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**ROCK ISLAND ARSENAL** 3.  
**RESEARCH & ENGINEERING DIVISION**  
**RESEARCH LABORATORIES**



**TECHNICAL NOTES**

MISCELLANEOUS ARTILLERY PROBLEMS

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# CONTENTS

<u>Section</u>	<u>page</u>
Derivation of the Moment-Area Method for Application to the Design of Control Rods.	1
Dynamic Analysis of Firing Platform.	8
Gas Flow Studies	16
Determine the Relationship Between Momentum Index (B) and Intrinsic Efficiency (I.E.)	24
Appendix A - Propellant Form Function	26
Appendix B - Thermochemistry of Propellant.	40

$l(t) = \text{centroid of } B(t)$

~~$B(t)$~~   $B(t) = \text{centroid of } R(t)$

$R(t) = \text{resisting force}$

## Derivation of the Moment-Area Method for Application to the Design of Control Rods

Given the differential equation for recoil

$$M\ddot{\psi} = B(t) - R(t) \quad (1)$$

Integrating

$$M\dot{\psi} = \int_a^t B(\tau) d\tau - \int_b^t R(\tau) d\tau \quad (2)$$

where  $B(t)$  is applied when  $t = a$ ; and  $R(t)$  is applied when  $t = b$ . Define

$$\int_a^t B(\tau) d\tau = G(t) \quad \int_b^t R(\tau) d\tau = H(t) \quad (3)$$

so

$$M\dot{\psi} = G(t) - H(t) \quad (4)$$

Integrating

$$M\psi = \int_a^t G(\tau) d\tau - \int_b^t H(\tau) d\tau \quad (5)$$

Now, integrating by parts

$$M\psi = \tau G(\tau) \Big|_a^t - \int_a^t \tau \frac{d}{d\tau} G(\tau) d\tau - \tau H(\tau) \Big|_b^t + \int_b^t \tau \frac{d}{d\tau} H(\tau) d\tau$$

or

$$M\psi = tG(t) - \int_a^t \tau B(\tau) d\tau - tH(t) + \int_b^t \tau R(\tau) d\tau$$

Rewrite as

$$M\psi = tG(t) - \frac{\int_a^t B(\tau) d\tau}{\int_a^t B(\tau) d\tau} \cdot \int_a^t \tau B(\tau) d\tau - tH(t) + \frac{\int_b^t R(\tau) d\tau}{\int_b^t R(\tau) d\tau} \cdot \int_b^t \tau R(\tau) d\tau \quad (6)$$

Define

$$\alpha(t) = \frac{\int_a^t \tau B(\tau) d\tau}{\int_a^t B(\tau) d\tau} \quad \beta(t) = \frac{\int_b^t \tau R(\tau) d\tau}{\int_b^t R(\tau) d\tau}$$

Note that  $\alpha(t)$  and  $\beta(t)$  are expressions for the centroids of the areas generated by the functions  $B(t)$  and  $R(t)$ , at time  $t$ , measured from  $t = 0$ .

Now, write equation 6 as

$$M\psi = tG(t) - \alpha(t) \cdot G(t) - tH(t) + \beta(t) \cdot H(t) \quad (7)$$

$R(t)$  = resistance to recoil

$t_R$  = recoil time

$b$  = time when  $R(t)$  is applied

$Q$  = centroid of  $B(t)$  curve

$p$  = centroid of  $R(t)$  curve

or

$$Mx = (t - \alpha) \int_0^t B(\tau) d\tau - (t - \beta) \int_0^t R(\tau) d\tau \quad (8)$$

At the end of recoil ( $x=L$ ) and ( $t=t_r$ ) the following relations hold

$$\int_0^{t_r} B(\tau) d\tau = \int_0^{t_r} R(\tau) d\tau = I \quad (\text{a constant}) \quad (9)$$

$$ML = [(t_r - \alpha) - (t_r - \beta)] I \quad (10)$$

or

$$\beta(t_r) = \frac{ML}{I} + \alpha(t_r) \quad (11)$$

Assume that the shape of the "total resistance to recoil" force,  $R(t)$ , is known. Then recalling

$$\beta(t) = \frac{\int_0^t \tau R(\tau) d\tau}{\int_0^t R(\tau) d\tau}$$

write

$$\beta(t_r) \cdot \int_0^{t_r} R(\tau) d\tau = ML + I\alpha = \int_0^{t_r} \tau R(\tau) d\tau$$

Now, from the two relations

$$I = \int_0^{t_r} R(\tau) d\tau \quad (12)$$

$$ML + I\alpha = \int_0^{t_r} \tau R(\tau) d\tau \quad (13)$$

the magnitude of  $R(t)$  and  $t_r$  may be determined.

CASE 1:  $a = b = 0$ ,  $t_r > t_0$  where  $t_0$  is the time  $B(t)$  acts.

Assume the following data to be known.

$$M = 3.8 \text{ * } 50\% \text{ in} \quad I = 200\theta \text{ * } \text{sec} \quad \alpha = 0.007 \text{ sec} \quad L = 4\theta \text{ in.}$$

**EXAMPLE 1:**

Assume a continuous function for  $R(t)$ , namely

$$R(t) = R_0 \sin \frac{\pi}{t_r} t$$

then

$$I = R_0 \int_0^{t_r} \sin \frac{\pi}{t_r} t dt = -R_0 \frac{t_r}{\pi} \cos \frac{\pi}{t_r} t \Big|_0^{t_r} = \frac{2R_0 t_r}{\pi}$$

$$\begin{aligned} ML + I\alpha &= R_0 \int_0^{t_r} t \sin \frac{\pi}{t_r} t dt = R_0 \left[ t \left( -\frac{t_r}{\pi} \right) \cos \frac{\pi}{t_r} t \right]_0^{t_r} + R_0 \frac{t_r}{\pi} \int_0^{t_r} \cos \frac{\pi}{t_r} t dt \\ &= R_0 \frac{t_r^2}{\pi} + R_0 \left( \frac{t_r}{\pi} \right)^2 \sin \frac{\pi}{t_r} t \Big|_0^{t_r} \\ &= R_0 \cdot \frac{t_r^2}{\pi} \end{aligned}$$

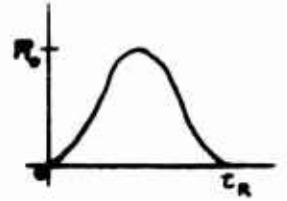
thus

$$\pi I = 2R_0 t_R \quad \pi(ML + I\alpha) = R_0 t_R^2$$

or

$$R_0 = \frac{\pi I^2}{4(ML + I\alpha)} = 16,100$$

$$t_R = \frac{2}{\pi} (ML + I\alpha) = 0.196$$



By the Moment-Area Method (eq. 10)

$$ML = [(t_R - \alpha) - (t_R - \beta)] I$$

Obviously  $\beta = \frac{1}{2} t_R$ , therefore

$$ML = [(t_R - \alpha) - (t_R - \frac{1}{2} t_R)] I$$

or

$$t_R = \frac{2}{I} (ML + I\alpha) = 0.196$$

as before.

### EXAMPLE 2:

Using the same data and assuming

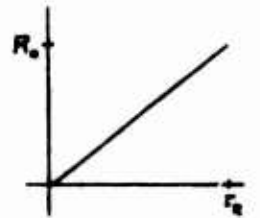
$$R(t) = R_0 \frac{t}{t_R}$$

then

$$I = \frac{R_0}{t_R} \int_0^{t_R} t dt = \frac{1}{2} R_0 t_R$$

and

$$ML + I\alpha = \frac{R_0}{t_R} \int_0^{t_R} t^2 dt = \frac{1}{3} R_0 t_R^2$$



thus

$$2I = R_0 t_R$$

$$3(ML + I\alpha) = R_0 t_R^2$$

or

$$R_0 = \frac{4I^2}{3(ML + I\alpha)} = 27,400$$

$$t_R = \frac{3(ML + I\alpha)}{2I} = 0.147$$

### EXAMPLE 3:

Using the same data and assuming

$$R(t) = R_0 \left(1 - \frac{t}{t_R}\right)$$

now

$$I = \int_0^{t_R} R_0 \left(1 - \frac{t}{t_R}\right) dt = R_0 \left(t - \frac{t^2}{2t_R}\right) \Big|_0^{t_R} = \frac{1}{2} R_0 t_R$$

and

$$ML + I\alpha = \int_0^{t_R} t \cdot R_0 \left(1 - \frac{t}{t_R}\right) dt = R_0 \left(\frac{t^2}{2} - \frac{t^3}{3t_R}\right) \Big|_0^{t_R} = \frac{1}{6} R_0 t_R^2$$

thus

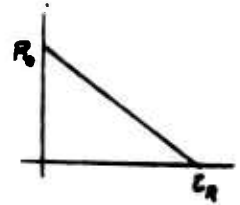
$$R_0 t_R = 2I$$

$$R_0 t_R^2 = 6(ML + I\alpha)$$

or

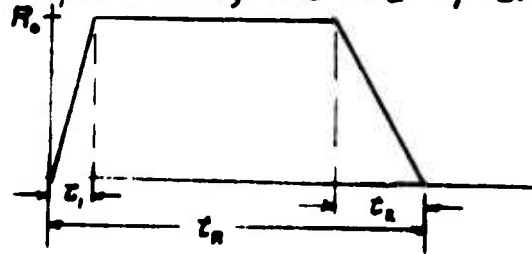
$$R_0 = \frac{2I^2}{3(ML + I\alpha)} = 13,680$$

$$t_R = \frac{3(ML + I\alpha)}{I} = 0.294$$



### EXAMPLE 4

Given the same data with  $R(t)$  defined by the accompanying figure. Also, assume  $t_1 = 0.005$  and  $t_2 = 0.010$ .



Now

$$\begin{aligned} I &= \int_0^{t_R} R(t) dt = \int_0^{t_1} \frac{R_0}{t_1} t dt + \int_{t_1}^{t_1+t_2} R_0 dt + \int_{t_1+t_2}^{t_R} R_0 \left(1 - \frac{t-t_1-t_2}{t_2}\right) dt \\ &= \frac{1}{2} R_0 t_1 + R_0 (t_2 - t_1) + \frac{1}{2} R_0 t_2 \end{aligned}$$

and

$$\begin{aligned} (ML + I\alpha) &= \int_0^{t_R} t R(t) dt = \int_0^{t_1} \frac{R_0}{t_1} t^2 dt + \int_{t_1}^{t_1+t_2} R_0 t dt + \int_{t_1+t_2}^{t_R} R_0 \cdot \frac{t_2-t}{t_2} t dt \\ &= \frac{1}{3} R_0 t_1^2 + \frac{1}{2} R_0 [(t_1+t_2)^2 - t_1^2] + R_0 t_2 \left(\frac{1}{2} t_2 - \frac{1}{3} t_2\right) \end{aligned}$$

or

$$I = \frac{1}{2} R_0 (2t_R - t_1 - t_2)$$

$$(ML + Id) = \frac{R_0}{6} [3t_R(t_R - t_2) - t_1^2 + t_2^2]$$

Note: the same results would be obtained by the moment-area method of equation B. That is:

$$ML = (t_R - \alpha) \int_0^{t_R} B(\tau) d\tau - (t_R - \beta) \int_0^{t_R} R(\tau) d\tau$$

or

$$ML = (t_R - \alpha) I - \left\{ \frac{1}{2} R_0 t_1 \cdot (t_R - \frac{2}{3} t_1) + R_0 (t_R - t_2 - t_1) \left[ t_2 + \frac{1}{2} (t_R - t_1 - t_2) \right] + \frac{1}{2} R_0 t_2 \cdot \frac{2}{3} t_2 \right\}$$

but

$$I = \frac{1}{2} R_0 (2t_R - t_1 - t_2)$$

thus

$$ML + Id = \frac{1}{2} R_0 t_R (2t_R - t_1 - t_2) - \left\{ \frac{1}{2} R_0 t_1 \cdot (t_R - \frac{2}{3} t_1) + R_0 (t_R - t_2 - t_1) \left[ t_2 + \frac{1}{2} (t_R - t_1 - t_2) \right] + \frac{1}{2} R_0 t_2 \cdot \frac{2}{3} t_2 \right\}$$

Expanding and collecting terms

$$ML + Id = \frac{R_0}{6} [3t_R(t_R - t_2) - t_1^2 + t_2^2]$$

as before.

Solving for  $t_R$  and  $R_0$

$$t_R = \frac{6(ML + Id) + 3It_2 + \sqrt{[6(ML + Id) + 3It_2]^2 - 12I[I(t_2^2 - t_1^2) + 3(ML + Id)(t_1 + t_2)]}}{6I}$$

So

$$t_R = 0.203 \quad R_0 = 10,280$$

Note the simplicity if  $t_1 = t_2$ . (Assume  $t_1 = t_2 = 0.010$ )

$$I = R_0 (t_R - t_1)$$

$$(ML + Id) = \frac{1}{2} R_0 t_R (t_R - t_1)$$

So

$$t_R = \frac{2(ML + Id)}{I} = 0.2016$$

$$R_0 = \frac{I^2}{2(ML + Id) - It_1} = 10,250$$

CASE 2:  $a \neq 0, b = 0$   $t_R - a > t_e$  Assume the following data to be known. (Firing-out-of-battery)

$$M = 3.8 \text{ } \# \text{ / in} \quad I = 2009 \text{ } \# \text{ / sec} \quad z = 0.007 \text{ sec} \quad L = 0$$

EXAMPLE 1

Define  $\alpha = -L'$  when  $t = a$ . Note:  $\alpha = a + \bar{z}$

Now

$$M\ddot{x} = -R_0$$

$$M\dot{x} = -R_0 t$$

$$Mx = -\frac{1}{2} R_0 t^2$$

At the end of the stroke

$$ML + I\alpha = \frac{1}{2} R_0 t_R^2 \quad I = R_0 t_R$$

thus

$$I\alpha = \frac{1}{2} R_0 t_R^2 = \frac{1}{2} I t_R$$

or

$$t_R = 2\alpha = 2(a + \bar{z})$$

$$R_0 = \frac{I}{2\alpha} = \frac{I}{2(a + \bar{z})}$$

Now

$$ML' = \frac{1}{2} R_0 a^2 = \frac{1}{2} \frac{I a^2}{2(a + \bar{z})} \quad (L' = 48")$$

so

$$4ML'(a + \bar{z}) = I a^2$$

or

$$I a^2 - 4ML'a - 4ML'\bar{z} = 0$$

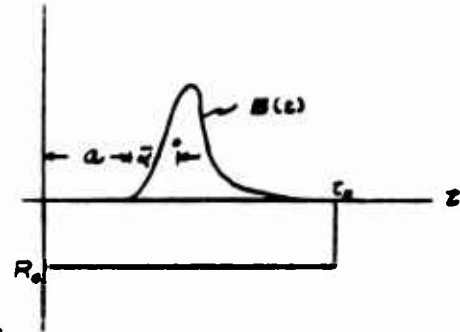
So

$$a = \frac{2ML'}{I} \left\{ 1 + \sqrt{1 + \frac{I\bar{z}}{ML'}} \right\}^{\frac{1}{2}} = 0.370$$

and

$$\alpha = 0.377 \quad t_R = 0.754 \quad R = 2660$$

Finally, note  $\dot{x} = -259$  "/sec when  $t = a$ .



### EXAMPLE 2

Define  $\dot{x} = -V = -259$  "/sec when  $t = a$ . Data is unchanged.

