in a range of continuous-wave excitation power are shown in Figure 7 (a). The carrier lifetime τ corresponding to the cut-off frequency (-3dB point) was obtained by fitting the response in the entire frequency range to the dependence $PL_{\omega} \propto [1 + (2\pi f\tau)^2]^{-1/2}$. A 250 ns carrier lifetime under low excitation condition was measured. The excess carrier concentration was estimated to be in the range $(2-4) \times 10^{15} \text{ cm}^{-3}$ at the excitation power in the range of 0.5-1 W/cm².

Fig. 7. Carrier lifetime measurements for the 1- μ m thick InAs_{0.8}Sb_{0.2} layer grown on AlGaInSb buffer on a GaSb substrate. (a) PL response spectra for excitation levels of 0.8, 1, 1.4, 1.31 μ m, and the excitation power used for carrier lifetime measurements.

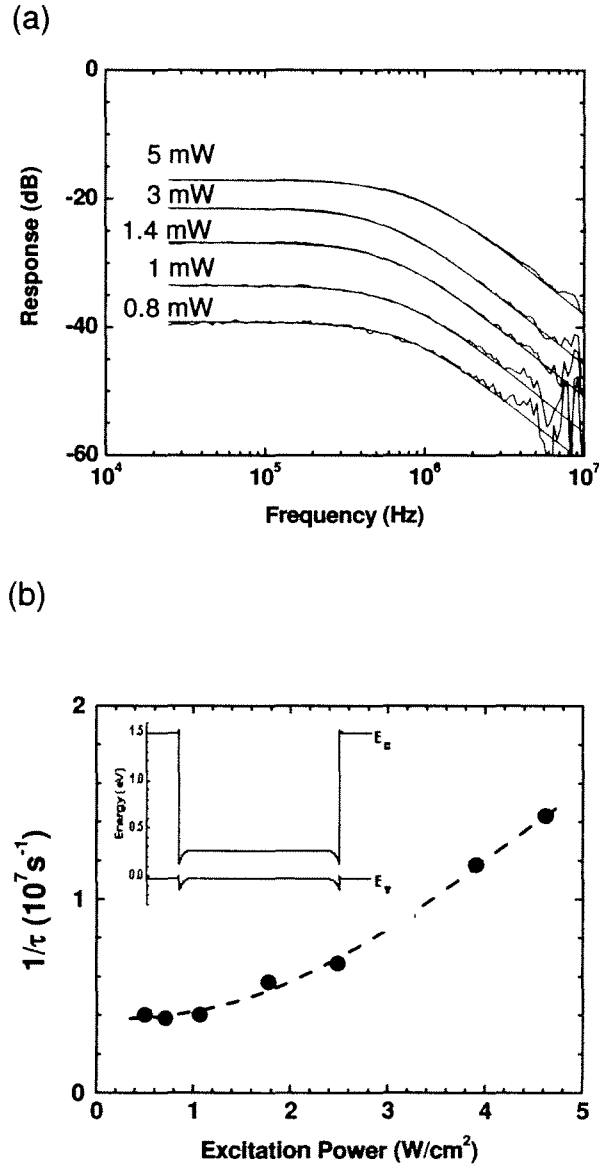


Fig. 7. Carrier lifetime measurements at $T = 77$ K on a $1\text{-}\mu\text{m}$ -thick $\text{InAs}_{0.8}\text{Sb}_{0.2}$ layer grown on an AlGaInSb buffer on a GaSb substrate. The PL responses and fits are presented for continuous wave excitation power levels of 0.8, 1, 1.4, 3 and 5 mW from bottom to top, respectively. The PL was excited at the wavelength of $1.31\text{ }\mu\text{m}$, and the excitation area was $2 \times 10^{-3} \text{ cm}^2$ FWHM (left). The reciprocal carrier lifetime is plotted versus continuous-wave excitation power density (right). A schematic band diagram of the InAsSb heterostructure used for carrier lifetime measurements is shown in the inset.

Conclusion

In summary, we conclude that growing compositionally graded buffers (Ga(Al)InSb on GaSb substrates) with a strained but unrelaxed top layer allows the fabrication of bulk $\text{InAs}_{1-x}\text{Sb}_x$ layers ($0.5\sim 1.5\ \mu\text{m}$ thick). These films have characteristics that are promising for the development of IR detectors operating within the spectral range from 5 to $12\ \mu\text{m}$. The critical element of the technology is the control of the in-plane lattice constant of the topmost section of the buffer. The in-plane lattice constant of this layer must be equal to the lattice constant of $\text{InAs}_{1-x}\text{Sb}_x$ with given x . The unrelaxed $\text{InAs}_{1-x}\text{Sb}_x$ epitaxial layers grown on top of such buffers demonstrated photoluminescence in the spectral range from 5.2 to $9.4\ \mu\text{m}$ within the temperature range of $77\sim 150\ \text{K}$. The carrier lifetime of $250\ \text{ns}$ was obtained at $T = 77\ \text{K}$ for structure consisting $\text{InAs}_{0.8}\text{Sb}_{0.2}$ epi-layers.

Acknowledgement

This work was supported by the National Science Foundation, Grant DMR0710154 and Army Research Office, Grant W911NF11-1-0109. The authors would like to thank IQE, Inc. for providing the superlattice samples.

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We discuss the lasers. Broad-area $375\ \text{A}/\text{cm}^2$ and the suppression incorporate a p suppress lasing feedback. A co output in a sin; single-mode op tuning range is produces up to the strongest ab

Keywords: Inter

1. Introduction

In 2002, a quantum well for the $3\text{--}5\ \mu\text{m}$ n room temperature have recently achieved increasingly more conduction-band quantum-well lasers only to wavelengths

The interband based on InAs e Auger non-radiative quantum cascade is an interband p conduction and v electron and hole back to the con wavelengths when voltage can be successfully cascade architecture