Now

$$-MV = -R_0 a = -\frac{Ia}{2(a+d)}$$

or

$$2MVa + 2MVd = Ia$$

so

$$a = \frac{2MVd}{I - 2MV} = 0.339$$

thus

$$\alpha = 0.346 \quad t_R = 0.692 \quad R = 2990$$

Note the sensitivity of this method. A change of 1 "/sec means a change of approximately 500 pounds. Or assume  $V = +258$  "/sec and compute  $L$ . ( $R = 2990$ )

$$(ML + I\alpha) = \frac{1}{2} I t_R$$

$$MV = R_0 a$$

$$I = R_0 t_R$$

$$\alpha = a + z$$

Solving

$$t_R = 0.692 \quad a = 0.328 \quad \alpha = 0.335 \quad L = 5.9''$$

CASE 3  $a = b = 0 \quad t_R < t_e$

Now, if  $t_R < t_e$ ; then, the equations thus far developed do not hold and it is necessary to utilize a computer for solution.

### EXAMPLE 1

Assume  $R(t)$  to be constant, and

$$M = 3.8 \text{ }^{\circ}\text{sec}^2/\text{in} \quad L = 7'' \quad I = 2009 \text{ }^{\circ}\text{sec} \quad \alpha = 0.007\text{sec}$$

then

$$I = R t_R$$

$$ML + I\alpha = \frac{1}{2} R t_R^2$$

so

$$t_R = \frac{2(ML + I\alpha)}{I} = 0.0405$$

but  $t_e = 0.088$ , thus this solution is invalid. By computer

$$t_R = 0.0393 \quad R = 50,700 \quad \alpha = 0.0063 \quad I = 1996 \text{ (when } t = t_R)$$

# DYNAMIC ANALYSIS OF FIRING PLATFORM

## 1. Introduction

Figure 1 depicts a 105 mm Howitzer mounted on a floating raft. The stability of the structure is to be determined under firing conditions.

It will be assumed that the entire structure is rigid; also, the relative motion of the recoiling parts will be neglected. Hence the applied load,  $R(t)$ -rod pull, will be the only excitation force applied. Also, it is assumed that this load will be applied through the trunnions and parallel to the center line of the gun tube.

The primary resistance is offered by the buoyant forces of the displaced fluid. Translation is resisted by drag forces. Finally, springs and dashpots are introduced to simulate anchoring.

## 2. Coordinate Systems

Assume the body to have three degrees of freedom namely  $\xi, \eta, \phi$ .

Define a coordinate system ( $O$ - $XYZ$ ) at the overall c.g. in the pre-firing condition

Define a coordinate system ( $O'$ - $X'Y'Z'$ ) with its origin fixed in the raft (at the c.g.) and at all times parallel to  $O$ - $XYZ$ .

Define a coordinate system ( $O'$ - $X''Y''Z''$ ) fixed in the raft.

Define a coordinate system ( $O''$ - $\Xi H Z$ ) with its origin at the trunnion and at all times parallel to  $O'$ - $X''Y''Z''$ .

Finally, define a coordinate system ( $O''$ - $\Xi' H' Z'$ ) such that  $O''$ - $H'$  makes a constant angle of  $-\gamma$  with  $O''$ - $H$  at all times at all times (in the  $O''$ - $\Xi' H$  plane)

Define the coordinates of  $O'$  in  $O$ - $XYZ$  as  $(\xi, \eta, 0)$ .  
Define the coordinates of  $O''$  in  $O'$ - $X''Y''Z''$  as  $(a, b, 0)$ .

# DYNAMIC ANALYSIS OF FIRING PLATFORM

TIME-ACTION OF BASE CONSIDERING BUOYANT AND VISCIOUS DRAG FORCES ON PLATFORM

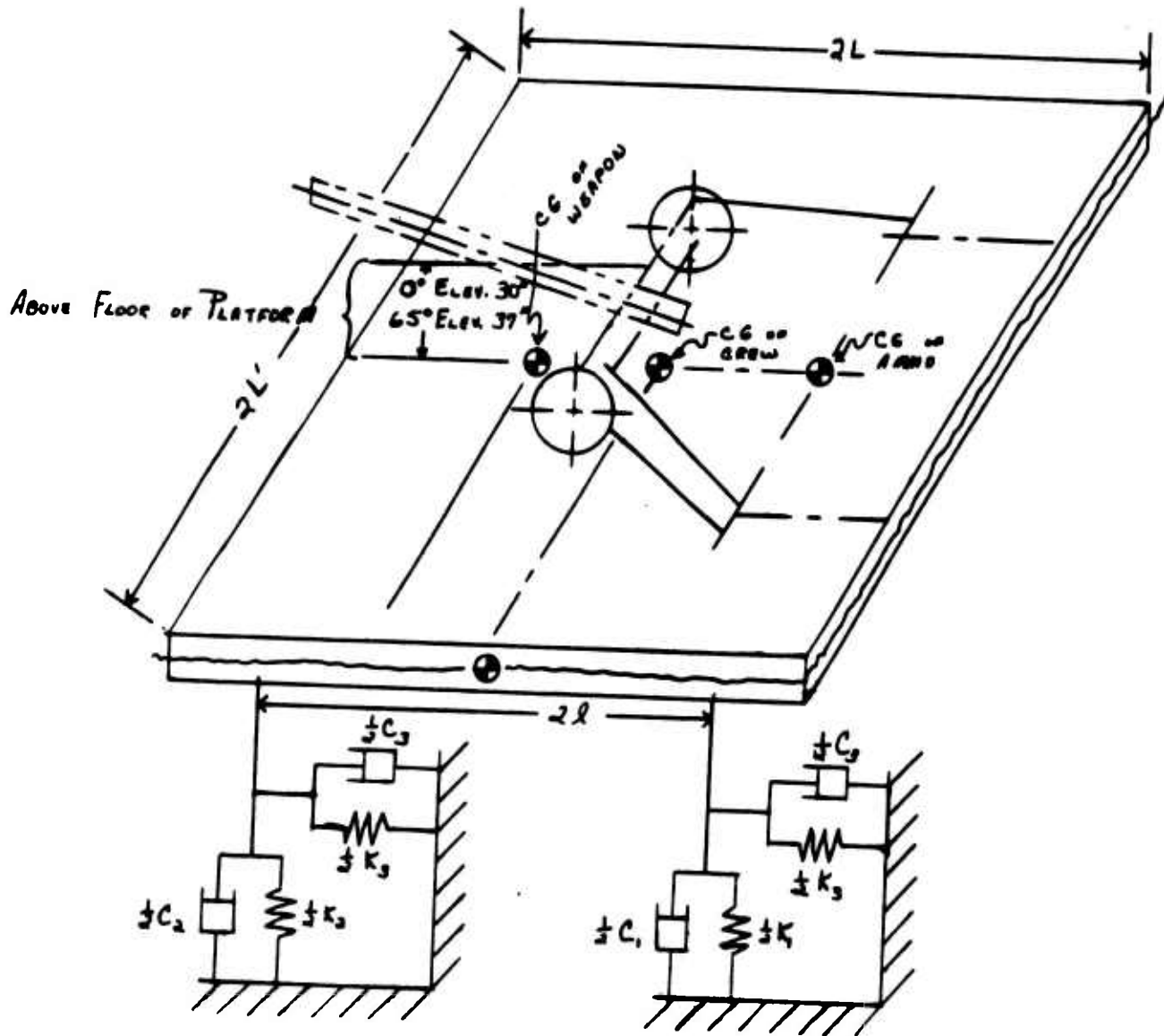


fig 1.

So

$$\begin{bmatrix} X \\ Y \\ Z \end{bmatrix} = \begin{bmatrix} X' \\ Y' \\ Z' \end{bmatrix} + \begin{bmatrix} \xi \\ \eta \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} X' \\ Y' \\ Z' \end{bmatrix} = \begin{bmatrix} \cos \phi & -\sin \phi & 0 \\ \sin \phi & \cos \phi & 0 \\ 0 & 0 & 1 \end{bmatrix} \cdot \begin{bmatrix} X'' \\ Y'' \\ Z'' \end{bmatrix}$$

$$\begin{bmatrix} X'' \\ Y'' \\ Z'' \end{bmatrix} = \begin{bmatrix} \bar{X} \\ H \\ Z \end{bmatrix} + \begin{bmatrix} a \\ b \\ 0 \end{bmatrix}$$

and

$$\begin{bmatrix} \bar{X} \\ H \\ Z \end{bmatrix} = \begin{bmatrix} \cos \bar{\gamma} & \sin \bar{\gamma} & 0 \\ -\sin \bar{\gamma} & \cos \bar{\gamma} & 0 \\ 0 & 0 & 1 \end{bmatrix} \cdot \begin{bmatrix} \bar{H}' \\ H' \\ Z' \end{bmatrix} \quad \text{where } \bar{\gamma} = -\gamma$$

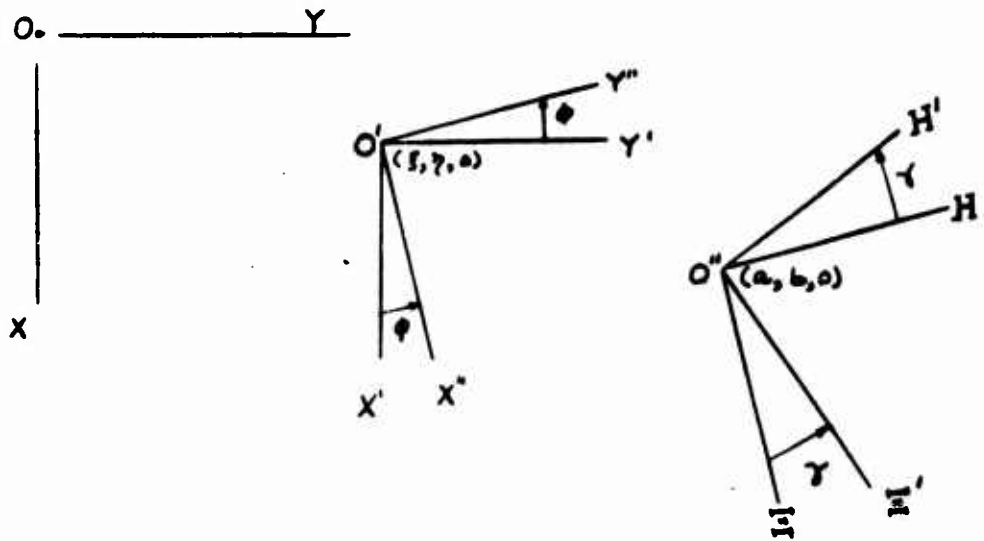


Fig. 2

### 3. Applied Load

Since  $R(t)$ , the applied load, acts in the direction of  $O''H'$ ,

$$\begin{bmatrix} R_x \\ R_y \\ R_z \end{bmatrix} = \begin{bmatrix} \cos \phi & -\sin \phi & 0 \\ \sin \phi & \cos \phi & 0 \\ 0 & 0 & 1 \end{bmatrix} \cdot \begin{bmatrix} \cos \bar{\gamma} & \sin \bar{\gamma} & 0 \\ -\sin \bar{\gamma} & \cos \bar{\gamma} & 0 \\ 0 & 0 & 1 \end{bmatrix} \cdot \begin{bmatrix} 0 \\ R(t) \\ 0 \end{bmatrix}$$

So

$$R_x = R(t) \sin(\bar{\gamma} - \phi) \quad R_y = R(t) \cos(\bar{\gamma} - \phi)$$

where  $\bar{\gamma}$  is the angle of elevation.

The moment about the mass center due to  $R(t)$  is

$$M_p = R(t) [a \cos \bar{\gamma} - b \sin \bar{\gamma}]$$

#### 4. Buoyant Forces

The buoyant force is equal to the weight of the displaced fluid and its line of action is vertical through the center of mass of the displaced fluid.

Assume a deflected position of the fluid mass as shown in Fig. 3. Now

$$r_A = r_1 \cos(\alpha_1 + \phi) \quad y_A = r_1 \sin(\alpha_1 + \phi)$$

$$r_B = r_2 \cos(\alpha_2 + \phi) \quad y_B = r_2 \sin(\alpha_2 + \phi)$$

$$r_C = r_3 \cos(\alpha_3 - \phi) \quad y_C = r_3 \sin(\alpha_3 - \phi)$$

$$r_D = r_4 \cos(\alpha_4 - \phi) \quad y_D = r_4 \sin(\alpha_4 - \phi)$$

Vertical depth of water line =  $h - \xi$

$$FH = r_4 \cos(\alpha_4 - \phi) - (h - \xi) \quad EI = r_1 \cos(\alpha_1 + \phi) - (h - \xi)$$

Also

$$r_E = h - \xi \quad y_E = r_1 \sin(\alpha_1 + \phi) - [r_1 \cos(\alpha_1 + \phi) - (h - \xi)] \tan \phi$$

let

$$A = r_1 \cos \alpha_1 = h + \xi_0 \quad B = r_1 \sin \alpha_1 = L + g$$

so

$$r_A = r_1 \cos(\alpha_1 + \phi) = r_1 \cos \alpha_1 \cos \phi - r_1 \sin \alpha_1 \sin \phi = A \cos \phi - B \sin \phi$$

$$r_E = h - \xi$$

$$y_A = r_1 \sin(\alpha_1 + \phi) = r_1 \sin \alpha_1 \cos \phi + r_1 \cos \alpha_1 \sin \phi = A \sin \phi + B \cos \phi$$

$$y_E = r_1 \sin(\alpha_1 + \phi) - [r_1 \cos(\alpha_1 + \phi) - (h - \xi)] \tan \phi$$

$$= A \sin \phi + B \cos \phi - [A \cos \phi - B \sin \phi - (h - \xi)] \tan \phi$$

$$= A \sin \phi + B \cos \phi - A \sin \phi + B \sin \phi \tan \phi + (h - \xi) \tan \phi$$

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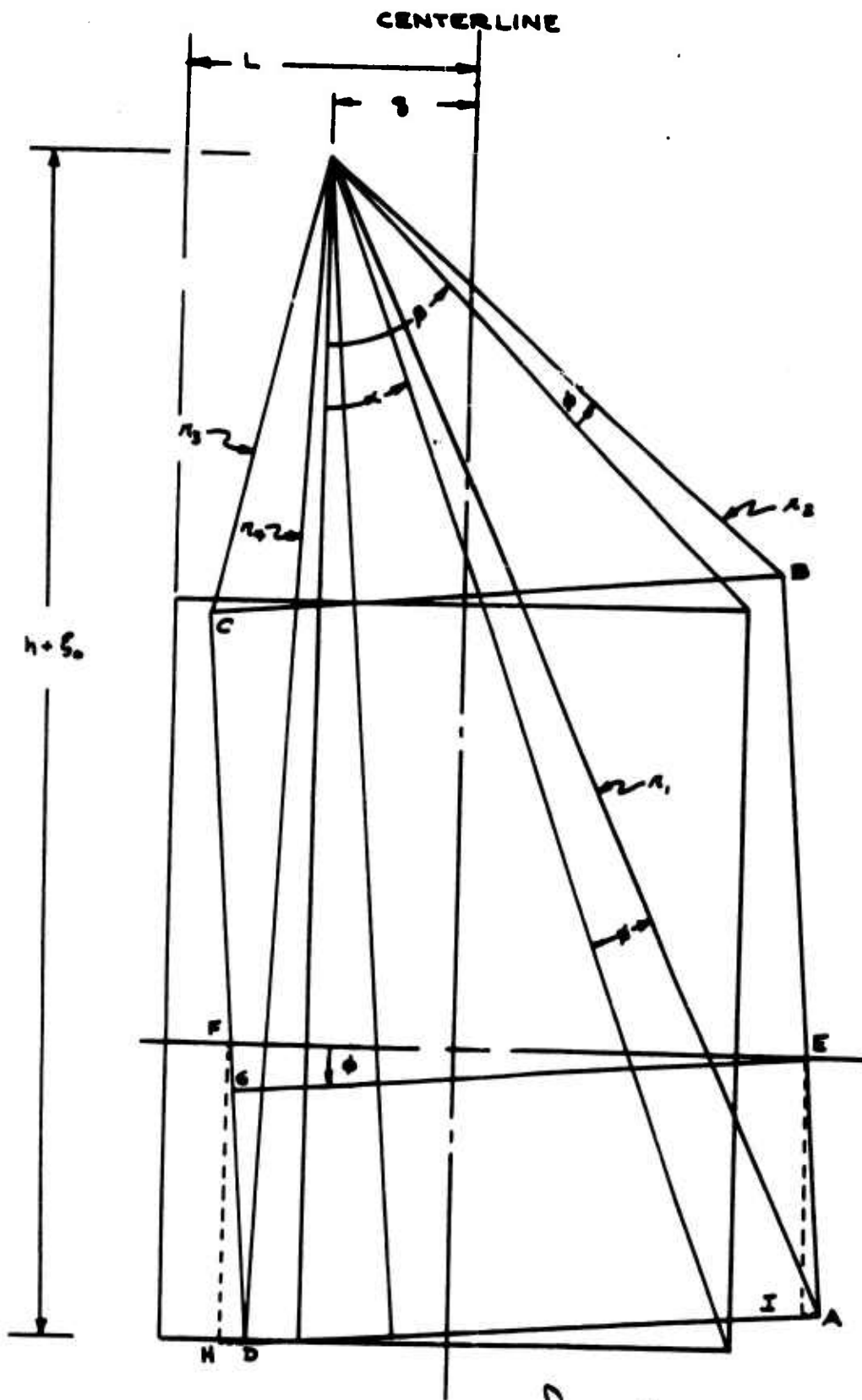


Fig 3.

or

$$y_E = B \sec \phi + (h - \xi) \tan \phi$$

Now

$$DA = 2L \quad EA = \{(x_A - x_E)^2 + (y_A - y_E)^2\}^{\frac{1}{2}} = A - B \tan \phi - (h - \xi) \sec \phi$$

Define the coordinate system  $E-x'y'$  by

$$x = x' \cos \phi - y' \sin \phi + (h - \xi)$$

$$y = x' \sin \phi + y' \cos \phi + y_E$$

so

$$\begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} \cos \phi & -\sin \phi \\ \sin \phi & \cos \phi \end{bmatrix} \begin{bmatrix} x' \\ y' \end{bmatrix} + \begin{bmatrix} h - \xi \\ y_E \end{bmatrix}$$

Now

$$x'_A = \bar{EA}$$

$$y'_A = 0$$

$$x'_B = \bar{EA}$$

$$y'_B = -2L$$

$$x'_E = 0$$

$$y'_E = 0$$

$$x'_F = -2L \tan \phi$$

$$y'_F = -2L$$

$$x'_G = 0$$

$$y'_G = -2L$$

thus

$$x'_\Delta = \frac{1}{3} \bar{FG} = \frac{1}{3} (-2L \tan \phi)$$

$$= -\frac{2}{3} L \tan \phi$$

$$y'_\Delta = \frac{2}{3} \bar{GE} = \frac{2}{3} (-2L)$$

$$= -\frac{4}{3} L$$

$$x'_\square = \frac{1}{2} \bar{EA}$$

$$y'_\square = \frac{1}{2} \bar{GE} = \frac{1}{2} (-2L)$$

$$= -L$$

So

$$y_\Delta = -\frac{2}{3} L \sin \phi \tan \phi - \frac{4}{3} L \cos \phi + y_E$$

$$y_\square = \frac{1}{2} \bar{EA} \sin \phi - L \cos \phi + y_E$$

Also

$$\begin{aligned}
 \text{Area}_\Delta &= \frac{1}{2}(\overline{GE} \cdot \overline{GF}) \\
 &= \frac{1}{2}(-2L \tan\phi)(-2L) \\
 &= 2L^2 \tan\phi \\
 \text{Area}_\square &= \overline{DA} \cdot \overline{FA} \\
 &= 2L \cdot \overline{FA} \\
 &= 2L(A - B \tan\phi - [h - \xi] \sec\phi)
 \end{aligned}$$

and

$$\text{C.G.} = \frac{\text{Area}_\Delta \cdot y_\Delta + \text{Area}_\square \cdot y_\square}{\text{Area}_\Delta + \text{Area}_\square}$$

The buoyant force is

$$\begin{aligned}
 F(\phi) &= 2L'(\text{Area}_\Delta + \text{Area}_\square)\sigma \\
 &= 4LL' [A + (L - B)\tan\phi - (h - \xi)\sec\phi]\sigma
 \end{aligned}$$

and the moment due to the force is

$$\begin{aligned}
 M(\phi) &= F(\phi) \cdot (\text{C.G.}) = 2L' [\text{Area}_\Delta \cdot y_\Delta + \text{Area}_\square \cdot y_\square] \sigma \\
 &= 2L'\sigma \left\{ 2L^2 \tan\phi \left[ -\frac{2}{3}L \sin\phi \tan\phi - \frac{1}{3}L \cos\phi + B \sec\phi + (h - \xi)\tan\phi \right] \right. \\
 &\quad \left. + 2L(A - B \tan\phi - [h - \xi]\sec\phi) \left( \frac{1}{2}[A - B \tan\phi - (h - \xi)\sec\phi] \sin\phi \right. \right. \\
 &\quad \left. \left. - L \cos\phi + B \sec\phi + (h - \xi)\tan\phi \right) \right\}
 \end{aligned}$$

$$\begin{aligned}
 M(\phi) &= 4LL' \left\{ -\frac{2}{3}L^2 \tan^2\phi \sin\phi - \frac{1}{3}L^2 \cos\phi \tan\phi + BL \tan\phi \sec\phi \right. \\
 &\quad \left. + (h - \xi)L \tan^2\phi + \frac{1}{2}[A - B \tan\phi - (h - \xi)\sec\phi]^2 \sin\phi \right. \\
 &\quad \left. + [A - B \tan\phi - (h - \xi)\sec\phi] [-L \cos\phi + B \sec\phi + (h - \xi)\tan\phi] \right\}
 \end{aligned}$$

Expanding and reducing by Trigonometric identities.

$$\begin{aligned}
 M(\phi) &= 4LL' \left\{ (BL - \frac{2}{3}L^2 - \frac{B^2}{2})(1 + \sec^2\phi) \sin\phi + (L - B)[(h - \xi)\sec^2\phi - A \cos\phi] \right. \\
 &\quad \left. + \frac{1}{2} \sin\phi [A^2 - (h - \xi)^2 \sec^2\phi] \right\} \sigma
 \end{aligned}$$

Linearizing and substituting for A and B

$$F(\phi) = 4LL'(\xi + \xi_0 - g\phi)\sigma$$

$$M(\phi) = 4LL' \left\{ (\xi + \xi_0)(g + h\phi) + \frac{\phi}{2} [3(\xi_0^2 - \xi^2) - 2(L^2 + 3g^2)] \right\} \sigma$$

### 5. Drag Forces

The drag force  $D$  is given by

$$D = \frac{1}{2} C_D \rho V^2 2L' [\xi + \xi_0 - (L+g)\phi] \dot{\eta}^2$$

or

$$D = \frac{\sigma L' C_D}{g} [\xi + \xi_0 - (L+g)\phi] \dot{\eta}^2 \quad (\text{where } C_D \sim 1.7)$$

### 6. Spring & Damping Forces

The forces resisting motion in the  $\eta$ -direction  $F_\eta$  are,

$$F_\eta = K_3 \eta + C_3 \dot{\eta}$$

The forces resisting motion in the  $\xi$ -direction  $F_\xi$  are

$$F_\xi = K_1 [\xi + \xi_0 - (L+g)\phi] + K_2 [\xi + \xi_0 + (L-g)\phi] \\ + C_1 [\dot{\xi} - (L+g)\dot{\phi}] + C_2 [\dot{\xi} + (L-g)\dot{\phi}]$$

The moments resisting motion in the  $\phi$ -direction  $F_\phi$  are

$$F_\phi = -K_1 [\xi + \xi_0 - (L+g)\phi](L+g) + K_2 [\xi + \xi_0 + (L-g)\phi](L-g) \\ - C_1 [\dot{\xi} - (L+g)\dot{\phi}](L+g) + C_2 [\dot{\xi} + (L-g)\dot{\phi}](L-g)$$

### 7. Equations of Motion

$$M \ddot{\eta} = R(\bar{c}) \cos(\bar{\gamma} - \phi) - \frac{\sigma L' C_D}{g} [\xi + \xi_0 - (L+g)\phi] \dot{\eta}^2 - K_3 \eta - C_3 \dot{\eta}$$

$$M \ddot{\xi} = R(\bar{c}) \sin(\bar{\gamma} - \phi) - 4LL'\sigma (\xi + \xi_0 - g\phi) - K_1 [\xi + \xi_0 - (L+g)\phi] \\ - K_2 [\xi + \xi_0 + (L-g)\phi] - C_1 [\dot{\xi} - (L+g)\dot{\phi}] - C_2 [\dot{\xi} + (L-g)\dot{\phi}] + W$$

$$I \ddot{\phi} = R(\bar{c}) [a \cos \bar{\gamma} - b \sin \bar{\gamma}] + 4LL'\sigma \{ (\xi + \xi_0)(g+h\phi) + \frac{g}{2} [3(\xi^2 - \xi_0^2) - \eta^2 + 3g^2] \} \\ + K_1 [\xi + \xi_0 - (L+g)\phi](L+g) - K_2 [\xi + \xi_0 + (L-g)\phi](L-g) \\ + C_1 [\dot{\xi} - (L+g)\dot{\phi}](L+g) - C_2 [\dot{\xi} + (L-g)\dot{\phi}](L-g)$$

with initial condition on  $\phi$ , say  $\phi_0$

$$\phi_0 = \frac{\xi_0 [K_1(L+g) + 4gLL'\sigma - K_2(L-g)]}{-4LL'\sigma (\xi_0 h + \xi_0^2 - \frac{L^2 + 3g^2}{g}) + K_1(L+g) + K_2(L-g)}$$

where

$$\xi_0 = \frac{W}{K_1 + K_2 + 4LL'\sigma}$$

## GAS FLOW STUDIES

Problem 1:

Consider a reservoir of gas at pressure  $p_0$ , density  $\rho_0$ , and temperature  $T_0$ . Assume that the gas is to expand by an adiabatic process through a convergent-divergent nozzle to the atmosphere; and that the flow is "steady-state". Determine the mass rate of flow of the gas. Define the subscript  $t$  to refer to conditions at the throat and the subscript 1 to refer to conditions in the chamber.



The continuity equation for steady flow is given as

$$Q = \rho \mu \bar{a} = \rho_t \mu_t \bar{a}_t \quad (1)$$

where  $Q$  is the mass rate of flow,  $\bar{a}$  is the area of the nozzle, and  $\mu$  is the velocity of the gases.

The energy equation for steady flow between the reservoir and any nozzle section is

$$C_p T + \frac{1}{2} \mu^2 = C_p T_1 \quad (2)$$

where  $C_p$  is the specific heat at constant pressure.

Now, equations 3, 4, 5 hold for adiabatic flow of a perfect gas

$$C_p - C_v = R \quad (3)$$

$$C_p / C_v = \gamma \quad (4)$$

$$C_p = \frac{\gamma}{\gamma - 1} R \quad (5)$$

where  $R$  is the gas constant and  $C_v$  is the specific heat at constant volume

Substituting  $C_p$  from eq. 5 into eq. 2 and solving for  $\mu$ ,

$$\mu = \sqrt{\frac{2\gamma}{\gamma - 1} R (T_1 - T)} \quad (6)$$

Now, the equation of state is given by

$$\frac{P}{\rho} = RT \quad (7)$$

and for adiabatic expansion this becomes

$$\frac{P}{\rho^\gamma} = \text{constant} \quad (8)$$

Differentiating eq. 8

$$\frac{1}{\rho^\gamma} dP - P\gamma \frac{1}{\rho^{\gamma+1}} d\rho = 0$$

or

$$\frac{dP}{d\rho} = \gamma \frac{P}{\rho} \quad (9)$$

and substituting from eq. 7

$$\frac{dP}{d\rho} = \gamma RT \quad (10)$$

The local velocity of sound is given by the equation

$$c = \left( \frac{dP}{d\rho} \right)^{\frac{1}{2}} \quad (11)$$

Substituting from equation 10

$$c = \sqrt{\gamma RT} \quad (12)$$

Now, equation 6 may be rewritten as

$$\frac{\gamma-1}{2\gamma R} \mu^2 = T_1 - T$$

or as

$$\frac{T_1}{T} = 1 + \frac{\gamma-1}{2\gamma R} \mu^2$$

substituting from eq. 12

$$\frac{T_1}{T} = 1 + \frac{\gamma-1}{2} \frac{\mu^2}{c^2} \quad (13)$$

Denote  $\frac{\mu}{c}$  by  $M$  (Mach number) then eq. 13 becomes

$$\frac{T_1}{T} = 1 + \frac{\gamma-1}{2} M^2 \quad (14)$$

Assume that there is no heat loss from this process (isentropic flow). For this type of flow

$$\frac{P}{\rho_i} = \left( \frac{\rho}{\rho_i} \right)^\gamma \quad (15)$$

$$\frac{T}{T_1} = \left( \frac{P}{P_1} \right)^{\frac{\gamma-1}{\gamma}} \quad (16)$$

Substituting equations 15 and 16 into equation 14

$$\frac{P}{P_1} = \left(1 + \frac{\gamma-1}{2} M^2\right)^{\frac{\gamma}{\gamma-1}} \quad (17)$$

$$\frac{\rho}{\rho_1} = \left(1 + \frac{\gamma-1}{2} M^2\right)^{\frac{1}{\gamma-1}} \quad (18)$$

where  $P_1$ ,  $\rho_1$ , and  $T_1$  refer to conditions in the reservoir such that the flow is from  $P_1, \rho_1, T_1$  to  $P, \rho, T$ .

Now, due to the convergent-divergent nozzle, assume that the velocity of the gases at the throat is equal to the local velocity of sound; thus, at the throat the Mach number ( $M$ ) is unity. Hence equations 14, 17, and 18 become

$$\frac{T_1}{T_c} = \frac{\gamma+1}{2} \quad (19)$$

$$\frac{P_1}{P_c} = \left(\frac{\gamma+1}{2}\right)^{\frac{\gamma}{\gamma-1}} \quad (20)$$

$$\frac{\rho_1}{\rho_c} = \left(\frac{\gamma+1}{2}\right)^{\frac{1}{\gamma-1}} \quad (21)$$

From equation 6, evaluated at the throat

$$\mu_c^2 = \frac{2\gamma}{\gamma-1} R(T_1 - T_c)$$

Now from equation 19

$$T_c = T_1 \left(\frac{2}{\gamma+1}\right)$$

hence

$$\mu_c^2 = \frac{2\gamma}{\gamma-1} R T_1 \frac{\gamma-1}{\gamma+1}$$

or

$$\mu_c^2 = \frac{2\gamma}{\gamma+1} R T_1 \quad (22)$$

The flow equation may be derived from equation 1.

$$Q = \rho_c \mu_c a_c$$

Substituting from equation 7

$$Q = P_c (R T_c)^{-1} \mu_c a_c$$

substituting from equations 19, 20 and 22

$$Q = \left[\gamma \left(\frac{2}{\gamma+1}\right)^{\frac{\gamma}{\gamma-1}} \frac{1}{R T_1}\right]^{\frac{1}{2}} P_1 a_c \quad (23)$$

Define  $U_1 =$  volume of chamber, then

$$Q = -U_1 \dot{\rho}_1 \quad (24)$$

From equation 8

$$P_0 \rho_0^{-\gamma} = P_1 \rho_1^{-\gamma}$$

so

$$\rho_1 = \rho_0 \left( \frac{P_1}{P_0} \right)^{\frac{1}{\gamma}} \quad (25)$$

also from equation 7 and equation 25

$$RT_1 = \frac{P_1}{\rho_1} = P_1 \rho_0^{-1} P_0^{-\frac{\gamma-1}{\gamma}} P_1^{\frac{\gamma-1}{\gamma}} = \rho_0^{-1} P_0^{-\frac{\gamma-1}{\gamma}} P_1^{\frac{1}{\gamma}} \quad (26)$$

Differentiating equation 25

$$\dot{\rho}_1 = \rho_0 P_0^{-\frac{1}{\gamma}} \cdot \frac{1}{\gamma} P_1^{\frac{1-\gamma}{\gamma}} \dot{P}_1 \quad (27)$$

Substituting equation 27 into equation 24 and equation 26 into equation 23, and equating

$$-U_1 \rho_0 P_0^{-\frac{1}{\gamma}} \cdot \frac{1}{\gamma} P_1^{\frac{1-\gamma}{\gamma}} \dot{P}_1 = \left[ \gamma \left( \frac{2}{\gamma+1} \right)^{\frac{\gamma+1}{\gamma-1}} \rho_0 P_0^{-\frac{1}{\gamma}} P_1^{\frac{1-\gamma}{\gamma}} \right]^{\frac{1}{2}} P_1 a_c$$

or

$$P_1^{\frac{1-3\gamma}{2\gamma}} dP = - \frac{a_c}{U_1} \rho_0^{-\frac{1}{\gamma}} P_0^{-\frac{1}{2\gamma}} \gamma^{\frac{1}{2}} \left( \frac{2}{\gamma+1} \right)^{\frac{\gamma+1}{2(\gamma-1)}} dt$$

Integrating

$$P_1^{\frac{1-\gamma}{2\gamma}} = \frac{\gamma-1}{2\gamma} \frac{a_c}{U_1} \rho_0^{-\frac{1}{\gamma}} P_0^{-\frac{1}{2\gamma}} \gamma^{\frac{1}{2}} \left( \frac{2}{\gamma+1} \right)^{\frac{\gamma+1}{2(\gamma-1)}} t + C$$

when  $t=0$ ,  $P_1 = P_0$  thus

$$P_1 = \left\{ P_0^{\frac{1-\gamma}{2\gamma}} + \frac{a_c(\gamma-1)}{2U_1} \left[ \gamma \left( \frac{2}{\gamma+1} \right)^{\frac{\gamma+1}{\gamma-1}} \frac{P_0}{\rho_0} \right]^{\frac{1}{2}} P_0^{-\frac{1-\gamma}{2\gamma}} t \right\}^{\frac{2\gamma}{1-\gamma}}$$

or

$$P_1 = P_0 \left( 1 + \frac{t}{\phi} \right)^{\frac{-2\gamma}{\gamma-1}} \quad (28)$$

where

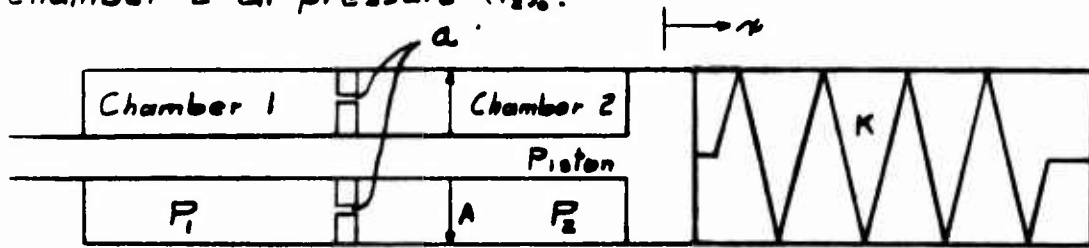
$$\phi = \frac{2U_1}{a_c(\gamma-1)} \left[ \frac{1}{\gamma} \left( \frac{\gamma+1}{2} \right)^{\frac{\gamma+1}{\gamma-1}} \frac{\rho_0}{P_0} \right]^{\frac{1}{2}} \quad (29)$$

also, it is useful to note

$$\rho_0 = P_0 (RT_0)^{-1} \quad (30)$$

Problem 2:

Define the flow of gas from Chamber 1 to Chamber 2 where Chamber 1 is initially under pressure ( $P_1$ ), and Chamber 2 at pressure ( $P_2$ ).



Define:

- $x$  = Translation of piston
- $U_1$  = Volume available to gas in Chamber 1
- $U_2 + Ax$  = Volume available to gas in Chamber 2
- $V$  = Velocity of gas flow thru  $a$  (per unit time)
- $Q$  = Mass rate of gas flow thru  $a$  (per unit time)
- $W$  = Initial weight of gas
- $w$  = Specific weight of gas
- $v$  = Specific volume of gas
- $\rho$  = Density of gas
- $P$  = Pressure
- $k$  = Ratio of specific heats
- $M$  = Mass of piston
- $A$  = Effective gas pressure area of piston
- $a$  = Orifice Area
- $K$  = Spring Rate
- $R$  = Gas Constant divided by average molecular weight
- $T$  = Temperature
- $g$  = Gravitational Constant

Note:

The subscript 1 refers to conditions in Chamber 1  
 The subscript 2 refers to conditions in Chamber 2

At time  $t$

$$w_1 = \frac{W_1 - \int_0^t gQ dt}{U_1} \quad (1)$$

$$w_2 = \frac{W_2 + \int_0^t gQ dt}{U_2 + Ax} \quad (2)$$

So

$$\rho_1 = \frac{w_1}{g} = \frac{\frac{w_1}{g} - \int_0^t Q dt}{U_1} \quad (3)$$

$$\rho_2 = \frac{w_2}{g} = \frac{\frac{w_2}{g} + \int_0^t Q dt}{U_2 + A\pi} \quad (4)$$

For an adiabatic expansion

$$P\pi^k = \text{Constant}$$

so for chamber 1,

$$(P_1)_0 (\pi_1)_0^k = P_1 \pi_1^k$$

or

$$\frac{P_1}{(P_1)_0} = \left(\frac{(\pi_1)_0}{\pi_1}\right)^k = \left(\frac{w_1}{(w_1)_0}\right)^k = \left(\frac{\rho_1}{(\rho_1)_0}\right)^k = \left\{ \frac{w_1 - g \int_0^t Q dt}{w_1} \right\}^k \quad (5)$$

Similarly

$$\frac{P_2}{(P_2)_0} = \left\{ \frac{[w_2 + g \int_0^t Q dt] U_2}{w_2 (U_2 + A\pi)} \right\}^k \quad (6)$$

Assuming that the length of the orifice is extremely short and no expansion of the gases occur within the orifice (incompressible flow); then,

$$V^2 = 2gc^2(P_1 - P_2)\pi_1 \quad (7)$$

where  $c$  is the coefficient of discharge. Also

$$Q = aV\rho_1 \quad (8)$$

The equation of motion is simply

$$M\ddot{x} = AP_2 - K(\pi + \pi_{or}) \quad (9)$$

where  $K\pi_{or} = A(P_2)_0$ .

Differentiating equations 3 and 4

$$Q = -U_1 \dot{\rho}_1 \quad (10)$$

$$Q = \frac{d}{dt} [\rho_2 (U_2 + A\pi)] \quad (11)$$

thus

$$-U_1 d\rho_1 = d[\rho_2 (U_2 + A\pi)]$$

integrating

$$-U_1 \rho_1 = \rho_2 (U_2 + A\pi) + C$$

When  $\pi = 0$

$$\rho_1 = (\rho_1)_0 = \frac{W_1}{gU_1} \quad \text{and} \quad \rho_2 = (\rho_2)_0 = \frac{W_2}{gU_2}$$

thus

$$-\frac{W_1}{g} = \frac{W_2}{g} + C$$

so

$$C = -\frac{W_1 + W_2}{g}$$

thus

$$\rho_1 = \frac{W_1 + W_2}{gU_1} - \rho_2 \frac{U_2 + AU_2}{U_1} \quad (12)$$

Substituting equation 10 into equation 8

$$-U_1 \dot{\rho}_1 = a \sqrt{\rho_1}$$

or by equation 7,

$$-U_1 \dot{\rho}_1 = a \rho_1 c [2gN_1 (\rho_1 - \rho_2)]^{\frac{1}{2}}$$

and by definition of  $N_1$  ( $gN_1 = 1/\rho_1$ ),

$$-U_1 \dot{\rho}_1 = ac \left[ \frac{2}{\rho_1} (\rho_1 - \rho_2) \right]^{\frac{1}{2}} \rho_1$$

or

$$-U_1 \dot{\rho}_1 = ac [2\rho_1 (\rho_1 - \rho_2)]^{\frac{1}{2}} \quad (13)$$

but

$$\rho_1 = B_1 \rho_1^{\frac{1}{2}} \quad \rho_2 = B_2 \rho_2^{\frac{1}{2}} \quad (14)$$

where (eq. 5)

$$B_1 = (\rho_1)_0 \left( \frac{gU_1}{W_1} \right)^{\frac{1}{2}} \quad B_2 = (\rho_2)_0 \left( \frac{gU_2}{W_2} \right)^{\frac{1}{2}} \quad (15)$$

Substituting eq's 14 into eq's 9 and 13 and rewriting eq. 12,

$$M\ddot{\rho} = AB_2 \rho_2^{\frac{1}{2}} - K(\pi + \pi_{sr}) \quad (16)$$

$$-U_1 \dot{\rho}_1 = ac [2\rho_1 (B_1 \rho_1^{\frac{1}{2}} - B_2 \rho_2^{\frac{1}{2}})]^{\frac{1}{2}} \quad (17)$$

$$\rho_1 = \frac{W_1 + W_2}{gU_1} - \rho_2 \frac{U_2 + AU_2}{U_1} \quad (18)$$

where

$$B_1 = (\rho_1)_0 \left( \frac{gU_1}{W_1} \right)^{\frac{1}{2}} \quad B_2 = (\rho_2)_0 \left( \frac{gU_2}{W_2} \right)^{\frac{1}{2}}$$

and  $P_1$  and  $P_2$  may be computed from equations 14. Equations 16, 17, 18 are the desired equations for the system.

Assume expansion of the gases exists in the orifice. Now, it is necessary to redefine the flow of gas through the orifice,  $a$ . The energy equation for an element of gas in the orifice is:

$$\frac{V^2}{2} = \frac{kRT_1}{k-1} \left(1 - \frac{T}{T_1}\right) \quad (19)$$

and the equation of state of the gas is

$$P = \rho RT \quad (20)$$

Thus one can write equation 19 as

$$V^2 = \frac{2k}{k-1} \left(\frac{P_1}{\rho_1} - \frac{P_2}{\rho_2}\right) \quad (21)$$

Now at the exit of the orifice and by use of eq. 14

$$V_2 = \left\{ \frac{2k}{k-1} (B_1 \rho_1^{k-1} - B_2 \rho_2^{k-1}) \right\}^{\frac{1}{2}} \quad (22)$$

and since

$$Q = a V_2 \rho_2 \quad (23)$$

then by eq. 10

$$-U_1 \rho_1 = a \rho_2 \left\{ \frac{2k}{k-1} (B_1 \rho_1^{k-1} - B_2 \rho_2^{k-1}) \right\}^{\frac{1}{2}} \quad (24)$$

Now the desired equations are

$$M \ddot{y} = A B_2 \rho_2^k - K(x + r_{or}) \quad (25)$$

$$-U_1 \rho_1 = a \rho_2 \left\{ \frac{2k}{k-1} (B_1 \rho_1^{k-1} - B_2 \rho_2^{k-1}) \right\}^{\frac{1}{2}} \quad (26)$$

$$\rho_1 = \frac{W_1 + W_2}{g U_1} - \rho_2 \frac{U_2 + A \dot{y}}{U_1} \quad (27)$$

## PROBLEM

Determine the relationship between momentum index (B) and intrinsic efficiency (I.E.)

## SOLUTION

Define

- $W_0$  = Weight of recoiling parts without brake
- $W_b$  = Weight of muzzle brake
- $W_0 + W_b$  = Weight of recoiling parts with brake
- $V_1$  = Maximum free recoil velocity (occurring at end of gas ejection period) of  $W_0 + W_b$
- $V_2$  = Maximum free recoil velocity (occurring at end of gas ejection period) of  $W_0$
- $v_1$  = Free recoil velocity at end of in-bore period of  $W_0 + W_b$
- $v_2$  = Free recoil velocity at end of in-bore period of  $W_0$
- $I_1$  = Impulse of in-bore period acting on recoiling parts
- $I$  = Total impulse acting on recoiling parts without brake.

By definition

$$B = 1 - \frac{(W_0 + W_b)(V_1 - v_1)}{W_0(V_2 - v_2)} \quad 1$$

$$I.E. = 1 - \frac{W_0 + W_b}{W_0} \cdot \frac{(W_0 + W_b)V_1^2}{W_0 V_2^2} = 1 - \left[ \frac{(W_0 + W_b)V_1}{W_0 V_2} \right]^2 \quad 2$$

but

$$(W_0 + W_b)v_1 = W_0 v_2$$

thus

$$B = 1 - \frac{(W_0 + W_b)(V_1 - v_1)}{W_0 V_2 - (W_0 + W_b)v_1}$$

or

$$B = 1 - \frac{(W_0 + W_b)V_1 - C}{W_0 V_2 - C} \quad 3$$

where

$$C = (W_0 + W_b)v_1 \quad 4$$

From eq. 2

$$(W_0 + W_0)V_1 = W_0 V_2 \sqrt{1 - I.E.}$$

Substituting this equation into eq. 3

$$B = 1 - \frac{W_0 V_2 \sqrt{1 - I.E.} - C}{W_0 V_2 - C}$$

or

$$(1 - B)(W_0 V_2 - C) = W_0 V_2 \sqrt{1 - I.E.} - C$$

expanding

$$W_0 V_2 - C - B W_0 V_2 + B C = W_0 V_2 \sqrt{1 - I.E.} - C$$

or

$$W_0 V_2 + B(C - W_0 V_2) = W_0 V_2 \sqrt{1 - I.E.}$$

Define

$$D = \frac{C}{W_0 V_2} = \frac{(W_0 + W_0)N_1}{W_0 V_2}$$

then

$$1 + B(D - 1) = \sqrt{1 - I.E.}$$

Squaring both sides

$$1 + 2B(D - 1) + B^2(D - 1)^2 = 1 - I.E.$$

thus

$$I.E. = -2B(D - 1) - B^2(D - 1)^2$$

or

$$I.E. = 2B(1 - D) - B^2(1 - D)^2$$

Let

$$K = (1 - D)B$$

then

$$I.E. = K(2 - K) \quad 5$$

where

$$K = \left[ 1 - \frac{(W_0 + W_0)N_1}{W_0 V_2} \right] B = \left( 1 - \frac{W_0 N_1}{W_0 V_2} \right) B = \left( 1 - \frac{I_1}{I} \right) B \quad 6$$

# APPENDIX A

## PROPELLANT FORM FUNCTION

### A1. Introduction

It is possible to express the volume of grain burned in several different ways. If we define  $N/C$  as the fraction of charge burned and  $f$  as the fraction of charge remaining then  $N/C$  is approximated by such expressions as  $N/C = 1 - f$  or  $N/C = K_0 + K_1 f + K_2 f^2$ . The exact equation is a cubic of the form  $N/C = k_0 + k_1 f + k_2 f^2 + k_3 f^3$ . The cubic arises since the change in volume is proportional to a length cubed. Since it is desirable to develop a ballistic system that will utilize a computer, the cubic equation presents no serious problem and undoubtedly adds a considerable amount of precision to the final result. Hence the propellant shapes to be considered here will be in the form of a cubic.

We will use the symbolism of Ref. 1, for this Appendix. Thus, we let

- $N/C$  = fraction of charge burned
- $N$  = amount of charge burned
- $C$  = initial amount of charge
- $V_0$  = initial volume of grain
- $V_x$  = volume remaining at any linear distance ( $x$ )
- $S_0$  = initial surface of grain
- $S_x$  = surface of grain remaining at any linear distance ( $x$ )
- $V_x/V_0$  = fraction of propellant remaining ( $N/C = 1 - V_x/V_0$ )
- $x$  = linear distance burned to flame front, perpendicular to burning surface
- $f$  = fraction of web remaining
- $S_x/S_0$  = progressivity
- $R, r$  = larger and smaller radii, respectively, of grain geometry
- $W$  = web, or thickness
- $L$  = length
- $D$  = width

Other symbols will be defined as needed.

## A2 Cord Propellant

From Figure A1 we can determine

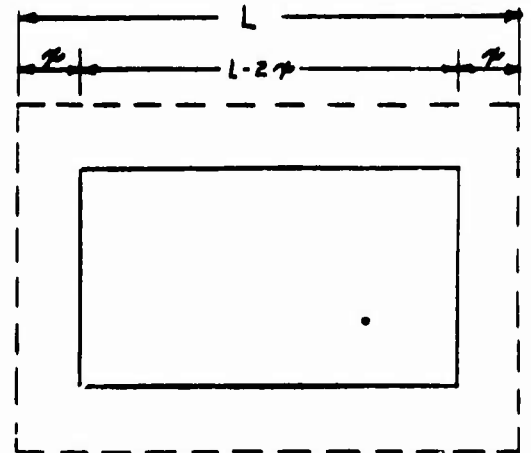
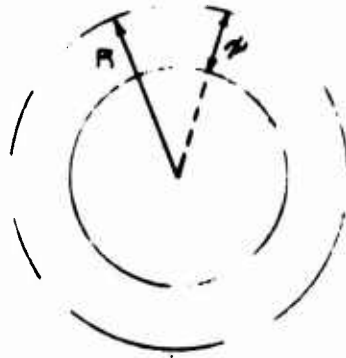
$$S_0 = 2\pi RL + 2\pi R^2 = 2\pi R(L+R)$$

$$\begin{aligned} S_r &= 2\pi(R-r)(L-2r) + 2\pi(R-r)^2 \\ &= 2\pi(RL - rL - 2Rr + 2r^2 + R^2 - 2Rr + r^2) \\ &= 2\pi[R(L+R) - (L+4R)r + 3r^2] \end{aligned}$$

thus

$$\frac{S_r}{S_0} = \frac{2\pi[R(L+R) - (L+4R)r + 3r^2]}{2\pi R(L+R)}$$

$$\frac{S_r}{S_0} = 1 - \frac{4R+L}{R^2+RL} r + \frac{3}{R^2+RL} r^2 \quad (A1)$$



also

Fig. A1

$$V_0 = \pi R^2 L$$

$$\begin{aligned} V_r &= \pi(R-r)^2(L-2r) = \pi(R^2 - 2Rr + r^2)(L-2r) \\ &= \pi(R^2L - 2R^2r - 2RrL + 4Rr^2 + Lr^2 - 2r^3) \\ &= \pi[R^2L - (2R^2 + 2RL)r + (4R+L)r^2 - 2r^3] \end{aligned}$$

so that

$$\frac{V_r}{V_0} = \frac{\pi[R^2L - (2R^2 + 2RL)r + (4R+L)r^2 - 2r^3]}{\pi R^2L}$$

$$\frac{V_r}{V_0} = 1 - \frac{2(R^2+RL)}{R^2L} r + \frac{4R+L}{R^2L} r^2 - \frac{2}{R^2L} r^3 \quad (A2)$$

Now

$$\frac{N}{C}(r) = 1 - \frac{V_0}{V_b} = 1 - 1 + \frac{2(R^2 + RL)}{R^2 L} r - \frac{4R + L}{R^2 L} r^2 + \frac{2}{R^2 L} r^3$$

$$\frac{N}{C}(r) = \frac{2(R^2 + RL)}{R^2 L} r - \frac{4R + L}{R^2 L} r^2 + \frac{2}{R^2 L} r^3 \quad (A3)$$

and since

$$r = R(1 - f)$$

write eq. A3 as:

$$\frac{N}{C}(f) = \frac{2(R^2 + RL)}{R^2 L} R(1 - f) - \frac{4R + L}{R^2 L} R^2(1 - f)^2 + \frac{2}{R^2 L} R^3(1 - f)^3$$

$$\frac{N}{C}(f) = \frac{2}{L}(R + L)(1 - f) - \frac{4R + L}{L}(1 - 2f + f^2) + \frac{2R}{L}(1 - 3f + 3f^2 - f^3)$$

$$\frac{N}{C}(f) = \frac{2R}{L} + 2 - \frac{2R}{L}f - 2f - \frac{4R + L}{L} + \frac{8R + 2L}{L}f - \frac{4R + L}{L}f^2 + \frac{2R}{L} - \frac{6R}{L}f + \frac{6R}{L}f^2 - \frac{2R}{L}f^3$$

$$\frac{N}{C}(f) = \left(\frac{2R}{L} + 2 - \frac{4R}{L} - 1 + \frac{2R}{L}\right) + \left(-\frac{2R}{L} - 2 + \frac{8R}{L} + 2 - \frac{6R}{L}\right)f + \left(\frac{6R}{L} - \frac{4R}{L} - 1\right)f^2 - \frac{2R}{L}f^3$$

or

$$\frac{N}{C}(f) = 1 + \frac{2R - L}{L} f^2 - \frac{2R}{L} f^3 \quad (A4)$$

### A3 Single-Perforated Propellant

From Fig. A2, we find

$$S_0 = 2\pi(R^2 - r^2) + 2\pi(R + r)L$$

$$S_p = 2\pi[(R - r)^2 - (r + r)^2] + 2\pi[(R - r)(L - 2r) + (r + r)(L - 2r)]$$

$$= 2\pi[R^2 - 2Rr + r^2 - r^2 - 2Rr - 2r^2 + RL - 2rL - 2Rr + 2r^2 + rL + rL - 2Rr - 2r^2]$$

$$= 2\pi[R^2 - r^2 + L(R + r)] - 8\pi(R + r)r$$

thus

$$\frac{S_p}{S_0} = \frac{2\pi[R^2 - r^2 + L(R + r) - 4(R + r)r]}{2\pi[R^2 - r^2 + L(R + r)]} = 1 - \frac{4(R + r)r}{R^2 - r^2 + L(R + r)}$$

$$\frac{S_p}{S_0} = 1 - \frac{4}{R - r + L} r \quad (A5)$$

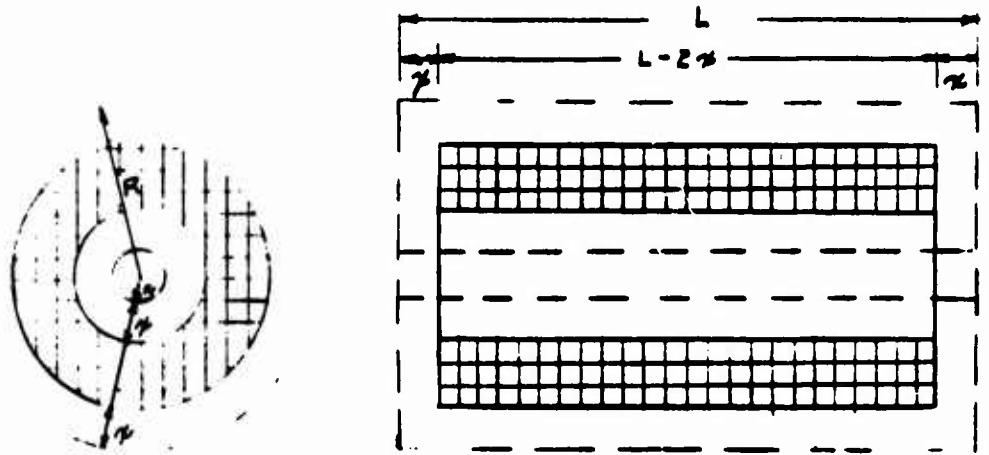


Fig. A2

also

$$V_0 = \pi(R^2 - r^2)L$$

$$\begin{aligned} V_1 &= \pi[(R-r)^2 - (r+r)^2](L-2r) \\ &= \pi(R^2 - 2Rr + r^2 - r^2 - 2Rr - r^2)(L-2r) \\ &= \pi(R^2L - 2R^2r - 2RLr + 4Rr^2 - Lr^2 + 2R^2r - 2RLr + 4Rr^2) \\ &= \pi L(R^2 - r^2) + 2\pi(r^2 - R^2 - RL - rL)r + 4\pi(R+r)r^2 \end{aligned}$$

so that

$$\frac{V_1}{V_0} = \frac{\pi L(R^2 - r^2) + 2\pi[(R+r)(r-R-L)]r + 4\pi(R+r)r^2}{\pi(R^2 - r^2)L}$$

$$\frac{V_1}{V_0} = 1 - \frac{2(R+r+L)r}{(R-r)L} + \frac{4r^2}{(R-r)L} \quad (A6)$$

Now,

$$\frac{N}{C}(r) = 1 - \frac{V_1}{V_0}$$

$$\frac{N}{C}(r) = \frac{2(R+r+L)r}{(R-r)L} - \frac{4r^2}{(R-r)L} \quad (A7)$$

and since

$$R-r - (r+r) = (R-r)r$$

$$R - r - 2x = (R - r)f$$

so

$$r = \frac{1}{2}(R - r)(1 - f)$$

Substituting this expression into eq. A7

$$\frac{N}{C}(f) = \frac{2(R - r + L)}{(R - r)L} \cdot \frac{1}{2}(R - r)(1 - f) - \frac{4}{(R - r)L} \cdot \frac{1}{4}(R - r)^2 \cdot (1 - f)^2$$

$$\frac{N}{C}(f) = \frac{R - r + L}{L}(1 - f) - \frac{R - r}{L}(1 - 2f + f^2)$$

$$\frac{N}{C}(f) = \frac{R - r + L}{L} - \frac{R - r + L}{L}f - \frac{R - r}{L} + \frac{2(R - r)}{L}f - \frac{R - r}{L}f^2$$

$$\frac{N}{C}(f) = \frac{1}{L}(R - r + L - R + r) - \frac{1}{L}(R - r + L - 2R + 2r)f - \frac{1}{L}(R - r)f^2$$

$$\frac{N}{C}(f) = 1 + \frac{R - r - L}{L}f - \frac{R - r}{L}f^2 \quad \text{(A8)}$$

#### A4 Seven Perforated Propellant

From Fig. A3

$$S_0 = 2\pi(R + 7r)L + 2\pi(R^2 - 7r^2)$$

$$S_x = 2\pi[(R - r) + 7(r + r)](L - 2x) + (R - r)^2 - 7(r + r)^2$$

$$S_x = 2\pi[RL - 2Rx - Lx + 2x^2 + 7rL + 7rL - 14rx - 14x^2 + R^2 - 2Rx + x^2 - 7r^2 - 14rx - 7x^2]$$

$$S_x = 2\pi[RL + 7rL + R^2 - 7r^2] + (-2R - L + 7L - 14r - 2R - 14r)x + (2 - 14 + 1 - 7)x^2$$

$$S_x = 2\pi\{[(R^2 - 7r^2) + L(R + 7r)] - [4(R + 7r) - 6L]x - 18x^2\}$$

thus

$$\frac{S_x}{S_0} = \frac{2\pi\{[(R^2 - 7r^2) + L(R + 7r)] - [4(R + 7r) - 6L]x - 18x^2\}}{2\pi[(R + 7r)L + (R^2 - 7r^2)]}$$

$$\frac{S_x}{S_0} = 1 - \frac{4(R + 7r) - 6L}{(R^2 - 7r^2) + L(R + 7r)}x - \frac{18x^2}{(R^2 - 7r^2) + L(R + 7r)} \quad \text{(A9)}$$

also

$$V_0 = \pi(R^2 - 7r^2)L$$

$$V_x = \pi[(R - r)^2 - 7(r + r)^2](L - 2x)$$

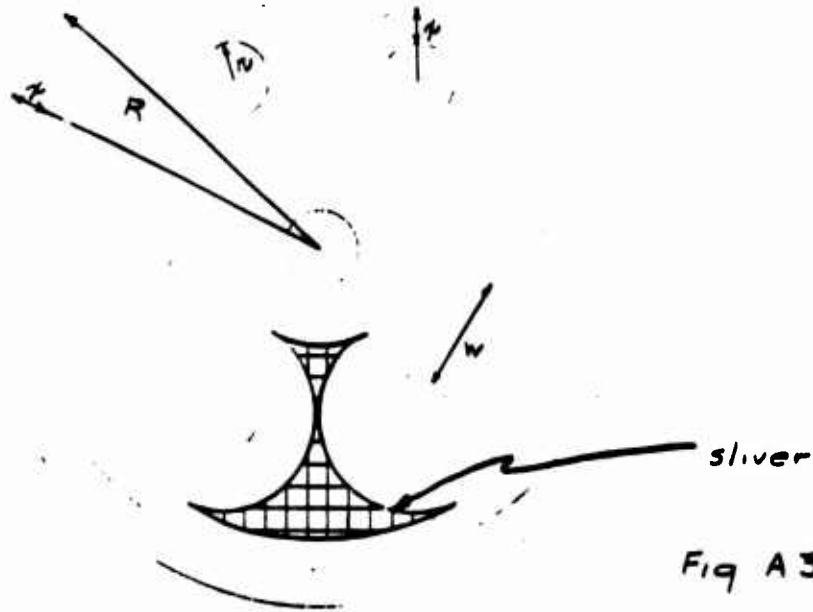
or

$$V_f = \pi [R^2 - 2Rr + r^2 - 7r^2 - 14Rr - 7r^2] (L - 2r)$$

$$V_f = \pi [R^2 L - 2RrL + r^2 L - 7r^2 L - 14RrL - 7r^2 L - 2R^2 r + 4Rr^2 - 2r^3 + 14r^2 r^2 + 28Rr^2 + 14r^3]$$

$$V_f = \pi [(R^2 L - 7r^2 L) + (-2RrL - 14RrL - 2R^2 r + 14r^2) r + (L - 7L + 4R + 28r) r^2 + 12r^3]$$

$$V_f = \pi \{ (R^2 - 7r^2) L - 2[(R + 7r)L + (R^2 - 7r^2)] r + [4(R + 7r) - 6L] r^2 + 12r^3 \}$$



$$\frac{V_f}{V_0} = \frac{\pi \{ (R^2 - 7r^2) L - 2[(R + 7r)L + (R^2 - 7r^2)] r + [4(R + 7r) - 6L] r^2 + 12r^3 \}}{\pi (R^2 - 7r^2) L}$$

$$\frac{V_f}{V_0} = 1 - \frac{2[(R + 7r)L + (R^2 - 7r^2)] r}{(R^2 - 7r^2) L} + \frac{4(R + 7r) - 6L}{(R^2 - 7r^2) L} r^2 + \frac{12}{(R^2 - 7r^2) L} r^3 \quad (A10)$$

Now

$$\frac{d}{C}(r) = 1 - \frac{V_f}{V_0}$$

$$\frac{d}{C}(r) = \frac{2[(R + 7r)L + (R^2 - 7r^2)] r}{(R^2 - 7r^2) L} - \frac{4(R + 7r) - 6L}{(R^2 - 7r^2) L} r^2 - \frac{12}{(R^2 - 7r^2) L} r^3 \quad (A11)$$

From Fig. A3, we note that before burning

$$R = 2W + 2r + r$$

or

$$W = \frac{1}{2}(R - 3r)$$

After burning of a distance,  $x$

$$Wf = \frac{1}{2}[(R - x) - 3(r + x)]$$

or

$$f = \frac{\frac{1}{2}[(R - x) - 3(r + x)]}{\frac{1}{2}(R - 3r)} = \frac{R - 3r - 4x}{R - 3r}$$

or

$$f = 1 - \frac{4x}{R - 3r}$$

hence

$$x = \frac{1}{4}(R - 3r)(1 - f)$$

Substituting this value of  $x$  into Eq. A11,

$$\begin{aligned} \frac{N}{C}(f) &= \frac{2[(R + 7r)L + (R^2 - 7r^2)]}{(R^2 - 7r^2)L} \cdot \frac{1}{4}(R - 3r)(1 - f) - \frac{4(R + 7r) - 6L}{(R^2 - 7r^2)L} \cdot \frac{1}{16}(R - 3r)^4(1 - f)^4 \\ &\quad - \frac{12}{(R^2 - 7r^2)L} \cdot \frac{1}{64}(R - 3r)^3(1 - f)^3 \end{aligned}$$

thus

$$\begin{aligned} (R^2 - 7r^2)L \cdot \frac{N}{C}(f) &= \frac{1}{2}(R^2 + 7rL + R^2 - 7r^2)(R - 3r)(1 - f) - \frac{1}{8}(4R + 28r - 6L)(R^2 - 6Rr + 9r^2)(1 - f)^4 \\ &\quad - \frac{3}{16}(R^3 - 9R^2 \cdot 3r + 3R \cdot 9r^2 - 27r^3)(1 - f)^3 \end{aligned}$$

$$\begin{aligned} (R^2 - 7r^2)L \cdot \frac{N}{C}(f) &= \frac{1}{2}(R^2L + 7rL^2 + R^2 - 7r^2)(R - 3r - Rf + 3rf) \\ &\quad - \frac{1}{8}(4R^3 + 28Rr - 6L)(R^2 - 6Rr + 9r^2)(1 - 2f + f^2) \\ &\quad - \frac{3}{16}(R^3 - 9R^2 \cdot 3r + 27Rr^2 - 27r^3)(1 - 3f + 3f^2 - f^3) \end{aligned}$$

$$\begin{aligned} (R^2 - 7r^2)L \cdot \frac{N}{C}(f) &= \frac{1}{2}(R^2L + 7rL^2 + R^2 - 7r^2)(R - 3r - Rf + 3rf) \\ &\quad - \frac{1}{8}(4R^3 - 24R^2r + 36Rr^2 + 28Rr^2 - 168Rr^2 + 28r^3 \\ &\quad - 6LR^2 + 36RrL - 54r^2L)(1 - 2f + f^2) \\ &\quad - \frac{3}{16}(R^3 - 9R^2r + 27Rr^2 - 27r^3)(1 - 3f + 3f^2 - f^3) \end{aligned}$$

Expanding

$$\begin{aligned}
 (R^2 - 7r^2)L \cdot \frac{N}{C}(f) &= \frac{1}{2}(R^2L - 3RrL - R^2L^2 + 3RrL^2 + 7RrL - 21r^2L - 7r^2R^2L^2 \\
 &+ 21r^2L^2 + R^3 - 3Rr^2 - R^2f + 3R^2r^2 - 7Rr^2 + 21r^3 + 7r^2R^2f - 21r^2f^2) \\
 &- \frac{1}{16}(4R^3 - 24R^2r + 36Rr^2 + 28r^2R^2 - 168Rr^2 + 252r^2 - 6LR^2 + 36RrL \\
 &- 54r^2L - 8R^2f + 48R^2r^2 - 72Rr^2f - 56r^2R^2f + 336Rr^2f^2 - 504r^2f^2 \\
 &+ 12LR^2f - 72RrL^2 + 108r^2L^2 + 4R^2f^2 - 24R^2r^2f^2 + 36Rr^2f^2 + 28Rr^2f^2 \\
 &- 168Rr^2f^2 + 252r^2f^2 - 6LR^2f^2 + 36RrL^2f^2 - 54r^2L^2f^2) - \frac{3}{16}(R^3 - 9R^2r \\
 &+ 27Rr^2 - 27r^2 - 3R^2f + 27R^2r^2 - 81Rr^2f + 81r^2f + 3R^2f^2 - 27R^2r^2f^2 \\
 &+ 81Rr^2f^2 - 81r^2f^2 - R^2f^3 + 9R^2r^2f^3 - 27Rr^2f^3 + 27r^2f^3)
 \end{aligned}$$

So

$$\begin{aligned}
 16(R^2 - 7r^2)L \cdot \frac{N}{C}(f) &= (8R^2L - 24RrL + 36RrL - 168r^2L + 8R^3 - 24Rr^2 \\
 &- 56Rr^2 + 168r^2 - 4R^3 + 24R^2r - 36Rr^2 - 28Rr^2 + 168Rr^2 - 252r^2 \\
 &+ 6LR^2 - 36RrL + 54r^2L - 3R^2 + 27R^2r - 81Rr^2 + 81r^2) + (-8R^2L \\
 &+ 24RrL - 56r^2L + 168r^2L - 8R^2 + 24R^2r + 56r^2R - 168r^2 + 8R^3 \\
 &- 48R^2r + 72Rr^2 + 56r^2R^2 - 336Rr^2 + 504r^2 - 12LR^2 + 72RrL \\
 &- 108r^2L + 9R^2 - 81R^2r + 243Rr^2 - 243r^2) f + (-4R^2 + 24R^2r \\
 &- 36Rr^2 - 28Rr^2 + 168Rr^2 - 252r^2 + 6LR^2 - 36RrL + 54r^2L - 9R^2 \\
 &+ 81R^2r - 243Rr^2 + 243r^2) f^2 + (3R^2 - 27R^2r + 81Rr^2 - 81r^2) f^3
 \end{aligned}$$

Thus

$$\begin{aligned}
 \frac{N}{C}(f) &= \frac{R^3 - R^2r - 5Rr^2 - 3r^2 + 14R^2L - 4RrL - 114r^2L}{16(R^2 - 7r^2)L} \\
 &+ \frac{9R^2 - 49R^2r + 33Rr^2 + 93r^2 - 20R^2L + 40RrL + 60r^2L}{16(R^2 - 7r^2)L} f \\
 &+ \frac{-18R^2 + 77R^2r - 111Rr^2 - 9r^2 + 6R^2L - 36RrL + 54r^2L}{16(R^2 - 7r^2)L} f^2 \\
 &+ \frac{3R^2 - 27R^2r + 81Rr^2 - 81r^2}{16(R^2 - 7r^2)L} f^3 \quad (A12)
 \end{aligned}$$

### A5 Spherical Grain

Without the benefit of a sketch, we can write

$$S_0 = 4\pi R^2$$

$$S_r = 4\pi(R-r)^2 = 4\pi(R^2 - 2Rr + r^2)$$

thus

$$\frac{S_r}{S_0} = \frac{4\pi(R^2 - 2Rr + r^2)}{4\pi R^2}$$

$$\frac{S_r}{S_0} = 1 - \frac{2}{R}r + \frac{1}{R^2}r^2 \quad (A13)$$

also

$$V_0 = \frac{4}{3}\pi R^3$$

$$V_r = \frac{4}{3}\pi(R-r)^3 = \frac{4}{3}\pi(R^3 - 3R^2r + 3Rr^2 - r^3)$$

so

$$\frac{V_r}{V_0} = \frac{\frac{4}{3}\pi(R^3 - 3R^2r + 3Rr^2 - r^3)}{\frac{4}{3}\pi R^3}$$

$$\frac{V_r}{V_0} = 1 - \frac{3}{R}r + \frac{3}{R^2}r^2 - \frac{1}{R^3}r^3 \quad (A14)$$

Now

$$\frac{N}{C}(r) = 1 - \frac{V_r}{V_0} = 1 - 1 + \frac{3}{R}r - \frac{3}{R^2}r^2 + \frac{1}{R^3}r^3$$

$$\frac{N}{C}(r) = \frac{3}{R}r - \frac{3}{R^2}r^2 + \frac{1}{R^3}r^3 \quad (A15)$$

Now

$$r = R(1-f)$$

so, write eq. A15 as

$$\frac{N}{C}(f) = \frac{3}{R} \cdot R(1-f) - \frac{3}{R^2} R^2(1-f)^2 + \frac{1}{R^3} R^3(1-f)^3$$

$$\begin{aligned} \frac{N}{C}(f) &= 3(1-f) - 3(1-2f+f^2) + 1-3f+3f^2-f^3 \\ &= 3-3f-3+6f-3f^2+1-3f+3f^2-f^3 \end{aligned}$$

$$\frac{N}{C}(f) = 1 - f^3 \quad (A16)$$

A6 Strip Propellant  
From Fig A4

$$S_0 = 2DL + 2LW + 2DW$$

$$S_0 = 2(DL + LW + DW)$$

and

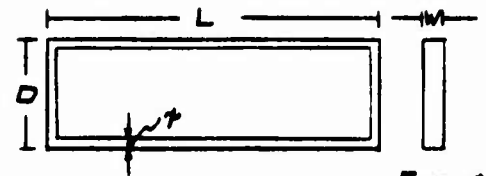


Fig A4

$$S_r = 2[(L-2r)(D-2r) + (L-2r)(W-2r) + (D-2r)(W-2r)]$$

$$S_r = 2[(LD + LW + DW) - 2r(D+L+W+D+W) + 12r^2]$$

$$S_r = 2(LD + LW + DW) - 8r(D+L+W) + 24r^2$$

so

$$\frac{S_r}{S_0} = \frac{2(LD + LW + DW) - 8r(D+L+W) + 24r^2}{2(DL + LW + DW)}$$

$$\frac{S_r}{S_0} = 1 - \frac{4(D+L+W)}{DL+LW+DW} r + \frac{12}{DL+LW+DW} r^2 \quad (A17)$$

Now,

$$V_0 = DLW$$

$$V_r = (D-2r)(L-2r)(W-2r) = (D-2r)(LW - 2rW - 2rL + 4r^2)$$

$$= DLW - 2rWD - 2rLD + 4r^2D - 2rLW + 4r^2W + 4r^2L - 8r^3$$

$$= DLW - 2r(WD + LD + LW) + 4r^2(D+W+L) - 8r^3$$

so

$$\frac{V_r}{V_0} = \frac{DLW - 2r(WD + LD + LW) + 4r^2(D+W+L) - 8r^3}{DLW}$$

$$\frac{V_r}{V_0} = 1 - \frac{2(WD + LD + LW)}{DLW} r + \frac{4(D+W+L)}{DLW} r^2 - \frac{8}{DLW} r^3 \quad (A18)$$

Now

$$\frac{r}{c}(\gamma) = 1 - \frac{V_r}{V_0} = \frac{2(WD + LD + LW)}{DLW} r - \frac{4(D+W+L)}{DLW} r^2 + \frac{8}{DLW} r^3$$

$$\frac{r}{c}(\gamma) = \frac{2(WD + LD + LW)}{DLW} r - \frac{4(D+W+L)}{DLW} r^2 + \frac{8}{DLW} r^3 \quad (A19)$$

but

$$W - 2r = fW$$

So

$$\kappa = \frac{1}{2}W(1-f)$$

Substituting into A19

$$\begin{aligned} \frac{N}{C}(f) &= \frac{2(WD+LD+LW)}{DLW} \cdot \frac{W}{2}(1-f) - \frac{4(D+W+L)}{DLW} \cdot \frac{W^2}{4}(1-f)^2 + \frac{8}{DLW} \frac{W^3}{8}(1-f)^3 \\ &= \frac{WD+LD+LW}{DL} (1-f) - \frac{W(D+W+L)}{DL} (1-2f+f^2) + \frac{W^2}{DL} (1-3f+3f^2-f^3) \\ &= \frac{WD+LD+LW}{DL} - \frac{WD+LD+LW}{DL} f - \frac{WD+W^2+WL}{DL} + \frac{2WD+2W^2+2WL}{DL} f \\ &\quad - \frac{WD+W^2+WL}{DL} f^2 + \frac{W^2}{DL} - \frac{3W^2}{DL} f + \frac{3W^2}{DL} f^2 - \frac{W^2}{DL} f^3 \end{aligned}$$

$$\begin{aligned} \frac{N}{C}(f) &= \frac{1}{DL} [WD+LD+LW - WD - W^2 - WL + W^2] \\ &\quad + \frac{1}{DL} [-WD - LD - LW + 2WD + 2W^2 + 2WL - 3W^2] f \\ &\quad + \frac{1}{DL} [-WD - W^2 - WL + 3W^2] f^2 - \frac{W^2}{DL} f^3 \\ \frac{N}{C}(f) &= 1 + \frac{LW+WD-LD-W^2}{DL} f + \frac{2W^2-WL-WD}{DL} f^2 - \frac{W^2}{DL} f^3 \quad (A20) \end{aligned}$$

### A7. Square Propellant

We may determine the desired equations by substituting  $D=L$  into equations A17, A18, A19, and A20. Thus:

$$\frac{S_x}{S_0} = 1 - \frac{4(2L+W)}{L^2+2LW} \kappa + \frac{12}{L^2+2LW} \kappa^2 \quad (A21)$$

$$\frac{V_x}{V_0} = 1 - \frac{2(2WL+L^2)}{L^2W} \kappa + \frac{4(2L+W)}{L^2W} \kappa^2 - \frac{8}{L^2W} \kappa^3 \quad (A22)$$

$$\frac{N}{C}(\kappa) = \frac{2(2WL+L^2)}{L^2W} \kappa - \frac{4(2L+W)}{L^2W} \kappa^2 + \frac{8}{L^2W} \kappa^3 \quad (A23)$$

$$\frac{N}{C}(f) = 1 + \frac{2LV-L^2-W^2}{L^2} f + \frac{2(W^2-WL)}{L^2} f^2 - \frac{W^2}{L^2} f^3 \quad (A24)$$

### A8. Cube Propellant

Here the desired equations may be determined from eq.s A21, A22, A23, and A24 by setting  $W=L$ .

Thus:

$$\frac{S_x}{S_0} = 1 - \frac{4}{L}x + \frac{4}{L^2}x^2 \quad (A25)$$

$$\frac{V_x}{V_0} = 1 - \frac{6}{L}x + \frac{12}{L^2}x^2 - \frac{8}{L^3}x^3 \quad (A26)$$

$$\frac{N}{C}(x) = \frac{6}{L}x - \frac{12}{L^2}x^2 + \frac{8}{L^3}x^3 \quad (A27)$$

$$\frac{N}{C}(f) = 1 - f^3 \quad (A28)$$

### A9. Inhibitors

The geometrical form factor for a given shape may be altered by use of an inhibitor. The inhibitor prevents burning at the surface which it covers. For example, consider an inhibitor placed on the outside surface of a single-perforated propellant cylinder but, not over the ends. Thus, the burning surfaces are the two ends and the inside surface. See Fig. A5.

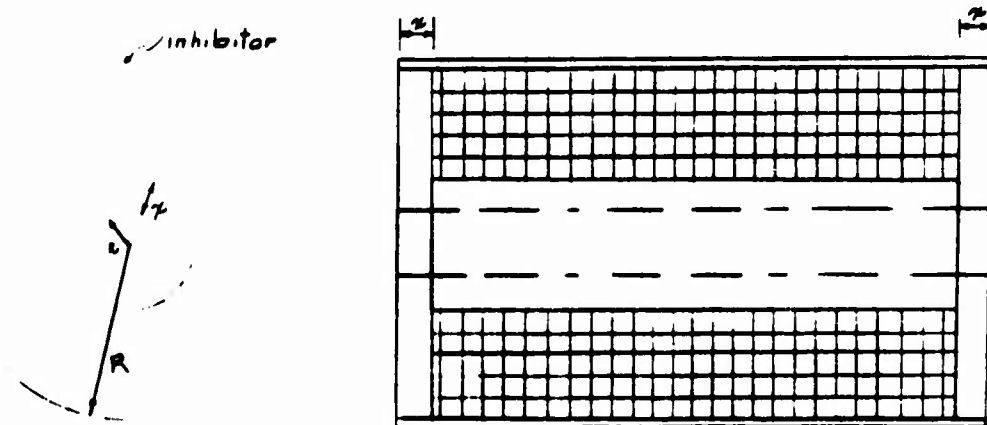


Fig. A5

The crosshatched section indicates the unburned propellant.

From Fig. A5, we find

$$S_0 = 2\pi(R^2 - r^2) + 2\pi rL$$

$$\begin{aligned} S_r &= 2\pi[R^2 - (r+\rho)^2] + 2\pi(r+\rho)(L-2\rho) \\ &= 2\pi[R^2 - r^2 - 2r\rho - \rho^2 + rL - 2r\rho + L\rho - 2\rho^2] \\ &= 2\pi[R^2 - r^2 + rL + \rho(L-4r) - 3\rho^2] \end{aligned}$$

thus

$$\frac{S_r}{S_0} = \frac{2\pi[R^2 - r^2 + rL + \rho(L-4r) - 3\rho^2]}{2\pi(R^2 - r^2 + rL)}$$

$$\frac{S_r}{S_0} = 1 + \frac{L-4r}{R^2 - r^2 + rL} \rho - \frac{3}{R^2 - r^2 + rL} \rho^2 \quad (A29)$$

also

$$V_0 = \pi(R^2 - r^2)L$$

$$\begin{aligned} V_r &= \pi[R^2 - (r+\rho)^2](L-2\rho) \\ &= \pi[r^2L - r^2L - 2r\rho L - \rho^2L - 2R^2\rho + 2r^2\rho + 4r\rho^2 + 2\rho^3] \\ &= \pi[(R^2 - r^2)L + \rho(2r^2 - 2R^2 - 2rL) + (4r - L)\rho^2 + 2\rho^3] \end{aligned}$$

so

$$\frac{V_r}{V_0} = \frac{\pi[(R^2 - r^2)L - 2(R^2 + rL - r^2)\rho + (4r - L)\rho^2 + 2\rho^3]}{\pi(R^2 - r^2)L}$$

$$\frac{V_r}{V_0} = 1 - \frac{2(R^2 + rL - r^2)}{(R^2 - r^2)L} \rho + \frac{4r - L}{(R^2 - r^2)L} \rho^2 + \frac{2}{(R^2 - r^2)L} \rho^3 \quad (A30)$$

Now

$$\frac{N}{C} = 1 - V_r/V_0$$

$$\frac{N}{C}(r) = \frac{2(R^2 + rL - r^2)}{(R^2 - r^2)L} r - \frac{4r - L}{(R^2 - r^2)L} r^2 - \frac{2}{(R^2 - r^2)L} r^3 \quad (A31)$$

Now

$$R - (r+\rho) = (R-r)\rho$$

$$\rho = (R-r)(1-f)$$

Substituting into A31

$$\frac{N}{C}(f) = \frac{2(R^2 + rL - r^2)}{(R^2 - r^2)L} (R-r)(1-f) - \frac{4r-L}{(R^2 - r^2)L} \cdot (R-r)^2(1-f)^2 - \frac{2}{(R^2 - r^2)L} (R-r)^3(1-f)^3$$

or

$$\frac{N}{C}(f) = \frac{2(R^2 + RL - L^2)}{(R+L)L} (1-f) - \frac{(4L-L)(R-L)}{(R+L)L} (1-f)^2 - \frac{2(R-L)^2}{(R+L)L} (1-f)^3$$

so

$$\begin{aligned}(R+L)L \cdot \frac{N}{C}(f) &= 2R^2 + 2RL - 2L^2 - (2R^2 + 2RL - 2L^2)f - (4RL - 4L^2 - LR + LN) \\ &+ 2(4RL - 4L^2 - LR + LN)f - (4RL - 4L^2 - LR + LN)f^2 - (2R^2 - 4RL + 2L^2) \\ &+ 3(2R^2 - 4RL + 2L^2)f - 3(2R^2 - 4RL + 2L^2)f^2 + (2R^2 - 4RL + 2L^2)f^3\end{aligned}$$

Combining terms

$$\begin{aligned}(R+L)L \cdot \frac{N}{C}(f) &= (2R^2 + 2RL - 2L^2 - \cancel{4RL} + \cancel{4L^2} + LR - LN - 2R^2 + 4\cancel{L^2} - 2L^2) \\ &+ (-2R^2 - 2\cancel{RL} + 2L^2 + 8RL - \cancel{8L^2} - 2LR + 2\cancel{LN} + 6R^2 - 12RL + 6\cancel{L^2})f \\ &+ (-4RL + 4L^2 + LR - LN - 6R^2 + 12RL - 6L^2)f^2 + 2(R-L)^2 f^3\end{aligned}$$

or

$$\begin{aligned}(R+L)L \cdot \frac{N}{C}(f) &= (R+L)L + (4R^2 - 2LR - 4RL)f \\ &+ (LR - LN - 6R^2 + 8RL - 2L^2)f^2 + 2(R-L)^2 f^3\end{aligned}$$

thus

$$\frac{N}{C}(f) = 1 + \frac{4R^2 - 2LR - 4RL}{(R+L)L} f + \frac{LR - LN - 6R^2 + 8RL - 2L^2}{(R+L)L} f^2 + \frac{2(R-L)^2}{(R+L)L} f^3 \quad (A32)$$

## APPENDIX B THERMOCHEMISTRY OF PROPELLANTS

### B1. The Atomic Composition of Propellant Constituents

We here compute the atomic composition (in gram atoms per gram of propellant) for the various chemical compounds of which the propellant is composed. Thus, we define:

- {C} = number of gram atoms of carbon per gram of propellant
- {H} = number of gram atoms of hydrogen per gram of propellant
- {N} = number of gram atoms of nitrogen per gram of propellant
- {O} = number of gram atoms of oxygen per gram of propellant

We will use M1 propellant as an example. Now, M1 propellant has constituents, by weight, of 85 parts Nitrocellulose (of which 13.15% is Nitrogen), 10 parts Dinitrotoluene, 5 parts Dibutylphthalate, 1 part Diphenylamine, 0.75 parts Ethyl Alcohol, and 0.50 parts Water. Thus:

#### M1 Propellant

Nitrocellulose (13.15% N)	85.00	83.13%
Dinitrotoluene	10.00	9.78%
Dibutylphthalate	5.00	4.89%
Diphenylamine	1.00	0.98%
Ethyl Alcohol	0.75	0.73%
Water	0.50	0.49%

The following atomic weights will be used:

$$C = 12.010 \quad H = 1.008 \quad N = 14.008 \quad O = 16.000$$

First consider one gram molecule of water ( $H_2O$ )

which consists of:

- 2 gram atoms of hydrogen
- 1 gram atom of oxygen

Thus

$$\begin{aligned} \text{One Gram Molecule of } H_2O &= (2 \times 1.008) + (1 \times 16.000) \\ &= 18.008 \text{ grams} \end{aligned}$$

So

$$\{H\} = \frac{2}{18.008} = 11106 \cdot 10^{-6} \tag{B1}$$

$$\{O\} = \frac{1}{18.008} = 5553 \cdot 10^{-6} \tag{B2}$$

For one gram molecule of ethyl alcohol ( $C_2H_5OH$ )  
 2 gram atoms of carbon  
 6 gram atoms of hydrogen  
 1 gram atom of oxygen

we have

$$\text{One Gram Molecule of } C_2H_5OH = (2 \times 12.010) + (6 \times 1.008) + (1 \times 16.000) \\ = 46.068 \text{ grams}$$

thus

$$\{C\} = \frac{2}{46.068} = 4341 \cdot 10^{-5} \quad (B)$$

$$\{H\} = \frac{6}{46.068} = 13024 \cdot 10^{-5} \quad (B1)$$

$$\{O\} = \frac{1}{46.068} = 2171 \cdot 10^{-5} \quad (B2)$$

For one gram molecule of diphenylamine ( $C_{12}H_{11}N$ )  
 12 gram atoms of carbon  
 11 gram atoms of hydrogen  
 1 gram atom of nitrogen

we have

$$\text{One Gram Molecule of } C_{12}H_{11}N = (12 \times 12.010) + (11 \times 1.008) + (1 \times 14.008) \\ = 169.216 \text{ grams}$$

thus

$$\{C\} = \frac{12}{169.216} = 7092 \cdot 10^{-5} \quad (B6)$$

$$\{H\} = \frac{11}{169.216} = 6501 \cdot 10^{-5} \quad (B7)$$

$$\{N\} = \frac{1}{169.216} = 591 \cdot 10^{-5} \quad (B8)$$

For one gram molecule of dibutylphthalate ( $C_{16}H_{22}O_4$ )  
 16 gram atoms of carbon  
 22 gram atoms of hydrogen  
 4 gram atoms of oxygen

we have

$$\text{One Gram Molecule of } C_{16}H_{22}O_4 = (16 \times 12.010) + (22 \times 1.008) + (4 \times 16.000) \\ = 278.336 \text{ grams}$$

thus

$$\{C\} = \frac{16}{278.336} = 5748 \cdot 10^{-5} \quad (B9)$$

$$\{H\} = \frac{22}{278.336} = 7904 \cdot 10^{-5} \quad (B10)$$

$$\{O\} = \frac{4}{278.336} = 1437 \cdot 10^{-5} \quad (B11)$$

Now, consider one gram molecule of dinitrotoluene ( $C_7H_6N_2O_4$ ) which consists of

- 7 gram atoms of carbon
- 6 gram atoms of hydrogen
- 2 gram atoms of nitrogen
- 4 gram atoms of oxygen

thus

$$\begin{aligned} \text{One Gram Molecule of } C_7H_6N_2O_4 &= (7 \times 12.01) + (6 \times 1.008) + (2 \times 14.008) + (4 \times 16.000) \\ &= 182.134 \text{ grams} \end{aligned}$$

So

$$\{C\} = \frac{7}{182.134} = 3843 \cdot 10^{-5} \quad (B12)$$

$$\{H\} = \frac{6}{182.134} = 3294 \cdot 10^{-5} \quad (B13)$$

$$\{N\} = \frac{2}{182.134} = 1098 \cdot 10^{-5} \quad (B14)$$

$$\{O\} = \frac{4}{182.134} = 2196 \cdot 10^{-5} \quad (B15)$$

Finally, we consider the more complicated problem involving the nitrocellulose. The cellulose molecule is written as  $C_6H_7O_2(OH)_3$ . On nitration, some of the (OH) groups are replaced by  $(ONO_2)$  groups. The resulting compound is  $C_6H_7O_2(OH)_{3-p}(ONO_2)_p$ , where  $p$  is the number of (OH) groups replaced by the  $(ONO_2)$  groups. Thus we have in one gram molecule of nitrocellulose  $C_6H_7O_2(OH)_{3-p}(ONO_2)_p$ :

- 6 gram atoms of carbon
- $10-p$  gram atoms of hydrogen
- $p$  gram atoms of nitrogen
- $5+2p$  gram atoms of oxygen

Thus:

$$\begin{aligned} \text{One Gram Molecule of } C_6H_7O_2(OH)_{3-p}(ONO_2)_p &= (6 \times 12.01) + (10-p) \times 1.008 + (p \times 14.008) + (5+2p) \times 16.000 \\ &= 162.14 + 45p \end{aligned}$$

So

$$\{C\} = \frac{6}{162.14 + 45p}$$

$$\{H\} = \frac{10-p}{162.14 + 45p}$$

$$\{N\} = \frac{p}{162.14 + 45p}$$

$$\{O\} = \frac{5+2p}{162.14 + 45p}$$

Define

$y$  = per cent of nitrogen

then since we have  $x$  gram atoms of nitrogen

$$y = \frac{14.008 x}{162.14 + 45x} \cdot 100$$

or

$$y = \frac{1400.8 x}{162.14 + 45x}$$

Now solving for  $x$

$$162.14 y + 45xy = 1400.8 x$$

or

$$x = \frac{162.14 y}{1400.8 - 45y}$$

Thus for  $y = 13.15$  we have

$$x = 2.635$$

So

$$\{C\} = \frac{6}{162.14 + 45x} = 2137 \cdot 10^{-5} \quad (B16)$$

$$\{H\} = \frac{10 - x}{162.14 + 45x} = 2624 \cdot 10^{-5} \quad (B17)$$

$$\{N\} = \frac{x}{162.14 + 45x} = 939 \cdot 10^{-5} \quad (B18)$$

$$\{O\} = \frac{5 + 2x}{162.14 + 45x} = 3659 \cdot 10^{-5} \quad (B19)$$

Now, for MI Propellant (eq's B1 thru B19)

$$\begin{aligned} \{C\} &= [(0.8313)(2137) + (0.0578)(2043) + (0.0489)(5748) + (0.0098)(7892) + (0.0073)(4541)] \cdot 10^{-5} \\ \{C\} &= 2535 \cdot 10^{-5} \quad (B20) \end{aligned}$$

$$\begin{aligned} \{H\} &= [(0.8313)(2624) + (0.0578)(3224) + (0.0489)(7804) + (0.0098)(6501) + (0.0073)(3024) + (0.0099)(1106)] \cdot 10^{-5} \\ \{H\} &= 3103 \cdot 10^{-5} \quad (B21) \end{aligned}$$

$$\begin{aligned} \{N\} &= [(0.8313)(939) + (0.0578)(698) + (0.0098)(581)] \cdot 10^{-5} \\ \{N\} &= 894 \cdot 10^{-5} \quad (B22) \end{aligned}$$

$$\begin{aligned} \{O\} &= [(0.8313)(3659) + (0.0578)(2136) + (0.0489)(4737) + (0.0098)(171) + (0.0073)(5553)] \cdot 10^{-5} \\ \{O\} &= 3370 \cdot 10^{-5} \quad (B23) \end{aligned}$$

## B2 Heat of Formation of the Propellant

The heats of formation of the products of combustion from graphite, hydrogen gas, nitrogen gas, and oxygen have been experimentally determined and are given by Table B1 for conditions at 300 °K

TABLE B1  
HEATS OF FORMATION OF THE PRODUCTS OF EXPLOSION  
In kilocalories per gram molecule for formation  
from graphite, gaseous oxygen, hydrogen, and nitrogen

Substance	Molecular Weight	Heats of Formation (T=300°K At constant volume)
CO	28.010	26.69
CO <sub>2</sub>	44.010	94.03
H <sub>2</sub> O liquid	18.016	67.43
H <sub>2</sub> O gas	18.016	57.50
CH <sub>4</sub>	16.042	17.27
NH <sub>3</sub>	17.032	10.38
NO	30.008	-21.53
OH	17.008	-9.31
N	14.008	-85.25
O	16.000	-58.75
H	1.008	-51.79

The technique of computing the heat of formation of the products of complete combustion for each propellant constituent (from graphite, hydrogen, nitrogen, oxygen) is illustrated below for M1 Propellant. Note: complete combustion implies the products CO<sub>2</sub> and H<sub>2</sub>O only, no CO, OH, etc. Thus

$$(CO_2) = \{C\} \quad (H_2O) = \frac{1}{2}\{H\} \quad (B24)$$

Define

HF<sub>PROD</sub> = Heat of Formation of the Products of Complete  
Combustion of a Propellant Constituent

From Table B1 we have 94030 cal./gm. as the heat of

TABLE B2

Constituent	Mol. Wt.	Atomic Composition in gm atoms/gm				Heat of Combustion cal/gm	HF of Products cal/gm	HF of Substance cal/gm
		C	H	N	O			
Nitrocellulose								
12.0 %N		0.02274	0.02933	0.00856	0.03608	2471	3127	656
12.2		0.02250	0.02879	0.00871	0.03617	2448	3087	642
12.4		0.02226	0.02826	0.00885	0.03626	2419	3046	627
12.6		0.02203	0.02772	0.00899	0.03635	2393	3005	612
12.8		0.02179	0.02718	0.00914	0.03643	2366	2965	599
13.0		0.02155	0.02664	0.00928	0.03652	2340	2923	585
13.15		0.02138	0.02624	0.00939	0.03658	2320	2894	574
13.2		0.02131	0.02610	0.00942	0.03661	2314	2884	570
13.25		0.02125	0.02596	0.00946	0.03663	2308	2873	563
13.4		0.02108	0.02556	0.00957	0.03669	2294	2843	549
Nitroglycerin $C_3H_5N_3O_9$	227.094	0.01321	0.02201	0.01321	0.03968	1635	1985	350
Nitroquaridine $CH_4N_4O_2$	104.074	0.00961	0.03843	0.03843	0.01921	2030	2199	169
Dinitrotoluene $C_7H_6N_2O_4$	182.134	0.03843	0.03294	0.01098	0.02196	4709	4724	15
Dibutylphthalate $C_{16}H_{32}O_4$	278.336	0.05748	0.07904	-	0.01437	7389	8070	681
Dimethylphthalate $C_{12}H_{14}O_4$	222.232	0.05598	0.06295	-	0.01795	5144	7200	2056
Diphenylamine $C_{12}H_{11}N$	169.216	0.07092	0.06501	0.00581	-	9068	8860	-208
Ethyl Centralite $C_{17}H_{29}N_2O$	268.346	0.06529	0.07453	0.00745	0.03726	8418	8652	234
Ethyl Alcohol $C_2H_5OH$	46.068	0.04341	0.13024	-	0.02171	7111	8473	1362
Water $H_2O$	18.008	-	0.11106	-	0.05553	-	3744	3744
2 Nitro di- phenylamine	214.216	0.03601	0.04668	0.00836	0.09336	6976	6840	-136
Graphite C	12.010	0.8326	-	-	-	-	-	0
Carbon Black C	12.010	0.8326	-	-	-	-	-	0

formation of  $\text{CO}_2$  and 67,430 cal./gm for  $\text{H}_2\text{O}$  (liquid). Thus

$$HF_{\text{PROD}} = 94030(\text{CO}_2) + 67430(\text{H}_2\text{O}) \quad (\text{B25})$$

Water ( $\text{H}_2\text{O}$ ) eq. B1

$$\{M\} = 0.11106$$

by B24

$$(\text{H}_2\text{O}) = 0.05553$$

by B25

$$HF_{\text{PROD}} = 3744 \text{ cal/gm.}$$

Ethyl Alcohol eq's B3 & B4

$$\{C\} = 0.04341 \quad \{M\} = 0.13024$$

by B24

$$(\text{CO}_2) = 0.04341 \quad (\text{H}_2\text{O}) = 0.06512$$

by B25

$$HF_{\text{PROD}} = 8473 \text{ cal/gm}$$

Diphenylamine eq's B6 & B7

$$\{C\} = 0.07092 \quad \{M\} = 0.06501$$

by B24

$$(\text{CO}_2) = 0.07092 \quad (\text{H}_2\text{O}) = 0.03250$$

by B25

$$HF_{\text{PROD}} = 8860 \text{ cal/gm}$$

Dibutylphthalate eq's B9 & B10

$$\{C\} = 0.05748 \quad \{M\} = 0.07904$$

by B24

$$(\text{CO}_2) = 0.05748 \quad (\text{H}_2\text{O}) = 0.03952$$

by B25

$$HF_{\text{PROD}} = 8070 \text{ cal/gm}$$

Dinitrotoluene eq's B12 & B13

$$\{C\} = 0.03843 \quad \{M\} = 0.03294$$

by B24

$$(\text{CO}_2) = 0.03843 \quad (\text{H}_2\text{O}) = 0.01647$$

by B25

$$HF_{\text{PROD}} = 4724 \text{ cal/gm}$$

Nitrocellulose (13.15% N) eq.'s B16 & B17

$$\{C\} = 0.02137 \quad \{H\} = 0.02624$$

by B24

$$(CO_2) = 0.02137 \quad (H_2O) = 0.01312$$

by B25

$$HF_{PROD} = 2894 \text{ cal/qm}$$

The heats of complete combustion of the propellant constituents (HC) have been experimentally determined.

For MI Propellant:

Water	HC = 0
Ethyl Alcohol	HC = 7111 cal/qm
Diphenylamine	HC = 9068 cal/qm
Dibutylphthalate	HC = 7389 cal/qm
Dinitrotoluene	HC = 4709 cal/qm
Nitrocellulose (13.15% N)	HC = 2320 cal/qm

The heat of formation of a propellant constituent ( $HF_{CONST}$ ) plus the heat of complete combustion of the propellant constituent (HC) is equal to the heat of formation of the products of complete combustion of the propellant constituent ( $HF_{PROD}$ ). So:

$$HF_{CONST} = HF_{PROD} - HC$$

Thus we have

Water	$HF_{CONST} = 3744 - 0 = 3744$
Ethyl Alcohol	$HF_{CONST} = 8473 - 7111 = 1362$
Diphenylamine	$HF_{CONST} = 8860 - 9068 = -208$
Dibutylphthalate	$HF_{CONST} = 8070 - 7389 = 681$
Dinitrotoluene	$HF_{CONST} = 4724 - 4709 = 15$
Nitrocellulose (13.15% N)	$HF_{CONST} = 2894 - 2320 = 574$

Now, we can determine the heat of formation of MI Propellant ( $HF_{MI}$ ) as:

$$HF_{MI} = (.8313)(574) + (.0978)(15) + (.0489)(681) - (.0098)(208) + (.0073)(1362) + (.0049)(3744)$$

$$HF_{MI} = 539.2 \text{ cal./qm.} \quad (B26)$$

### B3 Gas Pressure Due to the Products of Explosion

Define the equation of state of the gas as:

$$p = \frac{\eta RT}{V} \quad (B27)$$

where  $p$  is the gas pressure (atm),  $V$  the volume per unit mass of the gas ( $\text{cm}^3/\text{g}$ ),  $\eta$  the number of gram molecules in unit mass of gas,  $R$  the gas constant (82.06) and  $T$  the absolute temperature ( $^\circ\text{K}$ ).

For the conservation of atomic types

$$(N_2) = \frac{1}{2} \{N\} \quad (B28)$$

$$(CO) + (CO_2) = \{C\} \quad (B29)$$

$$(H_2) + (H_2O) = \frac{1}{2} \{H\} \quad (B30)$$

$$(CO) + 2(CO_2) + (H_2O) = \{O\} \quad (B31)$$

From which

$$\eta = \{C\} + \frac{1}{2} \{H\} + \frac{1}{2} \{N\} \quad (B32)$$

The equilibrium constant for the water gas reaction is

$$K = \frac{(CO)(H_2O)}{(CO_2)(H_2)} \quad (B33)$$

where  $K$  is given by the table

T °K	K
1000	0.7185
1200	1.406
1400	2.212
1600	3.043
1700	3.438
1800	3.832
1900	4.206
2000	4.574
2100	4.909
2200	5.235
2300	5.533
2400	5.821
2500	6.081
2600	6.331
2700	6.532

TABLE  
B3

T °K	K
2800	6.764
2900	6.949
3000	7.127
3100	7.281
3200	7.428
3300	7.562
3400	7.670
3500	7.764
3600	7.854
3700	7.923
3800	7.989
3900	8.037
4000	8.082

For MI propellant

$$\{C\} = 2535 \cdot 10^{-5} \quad \{H\} = 3103 \cdot 10^{-5} \quad \{N\} = 894 \cdot 10^{-5} \quad \{O\} = 3370 \cdot 10^{-5}$$

$$\gamma = 4534 \cdot 10^{-5}$$

$$\text{Heat of Formation of Propellant} = 538.2 \text{ cal./gm}$$

Define

$HF_{MI}$  = Heat of Formation of Propellant MI

$HF_{GAS}$  = Heat of Formation of Gases

IE = Internal Energy of Gases

HR = Heat of Reaction =  $HF_{GAS} - HF_{MI}$

It is necessary to determine T so that

$$HR - IE = 0$$

For MI Propellant, Assume  $T = 2300$  °K. Now, from Table B3 and equations B28, B29, B30, B31, and B33.

$$[CO_2] = 0.00250 \quad [CO] = 0.00285 \quad [H_2O] = 0.00585$$

$$[H_2] = 0.00966 \quad [N_2] = 0.00447$$

From Table B1

$$HR = (94,030)(0.00250) + (26,690)(0.00285) + (57,500)(0.00585) - 538.2 \\ = 643.1$$

From Table B4

$$IE = [(11.157)(0.00250) + (6.086)(0.00285) + (8.484)(0.00585) \\ + (5.577)(0.00966) + (6.020)(0.00447)] (T - 300) \\ = 594.8$$

So

$$HR - IE = 48.3$$

Hence the temperature (T) is incorrect.

T	[CO <sub>2</sub> ]	[CO]	[H <sub>2</sub> O]	[H <sub>2</sub> ]	[N <sub>2</sub> ]	HR	I.E.	HR-IE
2300	.00250	.00285	.00585	.00966	.00447	643.1	594.8	48.3
2400	.00243	.00282	.00582	.00958	.00447	642.4	629.0	13.4
2500	.00237	.00278	.00578	.00953	.00447	641.8	663.1	-21.3
2600	.00232	.00273	.00573	.00948	.00447	641.3	697.3	-56.0

and by straight line interpolation

$$T = 2440 \text{ } ^\circ\text{K}$$

which is the adiabatic flame temperature.

Also the force constant  $F$  is given by

$$F = \eta RT_0$$

$$F = 307,000 \text{ } \text{st.} \cdot \frac{\text{m}}{\text{m}}$$

TABLE B4

Mean Molecular Heats over the Temperature Range 300 °K to T °K  
 In calories per gram molecule per degree at constant volume

T °K	CO <sub>2</sub>	H <sub>2</sub> O	CO	H <sub>2</sub>	N <sub>2</sub>	OH	NO	O <sub>2</sub>
1000	9.409	6.823	5.403	5.055	5.326	5.136	5.592	5.751
1200	9.824	7.107	5.553	5.115	5.468	5.212	5.744	5.907
1400	10.165	7.388	5.684	5.189	5.597	5.298	5.869	6.037
1600	10.449	7.661	5.799	5.272	5.712	5.390	5.975	6.147
1700	10.577	7.797	5.852	5.318	5.766	5.439	6.024	6.198
1800	10.690	7.918	5.899	5.359	5.814	5.482	6.067	6.244
1900	10.798	8.044	5.945	5.405	5.861	5.529	6.110	6.290
2000	10.896	8.157	5.986	5.447	5.904	5.572	6.149	6.331
2100	10.990	8.273	6.012	5.492	5.945	5.617	6.186	6.373
2200	11.075	8.378	6.036	5.533	5.983	5.657	6.220	6.412
2300	11.157	8.484	6.086	5.577	6.020	5.700	6.252	6.452
2400	11.233	8.581	6.131	5.617	6.053	5.739	6.282	6.488
2500	11.306	8.678	6.162	5.659	6.086	5.780	6.310	6.525
2600	11.373	8.768	6.191	5.697	6.116	5.818	6.336	6.559
2700	11.438	8.857	6.218	5.736	6.146	5.856	6.361	6.594
2800	11.498	8.940	6.244	5.773	6.173	5.891	6.384	6.626
2900	11.556	9.022	6.269	5.811	6.199	5.927	6.406	6.659
3000	11.611	9.099	6.293	5.846	6.224	5.960	6.427	6.690
3100	11.664	9.174	6.315	5.882	6.248	5.993	6.448	6.722
3200	11.714	9.245	6.336	5.915	6.270	6.024	6.467	6.752
3300	11.762	9.315	6.357	5.948	6.292	6.053	6.486	6.782
3400	11.808	9.380	6.376	5.980	6.312	6.083	6.504	6.811
3500	11.853	9.444	6.395	6.012	6.332	6.115	6.521	6.840
3600	11.895	9.505	6.412	6.042	6.350	6.143	6.538	6.868
3700	11.936	9.565	6.429	6.073	6.368	6.171	6.554	6.895
3800	11.975	9.621	6.446	6.102	6.385	6.198	6.569	6.921
3900	12.013	9.677	6.462	6.131	6.402	6.225	6.585	6.947
4000	12.050	9.730	6.477	6.158	6.418	6.251	6.600	6.972

For monatomic gases, the mean molecular heats  
 are 2.980 throughout